# BLOW UP OF SOLUTIONS FOR A PETROVSKY TYPE EQUATION WITH LOGARITHMIC NONLINEARITY 

Jorge Ferreira, Nazli Irkil, Erhan Pişkin, Carlos Raposo, and Mohammad Shahrouzi


#### Abstract

This paper aims to investigate the initial boundary value problem of the nonlinear viscoelastic Petrovsky type equation with nonlinear damping and logarithmic source term. We derive the blow-up results by the combination of the perturbation energy method, concavity method, and differential-integral inequality technique.


## 1. Introduction

In this article, we are concerned with the following viscoelastic Petrovsky type equation with logarithmic nonlinearity,

$$
\begin{equation*}
u_{t t}+\Delta^{2} u-\int_{0}^{t} g(t-s) \triangle^{2} u(s) d s+\left|u_{t}\right|^{m-2} u_{t}=|u|^{p-2} u \ln |u| \tag{1.1}
\end{equation*}
$$

where $(x, t) \in \Omega \times \mathbb{R}^{+}$with initial data

$$
\begin{equation*}
u(x, 0)=u_{0}(x), u_{t}(x, 0)=u_{1}(x), x \in \Omega \tag{1.2}
\end{equation*}
$$

and boundary conditions

$$
\begin{equation*}
u(x, t)=\frac{\partial}{\partial \nu} u(x, t)=0, x \in \partial \Omega, t \geq 0 \tag{1.3}
\end{equation*}
$$

being $\Omega$ a bounded domain of $\mathbb{R}^{n}(n \geq 1)$ with smooth boundary $\partial \Omega$, and $v$ the unit outer normal to $\partial \Omega$.

The Petrovsky $[17,18]$ type of equation

$$
u_{t t}+\Delta^{2} u=f\left(x, t, u, u_{t}\right)
$$

originated from the study of beams and plates, and it can also be used in many branches of physics such as optics, geophysics, nuclear physics, and ocean acoustics. Many authors gave big attention to this problem for quite a long time. They made a lot of progress, as reported in $[1-3,6,8-10,13-15,21,22]$ with references therein. Memory damping has its origin in the mathematical

[^0]2020 Mathematics Subject Classification. Primary 35L20, 35L70.
Key words and phrases. Blow up, Petrovsky-type equation, logarithmic nonlinearity.
description of viscoelastic materials. The viscoelastic materials exhibit natural stabilization mechanism, which is due to the special property of these materials, to retain a memory of their past history. For nonlinearly damped viscoelastic equations with logarithmic source effect and time delay, see [16] with references therein.

Alabau-Boussouira et al. [1] discussed the initial-boundary value problem of a linear Petrovsky equation related to a plate model with memory in absence of source term,

$$
\begin{equation*}
u_{t t}+\Delta^{2} u-\int_{0}^{t} g(t-s) \triangle^{2} u(s) d s=0 \tag{1.4}
\end{equation*}
$$

and showed that the solution decays exponentially or polynomially as $t \rightarrow \infty$ if the initial data is sufficient small.

An energy source that acts in opposition to damping can cause destabilization in the model. The Petrovsky type models without memory and source term were discussed by Messaoudi [15],

$$
\begin{equation*}
u_{t t}+\Delta^{2} u+\left|u_{t}\right|^{m-2} u_{t}=|u|^{p-2} u \ln |u| \tag{1.5}
\end{equation*}
$$

where was established an existence result, and proved that the solution exists globally if $m \geq p$, however, if $m<p$ and the initial energy is negative, the solution blows up in finite time. Later, Chen and Zhou [6] proved that the solution of (1.5) blows up with positive initial energy. Moreover, they claimed that the solution blows up in finite time for vanishing initial energy under the condition $m=2$.

Li and Gao [10] considered (1.4) with memory and source term,

$$
\begin{equation*}
u_{t t}+\Delta^{2} u-\int_{0}^{t} g(t-s) \triangle^{2} u(s) d s+\left|u_{t}\right|^{m-2} u_{t}=|u|^{p-2} u \tag{1.6}
\end{equation*}
$$

They showed that the solution of the initial-boundary value problem of the model (1.6) blows up in finite time. In the case $m=2$, they studied the blowup of solution with a different method. In [13], Liu et al. studied the finite time blow-up for solutions of problem (1.6) with arbitrary high initial energy. Also, Liu et al. [14] obtained global nonexistence results of the problem (1.6) with a weak damping term for positive initial energy. Chen and Zhou [6] proved that the solution blows up with positive initial energy.

In recent years, a great deal of mathematical effort has been devoted to the study of nonlinear wave equations with logarithmic nonlinearity. The logarithmic source term has a wide range of applications of physics such as quantum field theory, nuclear physics, geophysics, optics $[4,5]$. The hyperbolic type equations with logarithmic source terms drew mathematicians' attention. The global existence, blow up, local existence, and asymptotic behavior of solutions were studied by many authors, see for instance $[3,5,7,9,11,12,19,20,23]$.

Gharabli et al. [3] considered the Petrovsky type equation with memory term and logarithmic nonlinearity given by

$$
\begin{equation*}
\left|u_{t}\right|^{\rho} u_{t t}+\Delta^{2} u+\Delta^{2} u_{t t}-\int_{0}^{t} g(t-s) \triangle^{2} u(s) d s=k u \ln u \tag{1.7}
\end{equation*}
$$

They established the local existence and general decay rate of the solutions of the problem (1.7). Later, in [2] was treated (1.7) in the absence $\Delta^{2} u_{t t}$ and in the presence of $u$ for $\rho=0$ and then, was proved the existence and general decay result under different suitable conditions.

As far as we know, in works about the Petrovsky type equation with viscoelastic and logarithmic nonlinearity, as seen from the above statements, there are not enough results on the blow-up problem (1.1) with positive initial energy. Hence, we intend to study finite time blow-up of the solution to the problem (1.1).

This work is written as follows. In Section 2, we prepare preliminaries, some notations. In the last section, we establish the blow-up results of the solution.

Throughout the present paper, we denote by $(\cdot, \cdot)$ the inner product of $L^{2}$ space. We use the standard Lebesgue space $L^{2}(\Omega)$ and Sobolev space $H_{0}^{2}(\Omega)$ with their usual scalar products and norms.

