# NEHARI MANIFOLD AND MULTIPLICITY RESULTS FOR A CLASS OF FRACTIONAL BOUNDARY VALUE PROBLEMS WITH $p$-LAPLACIAN 

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Abstract. In this work, we investigate the following fractional boundary value problems

$$
\left\{\begin{array}{l}
{ }_{t} D_{T}^{\alpha}\left(\left|{ }_{0} D_{t}^{\alpha}(u(t))\right|^{p-2}{ }_{0} D_{t}^{\alpha} u(t)\right) \\
=\nabla W(t, u(t))+\lambda g(t)|u(t)|^{q-2} u(t), t \in(0, T) \\
u(0)=u(T)=0
\end{array}\right.
$$

where $\nabla W(t, u)$ is the gradient of $W(t, u)$ at $u$ and $W \in C\left([0, T] \times \mathbb{R}^{n}, \mathbb{R}\right)$ is homogeneous of degree $r, \lambda$ is a positive parameter, $g \in C([0, T])$, $1<r<p<q$ and $\frac{1}{p}<\alpha<1$. Using the Fibering map and Nehari manifold, for some positive constant $\lambda_{0}$ such that $0<\lambda<\lambda_{0}$, we prove the existence of at least two non-trivial solutions.

## 1. Introduction

Fractional order models can be found to be more adequate than integer order models in some real world problems as fractional derivatives provide an excellent tool for the description of memory and hereditary properties of various materials and processes. The mathematical modeling of systems and processes in the fields of physics, chemistry, aerodynamics, electro dynamics of complex medium, polymer rheology, etc. involves derivatives of fractional order. As a consequence, the subject of fractional differential equations is gaining more importance and attention. There has been significant development in ordinary and partial differential equations involving both Riemann-Liouville and Caputo fractional derivatives. For details and examples, one can see the monographs $[2,11,15,18,19,21]$ and the papers $[1,6,10,14,29]$.

Recently, equations including both left and right fractional derivatives are discussed. Apart from their possible applications, equations with left and right derivatives is an interesting and new field in fractional differential equations

[^0]theory. In this topic, many results are obtained in dealing with the existence and multiplicity of solutions of nonlinear fractional differential equations by using techniques of nonlinear analysis, such as fixed point theory (including Leray-Schauder nonlinear alternative), topological degree theory (including coincidence degree theory) and comparison method (including upper and lower solutions and monotone iterative method), see $[4,12,30]$ and so on.

It should be noted that critical point theory and variational methods have also turned out to be very effective tools in determining the existence of solutions for integer order differential equations. The idea behind them is trying to find solutions of a given boundary value problem by looking for critical points of a suitable energy functional defined on an appropriate function space. In the last 30 years, the critical point theory has become a wonderful tool in studying the existence of solutions to differential equations with variational structures, we refer the reader to the books due to Mawhin and Willem [17], Rabinowitz [20], Schechter [22] and the references listed therein.

Motivated by the above classical works, in recent paper [13], the authors showed that critical point theory is an effective approach to deal with the existence of solutions for the following fractional boundary value problem

$$
\left\{\begin{array}{l}
{ }_{t} D_{T}^{\alpha}\left({ }_{0} D_{t}^{\alpha} u(t)\right)=\nabla W(t, u(t)), \quad t \in[0, T],  \tag{FBVP}\\
u(0)=u(T),
\end{array}\right.
$$

where $\alpha \in\left(\frac{1}{2}, 1\right),{ }_{t} D_{T}^{\alpha} u$ is the so called Riemann-Liouville fractional derivatives which is given by Definition $2.2, u \in \mathbb{R}^{n}, W \in C^{1}\left([0, T] \times \mathbb{R}^{n}, \mathbb{R}\right)$ and $\nabla W(t, u)$ is the gradient of $W(t, u)$ at $u$, which has been generalized in recent papers [16, 26-28, 31-33].

Note that the $\alpha$-order Riemann-Liouville fractional derivatives at time $t$ is not defined locally, it relies on the total effects of the commonly used integer derivative on the interval $[0, t]$. So it can be used to describe the variation of a system in which the instantaneous change rate depends on the past state, which is called the "memory" effect in a visualized manner [2]. In addition, as indicated in [3,23-25] the fractional theory has been applied to almost all fields of science including viscoelasticity and rheology, medicine and biology.

In this paper we want to contribute with the development of this new area on fractional differential equations theory. More precisely, the purpose of this work is to investigate the following fractional nonlinear Dirichlet problem

$$
\left\{\begin{array}{l}
{ }_{t} D_{T}^{\alpha}\left(\left|{ }_{0} D_{t}^{\alpha}(u(t))\right|^{p-2}{ }_{0} D_{t}^{\alpha} u(t)\right) \\
=\nabla W(t, u(t))+\lambda g(t)|u(t)|^{q-2} u(t), t \in(0, T) \\
u(0)=u(T)=0
\end{array}\right.
$$

where $\nabla W(t, u)$ is the gradient of $W(t, u)$ at $u$ and $W \in C\left([0, T] \times \mathbb{R}^{n}, \mathbb{R}\right)$ is homogeneous of degree $r, \lambda$ is a positive parameter, $g \in C([0, T]), 1<r<p<q$ and $\frac{1}{p}<\alpha<1$. We assume the following hypothesis:
$\left(\mathrm{H}_{1}\right) W:[0, T] \times \mathbb{R}^{n} \longrightarrow \mathbb{R}$ is homogeneous of degree $r$ that is

$$
W(t, s u)=s^{r} W(t, u)(s>0) \text { for all } t \in[0, T], u \in \mathbb{R}^{n}
$$

$\left(\mathrm{H}_{2}\right) W^{ \pm}(t, u)=\max ( \pm W(t, u), 0) \neq 0$ for all $u \neq 0$.
Note that, from $\left(\mathrm{H}_{1}\right), W(t, u)$ leads to the so-called Euler identity

$$
\begin{equation*}
u \nabla W(x, u)=r W(x, u) \tag{1.1}
\end{equation*}
$$

and

$$
\begin{equation*}
|W(x, u)| \leq K|u|^{r} \quad \text { for some constant } K>0 . \tag{1.2}
\end{equation*}
$$

Our main result is the following.
Theorem 1.1. Let $\frac{1}{p}<\alpha<1,1<r<p<q$ and assume that $W(t, u)$ satisfies the conditions $\left(\mathrm{H}_{1}\right)-\left(\mathrm{H}_{2}\right)$. Then there exists $\lambda_{0}>0$ such that for all $\lambda \in\left(0, \lambda_{0}\right),\left(P_{\lambda}\right)$ has at least two nontrivial solutions.
Remark 1.2. Recently, for $\lambda=0$, the authors in [7, 8] studied the existence and multiplicity of solutions for $\left(P_{\lambda}\right)$ (i.e., $\left(P_{0}\right)$ ) when the potential $W(t, u)$ is superquadratic or subquadratic at infinity. In our Theorem 1.1, we focus our attention on the case that the potential is of the form a combination of superquadratic term and subquadratic term. In addition, we do not need any assumption on the sign of the potential. Therefore, the recent related results are generalized and improved significantly.

## 2. Preliminaries

In this section, we give some background theory on the fractional calculus, in particular the Riemann-Liouville operators and results which will used throughout this paper. Let us start by introducing the definition of the fractional integral in the sense of Riemann-Liouville. We refer the reader to $[15,19]$ or other texts on basic fraction calculus.

Definition 2.1. Let $\alpha>0$ and $u$ be a function defined a.e. on $(a, b) \subset \mathbb{R}$ with values in $\mathbb{R}$. The left (resp. right) fractional integral in the sense of RiemannLiouville with inferior limit $a$ (resp. superior limit $b$ ) of order $\alpha$ of $u$ is given by

$$
{ }_{a} I_{t}^{\alpha} u(t)=\frac{1}{\Gamma(\alpha)} \int_{a}^{t}(t-s)^{\alpha-1} u(s) d s, \quad t \in[a, b]
$$

respectively

$$
{ }_{t} I_{b}^{\alpha} u(t)=\frac{1}{\Gamma(\alpha)} \int_{t}^{b}(t-s)^{\alpha-1} u(s) d s, \quad t \in[a, b]
$$

provided the right side is point-wise defined on $[a, b]$, where $\Gamma$ denotes Euler's Gamma function. If $u \in L^{1}(a, b)$, then ${ }_{a} I_{t}^{\alpha} u$ and $I_{b}^{\alpha} u$ are defined a.e. on $(a, b)$.

