# MODULI SPACES OF ORIENTED TYPE $\mathcal{A}$ MANIFOLDS OF DIMENSION AT LEAST 3 

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#### Abstract

We examine the moduli space of oriented locally homogeneous manifolds of Type $\mathcal{A}$ which have non-degenerate symmetric Ricci tensor both in the setting of manifolds with torsion and also in the torsion free setting where the dimension is at least 3 . These exhibit phenomena that is very different than in the case of surfaces. In dimension 3, we determine all the possible symmetry groups in the torsion free setting.


## 1. Introduction

Let $M$ be a smooth oriented manifold of dimension $m$; if $M \subset \mathbb{R}^{m}$, then the orientation will be given by $d x^{1} \wedge \cdots \wedge d x^{m}$. Let $\nabla$ be a connection on the tangent bundle of $M$. One says that $\mathcal{M}:=(M, \nabla)$ is torsion free if $\nabla_{\xi} \eta-\nabla_{\eta} \xi=$ $[\xi, \eta]$. Let $\vec{x}:=\left(x^{1}, \ldots, x^{m}\right)$ be a system of local coordinates on $M$. Adopt the Einstein convention and sum over repeated indices to expand $\nabla_{\partial_{x^{i}}} \partial_{x^{j}}=$ $\Gamma_{i j}{ }^{k} \partial_{x^{k}}$ in terms of the Christoffel symbols $\Gamma=\left(\Gamma_{i j}{ }^{k}\right)$; the condition that $\nabla$ is torsion free is then equivalent to the symmetry $\Gamma_{i j}{ }^{k}=\Gamma_{j i}{ }^{k}$. The importance of the torsion free condition lies in the following result which permits one to normalize the coordinate system so that only the second and higher order derivatives of the connection 1-form play a role. As we shall not be using this result, we shall omit the proof and refer the interested reader to, for example, Lemma 3.5 of [22] for the proof.
Theorem 1.1. $\mathcal{M}$ is torsion free if and only if for every point $P$ of $M$, there exist coordinates centered at $P$ so that $\Gamma_{i j}{ }^{k}(P)=0$.

Although much of Riemannian geometry involves the study of the LeviCivita connection, which is without torsion, in recent years connections which have torsion have played an important role in many developments. Papers have dealt with the theory of gravity with torsion [5, 25], B-metrics [18, 26], almost hypercomplex geometries [32], string theory [17, 24], spin geometries [3,27,28], contact geometries [2], almost product manifolds [34], non-integrable

[^0]geometries [1], the non-commutative residue for manifolds with boundary [37], Hermitian and anti-Hermitian geometry [33], CR geometry [14], generalized specialized holonomy [11], sigma models [16], Einstein-Weyl gravity at the linearized level [13], and Yang-Mills flow with torsion [20] to name just a few areas. Perhaps surprisingly, even the 2-dimensional case is of interest; connections on surfaces have been used to construct new examples of pseudo-Riemannian metrics without a corresponding Riemannian counterpart [9,10,12,31].

### 1.1. Local homogeneity

One says that $\mathcal{M}$ is locally homogeneous if given any two points $P$ and $Q$ of $M$, there exists the germ of a diffeomorphism $\Phi$ taking $P$ to $Q$ which preserves $\nabla$. Physically, the locally homogeneous setting is of particular interest as it corresponds to locally isotropic geometries. One has the following examples of homogeneous geometries:

Type $\mathcal{A}$. Let $M=\mathbb{R}^{m}$ and let $\Gamma \in\left(\mathbb{R}^{m}\right)^{*} \otimes\left(\mathbb{R}^{m}\right)^{*} \otimes \mathbb{R}^{m}$ be constant. The translation group $\mathbb{R}^{m}$ acts transitively on $M$ and preserves $\nabla$.
Type $\mathcal{B}$. Let $M=\mathbb{R}^{+} \times \mathbb{R}^{m-1}$ and let $\Gamma_{i j}{ }^{k}=\left(x^{1}\right)^{-1} C_{i j}{ }^{k}$ for $C \in\left(\mathbb{R}^{m}\right)^{*} \otimes$ $\left(\mathbb{R}^{m}\right)^{*} \otimes \mathbb{R}^{m}$ constant. The $a x+b$ group $\left(x^{1}, x^{2}, \ldots\right) \rightarrow\left(a x^{1}, a x^{2}+b^{2}, \ldots, a x^{m}+\right.$ $b^{m}$ ) for $a>0$ and $\vec{b}=\left(b^{2}, \ldots, b^{m}\right) \in \mathbb{R}^{m-1}$ acts transitively on $M$ and preserves $\nabla$.

Type $\mathcal{C}$. Let $\nabla$ be the Levi-Civita connection of a complete simply connected pseudo-Riemannian manifold $M$ of constant sectional curvature.

### 1.2. Two dimensional geometry

The examples given above provide a complete family of models for the locally homogeneous surfaces; any locally homogeneous surface admits a coordinate atlas modeled on one of these examples. These classes are not disjoint. No surface is both Type $\mathcal{A}$ and Type $\mathcal{C}$. However there are surfaces that are both Type $\mathcal{A}$ and Type $\mathcal{B}$ and there are surfaces which are both Type $\mathcal{B}$ and Type $\mathcal{C}$. We refer to Opozda [35] for a proof of the following result in the torsion free setting and to Arias-Marco and Kowalski [4] for the extension to the case of surfaces with torsion. We refer as well to $[15,23,29,30,36]$ for related work.

Theorem 1.2. Let $\mathcal{M}=(M, \nabla)$ be a locally homogeneous surface where $\nabla$ can have torsion. Then at least one of the following three possibilities hold that describe the local geometry:
$(\mathcal{A})$ There exists a coordinate atlas so the Christoffel symbols $\Gamma_{i j}{ }^{k}$ are constant.
$(\mathcal{B})$ There exists a coordinate atlas so the Christoffel symbols have the form $\Gamma_{i j}^{k}=\left(x^{1}\right)^{-1} C_{i j}{ }^{k}$ for $C_{i j}^{k}$ constant and $x^{1}>0$.
(C) $\nabla$ is the Levi-Civita connection of a metric of constant Gauss curvature.

### 1.3. The Ricci tensor

The curvature operator $R$ and the Ricci tensor $\rho$ of an arbitrary connection are given by setting

$$
R(\xi, \eta):=\nabla_{\xi} \nabla_{\eta}-\nabla_{\eta} \nabla_{\xi}-\nabla_{[\xi, \eta]} \text { and } \rho(\xi, \eta):=\operatorname{Tr}\{\sigma \rightarrow R(\sigma, \xi) \eta\} .
$$

In terms of local coordinates,

$$
\begin{align*}
R_{i j k}^{l} & =\partial_{x_{i}} \Gamma_{j k}^{l}-\partial_{x^{j}} \Gamma_{i k}{ }^{l}+\Gamma_{i n}{ }^{l} \Gamma_{j k}{ }^{n}-\Gamma_{j n}{ }^{l} \Gamma_{i k}{ }^{n},  \tag{1}\\
\rho_{j k} & =\partial_{x_{i}} \Gamma_{j k}{ }^{i}-\partial_{x^{j}} \Gamma_{i k}{ }^{i}+\Gamma_{i n}{ }^{i} \Gamma_{j k}{ }^{n}-\Gamma_{j n}{ }^{i} \Gamma_{i k}{ }^{n} .
\end{align*}
$$

Note that in this setting $\rho$ need not be symmetric.