## 2. Preliminaries

Firstly, we present the Sobolev's embedding inequality: suppose that $p$ is a constant such that $1 \leq p \leq \frac{2 n}{n-4}$ if $n \geq 5 ; p \geq 1$ if $n \leq 4$, then $H_{0}^{2}(\Omega) \hookrightarrow L^{p}(\Omega)$ continuously, and for $u \in H_{0}^{2}(\Omega)$

$$
\begin{equation*}
\|u\|_{p} \leq C\|\Delta u\|_{2} \tag{2.1}
\end{equation*}
$$

where $C$ denotes the best embedding constant.
With regard to problem (1.1), we assume that the parameter $p, m$ and the kernel function $g$ satisfy the following assumptions.
(A1) $2<p<\infty$, if $n \leq 4 ; 2<p<\frac{2(n-2)}{n-4}$, if $n \geq 5$.
(A2) $2 \leq m<\infty$, if $n \leq 4 ; 2 \leq m \leq \frac{2 n}{n-4}$, if $n \geq 5$.
$(A 3) g:[0,+\infty) \rightarrow[0,+\infty)$ is a nonincreasing and differentiable function satisfying

$$
\begin{equation*}
1-\int_{0}^{\infty} g(s) d s:=l>0 \tag{2.2}
\end{equation*}
$$

In addition to (2.2), $g$ satisfies, the inequalities

$$
\begin{gather*}
g(s) \geq 0, g^{\prime}(s) \leq 0 \\
\int_{0}^{\infty} g(s) d s<\frac{(p(1-b) / 2)-1}{(p(1-b) / 2)-1+(1 / 2 p(1-b))} \tag{2.3}
\end{gather*}
$$

We first state a local existence theorem. Using the Faedo-Galerkin method and the contraction mapping principle, the local existence result can also be proved in a similar way to those of [8].

Theorem 2.1. Assume that (A1) and (A3) hold. Then, for the initial data $u_{0} \in H_{0}^{2}(\Omega), u_{1} \in L^{2}(\Omega)$, there exists $T>0$, such that the problem (1.1) admits a unique local weak solution on $[0 ; T]$.

Remark 2.2. Condition (2.2) is necessary to guarantee the hyperbolicity and well-posedness of system (1.1).

In this part, we will introduce some material necessary in the proof of our main result (Theorem 3.1). For this purpose, we give some lemmas which will be used throughout this work.

Let us begin with defining the following total energy functional

$$
\begin{align*}
E(t)= & \frac{1}{2}\left\|u_{t}\right\|^{2}+\frac{1}{2}\left(1-\int_{0}^{t} g(s) d s\right)\|\Delta u\|^{2}+\frac{1}{2}(g \circ \Delta u)(t)  \tag{2.4}\\
& +\frac{1}{p^{2}}\|u\|_{p}^{p}-\frac{1}{p} \int_{\Omega} u^{p} \ln |u| d x
\end{align*}
$$

where

$$
\begin{equation*}
(g \circ \Delta u)(t)=\int_{0}^{t} g(t-s)\|\Delta u(s)-\Delta u(t)\|^{2} d s \tag{2.5}
\end{equation*}
$$

Lemma 2.3. Suppose that $(A 1)$ and $(A 2)$ hold and let $u(t)$ be the solution of the problem (1.1). Then the energy functional $E(t)$ is decreasing with respect to $t$ and

$$
\begin{equation*}
E^{\prime}(t)=\frac{1}{2}\left(g^{\prime} \circ \Delta u\right)(t)-\frac{1}{2} \int_{0}^{t} g(s) d s\|\Delta u(t)\|^{2}-\left\|u_{t}\right\|_{m}^{m} \leq 0 \tag{2.6}
\end{equation*}
$$

where

$$
\left(g^{\prime} \circ \Delta u\right)(t)=\int_{0}^{t} g^{\prime}(t-s) \int_{\Omega}|\Delta u(s)-\Delta u(t)|^{2} d x d t
$$

Proof. We multiply both sides of (1.1) by $u_{t}$ and then integrating from 0 to $t$ and using Green formula, we have

$$
E^{\prime}(t)=\frac{1}{2}\left(g^{\prime} \circ \Delta u\right)(t)-\frac{1}{2} \int_{0}^{t} g(s) d s\|\Delta u(t)\|^{2}-\left\|u_{t}\right\|_{m}^{m} \leq 0
$$

that is,
(2.7) $E(t)=\int_{0}^{t}\left[\frac{1}{2}\left(g^{\prime} \circ \Delta u\right)(t)-\frac{1}{2} \int_{0}^{t} g(s) d s\|\Delta u(t)\|^{2}-\left\|u_{t}\right\|_{m}^{m}\right]+E(0)$.

Lemma 2.4. Suppose that ( $A 1$ )-(A3) hold. We also define

$$
G(\alpha)=\frac{1}{2} \alpha^{2}-\frac{C_{1}^{p+1}}{p} \alpha^{p+1} \text { with } C / l^{1 / 2}=C_{1}
$$

and

$$
\alpha(t)=\left(l\|\Delta u(t)\|^{2}+(g \circ \Delta u)(t)\right)^{\frac{1}{2}} .
$$

Then

$$
\begin{equation*}
E(t) \geq G(\alpha(t)) \tag{2.8}
\end{equation*}
$$

Proof. We first note that, by (2.4), we obtain

$$
\begin{aligned}
E(t) & \geq \frac{1}{2}\left(1-\int_{0}^{t} g(s) d s\right)\|\Delta u\|^{2}+\frac{1}{2}(g \circ \Delta u)(t)+\frac{1}{p^{2}}\|u\|_{p}^{p}-\frac{1}{p} \int_{\Omega} u^{p} \ln |u| d x \\
& \geq \frac{1}{2} l\|\Delta u\|^{2}+\frac{1}{2}(g \circ \Delta u)(t)-\frac{1}{p}\|u\|_{p+1}^{p+1} \\
& \geq \frac{1}{2} l\|\Delta u\|^{2}+\frac{1}{2}(g \circ \Delta u)(t)-\frac{1}{p} C^{p+1}\|\Delta u\|_{2}^{p+1} \\
& \geq \frac{1}{2}\left(l\|\Delta u\|^{2}+(g \circ \Delta u)(t)\right)-\frac{1}{p} C_{1}^{p+1}\left(l\|\Delta u(t)\|^{2}+(g \circ \Delta u)(t)\right)^{\frac{p+1}{2}} \\
& =\frac{1}{2} \alpha^{2}-\frac{C_{1}^{p+1}}{p} \alpha^{p+1} \\
& =G(\alpha(t)) .
\end{aligned}
$$

Remark 2.5. We can easily see that $G$ is increasing for $0<\alpha<\lambda$, decreasing for $\alpha>\lambda, \lim _{\alpha \rightarrow \infty} G(\alpha) \rightarrow-\infty, G$ has a maximum at $\lambda=\left(\frac{p}{p+1}\right)^{\frac{1}{p-1}} C_{1}^{-\frac{p+1}{p-1}}$ and the maximum value is

$$
\begin{equation*}
G(\lambda)=\left(\frac{1}{2}\left(\frac{p}{p+1}\right)^{\frac{2}{p-1}}-\frac{1}{p}\left(\frac{p}{p+1}\right)^{\frac{p+1}{p-1}}\right) C_{1}^{-\frac{2(p+1)}{p-1}}=E_{1} \tag{2.9}
\end{equation*}
$$

