

SECOND COHOMOLOGY OF $\mathfrak{aff}(1)$ ACTING ON n -ARY DIFFERENTIAL OPERATORS

IMED BASDOURI, AMMAR DERBALI, AND SOUMAYA SAIDI

ABSTRACT. We compute the second cohomology of the affine Lie algebra $\mathfrak{aff}(1)$ on the dimensional real space with coefficients in the space $\mathcal{D}_{\underline{\lambda}, \mu}^n$ of n -ary linear differential operators acting on weighted densities where $\underline{\lambda} = (\lambda_1, \dots, \lambda_n)$. We explicitly give 2-cocycles spanning these cohomology.

1. Introduction

In mathematical deformation theory one studies how an object in a certain category of spaces can be varied in dependence on the points of a parameter space. In other words, deformation theory thus deals with the structure of families of objects like varieties, singularities, vector bundles, coherent sheaves, algebras or differentiable maps. Deformation problems appear in various areas of mathematics, in particular in algebra, algebraic and analytic geometry, and mathematical physics. Cohomology is a useful tool in Poisson Geometry, plays an important role in Deformation and Quantization Theory, and attracts more and more interest among algebraists.

In the theory of Lie groups, Lie algebras and their representation theory, a Lie algebra extension is an enlargement of a given Lie algebra \mathfrak{g} by another Lie algebra \mathfrak{h} . Extensions arise in several ways. There is the trivial extension obtained by taking a direct sum of two Lie algebras. Other types are the split extension and the central extension. Extensions may arise naturally, for instance, when forming a Lie algebra from projective group representations. Such a Lie algebra will contain central charges.

Starting with a polynomial loop algebra over finite-dimensional simple Lie algebra and performing two extensions, a central extension and an extension by a derivation, one obtains a Lie algebra which is isomorphic with an untwisted affine Kac-Moody algebra. Using the centrally extended loop algebra one may construct a current algebra in two spacetime dimensions. The Virasoro algebra is the universal central extension of the *Witt* algebra ([2]).

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Central extensions are needed in physics, because the symmetry group of a quantized system usually is a central extension of the classical symmetry group, and in the same way the corresponding symmetry Lie algebra of the quantum system is, in general, a central extension of the classical symmetry algebra (see [3]).

In 1967, Victor Kac and Robert Moody independently generalized the notion of classical Lie algebras, resulting in a new theory of infinite-dimensional Lie algebras, now called Kac-Moody algebras (see [9, 12]). They generalize the finite-dimensional simple Lie algebras and can often concretely be constructed as extensions. Kac-Moody algebras have been conjectured to be a symmetry groups of a unified superstring theory (see [8]). The centrally extended Lie algebras play a dominant role in quantum field theory, particularly in conformal field theory, string theory and in M -theory (see [6]). Central extensions are the simplest extensions of Lie algebras. They appear both in geometry and in physics. Thus, they play an important role in symplectic geometry [10, 13] and in various versions of quantization (see also [11]).

A large portion towards the end is devoted to background material for applications of Lie algebra extensions, both in mathematics and in physics, in areas where they are actually useful. A parenthetical link, (background material), is provided where it might be beneficial.

Lie algebra extensions are most interesting and useful for infinite-dimensional Lie algebras. The theory of group extensions and their interpretation in terms of cohomology is well known, see, e.g., [7]. Let \mathfrak{g} be a Lie algebra and M a \mathfrak{g} -module. The second space $H^2(\mathfrak{g}, M)$ classify the nontrivial extensions of the Lie algebra \mathfrak{g} by the module M :

$$0 \longrightarrow M \longrightarrow \cdot \longrightarrow \mathfrak{g} \longrightarrow 0,$$

the Lie structure on $\mathfrak{g} \oplus M$ being given by

$$[(g_1, \alpha), (g_2, \beta)]_{\mathfrak{g} \oplus M} = ([g_1, g_2], g_1 \cdot \beta - g_2 \cdot \alpha + \Omega(g_1, g_2)),$$

where Ω is a 2-cocycle with values in M .