### 1.4. Type $\mathcal{A}$ geometry

Let $S^{2}\left(\mathbb{R}^{m}\right)$ denote the space of symmetric 2-cotensors on $\mathbb{R}^{m} ; \sigma=\sigma_{i j} d x^{i} \otimes$ $d x^{j} \in S^{2}\left(\mathbb{R}^{m}\right)$ if and only if $\sigma_{i j}=\sigma_{j i}$. The natural parameter spaces with which we shall be working are defined by:

$$
\mathcal{W}(m):=\left(\mathbb{R}^{m}\right)^{*} \otimes\left(\mathbb{R}^{m}\right)^{*} \otimes \mathbb{R}^{m} \text { and } \mathcal{Z}(m):=S^{2}\left(\mathbb{R}^{m}\right) \otimes \mathbb{R}^{m}
$$

Since $\mathcal{Z}(m)$ is a subset of $\mathcal{W}(m)$, properties true on $\mathcal{W}(m)$ are often inherited by $\mathcal{Z}(m)$ and we shall thus often not mention $\mathcal{Z}(m)$ explicitly. If $\Gamma \in \mathcal{W}(m)$, then $\Gamma$ defines a Type $\mathcal{A}$ connection $\nabla^{\Gamma}$ on $\mathbb{R}^{m} ; \Gamma \in \mathcal{Z}(m) \subset \mathcal{W}(m)$ if and only if $\nabla^{\Gamma}$ is torsion free. If $\Gamma \in \mathcal{W}(m)$ and $\tilde{\Gamma} \in \mathcal{W}(m)$, introduce the equivalence relation $\Gamma \sim \tilde{\Gamma}$ if there exists the germ of an orientation preserving diffeomorphism $\Phi$ from a point $P$ in $\mathbb{R}^{m}$ to a point $\tilde{P}$ in $\mathbb{R}^{m}$ so that $\Phi^{*}\left(\nabla^{\tilde{\Gamma}}\right)=\nabla^{\Gamma}$; the precise points in question are irrelevant as the structures are homogeneous. Since the torsion free condition is preserved by diffeomorphism, $\sim$ defines an equivalence relation on $\mathcal{Z}(m)$ as well.

We say that $\mathcal{M}=(M, \nabla)$ is Type $\mathcal{A}$ if there is an atlas $\left\{\mathcal{U}_{\alpha}=\left(U_{\alpha}, \Gamma_{\alpha}\right), \Phi_{\alpha_{1} \alpha_{2}}\right\}$ where $\left\{U_{\alpha}, \Phi_{\alpha_{1} \alpha_{2}}\right\}$ forms an oriented coordinate atlas for $M$ (i.e., $\operatorname{det}\left(\Phi_{\alpha_{1} \alpha_{2}}\right)>$ $0)$ and where $\nabla$ is defined by $\Gamma_{\alpha} \in \mathcal{W}(m)$ on $U_{\alpha}$. The coordinate transformations $\left\{\Phi_{\alpha_{1} \alpha_{2}}\right\}$ satisfy the intertwining rule $\Phi_{\alpha_{1} \alpha_{2}} \nabla^{\Gamma_{\alpha_{2}}}=\nabla^{\Gamma_{\alpha_{1}}} \Phi_{\alpha_{1} \alpha_{2}}$. Note that $\mathcal{M}$ is torsion free if and only if $\Gamma_{\alpha} \in \mathcal{Z}(m)$ for all $\alpha$. The intertwining rule implies that $\Gamma_{\alpha_{i}} \sim \Gamma_{\alpha_{j}}$ for all $i$ and $j$. We wish study the moduli space of local isomorphism types. If $\tilde{\mathcal{M}}$ is another Type $\mathcal{A}$ manifold which is defined by an atlas $\left\{\left(\tilde{U}_{\beta}, \tilde{\Gamma}_{\beta}\right), \tilde{\Phi}_{\beta_{1}, \beta_{2}}\right\}$, then $\mathcal{M}$ and $\tilde{\mathcal{M}}$ are locally isomorphic if and only if $\Gamma_{\alpha} \sim \tilde{\Gamma}_{\beta}$ for all $\alpha, \beta$. The moduli space of such local isomorphism classes is then the quotient of $\mathcal{W}(m)$ by the equivalence relation $\sim$. We define:

$$
\mathfrak{W}^{+}(m):=\mathcal{W}(m) / \sim \text { and } \mathfrak{Z}^{+}(m):=\mathcal{Z}(m) / \sim .
$$

If $\Gamma \in \mathcal{W}(m)$, then Equation (1) shows that the Ricci tensor associated to $\Gamma$ is

$$
\begin{equation*}
\rho_{\Gamma, j k}=\Gamma_{i n}{ }^{i} \Gamma_{j k}{ }^{n}-\Gamma_{j n}{ }^{i} \Gamma_{i k}{ }^{n} . \tag{2}
\end{equation*}
$$

For generic $\Gamma \in \mathcal{W}(m), \rho_{\Gamma, j k} \neq \rho_{\Gamma, k j}$ so $\rho_{\Gamma}$ is in general not symmetric. One defines, therefore, the symmetric Ricci tensor by setting:

$$
\rho_{s, \Gamma}(\eta, \zeta):=\frac{1}{2}(\rho(\eta, \zeta)+\rho(\zeta, \eta)) \text {, i.e., } \rho_{s, \Gamma}:=\frac{1}{2}\left\{\rho_{\Gamma, j k}+\rho_{\Gamma, k j}\right\} d x^{j} \otimes d x^{k} \text {. }
$$

If $\Gamma \in \mathcal{Z}(m)$, then Equation (2) shows $\rho_{\Gamma, j k}=\rho_{\Gamma, k j}$ so the Ricci tensor is already symmetric and there is no need to symmetrize. We shall be interested in the case that $\rho_{s, \Gamma}$ is non-degenerate as this is the generic case, see Theorem 1.6 below. Let $\operatorname{sign}\left(\rho_{s, \Gamma}\right)=(p, q)$ be the signature of the symmetric Ricci tensor; there are $p$ timelike directions and $q$ spacelike directions. Thus $(p, q)=(m, 0)$ implies $\rho_{s, \Gamma}$ is negative definite while $(p, q)=(0, m)$ implies $\rho_{s, \Gamma}$ is positive definite. If $p+q=m$, i.e., $\rho_{s, \Gamma}$ is non-degenerate, set

$$
\begin{aligned}
& \mathcal{W}(p, q):=\left\{\Gamma \in \mathcal{W}(m): \operatorname{sign}\left(\rho_{s, \Gamma}\right)=(p, q)\right\}, \quad \mathfrak{W}^{+}(p, q):=\mathcal{W}(p, q) / \sim, \\
& \mathcal{Z}(p, q):=\left\{\Gamma \in \mathcal{Z}(m): \operatorname{sign}\left(\rho_{s, \Gamma}\right)=(p, q)\right\}, \\
& \mathfrak{Z}^{+}(p, q):=\mathcal{Z}(p, q) / \sim .
\end{aligned}
$$

### 1.5. Reduction to the action of the general linear group

Let $\mathrm{GL}^{+}(m, \mathbb{R})$ be the group of linear transformations of $\mathbb{R}^{m}$ which preserve the orientation, i.e., $\operatorname{det}(T)>0$. This group acts on the Christoffel symbols $\Gamma \in \mathcal{W}(m)$ of a Type $\mathcal{A}$ geometry by change of coordinates; two indices are down and one is up. If $\left\{e_{i}\right\}$ is a basis for $\mathbb{R}^{m}$ and if $T \in \operatorname{GL}(m, \mathbb{R})$, then

$$
(T \Gamma)\left(e_{i}, e_{j}, e^{k}\right):=\Gamma\left(T e_{i}, T e_{j}, T e^{k}\right) .
$$

One has the following observation [8] which shows that in fact one does not need to consider arbitrary diffeomorphisms in defining the moduli space if the symmetric Ricci tensor is non-degenerate as the diffeomorphisms $\Phi_{\alpha_{1} \alpha_{2}}$ in the atlas are affine. This (in principal) reduces the problem to one in group representation theory.