Now, we define the fractional derivative in the sense of Riemann-Liouville as follows.

Definition 2.2. Let $0<\alpha<1$. Then, the left (resp. right) fractional derivative in the sense of Riemann-Liouville with inferior limit $a$ (resp. superior limit $b$ ) of order $\alpha$ of $u$ is given by

$$
{ }_{a} D_{t}^{\alpha} u(t)=\frac{d}{d t}\left({ }_{a} I_{t}^{1-\alpha} u\right)(t), \forall t \in[a, b],
$$

respectively

$$
{ }_{t} D_{b}^{\alpha} u(t)=\frac{d}{d t}\left(I_{b}^{1-\alpha} u\right)(t), \forall t \in[a, b]
$$

provided that the right-hand side is point-wise defined.
Remark 2.3. From [15], if $u$ is an absolutely continuous function in $[a, b]$, then ${ }_{a} D_{t}^{\alpha} u$ and ${ }_{t} D_{b}^{\alpha} u$ are defined a.e. on $(a, b)$ and satisfy

$$
\begin{equation*}
{ }_{a} D_{t}^{\alpha} u(t)={ }_{a} I_{t}^{1-\alpha} u^{\prime}(t)+\frac{u(a)}{(t-a)^{\alpha} \Gamma(1-\alpha)} \tag{2.1}
\end{equation*}
$$

and

$$
\begin{equation*}
{ }_{t} D_{b}^{\alpha} u(t)={ }_{-} I_{b}^{1-\alpha} u^{\prime}(t)+\frac{u(b)}{(b-t)^{\alpha} \Gamma(1-\alpha)} . \tag{2.2}
\end{equation*}
$$

Moreover, if $u(a)=u(b)=0$, then ${ }_{a} D_{t}^{\alpha} u(t)={ }_{a} I_{t}^{1-\alpha} u^{\prime}(t)$ and ${ }_{t} D_{b}^{\alpha} u(t)=$ ${ }_{-} I_{b}^{1-\alpha} u^{\prime}(t)$. So in this case we have the equality of Riemann-Liouville fractional derivative and Caputo derivative defined by

$$
{ }_{a}^{c} D_{t}^{\alpha} u(t)={ }_{a} I_{t}^{1-\alpha} u^{\prime}(t)
$$

and

$$
{ }_{t}^{c} D_{b}^{\alpha} u(t)=-{ }_{t} I_{b}^{1-\alpha} u^{\prime}(t) .
$$

Consequently, one gets

$$
{ }_{a} D_{t}^{\alpha} u(t)={ }_{a}^{c} D_{t}^{\alpha} u(t)+\frac{u(a)}{(t-a)^{\alpha} \Gamma(1-\alpha)}
$$

and

$$
{ }_{t} D_{b}^{\alpha} u(t)={ }_{t}^{c} D_{b}^{\alpha} u(t)+\frac{u(b)}{(b-t)^{\alpha} \Gamma(1-\alpha)} .
$$

Next, we provide some properties concerning the left fractional operators of Riemann-Liouville. For more details we refer the reader to [5].

Proposition 2.4. For any $\alpha, \beta>0$ and any $u \in L^{1}(a, b)$, the following equality holds

$$
{ }_{a} I_{t}^{\alpha} \circ{ }_{a} I_{t}^{\beta} u={ }_{a} I_{t}^{\alpha+\beta} .
$$

From Proposition 2.4 and the equations (2.1) and (2.2), it is simple to deduce the following results concerning the composition between fractional integral and fractional derivative. That is, for any $0<\alpha<1$, if $u \in L^{1}(a, b)$ we have

$$
{ }_{a} D_{t}^{\alpha} \circ{ }_{a} I_{t}^{\alpha} u=u
$$

and if $u$ is absolutely continuous such that $u(a)=0$, then, one has

$$
{ }_{a} I_{t}^{\alpha} \circ{ }_{a} D_{t}^{\alpha} u=u
$$

Now, we presented an important result on the boundedness of the left fractional integral from $L^{p}(a, b)$ to $L^{p}(a, b)$.

Proposition 2.5. For any $\alpha>0$ and $p \geq 1$, ${ }_{a} I_{t}^{\alpha}$ is linear and continuous from $L^{p}(a, b)$ to $L^{p}(a, b)$. Moreover for all $u \in L^{p}(a, b)$, we have

$$
\left\|_{a} I_{t}^{\alpha} u\right\|_{p} \leq \frac{(b-a)^{\alpha}}{\Gamma(1+\alpha)}\|u\|_{p}
$$

In the same way, we give another classical result on the boundedness of the left fractional integral from $L^{p}(a, b)$ to $C_{a}(a, b)$ which completes Proposition 2.5 in the case $\frac{1}{p}<\alpha<1$, where $C_{a}(a, b):=\left\{u \in C(a, b): \lim _{t \rightarrow a^{+}} u(t)=0\right\}$.

Proposition 2.6. Let $0<\frac{1}{p}<\alpha<1$ and $p^{\prime}=\frac{p}{p-1}$ (the conjugate exponent of $p$ ). Then, for any $u \in L^{p}(a, b),{ }_{a} I_{t}^{\alpha} u$ is Hölder continuous on $(a, b]$ with exponent $\alpha-\frac{1}{p}>0$, moreover, $\lim _{t \rightarrow a^{+}}{ }_{a} I_{t}{ }^{\alpha} u(t)=0$. Consequently, ${ }_{a} I_{t}^{\alpha} u$ can be continuously extended by 0 in $t=a$. Finally, ${ }_{a} I_{t}^{\alpha} u \in C_{a}(a, b)$, and

$$
\begin{equation*}
\left\|_{a} I_{t}^{\alpha} u\right\|_{\infty} \leq \frac{(b-a)^{\alpha-\frac{1}{p}}}{\Gamma(\alpha)\left((a-1) p^{\prime}+1\right)^{\frac{1}{p^{\prime}}}}\|u\|_{p} \tag{2.3}
\end{equation*}
$$

Also, we will need the following formula for integration by parts, see (7) and (8) in [13].

Proposition 2.7. Let $0<\alpha<1$ and $p, q$ are such that

$$
p \geq 1, q \geq 1 \text { and } \frac{1}{p}+\frac{1}{q}<1+\alpha \text { or } p \neq 1, q \neq 1 \text { and } \frac{1}{p}+\frac{1}{q}=1+\alpha
$$

Then, for all $u \in L^{p}(a, b)$ and all $v \in L^{q}(a, b)$, one has

$$
\begin{equation*}
\int_{a}^{b} v(t)_{a} I_{t}^{\alpha} u(t) d t=\int_{a}^{b} u(t)_{a} I_{t}^{\alpha} v(t) d t \tag{2.4}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{a}^{b} u(t)_{a}^{c} D_{t}^{\alpha} v(t) d t=\left.v(t)_{t} I_{b}^{1-\alpha} u(t)\right|_{t=a} ^{t=b}+\int_{a}^{b} v(t)_{a} D_{t}^{\alpha} u(t) d t \tag{2.5}
\end{equation*}
$$

Moreover, if $v(a)=v(b)=0$, then, one gets

$$
\begin{equation*}
\int_{a}^{b} u(t)_{a} D_{t}^{\alpha} v(t) d t=\int_{a}^{b} v(t)_{a}^{c} D_{t}^{\alpha} u(t) d t \tag{2.6}
\end{equation*}
$$

## 3. Variational setting and main result

To show the existence of solutions to $\left(P_{\lambda}\right)$, we will use Nehari manifold and fibering maps theory. For this purpose we introduce some basic notations and results which are used in the proof of our main result.

As we say in Section 1, to deal with the existence of solutions for $\left(P_{\lambda}\right)$ or some simpler case for the nonlinear terms (for example, (FBVP)), the pioneer work is completed on some function space using the variational methods and critical point theory in [14]. Unfortunately, the authors in the recent work [8] pointed out that the fractional derivatives space $E_{0}^{\alpha, p}$ defined in [14] is problematic. Therefore, we choose the fractional Sobolev space $E_{0}^{\alpha, p}$ constructed in [8] (the same symbol as usual) in the sense of weak fractional derivatives.

For this purpose, we recall the definitions of left and right weak fractional derivatives.