In this paper we consider a natural class of “non-central” extensions of $\mathfrak{aff}(1)$, namely extensions by the modules $\mathcal{D}_{\lambda, \mu}$ of n -ary linear differential operators acting on weighted densities. The result is quite surprising: there exists a C_{n+k-1}^k extensions if and only if $\mu = \lambda_1 + \dots + \lambda_n + k$, where $k \in \mathbb{N}$.

We consider the one-parameter action of the Lie algebra of vector fields $\text{Vect}(\mathbb{R})$ by the *Lie derivative* on the space $C^\infty(\mathbb{R})$ of smooth functions on \mathbb{R} defined by:

$$(1.1) \quad L_{X \frac{d}{dx}}^\lambda (f) = Xf' + \lambda fX',$$

where $f \in C^\infty(\mathbb{R})$ and $X \frac{d}{dx} \in \text{Vect}(\mathbb{R})$ and where the superscript $'$ stands for $\frac{d}{dx}$. We denote by \mathcal{F}_λ the $\text{Vect}(\mathbb{R})$ -module structure on $C^\infty(\mathbb{R})$ defined by the action (1.1). Geometrically, \mathcal{F}_λ is the space weighted densities of weight λ on

\mathbb{R} :

$$\mathcal{F}_\lambda = \{f dx^\lambda, f \in C^\infty(\mathbb{R})\}.$$

Any differential operator A on \mathbb{R} can be viewed as the linear mapping $f(dx)^\lambda \mapsto (Af)(dx)^\mu$ from \mathcal{F}_λ to \mathcal{F}_μ (λ, μ in \mathbb{R}). Thus the space of differential operators is a $\text{Vect}(\mathbb{R})$ -module, denoted $\mathcal{D}_{\lambda,\mu}^1 := \text{Hom}_{\text{diff}}(\mathcal{F}_\lambda, \mathcal{F}_\mu)$. The $\text{Vect}(\mathbb{R})$ action is:

$$(1.2) \quad L_X^{\lambda,\mu}(A) = L_X^\mu \circ A - A \circ L_X^\lambda.$$

More generally we consider the structures $\text{Vect}(\mathbb{R})$ -modules on the space $\mathcal{D}_{\Delta,\mu}^n$ of n -linear differential operators: $A : \mathcal{F}_{\lambda_1} \otimes \cdots \otimes \mathcal{F}_{\lambda_n} \rightarrow \mathcal{F}_\mu$. The Lie algebra $\text{Vect}(\mathbb{R})$ acting on the space $\mathcal{D}_{\Delta,\mu}^n$ of n -ary Linear differential operators by:

$$(1.3) \quad L_X^{(\lambda_1, \dots, \lambda_n); \mu}(A) = L_X^\mu \circ A - A \circ L_X^{(\lambda_1, \dots, \lambda_n)},$$

where $L_X^{(\lambda_1, \dots, \lambda_n)}$ is the Lie derivative defined by the Leibniz rule:

$$(1.4) \quad \begin{aligned} L_X^{(\lambda_1, \dots, \lambda_n)}(\Phi_1 \otimes \Phi_2 \otimes \cdots \otimes \Phi_n) \\ = L_X^{\lambda_1}(\Phi_1) \otimes \Phi_2 \otimes \cdots \otimes \Phi_n + \cdots + \Phi_1 \otimes \cdots \otimes \Phi_{n-1} \otimes L_X^{\lambda_n}(\Phi_n). \end{aligned}$$

2. Definitions and notations

In this section, we recall the main definitions and facts related to the geometry of the space \mathbb{R} . We also recall some fundamental concepts from cohomology theory (see, e.g., [5, 7]).

2.1. Cohomology theory

Let us first recall some fundamental concepts from cohomology theory (see, e.g., [4, 5, 7, 14]). Let \mathfrak{g} be a Lie algebra acting on a vector space V . The space of n -cochains of \mathfrak{g} with values in V is the \mathfrak{g} -module

$$C^n(\mathfrak{g}, V) := \text{Hom}(\Lambda^n \mathfrak{g}; V).$$

The *coboundary operator* $\delta_n : C^n(\mathfrak{g}, V) \rightarrow C^{n+1}(\mathfrak{g}, V)$ is a \mathfrak{g} -map satisfying $\delta_n \circ \delta_{n-1} = 0$. The kernel of δ_n , denoted $Z^n(\mathfrak{g}, V)$, is the space of *n -cocycles*, among them, the elements in the range of δ_{n-1} are called *n -coboundaries*. We denote $B^n(\mathfrak{g}, V)$ the space of *n -coboundaries*.

