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# A LOCAL MINIMIZER IN SMECTIC LIQUID CRYSTALS

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ABSTRACT. In this paper, we consider a simplified energy functional for ferroelectric liquid crystals. With a special ansatz, we study existence of local minimizer with a certain symmetry.

# 1. Introduction

Let  $\Omega$  be bounded domain in  $\mathbf{R}^2$  which is occupied by a liquid crystal. Smectic Liquid crystal is described by the direction field  $\mathbf{n}$  and a complex vector field  $\psi$ . For more details about the governing energy, we refer the reader to [2, 3] and references are therein.

Throughout this paper, we assume that  $\psi = e^{iw}$  and  $\Omega = (-\pi, \pi) \times (-\pi h, \pi h)$  for small h > 0. We further suppose that **n** is a fixed unit vector and w is 2h- periodic. With the above conditions, we consider the simplest energy functional

(1.1) 
$$\mathcal{E}(w) = \int_{\Omega^+} \left\{ w_x^2 + (w_z + 1)^2 \right\} \, d\mathbf{x}$$

in a class of functions defined below. Here  $\Omega^+ = (-\pi, \pi) \times (0, \pi h)$ . Let

$$\mathcal{M} = \{ w : \Omega^+ \to \mathbf{R} : \Delta w = \delta_{z_1} - \delta_{z_2} \text{ in } \Omega^+, \frac{\partial w}{\partial \nu} = 0 \text{ on } \partial \Omega^+, z_1, z_2 \in \Omega \}.$$

From now on, we study minimization problem of the energy functional

$$\mathcal{E}(w) = \int_{\Omega^+} \left\{ w_x^2 + (w_z + 1)^2 \right\} \, d\mathbf{x} = \int_{\Omega^+} \left\{ |\nabla w|^2 + 2w_z \right\} \, d\mathbf{x} + 2\pi^2 h.$$

in  $\mathcal{M}$ .

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Since each function  $w \in \mathcal{M}$  depends on two points  $z_1, z_2 \in \Omega$ , the minimization problem is equivalent to find a pair of  $z_1$  and  $z_2$  in  $\Omega$  which minimizes  $\mathcal{E}(w)$ . For a given value h, let V be such that

$$\Delta V = \delta_{(0,0)} - \frac{1}{4\pi^2 h} \text{ in } \Omega,$$
$$\frac{\partial V}{\partial \nu} = 0 \text{ on } \partial \Omega.$$

It can be shown that for  $w \in \mathcal{M}$  the energy  $\mathcal{E}(w)$  can be written as  $-V(2z_2) + 2V(z_2 - z_1) - V(2z_1) + 2V(z_1 + z_2) + 2(z_1 - z_2) + 2\pi^2 h.$ 

For simplicity we assume that  $0 < z_1 < \frac{\pi h}{2}$ ,  $z_1 < z_2 < \pi h$ . Let  $\Gamma(z_1, z_2) = -V(2z_2) + 2V(z_2 - z_1) - V(2z_1) + 2V(z_1 + z_2) + 2(z_1 - z_2)$ . Then we have the following Lemmas.

LEMMA 1. For  $h > \frac{2}{\pi^2}$ ,  $\Gamma(d, \pi h - d)$  achieves its local minimum.

*Proof.* We first write V in terms of singular and regular parts as follows:

(1.2) 
$$V = \frac{1}{2\pi} \log |\mathbf{x}| - U, \ \mathbf{x} = (x, z),$$

(1.3) 
$$\Delta U = \frac{1}{4\pi^2 h} \text{ in } \Omega,$$

(1.4) 
$$\frac{\partial U}{\partial \nu} = \frac{\partial}{\partial \nu} \left( \frac{1}{2\pi} \log |\mathbf{x}| \right) \text{ on } \partial \Omega.$$

We get

(1.5) 
$$\Gamma(d, \pi h - d) = \frac{1}{\pi} \left( \log(\pi h - 2d) - \log(2d) \right) + 4d + 2 \left[ U(0, 2d) - U(0, \pi h - 2d) \right].$$