Definition 3.1. Let $0<\alpha \leq 1, u, v \in L^{1}([0, T], \mathbb{R})$, if

$$
\int_{0}^{T} \psi(t) v(t) d t=\int_{0}^{T} u(t)\left({ }_{t} D_{T}^{\alpha} \psi\right)(t) d t, \quad \forall \psi \in C_{0}^{\infty}([0, T], \mathbb{R})
$$

then $v$ is called as the left weak fractional derivative, and denoted by ${ }_{0} \dot{D}_{t}^{\alpha} u$; if

$$
\int_{0}^{T} \psi(t) v(t) d t=\int_{0}^{T} u(t)\left({ }_{t} D_{t}^{\alpha} \psi\right)(t) d t, \quad \forall \psi \in C_{0}^{\infty}([0, T], \mathbb{R})
$$

then $v$ is called as the right weak fractional derivative, and denoted by ${ }_{0} \dot{D}_{T}^{\alpha} u$.
Based on Definition 3.1, we can introduce the appropriate function space corresponding to $\left(P_{\lambda}\right)$. In fact, we only use the definition of left weak fractional derivative. For the simplicity and to be consistent with the classical notation, we still denote by ${ }_{0} D_{t}^{\alpha} u$ the left weak fractional derivative of $u$ in what follows. The set of all functions $u \in C^{\infty}([0, T], \mathbb{R})$ with $u(0)=u(T)=0$ is denoted by $C_{0}^{\infty}([0, T], \mathbb{R})$. For $\alpha>0$ we define the weak fractional derivative space $E_{0}^{\alpha, p}$ as the closure of $C_{0}^{\infty}([0, T], \mathbb{R})$ under the norm

$$
\begin{equation*}
\|u\|_{\alpha, p}=\left(\|u\|_{p}^{p}+\left\|_{0} D_{t}^{\alpha} u\right\|_{p}^{p}\right)^{\frac{1}{p}} \tag{3.1}
\end{equation*}
$$

Note that, as pointed out in [8], if we define the space $E_{0}^{\alpha, p}$ as in [14], taking a Cauchy sequence $\left\{u_{n}\right\} \subset C_{0}^{\infty}([0, T], \mathbb{R})$ with respect to the norm $\|\cdot\|_{\alpha, p}$, one has

$$
u_{n} \rightarrow u_{0}, \quad{ }_{0} D_{t}^{\alpha} u_{n} \rightarrow v_{0} \quad \text { in } \quad L^{p}([0, T], \mathbb{R}) .
$$

Unfortunately, ${ }_{0} D_{t}^{\alpha} u_{0}$ may not exist. Even if ${ }_{0} D_{t}^{\alpha} u_{0}$ exists, ${ }_{0} D_{t}^{\alpha} u_{0}$ may not be equal to $v_{0}$. That is, if we define the space as in [14], the space is not complete. Therefore, we could not choose it as our variational space. To fix this gap, the authors in [8] introduced the definition of weak fractional derivative operator (see [8] for the details). In other words, they optimized the completeness of the fractional derivative space. Thus, we can define the variational framework in the space chosen in our paper for the problem $\left(P_{\lambda}\right)$.

Lemma 3.2 ([8, Corollary 3.6 and Remark 3.10]). Let $0<\alpha \leq 1$ and $1<p<$ $\infty$. For all $u \in E_{0}^{\alpha, p}$, we have

$$
\begin{equation*}
\|u\|_{p} \leq \frac{T^{\alpha}}{\Gamma(\alpha+1)}\left\|_{0} D_{t}^{\alpha} u\right\|_{p} \tag{3.2}
\end{equation*}
$$

Moreover, if $\alpha>\frac{1}{p}$ and $\frac{1}{p}+\frac{1}{p^{\prime}}=1$, then

$$
\begin{equation*}
\|u\|_{\infty} \leq \frac{T^{\alpha-\frac{1}{p}}}{\Gamma(\alpha)\left((\alpha-1) p^{\prime}+1\right)^{\frac{1}{p^{\prime}}}}\left\|_{0} D_{t}^{\alpha} u\right\|_{p} \tag{3.3}
\end{equation*}
$$

According to (3.2), we can consider $E_{0}^{\alpha, p}$ with respect to the equivalent norm

$$
\begin{equation*}
\|u\|=\left\|_{0} D_{t}^{\alpha} u\right\|_{p} \tag{3.4}
\end{equation*}
$$

Lemma 3.3 ([8, Theorem 3.11]). Let $0<\alpha \leq 1$, and $1<p<\infty$. Assume that $\alpha>\frac{1}{p}$ and the sequence $\left\{u_{n}\right\} \rightharpoonup u$ weakly in $E_{0}^{\alpha, p}$. Then, $\left\{u_{n}\right\} \rightarrow u$ in $C([0, T], \mathbb{R})$, that is

$$
\left\|u_{n}-u\right\|_{\infty} \rightarrow 0 \quad \text { as } n \rightarrow \infty
$$

We say that $u \in E_{0}^{\alpha, p}$ is a solution to the problem $\left(P_{\lambda}\right)$, if $u$ satisfies the following equality

$$
\begin{array}{r}
\left.\int_{0}^{T}{ }_{0} D_{t}^{\alpha} u(t)\right|^{p-2}\left({ }_{0} D_{t}^{\alpha} u(t),{ }_{0} D_{t}^{\alpha} v(t)\right) d t-\int_{0}^{T}(\nabla W(t, u(t)), v(t)) d t \\
-\lambda \int_{0}^{T} g(t)|u(t)|^{q-2}(u(t), v(t)) d t=0 \quad \text { for any } v \in E_{0}^{\alpha, p}
\end{array}
$$

Therefore, associated to the problem $\left(P_{\lambda}\right)$, we define the functional

$$
\begin{equation*}
J_{\lambda}(u)=\frac{1}{p}\|u\|^{p}-\frac{1}{r} \int_{0}^{T} W(t, u(t)) d t-\frac{\lambda}{q} \int_{0}^{T} g(t)|u|^{q} d t \tag{3.5}
\end{equation*}
$$

We need to show that the following lemma holds.
Lemma 3.4. (i) The functional $J_{\lambda}$ is well defined on $E_{0}^{\alpha, p}$.
(ii) The functional $J_{\lambda}$ is of class $C^{1}\left(E_{0}^{\alpha, p}, \mathbb{R}\right)$ and for all $u, v \in E_{0}^{\alpha, p}$ we have $\left\langle J_{\lambda}^{\prime}(u), v\right\rangle=\left.\int_{0}^{T}{ }_{0} D_{t}^{\alpha} u(t)\right|^{p-2}\left({ }_{0} D_{t}^{\alpha} u(t),{ }_{0} D_{t}^{\alpha} v(t)\right) d t$

$$
\begin{equation*}
-\int_{0}^{T}(\nabla W(t, u(t)), v(t)) d t-\lambda \int_{0}^{T} g(t)|u(t)|^{q-2}(u(t), v(t)) d t \tag{3.6}
\end{equation*}
$$

Proof. (i) From the continuous embedding and the Hölder inequality, we have

$$
\begin{aligned}
J_{\lambda}(u) & \leq \frac{1}{p}\|u\|^{p}+\frac{1}{r} \int_{0}^{T} W(t, u(t)) d t+\frac{\lambda}{q} \int_{0}^{T} g(t)|u|^{q} d t \\
& \leq \frac{1}{p}\|u\|^{p}-\frac{K}{r}\|u\|_{r}^{r}-\frac{\lambda\|g\|_{\infty}}{q}\|u\|_{q}^{q} \\
& \leq \frac{1}{p}\|u\|^{p}+c_{1}\|u\|^{r}+c_{2}\|u\|^{q}
\end{aligned}
$$

which implies that $J_{\lambda}$ is well defined on $E_{0}^{\alpha, p}$.
(ii) Put $G(u)=\left.\left.\frac{1}{p}\right|_{0} D_{t}^{\alpha} u\right|^{p}-\frac{1}{r} W(t, u(t))-\frac{\lambda}{q} g(t)|u|^{q}$. Then, we can easily show that for all $u, v \in E_{0}^{\alpha, p}$ and for almost every $t \in[0, T]$

$$
\begin{aligned}
& \lim _{s \rightarrow 0} \frac{G(u(t)+s v(t))-G(u(t))}{s} \\
= & \left|{ }_{0} D_{t}^{\alpha} u(t)\right|^{p-2}\left({ }_{0} D_{t}^{\alpha} u(t),{ }_{0} D_{t}^{\alpha} v(t)\right)-(\nabla W(t, u(t)), v(t)) \\
& -\lambda g(t)|u(t)|^{q-2}(u(t), v(t))
\end{aligned}
$$