Lemma 2.6. Assume that (A1)-(A3) hold. Suppose that the solution $u$ of problem (1.1) satisfies

$$
\begin{equation*}
E(0)<E_{1}, \alpha(0)=l^{\frac{1}{2}}\left\|\Delta u_{0}\right\|^{2}>\lambda=\left(\frac{p}{p+1}\right)^{\frac{1}{p-1}} C_{1}^{-\frac{p+1}{p-1}} . \tag{2.10}
\end{equation*}
$$

Then there is a constant $\lambda_{1}>\lambda$ such that

$$
\alpha(t)=\left(l\|\Delta u(t)\|^{2}+(g \circ \Delta u)(t)\right)^{\frac{1}{2}} \geq \lambda_{1}
$$

and

$$
\begin{equation*}
\|u(t)\|_{p+1} \geq C_{1} \lambda_{1} . \tag{2.11}
\end{equation*}
$$

Proof. Since $E(0)<E_{1}$ and $G(\alpha)$ is a continuous function, there exist $\lambda_{1}^{\prime}$ and $\lambda_{1}$ with $\lambda_{1}^{\prime}<\lambda<\lambda_{1}$ such that

$$
G\left(\lambda_{1}^{\prime}\right)=G\left(\lambda_{1}\right)=E(0),
$$

which combined with Lemma 2.4 gives

$$
\begin{equation*}
G(\alpha(0))<E(0)=G\left(\lambda_{1}\right) . \tag{2.12}
\end{equation*}
$$

By (2.9)-(2.12) and Remark 2.5, we deduce

$$
\begin{equation*}
\alpha(0) \geq \lambda_{1} . \tag{2.13}
\end{equation*}
$$

Now we prove the first conclusion. If not, then there exists $t_{0}>0$ such that

$$
\alpha^{2}\left(t_{0}\right)=\left(l\left\|\Delta u\left(t_{0}\right)\right\|^{2}+(g \circ \Delta u)\left(t_{0}\right)\right)<\lambda_{1}^{2} .
$$

Case1. If $\lambda_{1}^{\prime}<\alpha\left(t_{0}\right)<\lambda_{1}$, according to Lemma 2.3 and Remark 2.5, we know $G\left(\alpha\left(t_{0}\right)\right)>E(0) \geq E\left(t_{0}\right)$, which contradicts to Lemma 2.4.

Case 2. If $\alpha\left(t_{0}\right)<\lambda_{1}^{\prime}$, then $\alpha\left(t_{0}\right)<\lambda_{1}^{\prime}<\lambda_{1}$. Set

$$
h(t)=\alpha(t)-\frac{\lambda_{1}^{\prime}+\lambda_{1}}{2} .
$$

Clearly, $h(t)$ is a continuous function, $h\left(t_{0}\right)<0$ and $h(0)>0$ by applying (2.13). Hence, there exists $t_{1} \in\left(0, t_{0}\right)$ such that $h\left(t_{1}\right)=0$, that is $\alpha\left(t_{1}\right)=\frac{\lambda_{1}^{\prime}+\lambda_{1}}{2}$ which implies

$$
G\left(\alpha\left(t_{1}\right)\right)>E(0) \geq E\left(t_{1}\right)
$$

This contradicts to Lemma 2.4.
To establish (2.11), by using the definition of energy function, we obtain

$$
E(0)+\frac{1}{p} \int_{\Omega} u^{p} \ln |u| d x \geq \frac{1}{2}\left(1-\int_{0}^{t} g(s) d s\right)\|\Delta u\|^{2}+\frac{1}{2}(g \circ \Delta u)(t) .
$$

Consequently, we obtain

$$
\begin{align*}
\frac{1}{p}\|u(t)\|_{p+1}^{p+1} \geq & \frac{1}{p} \int_{\Omega} u^{p} \ln |u| d x \geq \frac{1}{2}\left(1-\int_{0}^{t} g(s) d s\right)\|\Delta u\|^{2}  \tag{2.14}\\
& +\frac{1}{2}(g \circ \Delta u)(t)-E(0) \\
\geq & \frac{1}{2}\left(l\|\Delta u\|^{2}+(g \circ \Delta u)(t)\right)-E(0) \\
\geq & \frac{1}{2} \lambda_{1}^{2}-G\left(\lambda_{1}\right)=\frac{C_{1}^{p+1}}{p} \lambda_{1}^{p+1}
\end{align*}
$$

The proof is completed.
We give some lemmas which will be used in our proof. For proofs of Lemmas 2.7-2.10, we refer the readers to Kafini and Messaoudi [9].

Lemma 2.7. Suppose that (A1)-(A2) hold. There exists a positive constant depending on $\Omega$ only such that

$$
\begin{equation*}
\left(\int_{\Omega} u^{p} \ln |u| d x\right)^{\frac{s}{p}} \leq C\left[\int_{\Omega} u^{p} \ln |u| d x+\|\Delta u\|_{2}^{2}\right] \tag{2.15}
\end{equation*}
$$

for any $u \in L^{p+1}(\Omega)$ and $2 \leq s \leq p$, provided that $\int_{\Omega} u^{p} \ln |u| d x \geq 0$.
Lemma 2.8. Suppose that (A1)-(A2) hold. There exists a positive constant depending on $\Omega$ only such that

$$
\begin{equation*}
\|u\|_{p}^{p} \leq C\left[\int_{\Omega} u^{p} \ln |u| d x+\|\Delta u\|_{2}^{2}\right] \tag{2.16}
\end{equation*}
$$

for any $u \in L^{p}(\Omega)$, provided that $\int_{\Omega} u^{p} \ln |u| d x \geq 0$.