Theorem 1.3. Let $\mathcal{U}_{\alpha}=\left\{\left(U_{\alpha}, \Gamma_{\alpha}\right), \Phi_{\alpha_{1} \alpha_{2}}\right\}$ be an oriented Type $\mathcal{A}$ atlas on a Type $\mathcal{A}$ manifold $\mathcal{M}$. Assume that the Ricci tensor $\rho_{s, \mathcal{M}}$ is non-degenerate.
(1) $\Phi_{\alpha_{1} \alpha_{2}} \vec{x}_{\alpha_{2}}=A_{\alpha_{1} \alpha_{2}} \vec{x}_{\alpha_{2}}+\vec{b}_{\alpha_{1} \alpha_{2}}$ where $A_{\alpha_{1} \alpha_{2}} \in \mathrm{GL}^{+}(m, \mathbb{R})$ and $\vec{b}_{\alpha_{1} \alpha_{2}} \in$ $\mathbb{R}^{m}$.
(2) $\mathfrak{W}^{+}(p, q)=\mathcal{W}(p, q) / \mathrm{GL}^{+}(m, \mathbb{R})$ and $\mathfrak{Z}^{+}(p, q)=\mathcal{Z}(p, q) / \mathrm{GL}^{+}(m, \mathbb{R})$.

Proof. The symmetric Ricci tensor is an invariantly defined pseudo-Riemannian metric on $\mathcal{M}$ which is preserved by the Type $\mathcal{A}$ coordinate transformations $\Phi_{\alpha_{1} \alpha_{2}}$. Since $\Gamma$ is constant, the components of $\rho_{s, \Gamma}$ are constant on $U_{\alpha}$ for any $\alpha$. Thus $\rho_{s, \Gamma}$ is flat and the coordinate transformations have the form given. This establishes Assertion 1; as translations do not change $\Gamma$, only the action of $\mathrm{GL}^{+}(m, \mathbb{R})$ is relevant in examining the moduli spaces. Assertion 2 now follows.

Remark 1.4. We note that Theorem 1.3 fails if we do not assume the Ricci tensor is non-degenerate. We refer, for example, to [8] for a further discussion of this point in the torsion free setting when $m=2$.

The remainder of this paper is devoted to the study of these geometries. Since case of surfaces is dealt with in $[7,8,21]$, we shall assume for the remainder of this paper that $m=p+q \geq 3$; there are phenomena in this setting not found in the case $m=2$.

### 1.6. Principal bundles

Let $G$ be a Lie group which acts smoothly on a manifold $N$. Let $G_{P}:=\{g \in$ $G: g P=P\}$ be the isotropy group of the action. The action is said to be fixed point free if $G_{P}=\{\mathrm{id}\}$ for all $P$. The action is said to be proper if given points $P_{n} \in N$ and $g_{n} \in G$ with $P_{n} \rightarrow P \in N$ and $g_{n} P_{n} \rightarrow \tilde{P} \in N$, we can choose a convergent subsequence so $g_{n_{k}} \rightarrow g \in G$. We refer to $[6,19]$ for the proof of the following result; see also the discussion in [21].

Theorem 1.5. Let the action of a Lie group $G$ on a manifold $N$ be fixed point free, smooth, and proper. Then there is a natural smooth structure on the quotient space $N / G$ so that $G \rightarrow N \rightarrow N / G$ is a principal $G$ bundle.

### 1.7. Generic phenomena

Let $\mathfrak{P}_{m}=\mathfrak{P}_{m}(\Gamma)$ be a polynomial defined on $\mathcal{W}(m)$ which is divisible by $\operatorname{det}\left(\rho_{s, \Gamma}\right)$ and which doesn't vanish identically on $\mathcal{Z}(m)$. Let

$$
\begin{aligned}
\mathcal{W}\left(p, q ; \mathfrak{P}_{m}\right) & :=\left\{\Gamma \in \mathcal{W}(p, q): \mathfrak{P}_{m}(\Gamma) \neq 0\right\}, \\
\mathfrak{W}^{+}\left(p, q ; \mathfrak{P}_{m}\right) & :=\mathcal{W}\left(p, q ; \mathfrak{P}_{m}\right) / \mathrm{GL}^{+}(m, \mathbb{R}), \\
\mathcal{Z}\left(p, q ; \mathfrak{P}_{m}\right) & :=\left\{\Gamma \in \mathcal{Z}(p, q): \mathfrak{P}_{m}(\Gamma) \neq 0\right\}, \\
\mathfrak{Z}^{+}\left(p, q ; \mathfrak{P}_{m}\right) & :=\mathcal{Z}\left(p, q ; \mathfrak{P}_{m}\right) / \mathrm{GL}^{+}(m, \mathbb{R}) ;
\end{aligned}
$$

$\mathcal{W}\left(p, q ; \mathfrak{P}_{m}\right)$ and $\mathcal{Z}\left(p, q ; \mathfrak{P}_{m}\right)$ are open dense subsets of $\mathcal{W}(p, q)$ and $\mathcal{Z}(p, q)$, respectively. We will prove the following result in Section 2.

Theorem 1.6. There exists a polynomial $\mathfrak{P}_{m}$ so that $\mathcal{W}\left(p, q ; \mathfrak{P}_{m}\right)$ and $\mathcal{Z}(p, q$; $\left.\mathfrak{P}_{m}\right)$ are $\mathrm{GL}^{+}(m, \mathbb{R})$ invariant subsets on which $\mathrm{GL}^{+}(m, \mathbb{R})$ acts properly and without fixed points. Consequently, there are natural smooth structures on the moduli spaces $\mathfrak{W}^{+}\left(p, q ; \mathfrak{P}_{m}\right)$ and $\mathfrak{Z}^{+}\left(p, q ; \mathfrak{P}_{m}\right)$ so the projections $\mathcal{W}\left(p, q ; \mathfrak{P}_{m}\right) \rightarrow$ $\mathfrak{W}^{+}\left(p, q ; \mathfrak{P}_{m}\right)$ and $\mathcal{Z}\left(p, q ; \mathfrak{P}_{m}\right) \rightarrow \mathfrak{Z}^{+}\left(p, q ; \mathfrak{P}_{m}\right)$ are smooth principal $\mathrm{GL}^{+}(m, \mathbb{R})$ bundles.