From the standard elliptic PDEs [4], one can show that  $U(0, 2d) - U(0, \pi h - 2d)$  is a non-decreasing function. It suffices to show that  $\Phi(d) := \Gamma(d, \pi h - d) - 2[U(0, 2d) - U(0, \pi h - 2d)]$  has a local minimum. Note that

(1.6) 
$$\Phi'(d) = \frac{1}{\pi} \left( \frac{-2}{\pi h - 2d} - \frac{1}{d} \right) + 4 = 0,$$

It is clear that (1.6) has roots  $d_0 = \frac{2\pi h \pm \sqrt{4\pi^2 h^2 - 8h}}{8}$  if  $h > \frac{2}{\pi^2}$ . Since  $\lim_{d \uparrow \frac{\pi h}{2}} \Phi(d) = -\infty$  and  $\lim_{d \downarrow 0} \Phi(d) = \infty$ ,  $\Gamma(d, \pi h - d)$  has a local minimum.

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LEMMA 2.  $\Gamma$  is symmetric with respect to the line  $l: z_1 + z_2 = \pi h$ and the minimum point obtained in lemma 1 is also a critical point of  $\Gamma$ . Moreover,

(1) if  $z_1$  is fixed, then

$$\lim_{z_2 \downarrow z_1} \Gamma(z_1, z_2) = -\infty, \text{ and } \lim_{z_2 \uparrow \pi h} \Gamma(z_1, z_2) = \infty.$$

(2) if  $z_2$  is fixed, then

$$\lim_{z_1 \uparrow z_2} \Gamma(z_1, z_2) = -\infty, \text{ and } \lim_{z_1 \downarrow 0} \Gamma(z_1, z_2) = \infty.$$

*Proof.* Let (z, w) and  $(\tilde{z}, \tilde{w})$  be symmetric with respect to l. Then we get

(1.7) 
$$\tilde{z} - \tilde{w} = z - w, \ z + w = 2\pi h - (\tilde{z} + \tilde{w}).$$

Since V is symmetric with respect to the origin and  $2\pi h$ -periodic, we obtain

$$\begin{split} \Gamma(\tilde{z},\tilde{w}) &= -V(2\tilde{w}) + 2V(\tilde{w}-\tilde{z}) - V(2\tilde{z}) + 2V(\tilde{z}+\tilde{w}) + 2(\tilde{z}-\tilde{w}) \\ &= -V(2\pi h - 2z) + 2V(w-z) - V(2\pi h - 2w) \\ &+ 2V(2\pi h - z - w) + 2(z-w) \\ &= -V(2z) + 2V(w-z) - V(2w) + 2V(z+w) + 2(z-w). \end{split}$$

It also follows that  $(z_1, z_2) = (d, \pi h - d)$  is a critical point if d is a local minimum point of  $\Gamma(d, \pi h - d)$ .

From Lemma 1 and Lemma 2, we conclude the following theorem

THEOREM 3. If  $h > \frac{2}{\pi^2}$  and  $\Gamma(d, \pi h - d)$  has a local minimum at  $d = d_0$ , then  $(d_0, \pi h - d_0)$  is also local minimum of  $\Gamma$  in the domain D, where

$$D = \left\{ (z_1, z_2) : 0 < z_1 < \frac{\pi h}{2}, z_1 < z_2 < \pi h \right\}.$$

*Proof.* Using  $u = z_1 + z_2$ ,  $v = z_2 - z_1$ , we express  $\Gamma$  in terms of u and v as

$$\Gamma(z_1, z_2) = 2V(u) + 2V(v) - V(u - v) - V(u + v) - 2v := W(u, v).$$

Since  $\Gamma(d, \pi h - d)$  has a local minimum at  $d_0$ , then we get

$$\frac{\partial^2}{\partial u^2} W(d_0, \pi h - 2d_0) = -V''(2d_0) + V''(\pi h - 2d_0) > 0$$

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We compute

$$\frac{\partial^2}{\partial v^2} W(d_0, \pi h - 2d_0) = -V''(2d_0) + V''(\pi h).$$

By symmetry of V, we obtain

$$\frac{\partial^2}{\partial u \partial v} W(d_0, \pi h - 2d_0) = -V''(2\pi h - 2d_0) + V''(2d_0) = 0.$$