Thus, the following estimate holds.
Corollary 2.9. Let the assumptions of Lemma 2.7 and $m<p$ hold. Using the fact that $\|u\|_{m}^{m} \leq C\|u\|_{p}^{m} \leq C\left(\|u\|_{p}^{p}\right)^{\frac{m}{p}}$. Then we obtain the following

$$
\begin{equation*}
\|u\|_{m}^{m} \leq C\left[\left(\int_{\Omega} u^{p} \ln |u| d x\right)^{\frac{m}{p}}+\|\Delta u\|^{\frac{2 m}{p}}\right] \tag{2.17}
\end{equation*}
$$

Lemma 2.10. Suppose that (A1)-(A2) hold. The there exists a positive constant depending on $\Omega$ only such that

$$
\begin{equation*}
\|u\|_{p}^{s} \leq C\left[\|u\|_{p}^{p}+\|\Delta u\|_{2}^{2}\right] \tag{2.18}
\end{equation*}
$$

for any $u \in L^{p}(\Omega)$ and $2 \leq s \leq p$.
We define

$$
\begin{equation*}
H(t)=E_{1}-E(t) \tag{2.19}
\end{equation*}
$$

and we will use throughout this paper. As a result of Lemma 2.10 and definition of $H(t)$ and $E(t)$, we have the following estimative:
Lemma 2.11. Let $u$ be a solution of (1.1). Suppose that (A1) and (A2) hold. Then we get

$$
\begin{equation*}
\|u\|_{p}^{s} \leq C\left(-H(t)-\left\|u_{t}\right\|^{2}-(g \circ \Delta u)(t)+\|u\|_{p+1}^{p+1}+\|u\|_{p}^{p}\right) \tag{2.20}
\end{equation*}
$$

for any $u \in L^{p}(\Omega)$ and $2 \leq s \leq p$.
Proof. Using (2.2) and definition of $E(t)$, we obtain

$$
\begin{aligned}
\frac{1}{2} l\|\Delta u\|^{2} & \leq \frac{1}{2}\left(1-\int_{0}^{\infty} g(s) d s\right)\|\Delta u\|^{2} \\
& \leq E(t)-\frac{1}{2}\left(\left\|u_{t}\right\|^{2}+(g \circ \Delta u)(t)+\frac{2}{p^{2}}\|u\|_{p}^{p}\right)+\frac{1}{p} \int_{\Omega} u^{p} \ln |u| d x \\
& \leq E_{1}-H(t)-\frac{1}{2}\left(\left\|u_{t}\right\|^{2}+(g \circ \Delta u)(t)+\frac{2}{p^{2}}\|u\|_{p}^{p}\right)+\frac{1}{p}\|u\|_{p+1}^{p+1}
\end{aligned}
$$

Using Remark 2.5 and Lemma 2.6, we note that

$$
\left(\frac{p+1}{p}\right)^{\frac{p+1}{p-1}}\|u\|_{p+1}^{p+1} \geq C_{1}^{-\frac{2(p+1)}{p-1}}
$$

and

$$
E_{1}=\left(\frac{p}{p+1}\right)^{\frac{p+1}{p-1}}\left[\frac{1}{2} \frac{(p+1)}{p}-\frac{1}{p}\right] .
$$

Consequently, we obtain

$$
\begin{equation*}
E_{1} \leq\left[\frac{1}{2}\left(\frac{p+1}{p}\right)-\frac{1}{p}\right]\|u\|_{p+1}^{p+1} \tag{2.21}
\end{equation*}
$$

Then, a combination of Lemma 2.10 and (2.21) leads to

$$
\begin{aligned}
\|u\|_{p}^{s} & \leq \frac{p+1}{2 p}\|u\|_{p+1}^{p+1}-H(t)-\frac{1}{2}\left\|u_{t}\right\|^{2}-\frac{1}{2}(g \circ \Delta u)(t)-\frac{1}{p^{2}}\|u\|_{p}^{p} \\
& \leq C\left(-H(t)-\left\|u_{t}\right\|^{2}-(g \circ \Delta u)(t)+\|u\|_{p+1}^{p+1}+\|u\|_{p}^{p}\right)
\end{aligned}
$$

Finally we give the desired result.

## 3. Blow up result

In this part, we state and prove a blow up result for problem (1.1) in finite time with $E(0)<E_{1}$.

Theorem 3.1. Assume that $(A 1),(A 2)$ and $m<p$ hold. Assume further that $g$ satisfies (2.2), (2.3). Then any solution of (1.1) with initial data blows up in finite time provided that $E(0)<E_{1}$.
Proof. Using (2.4), Lemma 2.3, (2.19) and $E(0)<E_{1}$, we get

$$
\begin{aligned}
0<E_{1}-E(0) & =H(0) \\
\leq H(t)= & E_{1}-\frac{1}{2}\left\|u_{t}\right\|^{2}-\frac{1}{2}\left(1-\int_{0}^{t} g(s) d s\right)\|\Delta u\|^{2} \\
& -\frac{1}{2}(g \circ \Delta u)(t)-\frac{1}{p^{2}}\|u\|_{p}^{p}+\frac{1}{p} \int_{\Omega} u^{p} \ln |u| d x .
\end{aligned}
$$

By (2.9) and Lemma 2.6, we note that

$$
\begin{align*}
& E_{1}-\frac{1}{2}\left[\left\|u_{t}\right\|^{2}+\left(1-\int_{0}^{t} g(s) d s\right)\|\Delta u\|^{2}+(g \circ \Delta u)(t)+\frac{2}{p^{2}}\|u\|_{p}^{p}\right]  \tag{3.1}\\
< & E_{1}-\frac{1}{2} \lambda^{2}=-\frac{1}{p}\left(\frac{p}{p+1}\right)^{\frac{p+1}{p-1}} C_{1}^{-\frac{2(p+1)}{p-1}}<0
\end{align*}
$$

for $\forall t \geq 0$, which implies

$$
\begin{equation*}
0<H(0) \leq H(t) \leq \frac{1}{p} \int_{\Omega} u^{p} \ln |u| d x \tag{3.2}
\end{equation*}
$$

Let us define the following function

$$
\begin{equation*}
L(t)=H^{1-\beta}(t)+\varepsilon \int_{\Omega} u u_{t} d x \tag{3.3}
\end{equation*}
$$

where $\beta$ is a suitable positive constant which will be determined later and for

$$
\begin{equation*}
\frac{2(p-m)}{(m-1) p^{2}}<\beta<\frac{p-m}{(m-1) p}<1 \tag{3.4}
\end{equation*}
$$

Now, differentiating $L(t)$ with respect to $t$ and using equation (1.1) we have

$$
\begin{aligned}
L^{\prime}(t) & =(1-\beta) H^{-\beta}(t) H^{\prime}(t)+\varepsilon \int_{\Omega}\left|u_{t}\right|^{2} d x+\int_{\Omega} u u_{t t} d x \\
& \geq(1-\beta) H^{-\beta}(t)\left[\left\|u_{t}\right\|_{m}^{m}-\frac{1}{2}\left(g^{\prime} \circ \Delta u\right)(t)+\frac{1}{2} \int_{0}^{t} g(s) d s\|\Delta u(t)\|^{2}\right]+\varepsilon\left\|u_{t}\right\|^{2}
\end{aligned}
$$

$$
\begin{aligned}
& +\varepsilon \int_{\Omega} u\left(-\Delta^{2} u+\int_{0}^{t} g(t-s) \triangle^{2} u(s) d s-\left|u_{t}\right|^{m-2} u_{t}+|u|^{p-2} u \ln |u|\right) d x \\
\geq & (1-\beta) H^{-\beta}(t)\left\|u_{t}\right\|_{m}^{m}+\varepsilon\left\|u_{t}\right\|^{2}-\varepsilon\|\Delta u\|^{2} \\
& +\varepsilon \int_{0}^{t} g(t-s) \int_{\Omega} \Delta u(t) \triangle u(s) d x d s-\varepsilon \int_{\Omega}\left|u_{t}\right|^{m-2} u_{t} u d x+\varepsilon \int_{\Omega} u^{p} \ln |u| d x .
\end{aligned}
$$