### 1.8. Results concerning the isotropy subgroup

Let $G_{\Gamma}^{+}$be the group of orientation preserving symmetries of $\left(\mathbb{R}^{m}, \nabla^{\Gamma}\right)$; $G_{\Gamma}^{+}:=\left\{T \in \mathrm{GL}^{+}(m, \mathbb{R}): T \Gamma=\Gamma\right\}$. We will prove the following result in Section 3.

## Theorem 1.7.

(1) Let $\Gamma_{n} \in \mathcal{W}(p, q)$ satisfy $\Gamma_{n} \rightarrow \Gamma \in \mathcal{W}(p, q)$. If $\operatorname{dim}\left\{G_{\Gamma_{n}}^{+}\right\} \geq 1$, then $\operatorname{dim}\left\{G_{\Gamma}^{+}\right\} \geq 1$.
(2) There exists $c(m)$ so that if $\Gamma \in \mathcal{W}(p, q)$ and if no element of $G_{\Gamma}^{+}$ has infinite order, then every element in $G_{\Gamma}^{+}$has order at most $c(m)$. Furthermore, $\lim _{m \rightarrow \infty} c(m)=\infty$.

### 1.9. Definite symmetric Ricci tensor

## Let

$\tilde{\mathcal{W}}(p, q):=\left\{\Gamma \in \mathcal{W}(p, q): G_{\Gamma}^{+}=\{\operatorname{id}\}\right\}, \quad \tilde{\mathfrak{W}}^{+}(p, q):=\tilde{\mathcal{W}}(p, q) / \mathrm{GL}^{+}(m, \mathbb{R})$,

$$
\tilde{\mathcal{Z}}(p, q):=\left\{\Gamma \in \mathcal{Z}(p, q): G_{\Gamma}^{+}=\{\operatorname{id}\}\right\}, \quad \tilde{\mathfrak{Z}}^{+}(p, q):=\tilde{\mathcal{Z}}(p, q) / \mathrm{GL}^{+}(m, \mathbb{R})
$$

If $(p, q) \in\{(0, m),(m, 0)\}$ so $\rho_{s, \Gamma}$ is definite, then one need not consider the generic situation but can simply exclude the fixed point sets and work directly with the sets $\tilde{\mathcal{W}}(p, q)$ and $\tilde{\mathcal{Z}}(p, q)$. We shall prove the following result in Section 4.

Theorem 1.8. Let $(p, q) \in\{(m, 0),(0, m)\}$. Then:
(1) The action of $\mathrm{GL}^{+}(m, \mathbb{R})$ on $\mathcal{W}(p, q)$ and on $\mathcal{Z}(p, q)$ is proper.
(2) $\tilde{\mathcal{W}}(p, q)$ and $\tilde{\mathcal{Z}}(p, q)$ are open dense subsets of $\mathcal{W}(p, q)$ and $\mathcal{Z}(p, q)$, respectively.
(3) One can define natural smooth structures on the associated moduli spaces $\tilde{\mathfrak{W}}^{+}(p, q)$ and $\tilde{\mathfrak{Z}}^{+}(p, q)$ so that $\tilde{\mathcal{W}}(p, q) \rightarrow \tilde{\mathfrak{W}}^{+}(p, q)$ and $\tilde{\mathcal{Z}}(p, q) \rightarrow$ $\tilde{\mathfrak{Z}}^{+}(p, q)$ are smooth principal $\mathrm{GL}^{+}(m, \mathbb{R})$ bundles.

### 1.10. The higher signature setting

In Section 5, we will prove the following result which shows that Assertion 1 of Theorem 1.8 fails in the higher signature setting.

Theorem 1.9. Let $p \geq 1$, let $q \geq 1$, and let $p+q \geq 3$. There exists $\Gamma \in \mathcal{Z}(p, q)$ so that $G_{\Gamma}^{+}$is not compact. Consequently, the action of $\mathrm{GL}^{+}(m, \mathbb{R})$ on $\mathcal{Z}(p, q)$ or on $\mathcal{W}(p, q)$ is not proper.

### 1.11. Two dimensional geometry

The two dimensional setting is relatively easy to examine. Since it informs many of the constructions we will employ in the 3-dimensional setting, it seems worth while discussing it in a bit of detail; a construction which will be used in the proof of Theorem 1.7(2) will renter in analysis of the 2-dimensional setting. We introduce the following basic structure:

Definition 1.10. Let $\Gamma_{2}$ be the structure
$\Gamma_{11}{ }^{1}=\frac{1}{\sqrt{2}}, \quad \Gamma_{11}^{2}=0, \quad \Gamma_{12}{ }^{1}=0, \quad \Gamma_{12}^{2}=-\frac{1}{\sqrt{2}}, \quad \Gamma_{22}^{1}=-\frac{1}{\sqrt{2}}, \quad \Gamma_{22}^{2}=0$.
We obtain $\rho_{\Gamma}=\operatorname{diag}(-1,-1)$.
The following result was proved in [8] using different methods; we give a different proof in Section 6 to introduce arguments we will use subsequently.

Theorem 1.11. Adopt the notation established above.
(1) Let $(p, q)=(1,1)$ or $(p, q)=(0,2)$. Then the action of $\mathrm{GL}^{+}(2, \mathbb{R})$ on $\mathcal{Z}(p, q)$ is fixed point free and proper. Thus $\mathcal{Z}(p, q) \rightarrow \mathfrak{Z}^{+}(p, q)$ is a principal $\mathrm{GL}^{+}(2, \mathbb{R})$ bundle over a real analytic surface.
(2) Let $(p, q)=(2,0)$. If $G_{\Gamma}^{+} \neq\{\mathrm{id}\}$, then $G_{\Gamma}^{+}=\mathbb{Z}_{3}$ and $\Gamma$ is isomorphic to the structure $\Gamma_{2}$ of Definition 1.10. $\mathrm{GL}^{+}(2, \mathbb{R})$ acts properly on $\mathcal{Z}(2,0)$ and $\mathcal{Z}(2,0) \rightarrow \mathfrak{Z}_{2,0}^{+}-\left[\Gamma_{2}\right]$ is a principal $\mathrm{GL}(2, \mathbb{R})$ bundle over a real analytic surface once we remove the exceptional orbit corresponding to $\Gamma_{2}$.