By definition, the n^{th} cohomolgy space is the quotient space

$$H^n(\mathfrak{g}, V) = Z^n(\mathfrak{g}, V)/B^n(\mathfrak{g}, V).$$

We will only need the formula of δ_n (which will be simply denoted δ) in degrees 0, 1 and 2: for $\Xi \in C^0(\mathfrak{g}, V) = V$, $\delta\Xi(x) := x \cdot \Xi$, and for $\Lambda \in C^1(\mathfrak{g}, V)$,

$$(2.5) \quad \delta(\Lambda)(x, y) := g \cdot \Lambda(y) - y \cdot \Lambda(x) - \Lambda([x, y]) \quad \text{for any } x, y \in \mathfrak{g},$$

for $\Omega \in C^2(\mathfrak{g}, V)$,

$$(2.6) \quad \delta(\Omega)(x, y, z) := x \cdot \Omega(y, z) - y \cdot \Omega(x, z) + z \cdot \Omega(x, y)$$

$$- \Omega([x, y], z) + \Omega([x, z], y) - \Omega([y, z], x),$$

where $x, y, z \in \mathfrak{g}$.

2.2. Lie algebra $\mathfrak{aff}(1)$

The Lie algebra $\mathfrak{aff}(1)$ is realized as a subalgebra of the Lie algebra $\text{Vect}(\mathbb{R})$ (see [1]):

$$\mathfrak{aff}(1) = \text{Span}(X_1 = \frac{d}{dx}, X_x = x \frac{d}{dx}).$$

The commutation relations are

$$[X_1, X_x] = X_1, \quad [X_x, X_x] = 0, \quad [X_1, X_1] = 0.$$

2.3. The space of tensor densities on \mathbb{R}

The Lie algebra, $\text{Vect}(\mathbb{R})$, of vector fields on \mathbb{R} naturally acts, by the Lie derivative, on the space

$$\mathcal{F}_\lambda = \{f dx^\lambda : f \in C^\infty(\mathbb{R})\},$$

of weighted densities of degree λ . The Lie derivative L_X^λ of the space \mathcal{F}_λ along the vector field $X \frac{d}{dx}$ is defined by

$$(2.7) \quad L_X^\lambda = X \partial_x + \lambda X',$$

where $X, f \in C^\infty(\mathbb{R})$ and $X' := \frac{dX}{dx}$. More precisely, for all $f dx^\lambda \in \mathcal{F}_\lambda$, we have

$$L_X^\lambda(f dx^\lambda) = (X f' + \lambda f X') dx^\lambda.$$

In the paper, we restrict ourselves to the Lie algebra $\mathfrak{aff}(1)$ which is isomorphic to the Lie subalgebra of $\text{Vect}(\mathbb{R})$ spanned by

$$\{X_1, X_x\}.$$

2.4. The space of n -ary linear differential operators as a $\mathfrak{aff}(1)$ -module

The space of n -ary linear differential operators is a $\text{Vect}(\mathbb{R})$ -module, denoted

$$\mathcal{D}_{\underline{\lambda}, \mu}^n := \text{Hom}_{\text{diff}}(\mathcal{F}_{\lambda_1} \otimes \cdots \otimes \mathcal{F}_{\lambda_n}, \mathcal{F}_\mu).$$

The $\text{Vect}(\mathbb{R})$ action is:

$$(2.8) \quad L_X^{\underline{\lambda}, \mu}(A) = L_X^\mu \circ A - A \circ L_X^{\underline{\lambda}},$$

where $\underline{\lambda} = (\lambda_1, \dots, \lambda_n)$ and $L_X^{(\lambda_1, \dots, \lambda_n)}$ is the Lie derivative on $\mathcal{F}_{\lambda_1} \otimes \cdots \otimes \mathcal{F}_{\lambda_n}$ defined by the Leibnitz rule:

$$\begin{aligned} & L_X^{\underline{\lambda}}(f_1 dx^{\lambda_1} \otimes \cdots \otimes f_n dx^{\lambda_n}) \\ &= L_X^{\lambda_1}(f_1) \otimes \cdots \otimes f_n dx^{\lambda_n} + \cdots + f_1 dx^{\lambda_1} \otimes \cdots \otimes L_X^{\lambda_n}(f_n dx^{\lambda_n}). \end{aligned}$$

3. The space $H_{\text{diff}}^2(\mathfrak{aff}(1), \mathcal{D}_{\underline{\lambda}, \mu}^3)$

The first main result of paper is the following.