That is

$$
\begin{align*}
L^{\prime}(t) \geq & (1-\beta) H^{-\beta}(t)\left\|u_{t}\right\|_{m}^{m}+\varepsilon\left\|u_{t}\right\|^{2}-\varepsilon\|\Delta u\|^{2}  \tag{3.5}\\
& +\varepsilon \int_{0}^{t} g(t-s)\|\Delta u(t)\|^{2} \\
& +\varepsilon \int_{0}^{t} g(t-s) \int_{\Omega} \Delta u(t)[\triangle u(s)-\triangle u(t)] d x d s \\
& -\varepsilon \int_{\Omega}\left|u_{t}\right|^{m-2} u_{t} u d x+\varepsilon \int_{\Omega} u^{p} \ln |u| d x .
\end{align*}
$$

In view of Young's inequality, for any $\delta>0$ we obtain

$$
\begin{align*}
& \left|\int_{0}^{t} g(t-s) \int_{\Omega} \Delta u(t)[\triangle u(s)-\triangle u(t)] d x d s\right|  \tag{3.6}\\
\leq & \delta \int_{0}^{t} g(t-s)\|\triangle u(s)-\triangle u(t)\|^{2} d s+\frac{1}{4 \delta} \int_{0}^{t} g(s) d s\|\Delta u(t)\|^{2} \\
= & \delta(g \circ \Delta u)(t)+\frac{1}{4 \delta} \int_{0}^{t} g(s) d s\|\Delta u(t)\|^{2} .
\end{align*}
$$

Taking (3.6) into (3.5) yields that

$$
\begin{align*}
L^{\prime}(t) \geq & (1-\beta) H^{-\beta}(t)\left\|u_{t}\right\|_{m}^{m}+\varepsilon\left\|u_{t}\right\|^{2}-\varepsilon \int_{\Omega}\left|u_{t}\right|^{m-2} u_{t} u d x  \tag{3.7}\\
& -\varepsilon\left(1-\int_{0}^{t} g(s) d s\right)\|\Delta u(t)\|^{2}+\varepsilon \int_{\Omega} u^{p} \ln |u| d x \\
& -\varepsilon \delta(g \circ \Delta u)(t)-\frac{\varepsilon}{4 \delta} \int_{0}^{t} g(s) d s\|\Delta u(t)\|^{2} .
\end{align*}
$$

We use the definition of $H(t)$ to substitute for $\int_{\Omega} u^{p} \ln |u| d x$ for $0<b<\frac{p-1}{p}$. Therefore (3.7) takes form

$$
\begin{aligned}
L^{\prime}(t) \geq & (1-\beta) H^{-\beta}(t)\left\|u_{t}\right\|_{m}^{m}+\varepsilon\left(1+\frac{p(1-b)}{2}\right)\left\|u_{t}\right\|^{2} \\
& +\varepsilon p(1-b) H(t)-\varepsilon p(1-b) E_{1}+\varepsilon \frac{(1-b)}{p}\|u\|_{p}^{p} \\
& +\varepsilon\left[\left(\frac{p(1-b)}{2}-1\right)-\left(\frac{p(1-b)}{2}-1+\frac{1}{4 \delta}\right) \int_{0}^{\infty} g(s) d s\right]\|\Delta u(t)\|^{2} \\
& +\varepsilon\left(\frac{p(1-b)}{2}-\delta\right)(g \circ \Delta u)(t)+\varepsilon b \int_{\Omega} u^{p} \ln |u| d x-\varepsilon \int_{\Omega}\left|u_{t}\right|^{m-2} u_{t} u d x
\end{aligned}
$$

for a positive number $\delta, 0<\delta<\frac{p(1-b)}{2}$. Recalling (2.3) and $p>2$, the inequality above reduces to

$$
\begin{align*}
L^{\prime}(t) \geq & (1-\beta) H^{-\beta}(t)\left\|u_{t}\right\|_{m}^{m}+\varepsilon\left(1+\frac{p(1-b)}{2}\right)\left\|u_{t}\right\|^{2}  \tag{3.8}\\
& +\varepsilon p(1-b) H(t)-\varepsilon p(1-b) E_{1} \\
& +\varepsilon \frac{(1-b)}{p}\|u\|_{p}^{p}+\varepsilon \mu_{1}\|\Delta u(t)\|^{2}+\varepsilon \mu_{2}(g \circ \Delta u)(t) \\
& +\varepsilon b \int_{\Omega} u^{p} \ln |u| d x-\varepsilon \int_{\Omega}\left|u_{t}\right|^{m-2} u_{t} u d x .
\end{align*}
$$

At this point we choose $0<b<\frac{p-1}{p}$. This implies that

$$
\frac{p(1-b)}{2}-1+\frac{1}{4 \delta}>\frac{p(1-b)}{2}-1+\frac{1}{2 p(1-b)}=\frac{[p(1-b)-1]^{2}}{2 p(1-b)}>0
$$

Thus if we choose $b$ small enough, then we get from (2.3)

$$
\begin{gathered}
\mu_{1}=\left(\frac{p(1-b)}{2}-1\right)-\left(\frac{p(1-b)}{2}-1+\frac{1}{4 \delta}\right) \int_{0}^{\infty} g(s) d s>0 \\
\mu_{2}=\frac{p(1-b)}{2}-\delta>0
\end{gathered}
$$