So, from the Lagrange mean value theorem, (1.1) and (1.2), there exists a real number $\theta$ such that $|\theta| \leq|s|$ and

$$
\begin{align*}
& \frac{G(u(t)+s v(t))-G(u(t))}{s}  \tag{3.7}\\
= & \left|{ }_{0} D_{t}^{\alpha}(u(t)+\theta v(t))\right|^{p-2}\left({ }_{0} D_{t}^{\alpha}(u(t)+\theta v(t)),{ }_{0} D_{t}^{\alpha} v(t)\right) \\
& -(\nabla W(t,(u(t)+\theta v(t))), v(t))-\lambda g(t)|(u(t)+\theta v(t))|^{q-2}(u(t)+\theta v(t), v(t)) \\
\leq & \left|{ }_{0} D_{t}^{\alpha}(u(t)+\theta v(t))\right|^{p-1}\left|{ }_{0} D_{t}^{\alpha} v(t)\right| \\
& +r K|u(t)+\theta v(t)|^{r-1}|v(t)|+\lambda|g(t)||(u(t)+\theta v(t))|^{q-1}|v(t)| \\
\leq & \left|{ }_{0} D_{t}^{\alpha} u(t)\right|^{p-1}\left|{ }_{0} D_{t}^{\alpha} v(t)\right|+\left|{ }_{0} D_{t}^{\alpha} v(t)\right|^{p} \\
& +r K|u(t)|^{r-1}|v(t)|+r K|v(t)|^{r}+\lambda|g(t)||u(t)|^{q-1}|v(t)|+\lambda|g(t)||v(t)|^{q-1}
\end{align*}
$$

On the other hand, from the Hölder inequality, we get

$$
\left.\int_{0}^{T}\left|{ }_{0} D_{t}^{\alpha} u(t)\right|^{p-1}\right|_{0} D_{t}^{\alpha} v(t)\left|d t \leq\left\|\left.\left.\right|_{0} D_{t}^{\alpha} u\right|^{p-1}\right\|_{\frac{p}{p-1}}\left\|_{0} D_{t}^{\alpha} v\right\|_{p}\right.
$$

and

$$
\int_{0}^{T}|u(t)|^{\sigma-1}|v(t)| d t \leq\left\||u|^{\sigma-1}\right\|_{\frac{\sigma}{\sigma-1}}\|v\|_{\sigma} \quad \text { for } \quad \sigma=r \text { or } q
$$

Since $g$ is bounded, then, from the above inequalities, one concludes that the expression 3.7 is in $L^{1}([0, T])$. Therefore, by the dominated convergence theorem, we have

$$
\begin{aligned}
& \lim _{s \rightarrow 0} \frac{J_{\lambda}(u+s v)-J_{\lambda}(u)}{s} \\
= & \left.\int_{0}^{T}\right|_{0} D_{t}^{\alpha}\left(\left.u(t)\right|^{p-2}\left({ }_{0} D_{t}^{\alpha} u(t),{ }_{0} D_{t}^{\alpha} v(t)\right) d t\right. \\
& -\int_{0}^{T}(\nabla W(t, u(t)), v(t))-\lambda g(t)|u(t)|^{q-2}(u(t), v(t)) d t
\end{aligned}
$$

That is, $J_{\lambda}$ is Gâteaux differentiable.
In what follows, it is sufficient to prove that the Gâteaux derivative of $J_{\lambda}$ is continuous. It is similar to the one in Avci et al. [3], therefore we omit it.

We deduce from Lemma 3.4 and equation (1.1) that

$$
\left\langle J_{\lambda}^{\prime}(u), u\right\rangle=\int_{0}^{T}\left|{ }_{0} D_{t}^{\alpha} u(t)\right|^{p} d t-\int_{0}^{T} W(t, u(t)) d t-\lambda \int_{0}^{T} g(t)|u(t)|^{q} d t .
$$

It is easy to see that the energy functional $J_{\lambda}$ is not bounded below on the space $E_{0}^{\alpha, p}$, but it is bounded below on a suitable subset of $E_{0}^{\alpha, p}$. In order to investigate the problem $\left(P_{\lambda}\right)$, we define the constraint set

$$
\mathcal{N}_{\lambda}:=\left\{u \in E_{0}^{\alpha, p} \backslash\{0\}:\left\langle J_{\lambda}^{\prime}(u), u\right\rangle=0\right\}
$$

Note that $\mathcal{N}_{\lambda}$ contains every nonzero solution of $\left(P_{\lambda}\right)$, and $u \in \mathcal{N}_{\lambda}$ if and only if

$$
\begin{equation*}
\|u\|^{p}-\int_{0}^{T} W(t, u(t)) d t-\lambda \int_{0}^{T} g(t)|u(t)|^{q} d t=0 \tag{3.8}
\end{equation*}
$$

To obtain the existence of solutions, we split $\mathcal{N}_{\lambda}$ into three parts: corresponding to local minima, local maxima and points of inflection, are measurable sets defined as follows:
$\mathcal{N}_{\lambda}^{+}=\left\{u \in \mathcal{N}_{\lambda}:(p-1)\|u\|^{p}-(r-1) \int_{0}^{T} W(t, u(t)) d t-\lambda(q-1) \int_{0}^{T} g(t)|u|^{q} d t>0\right\}$,
$\mathcal{N}_{\lambda}^{-}=\left\{u \in \mathcal{N}_{\lambda}:(p-1)\|u\|^{p}-(r-1) \int_{0}^{T} W(t, u(t)) d t-\lambda(q-1) \int_{0}^{T} g(t)|u|^{q} d t<0\right\}$,
$\mathcal{N}_{\lambda}^{0}=\left\{u \in \mathcal{N}_{\lambda}:(p-1)\|u\|^{p}-(r-1) \int_{0}^{T} W(t, u(t)) d t-\lambda(q-1) \int_{0}^{T} g(t)|u|^{q} d t=0\right\}$.
Next, we present some important properties of $\mathcal{N}_{\lambda}^{+}, \mathcal{N}_{\lambda}^{-}$and $\mathcal{N}_{\lambda}^{0}$. Let $\bar{p}$ be such that $\frac{1}{p}+\frac{1}{\bar{p}}=1$ and put

$$
\mu_{0}=\frac{(p-r)(\Gamma(\alpha))^{q}((\alpha-1) \bar{p}+1)^{\frac{q}{p}}}{(q-r)\|g\|_{\infty} T^{1+q\left(\alpha-\frac{1}{p}\right)}}\left(\frac{(q-p)(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}}{K(q-r) T^{1+r\left(\alpha-\frac{1}{p}\right)}}\right)^{\frac{q-p}{p-r}}
$$

Then, we have the following crucial result.
Lemma 3.5. If $\lambda \in\left(0, \mu_{0}\right)$, then $\mathcal{N}_{\lambda}^{0}=\emptyset$.
Proof. We proceed by contradiction to prove that $\mathcal{N}_{\lambda}^{0}=\emptyset$ for all $\lambda \in\left(0, \mu_{0}\right)$. Let us suppose that there exists $u_{0} \in \mathcal{N}_{\lambda}^{0}$. Then, from (3.8) we obtain

$$
\begin{equation*}
(p-r)\left\|u_{0}\right\|^{p}-\lambda(q-r) \int_{0}^{T} g(t)\left|u_{0}\right|^{q} d t=0 \tag{3.9}
\end{equation*}
$$

and

$$
\begin{equation*}
(q-p)\left\|u_{0}\right\|^{p}-(q-r) \int_{0}^{T} W\left(t, u_{0}(t)\right) d t=0 \tag{3.10}
\end{equation*}
$$

From (3.3) and (3.9), one has

$$
\begin{equation*}
\left\|u_{0}\right\| \geq\left(\frac{(p-r)(\Gamma(\alpha))^{q}((\alpha-1) \bar{p}+1)^{\frac{q}{p}}}{\lambda(q-r)\|g\|_{\infty} T^{1+q\left(\alpha-\frac{1}{p}\right)}}\right)^{\frac{1}{q-p}} \tag{3.11}
\end{equation*}
$$

On the other hand, from (1.2), (3.3) and (3.10), one has

$$
\begin{equation*}
\left\|u_{0}\right\| \leq\left(\frac{K(q-r) T^{1+r\left(\alpha-\frac{1}{p}\right)}}{(q-p)(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}}\right)^{\frac{1}{p-r}} \tag{3.12}
\end{equation*}
$$

Combining (3.11) and (3.12) we obtain $\lambda \geq \mu_{0}$, which gives a contradiction. This completes the proof of Lemma 3.5.