Theorem 3.1. *The space $H_{\text{diff}}^2(\mathfrak{aff}(1), \mathcal{D}_{\underline{\lambda}, \mu}^3)$ has the following structure:*

$$(3.9) \quad H_{\text{diff}}^2(\mathfrak{aff}(1), \mathcal{D}_{\underline{\lambda}, \mu}^3) \simeq \begin{cases} \mathbb{R}^{C_{k+2}} & \text{if } \mu = \lambda_1 + \lambda_2 + \lambda_3 + k, \\ 0 & \text{otherwise,} \end{cases}$$

where $C_q^p = \frac{q!}{(q-p)!p!}$.

The following 2-cocycles span the corresponding cohomology spaces:

$$\begin{aligned} & \mathcal{C}\left(X \frac{d}{dx}, Y \frac{d}{dx}\right)(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3}) \\ &= \sum_{i+j=0}^{k-1} c_{i,j} (XY' - X'Y) f_1^{(i)} f_2^{(j)} f_3^{(i+j-k)} dx^{\lambda_1 + \lambda_2 + \lambda_3 + k}, \end{aligned}$$

where $c_{i,j}$ are constants, $f_1 dx^{\lambda_1} \in \mathcal{F}_{\lambda_1}$, $f_2 dx^{\lambda_2} \in \mathcal{F}_{\lambda_2}$, $f_3 dx^{\lambda_3} \in \mathcal{F}_{\lambda_3}$ and $X \frac{d}{dx}, Y \frac{d}{dx} \in \mathfrak{aff}(1)$.

We need the following lemma.

Lemma 3.2. *Let $b \in C^1(\mathfrak{aff}(1), \mathcal{D}_{\underline{\lambda}, \lambda_1 + \lambda_2 + \lambda_3 + k}^3)$ be defined as follows: for $X \frac{d}{dx} \in \mathfrak{aff}(1)$, $f_1 dx^{\lambda_1} \in \mathcal{F}_{\lambda_1}$, $f_2 dx^{\lambda_2} \in \mathcal{F}_{\lambda_2}$ and $f_3 dx^{\lambda_3} \in \mathcal{F}_{\lambda_3}$*

$$(3.10) \quad \begin{aligned} b(X)(f_1, f_2, f_3) &= \sum_{i+j+n=k} c_{i,j,n} X f_1^{(i)} f_2^{(j)} f_3^{(n)} dx^\mu \\ &+ \sum_{i+j+n=k-1} \alpha_{i,j,n} X' f_1^{(i)} f_2^{(j)} f_3^{(n)} dx^\mu, \end{aligned}$$

where the coefficients $c_{i,j,n}$ and $\alpha_{i,j,n}$ are constants.

Then the map $\delta b : \mathfrak{aff}(1) \times \mathfrak{aff}(1) \rightarrow \mathcal{D}_{\underline{\lambda}, \lambda_1 + \lambda_2 + \lambda_3 + k}^3$ is given by

$$(3.11) \quad \begin{aligned} & \delta b(X, Y)(f_1, f_2, f_3) \\ &= L_X^{\underline{\lambda}, \lambda_1 + \lambda_2 + \lambda_3 + k} b(Y)(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3}) \\ &\quad - L_Y^{\underline{\lambda}, \lambda_1 + \lambda_2 + \lambda_3 + k} b(X)(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3}) \\ &\quad - b([X, Y])(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3}) \\ &= \sum_{i+j+n=k} c_{i,j,n} (\mu - \lambda_1 - \lambda_2 - \lambda_3 - k) (X'Y - XY') f_1^{(i)} f_2^{(j)} f_3^{(n)}. \end{aligned}$$

3.1. Proof of theorem

Any 2-cocycle $\mathcal{C} \in Z_{\text{diff}}^1(\mathfrak{aff}(1), \mathcal{D}_{\underline{\lambda}, \mu}^3)$ should be retain the following general form:

$$(3.12) \quad \mathcal{C}(X, Y)(f_1, f_2, f_3) = \sum_{i+j+n=k-1} c_{i,j,n} (XY' - X'Y) f_1^{(i)} f_2^{(j)} f_3^{(n)} dx^\mu,$$

where $c_{i,j,n}$ are constants.