To estimate the ninth term of (3.8), we again use Young's inequality

$$
A B \leq \frac{\delta^{r}}{r} A^{r}+\frac{\delta^{-q}}{q} A^{q}, A, B \geq 0 \text { for all } \delta>0, \frac{1}{r}+\frac{1}{q}=1
$$

with $r=m$ and $q=\frac{m}{m-1}$. So we get

$$
\int_{\Omega}\left|u_{t}\right|^{m-2} u_{t} u d x \leq \frac{\tau^{m}}{m}\|u\|_{m}^{m}+\frac{m-1}{m} \tau^{-\frac{m}{m-1}}\left\|u_{t}\right\|_{m}^{m}
$$

which yields, by substitution in (3.8),

$$
\begin{align*}
L^{\prime}(t) \geq & {\left[(1-\beta) H^{-\beta}(t)-\varepsilon \frac{m-1}{m} \tau^{-\frac{m}{m-1}}\right]\left\|u_{t}\right\|_{m}^{m}-\varepsilon \frac{\tau^{m}}{m}\|u\|_{m}^{m} }  \tag{3.9}\\
& -\varepsilon p(1-b) E_{1} \\
& +\varepsilon\left(1+\frac{p(1-b)}{2}\right)\left\|u_{t}\right\|^{2}+\varepsilon \mu_{1}\|\Delta u(t)\|^{2}+\varepsilon \mu_{2}(g \circ \Delta u)(t) \\
& +\varepsilon p(1-b) H(t)+\varepsilon \frac{(1-b)}{p}\|u\|_{p}^{p}+\varepsilon b \int_{\Omega} u^{p} \ln |u| d x .
\end{align*}
$$

Of course (3.9) holds even if $\tau$ is time dependent since the integral is taken over the x-variable. Therefore by choosing $\tau$ so that $\tau^{-\frac{m}{m-1}}=M_{1} H^{-\beta}(t)$, for $M_{1}$ to be specified later, and substituting in (3.9), we get

$$
\begin{equation*}
L^{\prime}(t) \geq\left[(1-\beta)-\varepsilon \frac{m-1}{m} M_{1}\right] H^{-\beta}(t)\left\|u_{t}\right\|_{m}^{m} \tag{3.10}
\end{equation*}
$$

$$
\begin{aligned}
& -\varepsilon \frac{\left(M_{1}\right)^{1-m}}{m} H^{\beta(m-1)}(t)\|u\|_{m}^{m}-\varepsilon p(1-b) E_{1} \\
& +\varepsilon\left(1+\frac{p(1-b)}{2}\right)\left\|u_{t}\right\|^{2}+\varepsilon \mu_{1}\|\Delta u(t)\|^{2}+\varepsilon \mu_{2}(g \circ \Delta u)(t) \\
& +\varepsilon p(1-b) H(t)+\varepsilon \frac{(1-b)}{p}\|u\|_{p}^{p}+\varepsilon b \int_{\Omega} u^{p} \ln |u| d x
\end{aligned}
$$

By exploiting (3.2), Corollary 2.9 and Young's inequality, we have
(3.11) $\quad H^{\beta(m-1)}\|u\|_{m}^{m}$

$$
\begin{aligned}
& \leq\left(\frac{1}{p} \int_{\Omega} u^{p} \ln |u| d x\right)^{\beta(m-1)}\|u\|_{m}^{m} \\
& \leq C\left(\int_{\Omega} u^{p} \ln |u| d x\right)^{\beta(m-1)}\left[\left(\int_{\Omega} u^{p} \ln |u| d x\right)^{\frac{m}{p}}+\|\Delta u\|^{\frac{2 m}{p}}\right] \\
& \leq C\left[\left(\int_{\Omega} u^{p} \ln |u| d x\right)^{\beta(m-1)+\frac{m}{p}}+\left(\int_{\Omega} u^{p} \ln |u| d x\right)^{\beta(m-1)}\|\Delta u\|^{\frac{2 m}{p}}\right] \\
& \leq C\left[\left(\int_{\Omega} u^{p} \ln |u| d x\right)^{\beta(m-1)+\frac{m}{p}}+\left(\int_{\Omega} u^{p} \ln |u| d x\right)^{\beta(m-1) \frac{p}{p-m}}+\|\Delta u\|^{2}\right]
\end{aligned}
$$

From (3.4)

$$
2<\beta p(m-1)+m \leq p \text { and } 2<\frac{\beta(m-1) p^{2}}{p-m} \leq p
$$

Making using of Lemma 2.7, (3.11) yields that

$$
\begin{equation*}
H^{\beta(m-1)}\|u\|_{m}^{m} \leq C\left[\int_{\Omega} u^{p} \ln |u| d x+\|\Delta u\|^{2}\right] \tag{3.12}
\end{equation*}
$$

Combining (3.12) and (3.10) we arrive at

$$
\begin{align*}
L^{\prime}(t) \geq & {\left[(1-\beta)-\varepsilon \frac{m-1}{m} M_{1}\right] H^{-\beta}(t)\left\|u_{t}\right\|_{m}^{m} }  \tag{3.13}\\
& +\varepsilon\left(1+\frac{p(1-b)}{2}\right)\left\|u_{t}\right\|^{2} \\
& -\varepsilon p(1-b) E_{1}+\varepsilon \mu_{2}(g \circ \Delta u)(t) \\
& +\varepsilon\left(\mu_{1}-\frac{\left(M_{1}\right)^{1-m}}{m} C\right)\|\Delta u(t)\|^{2}+\varepsilon p(1-b) H(t) \\
& +\varepsilon \frac{(1-b)}{p}\|u\|_{p}^{p}+\varepsilon\left(b-\frac{\left(M_{1}\right)^{1-m}}{m} C\right) \int_{\Omega} u^{p} \ln |u| d x .
\end{align*}
$$

Noting that

$$
\begin{equation*}
H(t) \geq E_{1}-\frac{1}{2}\left\|u_{t}\right\|^{2}-\frac{1}{2}\|\Delta u\|^{2}-\frac{1}{2}(g \circ \Delta u)(t) \tag{3.14}
\end{equation*}
$$

$$
-\frac{1}{p}\|u\|_{p}^{p}+\frac{1}{p} \int_{\Omega} u^{p} \ln |u| d x .
$$

Inserting (3.14) into (3.13) and taking $p=p-2 a+2 a$, with

$$
\frac{p}{2}>a>\max \left\{\frac{p\left(1-2 \mu_{1}\right)}{2 l}, \frac{p\left(1-\mu_{2}\right)}{2}\right\}>1,
$$

$$
\begin{align*}
L^{\prime}(t) \geq & {\left[(1-\beta)-\varepsilon \frac{m-1}{m} M\right] H^{-\beta}(t)\left\|u_{t}\right\|_{m}^{m} }  \tag{3.15}\\
& +\varepsilon(1+a(1-b))\left\|u_{t}\right\|^{2} \\
& -2 \varepsilon(1-b) a\left[\frac{p+1}{2 p}-\frac{1}{p}\right]\|u\|_{p+1}^{p+1} \\
& +\varepsilon\left(\mu_{2}-\frac{(p-2 a)(1-b)}{2}\right)(g \circ \Delta u)(t) \\
& +\varepsilon\left(\mu_{1}-\frac{\left(M_{1}\right)^{1-m}}{m} C-\frac{(p-2 a)(1-b) l}{2}\right)\|\Delta u(t)\|^{2} \\
& +\varepsilon(1-b)(p-2 a) H(t)+\varepsilon \frac{(1-b)}{p}(1-p+2 a)\|u\|_{p}^{p} \\
& +\varepsilon\left(b-\frac{\left(M_{1}\right)^{1-m}}{m} C+\frac{(p-2 a)(1-b)}{p}\right) \int_{\Omega} u^{p} \ln |u| d x .
\end{align*}
$$