### 1.12. Three dimensional geometry

We now restrict to dimension $m=3$ and the torsion free setting in order to illustrate the possible isotropy subgroups. The examples where $\operatorname{dim}\left\{G_{\Gamma}^{+}\right\}>0$ form two families given in (1) and (2) below. The remaining structure groups are all finite and comprise one of the following $\left\{\mathbb{Z}_{3}, \mathbb{Z}_{2} \oplus \mathbb{Z}_{2}, \mathbb{Z}_{2}, s_{3}, a_{4}\right\}$; they appear in the two families given in (3) and (4) below. Thus one obtains that the constant $c(3)=3$ in Theorem 1.7. We will prove the following result in Section 7:

Theorem 1.12. Let $\Gamma \in \mathcal{Z}(p, q)$ for $p+q=3$. Assume $G_{\Gamma}^{+} \neq\{\mathrm{id}\}$. We can make a linear change of coordinates so that one of the following 4 possibilities holds:
(1) There exist $(a, b, c, d) \in \mathbb{R}^{4}$ so $G_{\Gamma}^{+}=\mathrm{SO}(1,1), \Gamma_{12}{ }^{3}=a, \Gamma_{13}{ }^{1}=b$, $\Gamma_{23}{ }^{2}=c, \Gamma_{33}{ }^{3}=d$, and $\rho=a d\left(e^{1} \otimes e^{2}+e^{2} \otimes e^{1}\right)+\left(-b^{2}+b d+c(-c+\right.$ d) $) e^{3} \otimes e^{3}$. We require $a d \neq 0$ and $-b^{2}+b d+c(-c+d) \neq 0$.
(2) There exist $(a, b, c, d) \in \mathbb{R}^{4}$ so $G_{\Gamma}^{+}=\mathrm{SO}(2), \Gamma_{11}{ }^{3}=a, \Gamma_{13}{ }^{1}=b, \Gamma_{13}{ }^{2}=$ $c, \Gamma_{22}{ }^{3}=a, \Gamma_{23}{ }^{1}=-c, \Gamma_{23}{ }^{2}=b, \Gamma_{33}{ }^{3}=d, \rho_{\Gamma}=\operatorname{diag}(a d, a d, 2(b d-$ $\left.b^{2}+c^{2}\right)$ ). We require $a d \neq 0$ and $b d-b^{2}+c^{2} \neq 0$.
(3) The group $G_{\Gamma}^{+}$is finite, there exists an element of order 3 in $G_{\Gamma}^{+}$, and there exist $(a, b, c, d) \in \mathbb{R}^{4}$ so $\Gamma_{11}{ }^{1}=1, \Gamma_{11}{ }^{3}=a, \Gamma_{12}^{2}=-1, \Gamma_{13}{ }^{1}=b$, $\Gamma_{13}{ }^{2}=c, \Gamma_{22}{ }^{1}=-1, \Gamma_{22}{ }^{3}=a, \Gamma_{23}{ }^{1}=-c, \Gamma_{23}{ }^{2}=b, \Gamma_{33}{ }^{3}=d$, and $\rho_{\Gamma}=(a d-2)\left(e^{1} \otimes e^{1}+e^{2} \otimes e^{2}+2\left(b d-b^{2}+c^{2}\right)\right) e^{3} \otimes e^{3}$. We require $a d-2 \neq 0$ and $-b^{2}+c^{2}+b d \neq 0$. We have $G_{\Gamma}^{+}=\mathbb{Z}_{3}$ except for the following exceptional structures which are given up to isomorphism by:
(a) $a=0, b=0, c=1, d=0$, and $G_{\Gamma}^{+}=s_{3}$.
(b) $c=0, a=b= \pm \frac{1}{\sqrt{2}}, d= \pm \sqrt{2}$, and $G_{\Gamma}^{+}=a_{4}$.
(4) The group $G_{\Gamma}^{+}$is finite and all elements of $G_{\Gamma}^{+}$have order 2. There are two structures up to isomorphism:
(a) $G_{\Gamma}^{+}=\mathbb{Z}_{2} \oplus \mathbb{Z}_{2}, \Gamma_{12}{ }^{3}=1, \Gamma_{13}{ }^{2}=1, \Gamma_{23}{ }^{1}=-1$, and $\rho=-2 e^{1} \otimes e^{1}+2 e^{2} \otimes e^{2}+2 e^{3} \otimes e^{3}$.
(b) $G_{\Gamma}^{+}=\mathbb{Z}_{2}, \Gamma_{i j}{ }^{k}=0$ unless the index 3 appears an odd number of times, $\Gamma_{11}^{3}=a, \Gamma_{12}^{3}=b, \Gamma_{13}^{1}=c, \Gamma_{13}^{2}=d, \Gamma_{21}^{3}=b, \Gamma_{22}^{3}=e$, $\Gamma_{23}{ }^{1}=f, \Gamma_{23}{ }^{2}=g, \Gamma_{31}{ }^{1}=c, \Gamma_{31}{ }^{2}=d, \Gamma_{32}{ }^{1}=f, \Gamma_{32}{ }^{2}=g$, $\Gamma_{33}{ }^{3}=h . \quad \rho_{11}=-2 b d+a(-c+g+h), \quad \rho_{12}=\rho_{21}=-d e-a f+b h$, $\rho_{33}=-c^{2}-2 d f+c h+g(-g+h)$. One requires $\operatorname{det}(\rho) \neq 0$.

### 1.13. The unoriented category

There are similar results in the unoriented category. One does not assume $\mathcal{M}$ is oriented and one replaces the structure group $\mathrm{GL}^{+}(m, \mathbb{R})$ by the full general linear group. Theorems 1.6-1.9 extend to this context with only the appropriate minor modifications of notation. The corresponding analysis of Theorem 1.12 in dimension 3 would become much more complicated and we have not attempted it for that reason nor have we considered torsion in these results for the same reason as our purpose was to be illustrative rather than exhaustive. We have chosen to work in the smooth category; however all the structures in question and the relevant morphisms are in fact real analytic.

## 2. Generic properties

We introduce the following tensors. Let
(3) $\omega:=\Gamma_{i j}{ }^{j} d x^{i}, \quad \rho_{1, \Gamma}=\Gamma_{i n}{ }^{i} \Gamma_{j k}{ }^{n} d x^{j} \otimes d x^{k}$, and $\rho_{2, \Gamma}=\Gamma_{j n}{ }^{i} \Gamma_{i k}{ }^{n} d x^{j} \otimes d x^{k}$.

Note that $\rho=\rho_{1}-\rho_{2}$. If $\rho_{s, \Gamma}$ is non-degenerate, then $\rho_{s, \Gamma}(\varepsilon):=\rho_{s, \Gamma}+\varepsilon \rho_{2, s, \Gamma}$ is invertible for small $\varepsilon$. Let $\varrho_{s, \Gamma}^{i \ell}(\varepsilon)$ be the components of the inverse matrix; this defines the dual symmetric non-degenerate 2 -tensor on $\left(\mathbb{R}^{m}\right)^{*}$. As $\rho_{s, \Gamma}(\varepsilon)$ is real analytic in $\varepsilon$, we sum over repeated indices to expand

$$
\rho_{s, \Gamma}^{i \ell}(\varepsilon) \Gamma_{i j}{ }^{j} e_{\ell}=\sum_{n=0}^{\infty} \xi_{\Gamma, n} \varepsilon^{n} \text { where } \xi_{\Gamma, n} \in \mathbb{R}^{m}
$$

We begin the proof of Theorem 1.6 with the following observation.
Lemma 2.1. There exists a polynomial $\mathfrak{P}_{m}=\mathfrak{P}_{m}(\Gamma)$ and an integer $\kappa_{m}$ so that:
(1) If $\Gamma \in \mathcal{W}(m)$ and if $\mathfrak{P}_{m}(\Gamma) \neq 0$, then
(a) $\rho_{s, \Gamma}$ is non-degenerate,
(b) $\mathcal{B}_{\Gamma}:=\left\{\xi_{\Gamma, 0}, \xi_{\Gamma, 1}, \ldots, \xi_{\Gamma, m-1}\right\}$ is a basis for $\mathbb{R}^{m}$.
(c) $G_{\Gamma}^{+}=\{\mathrm{id}\}$.
(2) If $T \in \mathrm{GL}^{+}(m, \mathbb{R})$ and if $\Gamma \in \mathcal{W}(m)$, then $\mathfrak{P}_{m}(\Gamma)=\operatorname{det}(T)^{\kappa(m)} \mathfrak{P}_{m}(T \Gamma)$.
(3) There exists $\Gamma \in \mathcal{Z}(m)$ so that $\mathfrak{P}_{m}(\Gamma) \neq 0$.