The 2-cocycle condition reads as follows: for all $f_1 dx^{\lambda_1} \in \mathcal{F}_{\lambda_1}$, $f_2 dx^{\lambda_2} \in \mathcal{F}_{\lambda_2}$, $f_3 dx^{\lambda_3} \in \mathcal{F}_{\lambda_3}$ and for all $X, Y \in \mathfrak{aff}(1)$, we have

$$(3.13) \quad L_X^{\lambda_1, \lambda_2, \lambda_3, \mu} \mathcal{C}(Y, Z)(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3}) \\ - \mathcal{C}([X, Y], Z) + \circlearrowleft(X, Y, Z) = 0,$$

where $\circlearrowleft(X, Y, Z)$ denotes summation over the cyclic permutation on X, Y, Z .

A direct computation, proves that we have

$$\begin{aligned} \Phi_1 &= L_X^{\lambda_1, \lambda_2, \lambda_3, \mu} \mathcal{C}(Y, Z)(f_1, f_2, f_3) \\ &= L_X^\mu \mathcal{C}(Y, Z)(f_1, f_2, f_3) - \mathcal{C}(Y, Z)(L_X^{\lambda_1}(f_1), f_2, f_3) \\ &\quad - \mathcal{C}(Y, Z)(f_1, L_X^{\lambda_2}(f_2), f_3) - \mathcal{C}(Y, Z)(f_1, f_2, L_X^{\lambda_3}(f_3)) \\ &= \sum_{i+j+n=k-1} c'_{i,j,n} X(YZ' - Y'Z) f_1^{(i)} f_2^{(j)} f_3^{(n)} \\ &\quad + \sum_{i+j+n=k-1} c_{i,j,n} (\mu - \lambda_1 - \lambda_2 - \lambda_3 - k + 1) X'(YZ' - Y'Z) f_1^{(i)} f_2^{(j)} f_3^{(n)}, \end{aligned}$$

$$\begin{aligned} \Phi_2 &= L_Y^{\lambda_1, \lambda_2, \lambda_3, \mu} \mathcal{C}(X, Z)(f_1, f_2, f_3) \\ &= \sum_{i+j+n=k-1} c'_{i,j,n} Y(XZ' - X'Z) f_1^{(i)} f_2^{(j)} f_3^{(n)} \\ &\quad + \sum_{i+j+n=k-1} c_{i,j,n} (\mu - \lambda_1 - \lambda_2 - \lambda_3 - k + 1) Y'(XZ' - X'Z) f_1^{(i)} f_2^{(j)} f_3^{(n)}, \end{aligned}$$

$$\begin{aligned} \Phi_3 &= L_Z^{\lambda_1, \lambda_2, \lambda_3, \mu} \mathcal{C}(X, Y)(f_1, f_2, f_3) \\ &= \sum_{i+j+n=k-1} c'_{i,j,n} Z(XY' - X'Y) f_1^{(i)} f_2^{(j)} f_3^{(n)} \\ &\quad + \sum_{i+j+n=k-1} c_{i,j,n} (\mu - \lambda_1 - \lambda_2 - \lambda_3 - k + 1) Z'(XY' - X'Y) f_1^{(i)} f_2^{(j)} f_3^{(n)}, \end{aligned}$$

$$\Psi_1 = \mathcal{C}([X, Y], Z) = \sum_{i+j+n=k-1} c_{i,j,n} Z'(XY' - X'Y) f_1^{(i)} f_2^{(j)} f_3^{(n)},$$

$$\Psi_2 = \mathcal{C}([X, Z], Y) = \sum_{i+j+n=k-1} c_{i,j,n} Y'(XZ' - X'Z) f_1^{(i)} f_2^{(j)} f_3^{(n)},$$

$$\Psi_3 = \mathcal{C}([Y, Z], X) = \sum_{i+j+n=k-1} c_{i,j,n} X'(YZ' - Y'Z) f_1^{(i)} f_2^{(j)} f_3^{(n)}.$$