Then by $\ln |u| \geq 1$,

$$
\int_{\Omega} u^{p} \ln |u| d x \geq\|u\|_{p}^{p}
$$

and (3.15) becomes

$$
\begin{aligned}
L^{\prime}(t) \geq & {\left[(1-\beta)-\varepsilon \frac{m-1}{m} M\right] H^{-\beta}(t)\left\|u_{t}\right\|_{m}^{m} } \\
& +\varepsilon(1+a(1-b))\left\|u_{t}\right\|^{2} \\
& -2 \varepsilon(1-b) a\left[\frac{p+1}{2 p}-\frac{1}{p}\right]\|u\|_{p+1}^{p+1} \\
& +\varepsilon\left(\mu_{2}-\frac{(p-2 a)(1-b)}{2}\right)(g \circ \Delta u)(t) \\
& +\varepsilon\left(\mu_{1}-\frac{\left(M_{1}\right)^{1-m}}{m} C-\frac{(p-2 a)(1-b) l}{2}\right)\|\Delta u(t)\|^{2} \\
& +\varepsilon(1-b)(p-2 a) H(t) \\
& +\varepsilon\left(b+\frac{1-b}{p}-\frac{\left(M_{1}\right)^{1-m}}{m} C\right)\|u\|_{p}^{p} .
\end{aligned}
$$

At this point, we take $M_{1}$ sufficiently large such that

$$
b+\frac{1-b}{p}-\frac{\left(M_{1}\right)^{1-m}}{m} C>0
$$

and

$$
\mu_{1}-\frac{\left(M_{1}\right)^{1-m}}{m} C-\frac{(p-2 a)(1-b) l}{2}>0 .
$$

Since $H(0)=E_{1}-E(0)>0$, and $M_{1}$ and $b$ are fixed, taking $\varepsilon$ small enough yields

$$
(1-\beta)-\varepsilon \frac{m-1}{m} M \geq 0
$$

and

$$
\begin{equation*}
L(0)=H^{1-\beta}(0)+\varepsilon \int_{\Omega} u_{0} u_{1} d x>0 \tag{3.16}
\end{equation*}
$$

Therefore, (3.15) takes the form

$$
\begin{equation*}
L^{\prime}(t) \geq \lambda\left[H(t)+\left\|u_{t}\right\|^{2}+\|\Delta u\|^{2}+(g \circ \Delta u)(t)+\|u\|_{p+1}^{p+1}+\|u\|_{p}^{p}\right] \tag{3.17}
\end{equation*}
$$ where $\lambda>0$ is the minimum of the coefficients of $H(t),\left\|u_{t}\right\|^{2},(g \circ \Delta u)(t)$, $\|\Delta u\|^{2},\|u\|_{p}^{p}$ and $\int_{\Omega} u^{p} \ln |u| d x$.

Consequently, we obtain

$$
\begin{equation*}
L(t)>L(0), t \geq 0 \tag{3.18}
\end{equation*}
$$

Now we estimate

$$
\left|\int_{\Omega} u u_{t} d x\right| \leq\|u\|\left\|u_{t}\right\| \leq C\|u\|_{p}\left\|u_{t}\right\|
$$

which implies

$$
\left|\int_{\Omega} u u_{t} d x\right|^{\frac{1}{1-\alpha}} \leq\|u\|_{p}^{\frac{1}{1-\alpha}}\left\|u_{t}\right\|^{\frac{1}{1-\alpha}}
$$

Applying Young's inequality we get

$$
\begin{equation*}
\left|\int_{\Omega} u u_{t} d x\right|^{1 /(1-\alpha)} \leq C\left[C\|u\|_{p}^{\frac{\mu}{1-\alpha}}\left\|u_{t}\right\|^{\frac{\kappa}{1-\alpha}}\right] \text { for } \frac{1}{\mu}+\frac{1}{\kappa}=1 \tag{3.19}
\end{equation*}
$$

To be able to use Lemma 2.11, we take $\kappa=2 /(1-\alpha)$, to get

$$
\mu=2(1-\alpha) /(1-2 \alpha)
$$

Therefore (3.19) has the form

$$
\left|\int_{\Omega} u u_{t} d x\right|^{1 /(1-\alpha)} \leq C\left[\left\|u_{t}\right\|^{2}+\|u\|_{p}^{s}\right]
$$

where $s=2 /(1-2 \alpha) \leq p$. By using Lemma 2.11 we get

$$
\left|\int_{\Omega} u u_{t} d x\right|^{1 /(1-\alpha)} \leq C\left[H(t)+\left\|u_{t}\right\|^{2}+\|\Delta u\|^{2}+(g \circ \Delta u)(t)+\|u\|_{p+1}^{p+1}+\|u\|_{p}^{p}\right] .
$$

On the other hand by $(a+b)^{p} \leq 2^{p-1}\left(a^{p}+b^{p}\right)$, we have

$$
\begin{align*}
L(t)^{\frac{1}{1-\alpha}} & =\left[H^{1-\alpha}(t)+\varepsilon \int_{\Omega} u u_{t} d x\right]^{\frac{1}{1-\alpha}}  \tag{3.20}\\
& \leq 2^{1 /(1-\alpha)}\left[H(t)+\left|\int_{\Omega} u u_{t} d x\right|^{\frac{1}{1-\alpha}}\right] \\
& \leq C\left[H(t)+\left\|u_{t}\right\|^{2}+\|\Delta u\|^{2}+(g \circ \Delta u)(t)+\|u\|_{p+1}^{p+1}+\|u\|_{p}^{p}\right]
\end{align*}
$$

By associating (3.17) and (3.20) we arrive at

$$
\begin{equation*}
L^{\prime}(t) \geq \xi L^{\frac{1}{1-\alpha}}(t) \tag{3.21}
\end{equation*}
$$

where $\xi$ is a positive constant.
Integration of (3.21) over $(0, t)$ we reach

$$
L^{\frac{\alpha}{1-\alpha}}(t) \geq \frac{1}{L^{-\frac{\alpha}{1-\alpha}}(0)-\frac{\xi \alpha t}{1-\alpha}}
$$

Acknowledgment. The authors would like to express their gratitude to the anonymous referee for their constructive comments and suggestions that allowed to improve this manuscript.