Lemma 3.6. If $\lambda \in\left(0, \mu_{0}\right)$, then $J_{\lambda}$ is coercive and bounded below on $\mathcal{N}_{\lambda}$.
Proof. Let $u \in \mathcal{N}_{\lambda}$. Then, using (1.2) and (3.3), we obtain

$$
\begin{equation*}
\int_{0}^{T} W(t, u(t)) d t \leq K \int_{0}^{T}|u(t)|^{r} d t \leq \frac{K T^{1+r\left(\alpha-\frac{1}{p}\right)}}{(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{\bar{p}}}}\|u\|^{r} \tag{3.13}
\end{equation*}
$$

Consequently, from (3.8), we obtain

$$
\begin{aligned}
J_{\lambda}(u) & =\frac{q-p}{q p}\|u\|^{p}-\frac{q-r}{r q} \int_{0}^{T} W(t, u(t)) d t \\
& \geq \frac{q-p}{q p}\|u\|^{p}-\frac{K(q-r) T^{1+r\left(\alpha-\frac{1}{p}\right)}}{q r(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}}\|u\|^{r} .
\end{aligned}
$$

Since $r<p<q, J_{\lambda}$ is coercive and bounded below on $\mathcal{N}_{\lambda}$. The proof of Lemma 3.6 is now completed.

Now as we know that the Nehari manifold is closely related to the behavior of the functions $\Phi_{u}:[0, \infty) \rightarrow \mathbb{R}$ defined as

$$
\Phi_{u}(s)=J_{\lambda}(s u)
$$

Such maps are called fibering maps and were introduced by Drabek and Pohozaev in [9]. For $u \in E_{0}^{\alpha, p}$, we define

$$
\Phi_{u}(s)=\frac{s^{p}}{p}\|u\|^{p}-\frac{s^{r}}{r} \int_{0}^{T} W(t, u(t)) d t-\lambda \frac{s^{q}}{q} \int_{0}^{T} g(t)|u(t)|^{q} d t
$$

then, we have

$$
\Phi_{u}^{\prime}(s)=s^{p-1}\|u\|^{p}-s^{r-1} \int_{0}^{T} W(t, u(t)) d t-\lambda s^{q-1} \int_{0}^{T} g(t)|u(t)|^{q} d t
$$

and

$$
\begin{align*}
\Phi_{u}^{\prime \prime}(s)= & (p-1) s^{p-2}\|u\|^{p}-(r-1) s^{r-2} \int_{0}^{T} W(t, u(t)) d t \\
& -\lambda(q-1) s^{q-2} \int_{0}^{T} g(t)|u(t)|^{q} d t \tag{3.14}
\end{align*}
$$

Then, it is easy to see that $s u \in \mathcal{N}_{\lambda}$ if and only if $\Phi_{u}^{\prime}(s)=0$ and in particular, $u \in \mathcal{N}_{\lambda}$ if and only if $\Phi_{u}^{\prime}(1)=0$.

Before studying the behavior of Nehari manifold using fibering maps, we introduce some notations

$$
\begin{aligned}
& \mathcal{W}^{ \pm}=\left\{u \in E_{0}^{\alpha, p} \backslash\{0\}: \int_{0}^{T} W(t, u(t)) d t \gtrless 0\right\}, \\
& \mathcal{W}^{0}=\left\{u \in E_{0}^{\alpha, p} \backslash\{0\}: \int_{0}^{T} W(t, u(t)) d t=0\right\}, \\
& \mathcal{G}^{ \pm}=\left\{u \in E_{0}^{\alpha, p} \backslash\{0\}: \int_{0}^{T} g(t)|u(t)|^{q} d t \gtrless 0\right\},
\end{aligned}
$$

and

$$
\mathcal{G}^{0}=\left\{u \in E_{0}^{\alpha, p} \backslash\{0\}: \int_{0}^{T} g(t)|u(t)|^{q} d t=0\right\}
$$

In what follows, we study the fibering map $\Phi_{u}$ according to the sign of $\int_{0}^{T} g(t)|u(t)|^{q} d t$ and $\int_{0}^{T} \nabla W(t, u(t)) d t$. For this purpose, we define $m_{u}:[0, \infty)$ $\rightarrow \mathbb{R}$ by

$$
\begin{equation*}
m_{u}(s)=s^{p-r}\|u\|^{p}-\lambda s^{q-r} \int_{0}^{T} g(t)|u(t)|^{q} d t \tag{3.15}
\end{equation*}
$$

Then, for $s>0$ we have

$$
\begin{align*}
\Phi_{u}^{\prime}(s) & =s^{p-1}\|u\|^{p}-s^{r-1} \int_{0}^{T} W(t, u(t)) d t-\lambda s^{q-1} \int_{0}^{T} g(t)|u(t)|^{q} d t \\
& =s^{r-1}\left(m_{u}(s)-\int_{0}^{T} W(t, u(t)) d t\right) \tag{3.16}
\end{align*}
$$

which implies that $s u \in \mathcal{N}_{\lambda}$ if and only if $s$ is a solution of the following equation

$$
m_{u}(s)=\int_{0}^{T} W(t, u(t)) d t
$$

Moreover, it is obvious that $m_{u}(0)=0$ and

$$
\begin{equation*}
m_{u}^{\prime}(s)=(p-r) s^{p-r-1}\|u\|^{p}-\lambda(q-r) s^{q-r-1} \int_{0}^{T} g(t)|u(t)|^{q} d t \tag{3.17}
\end{equation*}
$$

Lemma 3.7. If $u \in \mathcal{W}^{0} \cap \mathcal{G}^{0}$, then $\Phi_{u}$ has no critical point.
Proof. In this case $\Phi_{u}(0)=0$ and $\Phi_{u}^{\prime}(s)>0, \forall s>0$ which implies that $\Phi_{u}$ is strictly increasing and hence has no critical point.

Lemma 3.8. If $u \in \mathcal{W}^{0} \cap \mathcal{G}^{+}$, then $\Phi_{u}$ has a unique critical point which corresponds to a global maximum point. Moreover, there exists $s_{0}>0$ such that $s_{0} u \in \mathcal{N}_{\lambda}^{+}$and $J_{\lambda}\left(s_{0} u\right)<0$.

Proof. In this case, there exists a unique $\bar{s} \in(0, \infty)$ such that $m_{u}^{\prime}(\bar{s})=0$. In addition, $m_{u}^{\prime}(s)>0$ for $s \in(0, \bar{s})$ and $m_{u}^{\prime}(s)<0$ for $s \in(\bar{s}, \infty)$. Note that $m_{u}(0)=0$ and $m_{u}(s) \rightarrow-\infty$ as $s \rightarrow \infty$. So, for $u \in \mathcal{W}^{-}$, there exists a unique $s_{0}$ such that $m_{u}\left(s_{0}\right)=\int_{0}^{T} W(t, u(t)) d t$. Consequently, according to (3.16), we have $\Phi_{u}^{\prime}(s)>0$ for $0<s<s_{0}$, and $\Phi_{u}^{\prime}(s)<0$ for $s>s_{0}$. That is, $\Phi_{u}$ is increasing on $\left(0, s_{0}\right)$, decreasing on $\left(s_{0}, \infty\right)$. Therefore, $\Phi_{u}$ has exactly one critical point at $s_{0}$, which is a global maximum point. Thus, by (3.14), $s_{0} u \in \mathcal{N}_{\lambda}^{-}$.
Lemma 3.9. If $u \in \mathcal{W}^{+} \cap \mathcal{G}^{0}$, then $\Phi_{u}$ has a unique critical point which correspond to a global minimum point. Moreover, there exists $s_{1}>0$ such that $s_{1} u \in \mathcal{N}_{\lambda}^{+}$and $J_{\lambda}\left(s_{1} u\right)<0$.
Proof. In this case, it is easy to see that $m_{u}(0)=0$ and $m_{u}^{\prime}(s)>0, \forall s>0$, which implies that $m_{u}$ is strictly increasing. Since $u \in \mathcal{W}^{+}$, there exists a unique $s_{1}>0$ such that $m_{u}\left(s_{1}\right)=\int_{0}^{T} W(t, u(t)) d t$. This implies that $\Phi_{u}$ is decreasing on $\left(0, s_{1}\right)$, increasing on $\left(s_{1}, \infty\right)$ and $\Phi_{u}^{\prime}\left(s_{1}\right)=0$. Thus, $\Phi_{u}$ has exactly one critical point corresponding to global minimum point. Hence $s_{1} u \in \mathcal{N}_{\lambda}^{+}$. Moreover, since $J_{\lambda}(0)=0$, then we have $J_{\lambda}\left(s_{1} u\right)<0$.