Then, by considering the equation (3.13), we can write

$$(3.14) \quad L_X^{\lambda_1, \lambda_2, \lambda_3, \mu} \mathcal{C}(Y, Z)(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3})$$

$$-\mathcal{C}([X, Y], Z) + \circlearrowleft (X, Y, Z) = \sum_{t=1}^3 (\Phi_t - \Psi_t) = 0.$$

We show that the equation (3.14) is satisfied only if $c'_{i,j,n} = 0$ and $c_{i,j,n}(\mu - \lambda_1 - \lambda_2 - \lambda_3 - k) = 0$.

Any trivial 2-cocycle should be of the form

$$\begin{aligned} & L_X^{\lambda, \mu} b(Y)(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3}) - L_Y^{\lambda, \mu} b(X)(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3}) \\ & - b([X, Y])(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3}). \end{aligned}$$

By using Lemma 3.2, we have

$$\begin{aligned} (3.15) \quad & \delta b(X, Y)(f_1, f_2, f_3) \\ & = L_X^{\lambda, \lambda_1 + \lambda_2 + \lambda_3 + k} b(Y)(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3}) \\ & \quad - L_Y^{\lambda, \lambda_1 + \lambda_2 + \lambda_3 + k} b(X)(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3}) \\ & \quad - b([X, Y])(f_1 dx^{\lambda_1}, f_2 dx^{\lambda_2}, f_3 dx^{\lambda_3}) \\ & = \sum_{i+j+n=k} c_{i,j,n}(\mu - \lambda_1 - \lambda_2 - \lambda_3 - k)(X'Y - XY')f_1^{(i)}f_2^{(j)}f_3^{(n)}. \end{aligned}$$

4. The space $\mathbf{H}_{\text{diff}}^2(\mathfrak{aff}(1), \mathcal{D}_{\lambda, \mu}^n)$

The second main result of paper is the following.

Theorem 4.1. *The space $\mathbf{H}_{\text{diff}}^2(\mathfrak{aff}(1), \mathcal{D}_{\lambda, \mu}^n)$ has the following structure:*

$$(4.16) \quad \mathbf{H}_{\text{diff}}^2(\mathfrak{aff}(1), \mathcal{D}_{\lambda, \mu}^n) \simeq \begin{cases} \mathbb{R}^{C_{n+k-1}} & \text{if } \mu = \lambda_1 + \dots + \lambda_n + k, \\ 0 & \text{otherwise.} \end{cases}$$

The following 2-cocycles span the corresponding cohomology spaces:

$$\begin{aligned} & \mathcal{C}\left(X \frac{d}{dx}, Y \frac{d}{dx}\right)(f_1 dx^{\lambda_1}, \dots, f_n dx^{\lambda_n}) \\ & = \sum_{i_1 + \dots + i_n = 0}^{k-1} c_{i_1, \dots, i_n}(XY' - X'Y)f_1^{(i_1)} \dots f_n^{(i_n)} dx^{\lambda_1 + \dots + \lambda_n + k}, \end{aligned}$$

where c_{i_1, \dots, i_n} are constants, $f_i dx^{\lambda_i} \in \mathcal{F}_{\lambda_i}$ and $X \frac{d}{dx}, Y \frac{d}{dx} \in \mathfrak{aff}(1)$.

We need the following lemma.

Lemma 4.2. *Let $b \in C^1(\mathfrak{aff}(1), \mathcal{D}_{\lambda, \lambda_1 + \dots + \lambda_n + k}^n)$ be defined as follows: for $X \frac{d}{dx} \in \mathfrak{aff}(1)$, $f_i dx^{\lambda_i} \in \mathcal{F}_{\lambda_i}$.*

$$(4.17) \quad \begin{aligned} b(X)(f_1, f_2, \dots, f_n) & = \sum_{i_1 + \dots + i_n = k} c_{i_1, \dots, i_n} X f_1^{(i_1)} \dots f_n^{(i_n)} dx^\mu \\ & \quad + \sum_{i_1 + \dots + i_n = k-1} \alpha_{i_1, \dots, i_n} X' f_1^{(i_1)} \dots f_n^{(i_n)} dx^\mu, \end{aligned}$$

where the coefficients c_{i_1, \dots, i_n} and α_{i_1, \dots, i_n} are constants.