Proof. Clearly $\rho_{s, \Gamma}$ is non-degenerate if and only if $\operatorname{det}\left(\rho_{s, \Gamma}\right) \neq 0$. So we will make $\operatorname{det}\left(\rho_{s, \Gamma}\right)$ a factor of our polynomial to ensure that Assertion (1a) is valid. We apply Cramer's rule. Let $\tilde{\rho}_{s, \Gamma}$ be the matrix of cofactors of $\rho_{s, \Gamma}$; this is a matrix valued polynomial which is well defined for all $\Gamma \in \mathcal{Z}(m)$ such that if $\operatorname{det}\left(\rho_{s, \Gamma}\right) \neq 0$, then $\rho_{s, \Gamma}^{-1}=\operatorname{det}\left(\rho_{s, \Gamma}\right)^{-1} \tilde{\rho}_{s, \Gamma}$. Suppose $\rho_{s, \Gamma}$ is invertible. We use the Neumann series to expand

$$
\begin{aligned}
& \left(\rho_{s, \Gamma}+\varepsilon \rho_{2, s, \Gamma}\right)^{-1}=\left\{\left(\rho_{s, \Gamma}\left(\operatorname{id}+\varepsilon \rho_{s, \Gamma}^{-1} \rho_{2, s, \Gamma}\right)\right\}^{-1}=\left(\operatorname{id}+\varepsilon \rho_{s, \Gamma}^{-1} \rho_{2, s, \Gamma}\right)^{-1} \rho_{s, \Gamma}^{-1}\right. \\
= & \sum_{n=0}^{\infty}(-1)^{n}\left(\rho_{s, \Gamma}^{-1} \rho_{2, s, \Gamma}\right)^{n} \rho_{s, \Gamma}^{-1} \varepsilon^{n}=\sum_{n=0}^{\infty}(-1)^{n} \operatorname{det}\left(\rho_{s, \Gamma}\right)^{-n-1}\left(\tilde{\rho}_{s, \Gamma} \rho_{2, s, \Gamma}\right)^{n} \tilde{\rho}_{s, \Gamma} \varepsilon^{n} .
\end{aligned}
$$

We let $\tilde{\xi}_{n}:=\left(\tilde{\rho}_{s, \Gamma} \rho_{2, s, \Gamma}\right)^{n} \tilde{\rho}_{s, \Gamma} \omega$. We then have

$$
\begin{equation*}
\xi_{\Gamma, n}=(-1)^{n} \operatorname{det}\left(\rho_{s, \Gamma}\right)^{-n-1} \tilde{\xi}_{\Gamma, n} . \tag{4}
\end{equation*}
$$

If $\operatorname{det}\left(\rho_{s, \Gamma}\right) \neq 0$, then $\left\{\xi_{\Gamma, 0}, \ldots, \xi_{\Gamma, n}\right\}$ is a basis for $\mathbb{R}^{m}$ if and only if $\left\{\tilde{\xi}_{\Gamma, 0}, \ldots\right.$, $\left.\tilde{\xi}_{\Gamma, n}\right\}$ is a basis for $\mathbb{R}^{m}$. We define a polynomial $\mathfrak{P}_{m}$ which satisfies (1a) and (1b) by defining:

$$
\mathfrak{P}_{m}(\Gamma):=\operatorname{det}\left(\rho_{s, \Gamma}\right) \operatorname{det}\left(\tilde{\xi}_{0, \Gamma}, \ldots, \tilde{\xi}_{m, \Gamma}\right) .
$$

Since contracting an upper index against a lower index is invariant under the action of $\mathrm{GL}^{+}(m, \mathbb{R})$, Theorem 1.3 shows the tensors of Equation (3) are invariantly defined; if $T \in \mathrm{GL}^{+}(m, \mathbb{R})$ and $\Gamma \in \mathcal{W}(m)$, one has that:

$$
\begin{array}{lll}
\omega(T \Gamma)=T \omega(\Gamma), & \rho_{1}(T \Gamma)=T \rho_{1}(\Gamma), & \rho_{2}(T \Gamma)=T \rho_{2}(\Gamma) \\
\rho_{2, s}(T \Gamma)=T \rho_{2, s}(\Gamma), & T \xi_{\Gamma, i}=\xi_{T \Gamma, i}, & T \mathcal{B}_{\Gamma}=\mathcal{B}_{T \Gamma} . \tag{5}
\end{array}
$$

Let $T \in G_{\Gamma}^{+}$with $\mathfrak{P}_{m}(\Gamma) \neq 0$. Then $\mathcal{B}_{\Gamma}$ is a basis for $\mathbb{R}^{m}$. Since $T \mathcal{B}_{\Gamma}=\mathcal{B}_{T \Gamma}=$ $\mathcal{B}_{\Gamma}, T=\mathrm{id}$. Assertion (1c) now follows.

We may verify Assertion 2 as follows. We have

$$
\begin{equation*}
\operatorname{det}\left(\rho_{s, T \Gamma}\right)=\operatorname{det}\left(T \rho_{s, \Gamma}\right)=\operatorname{det}(T)^{2} \operatorname{det}\left(\rho_{s, \Gamma}\right) . \tag{6}
\end{equation*}
$$

In particular, if $\operatorname{det}\left(\rho_{s, \Gamma}\right)=0$, then $\operatorname{det}\left(\rho_{s, T \Gamma}\right)=0$ and Assertion 2 holds trivially. Suppose $\rho_{s, \Gamma}$ is non-singular. Let $c(m):=(1+2+\cdots+(m-1))+1$ and let $\kappa(m)=2 c(m)+m+2$. Since $T \xi_{\Gamma, i}=\xi_{T \Gamma, i}$, we may verify Assertion 2 by using Equation (4) and Equation (6) to compute:

$$
\begin{aligned}
\mathfrak{P}_{m}(T \Gamma) & =\operatorname{det}\left(\rho_{s, T \Gamma}\right) \operatorname{det}\left(\tilde{\xi}_{T \Gamma, 0}, \ldots, \tilde{\xi}_{T \Gamma, m-1}\right) \\
& =\operatorname{det}\left(\rho_{s, T \Gamma}\right)^{c(m)+1} \operatorname{det}\left(\xi_{T \Gamma, 0}, \ldots, \xi_{T \Gamma, m-1}\right) \\
& =\operatorname{det}(T)^{2 c(m)+2} \operatorname{det}\left(\rho_{s, \Gamma}\right)^{c(m)+1} \operatorname{det}\left(T \xi_{\Gamma, 0}, \ldots, T \xi_{\Gamma, m-1}\right) \\
& =\operatorname{det}(T)^{2 c(m)+m+2} \operatorname{det}\left(\rho_{s, \Gamma}\right)^{c(m)+1} \operatorname{det}\left(\xi_{\Gamma, 0}, \ldots, \xi_{\Gamma, m-1}\right) \\
& =\operatorname{det}(T)^{2 c(m)+m+2} \operatorname{det}\left(\rho_{s, \Gamma}\right) \operatorname{det}\left(\tilde{\xi}_{\Gamma, 0}, \ldots, \tilde{\xi}_{\Gamma, m-1}\right) \\
& =\operatorname{det}(T)^{2 c(m)+m+2} \mathfrak{P}_{m}(\Gamma) .
\end{aligned}
$$

We complete the proof by exhibiting torsion free Christoffel symbols in all dimensions where $\rho_{s, \Gamma}$ is non-degenerate and where $\left\{\Xi_{\Gamma, 0}, \ldots, \Xi_{\Gamma, m-1}\right\}$ are a basis for $\mathbb{R}^{m}$. We proceed by considering various cases.