Lemma 3.10. If $u \in \mathcal{W}^{+} \cap \mathcal{G}^{+}$, then there exists $\mu_{1}>0$ such that for $\lambda \in$ $\left(0, \mu_{1}\right), \Phi_{u}$ have a positive value and $\Phi_{u}$ has exactly two critical points which correspond to the local minimum and local maximum. Moreover, there exists $s_{2}>0$ such that $s_{2} u \in \mathcal{N}_{\lambda}^{+}$and $J_{\lambda}\left(s_{2} u\right)<0$.
Proof. Let $u \in E_{0}^{\alpha, p}$. As in above, we define

$$
M_{u}(s)=\frac{s^{p}}{p}\|u\|^{p}-\lambda \frac{s^{q}}{q} \int_{0}^{T} g(t)|u(t)|^{q} d t
$$

Then,

$$
M_{u}^{\prime}(s)=s^{p-1}\|u\|^{p}-\lambda s^{q-1} \int_{0}^{T} g(t)|u(t)|^{q} d t
$$

It is clear that $M_{u}$ attains its maximum value at $\widetilde{s}=\left(\frac{\|u\|^{p}}{\lambda \int_{0}^{T} g(t)|u(t)|^{q} d t}\right)^{\frac{1}{q-p}}$. Moreover,

$$
M_{u}(\widetilde{s})=\left(\frac{1}{p}-\frac{1}{q}\right)\left(\frac{\|u\|^{q}}{\lambda \int_{0}^{T} g(t)|u(t)|^{q} d t}\right)^{\frac{p}{q-p}}
$$

and

$$
M_{u}^{\prime \prime}(\widetilde{s})=(p-q) \frac{\|u\|^{\frac{p(q-2)}{q-p}}}{\left(\lambda \int_{0}^{T} g(t)|u(t)|^{q} d t\right)^{\frac{p-2}{q-p}}}<0
$$

From (3.3), we deduce that

$$
\begin{equation*}
M_{u}(\widetilde{s}) \geq \frac{q-p}{q p}\left(\frac{(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}}{\lambda\|g\|_{\infty} T^{1+q\left(\alpha-\alpha \frac{1}{p}\right)}}\right)^{\frac{p}{q-p}}:=\delta \tag{3.18}
\end{equation*}
$$

which is independent of $u$. We now show that there exists $\mu_{1}>0$ such that $\Phi_{u}(\widetilde{s})>0$. Using (1.2) and (3.3), we get

$$
\begin{aligned}
\frac{\widetilde{s}^{r}}{r} \int_{0}^{1} W(t, u(t)) d t & \leq \frac{K T^{1+r\left(\alpha-\frac{1}{p}\right)}\|u\|^{r}}{r(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}} \widetilde{s}^{r} \\
& =\frac{K T^{1+r\left(\alpha-\frac{1}{p}\right)}}{r(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}}\|u\|^{r}\left(\frac{\|u\|^{p}}{\lambda \int_{0}^{T} g(t)|u(t)|^{q} d t}\right)^{\frac{r}{q-p}} \\
& =\frac{K T^{1+r\left(\alpha-\frac{1}{p}\right)}}{r(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}}\left(\frac{\|u\|^{q}}{\lambda \int_{0}^{T} g(t)|u(t)|^{q} d t}\right)^{\frac{r}{q-p}} \\
& =\frac{K T^{1+r\left(\alpha-\frac{1}{p}\right)}}{r(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}}\left(\frac{p q}{q-p}\right)^{\frac{r}{p}}\left(M_{u}(\widetilde{s})\right)^{\frac{r}{p}}
\end{aligned}
$$

Thus

$$
\begin{aligned}
\Phi_{u}(\widetilde{s}) & =M_{u}(\widetilde{s})-\frac{\widetilde{s}^{r}}{r} \int_{0}^{T} W(t, u(t)) d t \\
& \geq M_{u}(\widetilde{s})-\frac{K T^{1+r\left(\alpha-\frac{1}{p}\right)}}{r(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}}\left(\frac{p q}{q-p}\right)^{\frac{r}{p}}\left(M_{u}(\widetilde{s})\right)^{\frac{r}{p}} \\
& \geq \delta^{\frac{r}{p}}\left(\delta^{\frac{p-r}{p}}-\frac{K T^{1+r\left(\alpha-\frac{1}{p}\right)}}{r(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}}\left(\frac{p q}{q-p}\right)^{\frac{r}{p}}\right)>0
\end{aligned}
$$

for $0<\lambda<\mu_{1}$, where $\delta$ is the constant given in (3.18) and
$\mu_{1}=\frac{(q-p)(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}}{q p\|g\|_{\infty} T^{1+q\left(\alpha-\frac{1}{p}\right)}}\left(\frac{r(\Gamma(\alpha))^{r}((\alpha-1) \bar{p}+1)^{\frac{r}{p}}}{K T^{1+r\left(\alpha-\frac{1}{p}\right)}}\left(\frac{q-p}{q p}\right)^{\frac{r}{p}}\right)^{\frac{q-p}{p-r}}$.
The same arguments used in the proof of Lemma 3.8 show that $\Phi_{u}$ has exactly two critical points which correspond to the local minimum and local maximum. Moreover, there exists $s_{2}>0$ such that $s_{2} u \in \mathcal{N}_{\lambda}^{+}$and $J_{\lambda}\left(s_{2} u\right)<0$. The proof of Lemma 3.10 is now completed.

Remark 3.11. In what follows, let us define $\lambda_{0}$ as

$$
\begin{equation*}
\lambda_{0}=\min \left(\mu_{0}, \mu_{1}\right) \tag{3.19}
\end{equation*}
$$

Note that if $0<\lambda<\lambda_{0}$, then all the above Lemmas hold true.
Lemma 3.12. Let $u$ be a local minimizer for $J_{\lambda}$ on subset $\mathcal{N}_{\lambda}^{+}$or $\mathcal{N}_{\lambda}^{-}$of $\mathcal{N}_{\lambda}$ such that $u \notin \mathcal{N}_{\lambda}^{0}$. Then $u$ is a critical point of $J_{\lambda}$.
Proof. Since $u$ is a minimizer for $J_{\lambda}$ under the constraint

$$
I_{\lambda}(u):=\left\langle J_{\lambda}^{\prime}(u), u\right\rangle=0 .
$$

Then, applying the theory of Lagrange multipliers, we get the existence of $\mu \in \mathbb{R}$ such that

$$
J_{\lambda}^{\prime}(u)=\mu I_{\lambda}^{\prime}(u)
$$

So, we get

$$
\left\langle J_{\lambda}^{\prime}(u), u\right\rangle=\mu\left\langle I_{\lambda}^{\prime}(u), u\right\rangle=\mu \Phi_{u}^{\prime \prime}(1)=0
$$

but $u \notin \mathcal{N}_{\lambda}^{0}$ and so $\Phi_{u}^{\prime \prime}(1) \neq 0$. Hence, $\mu=0$, which gives the proof of Lemma 3.12.