Then the map $\delta b : \mathbf{aff}(1)^{\otimes 2} \rightarrow \mathcal{D}_{\underline{\lambda}, \lambda_1 + \dots + \lambda_n + k}^n$ is given by

$$\begin{aligned}
(4.18) \quad \delta b(X, Y)(f_1, \dots, f_n) &= L_X^{\lambda_1, \dots, \lambda_n, \lambda_1 + \dots + \lambda_n + k} b(Y)(f_1 dx^{\lambda_1}, \dots, f_n dx^{\lambda_n}) \\
&\quad - L_Y^{\lambda_1, \dots, \lambda_n, \lambda_1 + \dots + \lambda_n + k} b(X)(f_1 dx^{\lambda_1}, \dots, f_n dx^{\lambda_n}) \\
&\quad - b([X, Y])(f_1 dx^{\lambda_1}, \dots, f_n dx^{\lambda_n}) \\
&= \sum_{i_1 + \dots + i_n = k} c_{i_1, \dots, i_n} (\mu - \lambda_1 - \dots - \lambda_n - k) (X'Y - XY') f_1^{(i_1)} \dots f_n^{(i_n)}.
\end{aligned}$$

4.1. Proof of theorem

Any 2-cocycle $\mathcal{C} \in Z_{\text{diff}}^1(\mathbf{aff}(1), \mathcal{D}_{\underline{\lambda}, \lambda_1 + \dots + \lambda_n + k}^n)$ should retain the following general form:

$$\begin{aligned}
(4.19) \quad \mathcal{C}(X, Y)(f_1, \dots, f_n) &= \sum_{i_1 + \dots + i_n = k-1} c_{i_1, \dots, i_n} (XY' - X'Y) f_1^{(i_1)} \dots f_n^{(i_n)} dx^\mu,
\end{aligned}$$

where c_{i_1, \dots, i_n} are constants.

The 2-cocycle condition reads as follows: for all $f_i dx^{\lambda_i} \in \mathcal{F}_{\lambda_i}$, and for all $X, Y \in \mathbf{aff}(1)$, we have

$$L_X^{\lambda_1, \dots, \lambda_n; \mu} \mathcal{C}(Y, Z)(f_1 dx^{\lambda_1}, \dots, f_n dx^{\lambda_n}) - \mathcal{C}([X, Y], Z) + \odot(X, Y, Z) = 0.$$

A direct computation, proves that the coefficient of the component $f_1^{(i_1)} \dots f_n^{(i_n)}$ in the 2-cocycle condition above is equal to

$$(4.20) \quad \mu - \lambda_1 - \dots - \lambda_n - k = 0.$$

Any trivial 2-cocycle should be of the form

$$\begin{aligned}
&L_X^{\lambda_1, \dots, \lambda_n; \mu} b(Y)(f_1 dx^{\lambda_1}, \dots, f_n dx^{\lambda_n}) - L_Y^{\lambda_1, \dots, \lambda_n; \mu} b(X)(f_1 dx^{\lambda_1}, \dots, f_n dx^{\lambda_n}) \\
&- b([X, Y])(f_1 dx^{\lambda_1}, \dots, f_n dx^{\lambda_n}).
\end{aligned}$$

By using Lemma 4.2, we have

$$\begin{aligned}
(4.21) \quad \delta b(X, Y)(f_1, \dots, f_n) &= L_X^{\lambda_1, \dots, \lambda_n; \lambda_1 + \dots + \lambda_n + k} b(Y)(f_1 dx^{\lambda_1}, \dots, f_n dx^{\lambda_n}) \\
&\quad - L_Y^{\lambda_1, \dots, \lambda_n; \lambda_1 + \dots + \lambda_n + k} b(X)(f_1 dx^{\lambda_1}, \dots, f_n dx^{\lambda_n}) \\
&\quad - b([X, Y])(f_1 dx^{\lambda_1}, \dots, f_n dx^{\lambda_n}) \\
&= \sum_{i_1 + \dots + i_n = k} c_{i_1, \dots, i_n} (\mu - \lambda_1 - \dots - \lambda_n - k) (X'Y - XY') f_1^{(i_1)} \dots f_n^{(i_n)}.
\end{aligned}$$