## References

[1] F. Alabau-Boussouira, P. Cannarsa, and D. Sforza, Decay estimates for second order evolution equations with memory, J. Funct. Anal. 254 (2008), no. 5, 1342-1372. https: //doi.org/10.1016/j.jfa.2007.09.012
[2] M. M. Al-Gharabli, A. Guesmia, and S. A. Messaoudi, Existence and a general decay results for a viscoelastic plate equation with a logarithmic nonlinearity, Commun. Pure Appl. Anal. 18 (2019), no. 1, 159-180. https://doi.org/10.3934/cpaa. 2019009
[3] M. M. Al-Gharabli and S. A. Messaoudi, Existence and a general decay result for a plate equation with nonlinear damping and a logarithmic source term, J. Evol. Equ. 18 (2018), no. 1, 105-125. https://doi.org/10.1007/s00028-017-0392-4
[4] I. Bialynicki-Birula and J. Mycielski, Wave equations with logarithmic nonlinearities, Bull. Acad. Polon. Sci. Sér. Sci. Math. Astronom. Phys. 23 (1975), no. 4, 461-466.
[5] T. Cazenave and A. Haraux, Équations d'évolution avec non linéarité logarithmique, Ann. Fac. Sci. Toulouse Math. (5) 2 (1980), no. 1, 21-51.
[6] W. Chen and Y. Zhou, Global nonexistence for a semilinear Petrovsky equation, Nonlinear Anal. 70 (2009), no. 9, 3203-3208. https://doi.org/10.1016/j.na.2008.04.024
[7] H. Di, Y. Shang, and Z. Song, Initial boundary value problem for a class of strongly damped semilinear wave equations with logarithmic nonlinearity, Nonlinear Anal. Real World Appl. 51 (2020), 102968, 22 pp. https://doi.org/10.1016/j.nonrwa. 2019. 102968
[8] T. G. Ha and S.-H. Park, Blow-up phenomena for a viscoelastic wave equation with strong damping and logarithmic nonlinearity, Adv. Difference Equ. 2020, Paper No. 235, 17 pp. https://doi.org/10.1186/s13662-020-02694-x
[9] M. Kafini and S. Messaoudi, Local existence and blow up of solutions to a logarithmic nonlinear wave equation with delay, Appl. Anal. 99 (2020), no. 3, 530-547. https: //doi.org/10.1080/00036811.2018.1504029
[10] F. Li and Q. Gao, Blow-up of solution for a nonlinear Petrovsky type equation with memory, Appl. Math. Comput. 274 (2016), 383-392. https://doi.org/10.1016/j.amc. 2015.11.018
[11] W. Lian and R. Xu, Global well-posedness of nonlinear wave equation with weak and strong damping terms and logarithmic source term, Adv. Nonlinear Anal. 9 (2020), no. 1, 613-632. https://doi.org/10.1515/anona-2020-0016
[12] G. Liu, The existence, general decay and blow-up for a plate equation with nonlinear damping and a logarithmic source term, Electron. Res. Arch. 28 (2020), no. 1, 263-289. https://doi.org/10.3934/era. 2020016
[13] L. Liu, F. Sun, and Y. Wu, Blow-up of solutions for a nonlinear Petrovsky type equation with initial data at arbitrary high energy level, Bound. Value Probl. 2019 (2019), Paper No. 15, 18 pp. https://doi.org/10.1186/s13661-019-1136-x
[14] L. Liu, F. Sun, and Y. Wu, Finite time blow-up for a nonlinear viscoelastic Petrovsky equation with high initial energy, Partial Differ. Equ. Appl. 1 (2020), no. 5, Paper No. 31, 18 pp. https://doi.org/10.1007/s42985-020-00031-1
[15] S. A. Messaoudi, Global existence and nonexistence in a system of Petrovsky, J. Math. Anal. Appl. 265 (2002), no. 2, 296-308. https://doi.org/10.1006/jmaa.2001.7697
[16] S.-H. Park, Blowup for nonlinearly damped viscoelastic equations with logarithmic source and delay terms, Adv. Difference Equ. 2021 (2021), Paper No. 316, 14 pp. https: //doi.org/10.1186/s13662-021-03469-8
[17] I. G. Petrovsky, Über das Cauchysche Problem für Systeme von partiellen Differentialgleichungen, Mat. sb. (Mosk.) 44 (1937), no. 5, 815-870.
[18] I. G. Petrowsky, Sur l'analyticité des solutions des systèmes d'équations différentielles, Rec. Math. N. S. [Mat. Sbornik] 5(47) (1939), 3-70.
[19] E. Pişkin and N. Irkıl, Blow up of the solution for hyperbolic type equation with logarithmic nonlinearity, Aligarh Bull. Math. 39 (2020), no. 1, 43-53.
[20] E. Pişkin and N. Irkıl, Existence and decay of solutions for a higher-order viscoelastic wave equation with logarithmic nonlinearity, Commun. Fac. Sci. Univ. Ank. Ser. A1. Math. Stat. 70 (2021), no. 1, 300-319.
[21] E. Pişkin and N. Polat, On the decay of solutions for a nonlinear Petrovsky equation, Math. Sci. Letters. 3 (2014), no. 1, 43-47.
[22] F. Tahamtani and M. Shahrouzi, Existence and blow up of solutions to a Petrovsky equation with memory and nonlinear source term, Bound. Value Probl. 2012 (2012), 50, 15 pp. https://doi.org/10.1186/1687-2770-2012-50
[23] Y. Ye, Global solution and blow-up of logarithmic Klein-Gordon equation, Bull. Korean Math. Soc. 57 (2020), no. 2, 281-294. https://doi.org/10.4134/BKMS.b190190

## Jorge Ferreira

Department of Exact Sciences
Federal Fluminense University
Volta Redonda - RJ, 27213-145, Brazil
Email address: jorge_ferreira@uff.br
Nazli Irkil
Department of Mathematics
Dicle University
Diyarbakir 1280, Turkey
Email address: nazliirkil@gmail.com

Erhan Pişkin
Department of Mathematics
Dicle University
Diyarbakir 1280, Turkey
Email address: episkin@dicle.edu.tr
Carlos Raposo
Department of Mathematics
Federal University of Bahia
SAlVador - BA, 40.170-110, Brazil
Email address: carlos.raposo@ufba.br
Mohammad Shahrouzi
Department of Mathematics
Jahrom University
Jahrom 74137-66171, Iran
Email address: mshahrouzi@jahromu.ac.ir


[^0]:    Received November 26, 2021; Revised March 2, 2022; Accepted June 16, 2022.