## 4. Proof of Theorem 1.1

Throughout this section, we assume that $\frac{1}{2}<\alpha<1$ and $1<r<p<q$. Let $\lambda_{0}$ be the constant given by (3.19). Then the proof of Theorem 1.1 is based on the following two Propositions.
Proposition 4.1. Assume that hypothesis of Theorem 1.1 are satisfied. Then, for all $0<\lambda<\lambda_{0}, J_{\lambda}$ achieves its minimum on $\mathcal{N}_{\lambda}^{+}$.
Proof. Since $J_{\lambda}$ is bounded below on $\mathcal{N}_{\lambda}$ and also on $\mathcal{N}_{\lambda}^{+}$, there exists a minimizing sequence $\left\{u_{k}\right\} \subset \mathcal{N}_{\lambda}^{+}$such that

$$
\lim _{k \rightarrow \infty} J_{\lambda}\left(u_{k}\right)=\inf _{u \in \mathcal{N}_{\lambda}^{+}} J_{\lambda}(u)
$$

As $J_{\lambda}$ is coercive on $\mathcal{N}_{\lambda},\left\{u_{k}\right\}$ is a bounded sequence in $E_{0}^{\alpha, p}$ up to a subsequence, there exists $u_{\lambda} \in E_{0}^{\alpha, p}$ such that

$$
u_{k} \rightharpoonup u_{\lambda} \quad \text { weakly in } E_{0}^{\alpha, p}
$$

Let $u \in E_{0}^{\alpha, p}$ such that $\int_{0}^{T} W(t, u(t)) d t>0$. Then, from Lemma 3.9 and Lemma 3.10, there exists $s_{1}>0$ such that $s_{1} u \in \mathcal{N}_{\lambda}^{+}$and $J_{\lambda}\left(s_{1} u\right)<0$. Hence,

$$
\inf _{u \in \mathcal{N}_{\lambda}^{+}} J_{\lambda}(u)<0
$$

Since $\left\{u_{k}\right\} \subset \mathcal{N}_{\lambda}$, we get

$$
J_{\lambda}\left(u_{k}\right)=\left(\frac{1}{p}-\frac{1}{r}\right)\left\|u_{k}\right\|^{p}-\left(\frac{1}{r}-\frac{1}{q}\right) \int_{0}^{T} W\left(t, u_{k}(t)\right) d t
$$

which yields that

$$
\left(\frac{1}{r}-\frac{1}{q}\right) \int_{0}^{T} W\left(t, u_{k}(t)\right) d t=\left(\frac{1}{p}-\frac{1}{q}\right)\left\|u_{k}\right\|^{p}-J_{\lambda}\left(u_{k}\right)
$$

Let $k$ go to infinity in the above equation, we get

$$
\begin{equation*}
\int_{0}^{T} W\left(t, u_{\lambda}(t)\right) d t>0 \tag{4.1}
\end{equation*}
$$

Now, we claim that $u_{k} \rightarrow u_{\lambda}$ strongly in $E_{0}^{\alpha, p}$. Otherwise, we have

$$
\begin{equation*}
\left\|u_{\lambda}\right\|^{p}<\lim \inf _{k \rightarrow \infty}\left\|u_{k}\right\|^{p} \tag{4.2}
\end{equation*}
$$

Since $\Phi_{u_{\lambda}}^{\prime}\left(s_{1}\right)=0$ it follows from (4.2) that $\Phi_{u_{k}}^{\prime}\left(s_{1}\right)>0$ for sufficiently large $k$. So, we must have $s_{1}>1$. However, $s_{1} u_{\lambda} \in \mathcal{N}_{\lambda}^{+}$and so

$$
J_{\lambda}\left(s_{1} u_{\lambda}\right)<J_{\lambda}\left(u_{\lambda}\right) \leq \lim _{k \rightarrow \infty} J_{\lambda}\left(u_{k}\right)=\inf _{u \in \mathcal{N}_{\lambda}^{+}} J_{\lambda}(u),
$$

which gives a contradiction. Thus,

$$
u_{k} \rightarrow u_{\lambda} \text { strongly in } E_{0}^{\alpha, p}
$$

which implies that $u_{\lambda} \in \mathbb{N}_{\lambda}=\mathcal{N}_{\lambda}^{+} \cup \mathcal{N}_{\lambda}^{0}$. In addition, it is easy to check by contradiction that $u_{\lambda} \in \mathcal{N}_{\lambda}^{+}$. Consequently, from (4.1), $u_{\lambda}$ is a nontrivial solution of $\left(P_{\lambda}\right)$.
Proposition 4.2. Assume that hypothesis of Theorem 1.1 are satisfied. Then, for all $0<\lambda<\lambda_{0}$, $J_{\lambda}$ achieves its minimum on $\mathcal{N}_{\lambda}^{-}$.
Proof. Let $u \in \mathcal{N}_{\lambda}^{-}$. Therefore, using the result in Lemma 3.10, we have the existence of $\mu_{1}>0$ such that $J_{\lambda}(u) \geq \mu_{1}$. So, there exists a minimizing sequence $\left\{v_{k}\right\} \subset \mathcal{N}_{\lambda}^{-}$such that

$$
\begin{equation*}
\lim _{k \rightarrow \infty} J_{\lambda}\left(v_{k}\right)=\inf _{u \in \mathcal{N}_{\lambda}^{-}} J_{\lambda}(u)>0 \tag{4.3}
\end{equation*}
$$

Moreover, since $J_{\lambda}$ is coercive, $\left\{v_{k}\right\}$ is a bounded sequence in $E_{0}^{\alpha, p}$ up to a sub-sequence, there exists $v_{\lambda} \in E_{0}^{\alpha, p}$ such that

$$
v_{k} \rightharpoonup v_{\lambda} \quad \text { weakly in } E_{0}^{\alpha, p} .
$$

Since $v_{k} \in \mathcal{N}_{\lambda}$, then we have

$$
\begin{equation*}
J_{\lambda}\left(v_{k}\right)+\left(\frac{1}{q}-\frac{1}{p}\right)\left\|v_{k}\right\|^{p}+\lambda\left(\frac{1}{q}-\frac{1}{r}\right) \int_{0}^{1} g(t)\left|v_{k}(t)\right|^{q} d t \tag{4.4}
\end{equation*}
$$

Let $k$ go to infinity in (4.4), it follows from (4.3) that

$$
\begin{equation*}
\int_{0}^{1} g(t)\left|v_{\lambda}(t)\right|^{q} d t>0 \tag{4.5}
\end{equation*}
$$

Hence, $v_{\lambda} \in \mathcal{G}^{+}$and so $\Phi_{v_{\lambda}}$ has a global maximum at some point $\widetilde{s}$. Consequently, $\widetilde{s} v_{\lambda} \in \mathcal{N}_{\lambda}^{-}$. On the other hand, $v_{k} \in \mathcal{N}_{\lambda}^{-}$implies that 1 is a global maximum point for $\Phi_{v_{k}}$, i.e.,

$$
\begin{equation*}
J_{\lambda}\left(\widetilde{s} v_{k}\right)=\Phi_{v_{k}}(\widetilde{s}) \leq \Phi_{v_{k}}(1)=J_{\lambda}\left(v_{k}\right) \tag{4.6}
\end{equation*}
$$

Now, as in the step 1, we claim that $v_{k} \rightarrow v_{\lambda}$. Assume it is not true, then

$$
\left\|v_{\lambda}\right\|^{p}<\lim \inf _{k \rightarrow \infty}\left\|v_{k}\right\|^{p},
$$

it follows from (4.6) that

$$
\begin{aligned}
J_{\lambda}\left(\widetilde{s} v_{\lambda}\right) & =\frac{\widetilde{s}^{p}}{p}\left\|v_{\lambda}\right\|^{p}-\frac{\widetilde{s}^{r}}{r} \int_{0}^{1} W\left(t, v_{\lambda}(t)\right) d t-\lambda \frac{\widetilde{s}^{q}}{q} \int_{0}^{1} g(t)\left|v_{\lambda}(t)\right|^{q} d t \\
& <\inf _{k \rightarrow \infty}\left(\frac{\widetilde{s}^{p}}{p}\left\|v_{k}\right\|^{p}-\frac{\widetilde{s}^{r}}{r} \int_{0}^{1} W\left(t, v_{k}(t)\right) d t-\lambda \frac{\widetilde{s}^{q}}{q} \int_{0}^{1} g(t)\left|v_{k}(t)\right|^{q} d t\right) \\
& \leq \lim _{k \rightarrow \infty} J_{\lambda}\left(\widetilde{s} v_{k}\right) \leq \lim _{k \rightarrow \infty} J_{\lambda}\left(v_{k}\right)=\inf _{u \in \mathcal{N}_{\lambda}^{-}} J_{\lambda}(u),
\end{aligned}
$$

which gives a contradiction. Hence, $v_{k} \rightarrow v_{\lambda}$ and so $v_{\lambda} \in \mathcal{N}_{\lambda}^{-} \cup \mathcal{N}_{\lambda}^{0}$, since $\mathcal{N}_{\lambda}^{0}=\emptyset$, then, $v_{\lambda}$ is a minimizer for $J_{\lambda}$ on $\mathcal{N}_{\lambda}^{-}$. On the other hand, from (4.5),
$v_{\lambda}$ is a nontrivial solution of problem $\left(P_{\lambda}\right)$. Finally, since $\mathcal{N}_{\lambda}^{-} \cap \mathcal{N}_{\lambda}^{+}=\emptyset, u_{\lambda}$ and $v_{\lambda}$ are distinct. That is the result of Theorem 1.1 holds true.