References

- [1] I. Basdouri, I. Laraiedh, and O. Ncib, *The linear $\mathfrak{aff}(n|1)$ -invariant differential operators on weighted densities on the superspace $\mathbb{R}^{1|n}$ and $\mathfrak{aff}(n|1)$ -relative cohomology*, Int. J. Geom. Methods Mod. Phys. **10** (2013), no. 4, 1320004, 9 pp.
- [2] G. G. A. Bäuerle and E. A. de Kerf, *Lie algebras. Part 1*, Studies in Mathematical Physics, **1**, North-Holland Publishing Co., Amsterdam, 1990.
- [3] G. G. A. Bäuerle, E. A. de Kerf, A. van Groesen, E. M. de Jager, and A. P. E. Ten Kroode, *Finite and infinite dimensional Lie algebras and their application in physics*. Studies in mathematical physics **7**, North-Holland, 1997.
- [4] G. Dito, M. Flato, D. Sternheimer, and L. Takhtajan, *Deformation quantization and Nambu mechanics*, Comm. Math. Phys. **183** (1997), no. 1, 1–22.
- [5] G. Dito and D. Sternheimer, *Deformation quantization: genesis, developments and metamorphoses*, IRMA Lectures in Mathematics and Theoretical Physics 1, Walter De Gruyter, Berlin 2002, 9–54.
- [6] J.-P. Francoise, G. L. Naber, and T. S. Tsun, *Encyclopedia of Mathematical Physics*. Current Algebra. ISBN 978-0-12-512666-3- via Science Direct, 2006.
- [7] D. B. Fuks, *Cohomology of infinite-dimensional Lie algebras*, Contemporary Soviet Mathematics, Consultants Bureau, New York, 1986.
- [8] P. Goddard, *Kac-Moody and Virasoro algebras in mathematical physics*, in VIIIth international congress on mathematical physics (Marseille, 1986), 390–401, World Sci. Publishing, Singapore, 1987.
- [9] V. G. Kac, *Infinite-dimensional Lie algebras*, third edition, Cambridge Uni. Press, Cambridge, 1990.
- [10] A. A. Kirillov, *Elementy teorii predstavlenii*, Nauka, Moscow 1972.
- [11] ———, *Geometric quantization*, in Current problems in mathematics. Fundamental directions, Vol. 4, 141–178, 291, Itogi Nauki i Tekhniki, Akad. Nauk SSSR, Vsesoyuz. Inst. Nauchn. i Tekhn. Inform., Moscow, 1985.
- [12] A. W. Knap, *Lie groups beyond an introduction*, second edition, Progress in Mathematics, **140**, Birkhäuser Boston, Inc., Boston, MA, 2002.
- [13] B. Kostant, *Quantization and unitary representations. I. Prequantization*, in Lectures in modern analysis and applications, III, 87–208. Lecture Notes in Math., 170, Springer, Berlin, 1970.
- [14] D. Sternheimer, J. Rawnsley, and S. Gutt, *Deformation theory and symplectic geometry*, Mathematical Physics Studies, **20**, Kluwer Academic Publishers Group, Dordrecht, 1997.

IMED BASDOURI
 DÉPARTEMENT DE MATHÉMATIQUES
 FACULTÉ DES SCIENCES DE GAFSA
 TUNISIE
 Email address: basdourimed@yahoo.fr

AMMAR DERBALI
 UNIVERSITÉ DE GAFSA
 FACULTÉ DES SCIENCES
 DÉPARTEMENT DE MATHÉMATIQUES
 TUNISIE
 Email address: ammar.derbali1@yahoo.fr

SOUMAYA SAIDI
UNIVERSITÉ DE GAFSA
FACULTÉ DES SCIENCES
DÉPARTEMENT DE MATHÉMATIQUES
TUNISIE
Email address: soumayasaidi24@yahoo.com