Case 1. Let $m=2$. We use the parametrization of [8] and define a torsion free tensor by setting $\Gamma(x)_{i j}{ }^{k}$

$$
\begin{array}{llll}
\Gamma(x)_{11}^{1}=x+\frac{1}{x}, & \Gamma(x)_{11}{ }^{2}=0, & \Gamma(x)_{12}^{1}=0 & \Gamma(x)_{21}{ }_{1}^{1}=0, \\
\Gamma(x)_{22}^{1}=x, & \Gamma(x)_{12}^{1}=x, & \Gamma(x)_{21}{ }^{1}=x, & \Gamma(x)_{22}^{2}=1 .
\end{array}
$$

We may then compute:
$\rho_{2, \Gamma}(x)=\left(\begin{array}{cc}2+\frac{1}{x^{2}}+2 x^{2} & x \\ x & 1+2 x^{2}\end{array}\right), \rho_{s, \Gamma}(x)=\left(\begin{array}{ll}1 & 0 \\ 0 & 1\end{array}\right), \omega=\binom{x+\frac{1}{x}}{1}$.

In particular $\rho_{s, \Gamma}$ is non-singular. As $\left(\rho_{s, \Gamma}+\varepsilon \rho_{2, s, \Gamma}\right)^{-1}=\mathrm{id}-\varepsilon \rho_{2, s, \Gamma}+O\left(\varepsilon^{2}\right)$, one has that

$$
\xi_{0}(x)=\binom{x+\frac{1}{x}}{1}, \quad \xi_{1}(x)=\binom{\frac{1}{x^{3}}+\frac{3}{x}+5 x+2 x^{3}}{2+3 x^{2}}
$$

Choose $x$ so $\xi_{0}(x)$ and $\xi_{1}(x)$ are linearly independent to complete the proof if $m=2$.

Case 2. Suppose $m=2 \bar{m}$ for $m \geq 2$. We let $\Gamma(\vec{x}):=\Gamma\left(x_{1}\right) \oplus \cdots \oplus \Gamma\left(x_{\bar{m}}\right)$. The structures decouple;

$$
\begin{aligned}
& \rho_{s, \Gamma}(\vec{x})=\rho_{s, \Gamma}\left(x_{1}\right) \oplus \cdots \oplus \rho_{s, \Gamma}\left(x_{\bar{m}}\right)=\mathrm{id}, \quad \rho_{2}(\vec{x})=\rho_{2}\left(x_{1}\right) \oplus \cdots \oplus \rho_{2}\left(x_{\bar{m}}\right) \\
& \left(\rho_{s, \Gamma}(\vec{x})+\varepsilon \rho_{2}(\vec{x})\right)^{-1}=\sum_{n=0}^{\infty}(-1)^{n} \varepsilon^{n} \rho_{2}(\vec{x})^{n} .
\end{aligned}
$$

Since $\rho_{s, \Gamma}=\mathrm{id}, \rho_{s, \Gamma}$ is non-singular. Let $V:=\operatorname{Span}_{0 \leq n \leq m-1}\left\{\rho_{2}(\vec{x})^{n} \omega(\vec{x})\right\}$. We must show $V=\mathbb{R}^{m}$. As the minimal polynomial of $\rho_{2}(\vec{x})$ has degree at most $m-1$, it is not necessary to truncate by taking $n \leq m-1$ and we have

$$
V=\operatorname{Span}_{0 \leq n}\left\{p_{2}(\vec{x})^{n} \omega(\vec{x})\right\}
$$

We assume $0<x_{1}<\cdots<x_{\bar{m}}$. As $n \rightarrow \infty$, the terms in $x_{\bar{m}}$ will dominate. We examine the final block

$$
\rho_{2}\left(\vec{x}_{\bar{m}}\right)^{n}=\left(2 x_{\bar{m}}^{2}\right)^{n}\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right)^{n}\binom{x_{\bar{m}}}{0}+O\left(x_{\bar{m}}^{2 n}\right) .
$$

The other blocks do not play a role so $\lim _{n \rightarrow \infty}\left\{\left(2 x_{\bar{m}}^{2}\right)^{-n} x_{\bar{m}}^{-1}\right\} \omega=e_{2 m-1}$. Since $V$ is a closed $\mathbb{Z}\left(\rho_{2}(\vec{x})\right)$ module, $e_{2 m-1} \in V$. Examining $\omega-\left(\frac{1}{x_{\bar{m}}}+x_{\bar{m}}\right) e_{2 m-1}$ and applying a similar argument to the last block yields as well $e_{2 m} \in V$. We can now work our way backwards through the blocks to see $V=\mathbb{R}^{m}$ as desired.
Case 3. Suppose $m=3+2 k$ for $k \geq 0$ is odd. Applying exactly the same asymptotic analysis as used in the even dimensional case, we are reduced to considering the case $m=3$. We set

$$
\begin{array}{lll}
\Gamma_{11}^{1}=2, & \Gamma_{22}^{2}=4, & \Gamma_{33}^{3}=2, \\
\Gamma_{11}^{3}=1, & \Gamma_{13}^{1}=1, & \Gamma_{31}^{1}=1, \\
\Gamma_{23}^{2}=1, & \Gamma_{32}{ }^{2}=1, & \Gamma_{22}^{3}=1
\end{array}
$$

We then compute:

$$
\rho_{2}=\left(\begin{array}{ccc}
6 & 0 & 2 \\
0 & 18 & 4 \\
2 & 4 & 6
\end{array}\right), \quad \rho_{s, \Gamma}=\left(\begin{array}{ccc}
2 & 0 & 0 \\
0 & 2 & 0 \\
0 & 0 & 2
\end{array}\right), \quad \omega=\left(\begin{array}{l}
2 \\
4 \\
4
\end{array}\right) .
$$

Let $A:=\frac{1}{2} \rho_{2}$. We then have

$$
\left(\rho_{s, \Gamma}+\varepsilon \rho_{2, s, \Gamma}\right)^{-1}=\frac{1}{2}(\mathrm{id}+A)^{-1}=\frac{1}{2} \sum_{n=0}^{\infty}(-1)^{n} \varepsilon^{n} A^{n} .
$$

We complete the proof by computing:

$$
\Xi_{0}=\left(\begin{array}{l}
1 \\
2 \\
2
\end{array}\right), \Xi_{1}=\left(\begin{array}{c}
5 \\
15 \\
11
\end{array}\right), \Xi_{2}=\left(\begin{array}{c}
26 \\
78 \\
68
\end{array}\right), \operatorname{det}\left(\begin{array}{ccc}
1 & 5 & 26 \\
2 & 15 & 78 \\
2 & 11 & 68
\end{array}\right)=54
$$

The proof of Theorem 1.6. We use Lemma 2.1 to see $\mathrm{GL}^{+}(m, \mathbb{R})$ preserves $\mathcal{W}\left(p, q ; \mathfrak{P}_{m}\right)$ and $\mathcal{Z}\left(p, q ; \mathfrak{P}_{m}\right)$ and that $\mathrm{GL}^{+}(m, \mathbb{R})$ acts without fixed points on these sets. Theorem 1.6 will then follow from Theorem 1.5 if we can show the action is proper.