## References

[1] R. P. Agarwal, M. Benchohra, and S. Hamani, Boundary value problems for fractional differential equations, Georgian Math. J. 16 (2009), no. 3, 401-411.
[2] O. P. Agrawal, J. A. Tenreiro Machado, and J. Sabatier, Fractional Derivatives and Their Application: Nonlinear Dynamics, Springer-Verlag, Berlin, 2004,
[3] R. Almeida, Analysis of a fractional SEIR model with treatment, Appl. Math. Lett. 84 (2018), 56-62. https://doi.org/10.1016/j.aml.2018.04.015
[4] Z. Bai and H. Lü, Positive solutions for boundary value problem of nonlinear fractional differential equation, J. Math. Anal. Appl. 311 (2005), no. 2, 495-505. https://doi. org/10.1016/j.jmaa.2005.02.052
[5] L. Bourdin, Existence of a weak solution for fractional Euler-Lagrange equations, J. Math. Anal. Appl. 399 (2013), no. 1, 239-251. https://doi.org/10.1016/j.jmaa. 2012. 10.008
[6] M. Chamekh, A. Ghanmi, and S. Horrigue, Iterative approximation of positive solutions for fractional boundary value problem on the half-line, Filomat 32 (2018), no. 18, 61776187.
[7] T. Chen and W. Liu, Solvability of fractional boundary value problem with p-Laplacian via critical point theory, Bound. Value Probl. 2016 (2016), Paper No. 75, 12 pp. https: //doi.org/10.1186/s13661-016-0583-x
[8] T. Chen, W. Liu, and H. Jin, Infinitely many weak solutions for fractional dirichlet problem with p-Laplacian, arXiv:1605.09238.
[9] P. Drábek and S. I. Pohozaev, Positive solutions for the p-Laplacian: application of the fibering method, Proc. Roy. Soc. Edinburgh Sect. A 127 (1997), no. 4, 703-726. https://doi.org/10.1017/S0308210500023787
[10] A. Ghanmi, M. Kratou, and K. Saoudi, A multiplicity results for a singular problem involving a Riemann-Liouville fractional derivative, Filomat 32 (2018), no. 2, 653-669. https://doi.org/10.2298/fil1802653g
[11] R. Hilfer, Applications of Fractional Calculus in Physics, World Scientific Publishing Co., Inc., River Edge, NJ, 2000. https://doi.org/10.1142/9789812817747
[12] W. Jiang, The existence of solutions to boundary value problems of fractional differential equations at resonance, Nonlinear Anal. 74 (2011), no. 5, 1987-1994. https://doi.org/ 10.1016/j.na.2010.11.005
[13] F. Jiao and Y. Zhou, Existence results for fractional boundary value problem via critical point theory, Internat. J. Bifur. Chaos Appl. Sci. Engrg. 22 (2012), no. 4, 1250086, 17 pp. https://doi.org/10.1142/S0218127412500861
[14] , Existence of solutions for a class of fractional boundary value problems via critical point theory, Comput. Math. Appl. 62 (2011), no. 3, 1181-1199. https://doi. org/10.1016/j.camwa.2011.03.086
[15] A. A. Kilbas, H. M. Srivastava, and J. J. Trujillo, Theory and Applications of Fractional Differential Equations, North-Holland Mathematics Studies, 204, Elsevier Science B.V., Amsterdam, 2006.
[16] W. Liu, M. Wang, and T. Shen, Analysis of a class of nonlinear fractional differential models generated by impulsive effects, Bound. Value Probl. 2017 (2017), Paper No. 175, 18 pp. https://doi.org/10.1186/s13661-017-0909-3
[17] J. Mawhin and M. Willem, Critical Point Theory and Hamiltonian Systems, Applied Mathematical Sciences, 74, Springer-Verlag, New York, 1989. https://doi.org/10. 1007/978-1-4757-2061-7
[18] K. S. Miller and B. Ross, An Introduction to the Fractional Calculus and Fractional Differential Equations, A Wiley-Interscience Publication, John Wiley \& Sons, Inc., New York, 1993.
[19] I. Podlubny, Fractional Differential Equations, Mathematics in Science and Engineering, 198, Academic Press, Inc., San Diego, CA, 1999.
[20] P. H. Rabinowitz, Minimax methods in critical point theory with applications to differential equations, CBMS Regional Conference Series in Mathematics, 65, Published for the Conference Board of the Mathematical Sciences, Washington, DC, 1986. https: //doi.org/10.1090/cbms/065
[21] S. G. Samko, A. A. Kilbas, and O. I. Marichev, Fractional Integrals and Derivatives, translated from the 1987 Russian original, Gordon and Breach Science Publishers, Yverdon, 1993.
[22] M. Schechter, Linking Methods in Critical Point Theory, Birkhäuser Boston, Inc., Boston, MA, 1999. https://doi.org/10.1007/978-1-4612-1596-7
[23] J. Vanterler da Costa Sousa, E. Capelas de Oliveira, and L. A. Magna, Fractional calculus and the ESR test, AIMS Mathematics 2 (2017), 692-705. https://doi.org/ 10.3934/Math. 2017.4.692
[24] J. Vanterler da C. Sousa, Magun N. N. dos Santos, L. A. Magna, and E. Capelas de Oliveira, Validation of a fractional model for erythrocyte sedimentation rate, Comput. Appl. Math. 37 (2018), no. 5, 6903-6919. https://doi.org/10.1007/s40314-018-07170
[25] D. Tavares, R. Almeida, and D. F. M. Torres, Combined fractional variational problems of variable order and some computational aspects, J. Comput. Appl. Math. 339 (2018), 374-388. https://doi.org/10.1016/j.cam.2017.04.042
[26] C. Torres Ledesma, Mountain pass solution for a fractional boundary value problem, J. Fract. Calc. Appl. 5 (2014), no. 1, 1-10.
[27] , Boundary value problem with fractional p-Laplacian operator, Adv. Nonlinear Anal. 5 (2016), no. 2, 133-146. https://doi.org/10.1515/anona-2015-0076
[28] C. E. Torres Ledesma and N. Nyamoradi, Impulsive fractional boundary value problem with p-Laplace operator, J. Appl. Math. Comput. 55 (2017), no. 1-2, 257-278. https: //doi.org/10.1007/s12190-016-1035-6
[29] W. Xie, J. Xiao, and Z. Luo, Existence of solutions for fractional boundary value problem with nonlinear derivative dependence, Abstr. Appl. Anal. 2014 (2014), Art. ID 812910, 8 pp. https://doi.org/10.1155/2014/812910
[30] S. Zhang, Existence of a solution for the fractional differential equation with nonlinear boundary conditions, Comput. Math. Appl. 61 (2011), no. 4, 1202-1208. https://doi. org/10.1016/j.camwa.2010.12.071
[31] , Solutions for a class of fractional boundary value problem with mixed nonlinearities, Bull. Korean Math. Soc. 53 (2016), no. 5, 1585-1596. https://doi.org/10. 4134/BKMS.b150857
[32] Z. Zhang and J. Li, Variational approach to solutions for a class of fractional boundary value problems, Electron. J. Qual. Theory Differ. Equ. 2015 (2015), No. 11, 10 pp. https://doi.org/10.14232/ejqtde.2015.1.11
[33] Y. Zhao and L. Tang, Multiplicity results for impulsive fractional differential equations with p-Laplacian via variational methods, Bound. Value Probl. 2017 (2017), Paper No. 123, 15 pp. https://doi.org/10.1186/s13661-017-0855-0

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