Now

$$V(\mathbf{x}) = \frac{1}{2\pi} \log |\mathbf{x}| - U(\mathbf{x}),$$
  

$$\Delta U = \frac{1}{4\pi^2 h},$$
  

$$\frac{\partial U}{\partial \nu} = \begin{cases} \frac{1}{2(\pi^2 + z^2)} & \text{if } x = \pm \pi, \\ \frac{h}{2(\pi^2 h^2 + x^2)} & \text{if } z = \pm \pi h \end{cases}$$

Let (z, w) be a point on the line passing through  $(d_0, \pi h - d_0)$  perpendicular to the line  $z_1 + z_2 = \pi h$ . Then  $z - d_0 = w - (\pi h - d_0)$ , and thus

$$w = z + \pi h - 2d_0, \quad w - z = \pi h - 2d_0,$$
  
$$z + w = 2z + \pi h - 2d_0, \quad 2w = 2z + 2(\pi h - 2d_0).$$

Since  $U(0, \cdot)$  is  $2\pi h$ -periodic and symmetric with respect to  $z = \pi h$ , we get

$$F(z) := U(0, 2z) + U(0, 2w) - 2U(z + w)$$
  
=  $U(0, 2z) + U(0, 2z + 2(\pi h - 2d_0)) - 2U(0, 2z + \pi h - 2d_0)$   
=  $U(0, 2z) + U(0, 2z - 4d_0)) - 2U(0, \pi h - (2z - 2d_0)).$ 

Note that  $d_0 < \frac{\pi h}{4}$ . For any z satisfying  $d_0 \leq z \leq \pi h$ ,

$$F'(z) = 2U_z(0, 2z) + 2U_z(0, 2z - 4d_0) + 4U_z(0, \pi h - (2z - 2d_0)).$$

Then

$$\lim_{z \downarrow d_0} F'(z) = 2U_z(0, 2d_0) + 2U_z(0, -2d_0) + 4U_z(0, \pi h)$$
$$= 4U_z(0, \pi h) > 0.$$

This implies that F is increasing on  $[d_0, d_0 + \delta_1)$  for some  $\delta_1 > 0$ . Similarly, F is decreasing on  $(d_0 - \delta_2, d_0]$  for some  $\delta_2 > 0$  by symmetry.

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Now,

$$\begin{split} \Gamma(z,w) = & F(z) - \frac{1}{2\pi} \left( \log |2z - 4d_0| + \log |2z| - 2\log |\pi h - (2z - 2d_0)| \right) \\ & + 2V(\pi h - 2d_0) - 2(\pi h - 2d_0). \end{split}$$

It suffices to show that  $G(z) = \Gamma(z, w) - F(z)$  has a local minimum at  $z = d_0$ . We calculate

$$G''(z) = \frac{1}{(z - 2d_0)^2} + \frac{1}{z^2} - \frac{8}{(2z - 2d_0 - \pi h)^2}.$$

and thus using  $d_0 < \frac{\pi h}{4}$  we get

$$G''(d_0) = \frac{2}{d_0^2} - \frac{8}{\pi^2 h^2} > \frac{24}{\pi^2 h^2} > 0.$$

Hence

$$\frac{\partial^2}{\partial v^2}W(d_0,\pi h - 2d_0) > 0.$$

By second derivative test, we see that  $\Gamma$  indeed achieves its local minimum at  $(d_0, \pi h - d_0)$ . This completes the proof.

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#### References

- P. G. de Gennes and J. Prost, *The Physics of Liquid Crystals*, Oxford University Press, 1993.
- [2] S. T. Lagerwall, Ferroelectric and Antiferroelectric Liquid Crystals, Wiley-VCH, 1999.
- [3] J. Park and M. C. Calderer, Analysis of nonlocal electrostatic effects in chiral smectic c liquid crystals, SIAM J. Appl. Math. 66 (2006), 2107-2126.
- [4] L. C. Evans, *Partial differential equations*, Graduate Studies in Mathematics, 19. American Mathematical Society, Providence, RI, 1998.

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