Given bases $\mathcal{B}$ and $\tilde{\mathcal{B}}$ for $\mathbb{R}^{m}$, let $T_{\mathcal{B}, \tilde{\mathcal{B}}}$ be the unique linear transformation taking $\mathcal{B}$ to $\tilde{\mathcal{B}}$. Let $\Gamma \in \mathcal{Z}\left(p, q ; \mathfrak{P}_{m}\right)$. By Equation (5), $T \mathcal{B}_{\Gamma}=\mathcal{B}_{T \Gamma}$ so $T=$ $T_{\mathcal{B}_{\Gamma}, \mathcal{B}_{T \Gamma}}$. Let $\left\{\Gamma_{k}, \Gamma, \tilde{\Gamma}\right\} \subset \mathcal{W}\left(p, q ; \mathfrak{P}_{m}\right)$ and $T_{k} \in \mathrm{GL}^{+}(m, \mathbb{R})$. Assume that $\Gamma_{k} \rightarrow \Gamma$ and $T \Gamma_{k} \rightarrow \tilde{\Gamma}$. This implies $\mathcal{B}_{\Gamma_{k}} \rightarrow \mathcal{B}_{\Gamma}$ and $\mathcal{B}_{T \Gamma_{k}} \rightarrow \mathcal{B}_{\tilde{\Gamma}}$. Consequently,

$$
T_{k}=T_{\mathcal{B}_{\Gamma_{k}}, \mathcal{B}_{T \Gamma_{k}}} \rightarrow T_{\mathcal{B}_{\Gamma}, \mathcal{B}_{\overparen{\Gamma}}}
$$

## 3. The proof of Theorem 1.7

### 3.1. The proof of Theorem $1.7(1)$

Suppose that

$$
\Gamma_{n} \in \mathcal{W}(p, q) \text { and } \Gamma_{n} \rightarrow \Gamma \in \mathcal{W}(p, q)
$$

Assume $\operatorname{dim}\left\{G_{\Gamma_{n}}^{+}\right\} \geq 1$. We must show $\operatorname{dim}\left\{G_{\Gamma}^{+}\right\} \geq 1$. Since $\operatorname{dim}\left\{G_{\Gamma_{n}}^{+}\right\} \geq 1$, we may find $0 \neq \xi_{n} \in \operatorname{gl}(m, \mathbb{R})$ so that $\exp \left(t \xi_{n}\right)$ defines a 1-parameter subgroup of $G_{\Gamma_{n}}^{+}$. Let $\|\cdot\|$ be a norm on $\operatorname{gl}(m, \mathbb{R})$. We may assume without loss of generality that $\left\|\xi_{n}\right\|=1$ and extract a convergent subsequence $\xi_{n} \rightarrow \xi$ with $\|\xi\|=1$. We then have by continuity that $\exp (t \xi)$ is a 1-parameter subgroup of $G_{\Gamma}^{+}$and thus, in particular, $\operatorname{dim}\left\{G_{\Gamma}^{+}\right\} \geq 1$.

### 3.2. The proof of Theorem $1.7(2)$

Fix $m \geq 3$ and let

$$
\vartheta=\left(\vartheta_{i j k}\right) \text { for } \vartheta \in\{0,1\} \text { and } 1 \leq i, j, k \leq m .
$$

Let $A_{\vartheta}$ be the Abelian group which is generated multiplicatively by indeterminates $\kappa_{1}, \ldots, \kappa_{m}$ subject to the relations $\kappa_{i} \kappa_{j}=\kappa_{k}$ whenever $\vartheta_{i j k}=1$. Let $\operatorname{Tor}\left(A_{\vartheta}\right)$ be the subgroup of $A$ consisting of all elements of finite order. Let

$$
c(m):=\max _{T \in \operatorname{Tor}(A)} \operatorname{order}(T) .
$$

Let $\Gamma \in \mathcal{Z}(p, q)$. Assume that no element of $G_{\Gamma}^{+}$has infinite order. Let $T \in G_{\Gamma}^{+}$. Since $T$ has finite order, there exists a complex basis $\left\{f_{1}, \ldots, f_{m}\right\}$ for $\mathbb{C}$ so that $T f_{i}=\lambda_{i} f_{i}$ where $\left|\lambda_{i}\right|=1$. If we express $\lambda_{i}=e^{\sqrt{-1} \theta_{i}}$, then the $\theta_{i}$ are the rotation angles of $T$ regarded as a real map and the eigenvalues $\lambda_{i}$ occur in conjugate pairs for $\theta_{i} \notin\{0, \pi\}$. We have $T \Gamma_{i j}{ }^{k}=\lambda_{i} \lambda_{j} \lambda_{k}^{-1} \Gamma_{i j}{ }^{k}$. Set $\vartheta_{\Gamma, i j k}=1$ if $\Gamma_{i j}{ }^{k} \neq 0$ and set $\vartheta_{J, i j k}=0$ otherwise. Then $\vec{\lambda}$ can be regarded as an element
of $A_{\vartheta}$. Furthermore, if $A_{\vartheta}$ has an element of infinite order, then there exists $\vec{\lambda}$ where the eigenvalues occur suitably in conjugate pairs so $T_{\lambda} \Gamma=\Gamma$ and $T_{\lambda}$ has infinite order. Since this is false, $A_{\vartheta(\Gamma)}$ is finite and thus order $\left(T_{\lambda}\right) \leq c(m)$ as desired.

To show that $\lim _{m \rightarrow \infty} c(m)=\infty$, we will construct a family of Type $\mathcal{A}$ connections $\Gamma_{3 \ell}$ on $\mathbb{R}^{3 \ell}$ so that there exists an element $T \in G_{\Gamma_{3 \ell}}^{+}$which has order $2^{\ell}-1$ and such that there is no element of infinite order in $G_{\Gamma_{3 \ell}}^{+}$. We shall work in the torsion free setting so there is no need to symmetrize. Recall that

$$
\rho=\rho_{1}-\rho_{2} \text { where } \rho_{1 ; j k}:=\Gamma_{i n}{ }^{i} \Gamma_{j k}{ }^{n} \text { and } \rho_{2 ; j k}:=\Gamma_{j n}{ }^{i} \Gamma_{i k}{ }^{n} .
$$

Central to our construction is a 3 -dimensional example. Let $\left\{e_{1}, e_{2}, e_{3}\right\}$ be the standard basis for $\mathbb{R}^{3}$. Introduce a complex basis

$$
\begin{array}{lll}
f_{1}:=e_{1}+\sqrt{-1} e_{2}, & f_{2}:=e_{1}-\sqrt{-1} e_{2}, & f_{3}:=e_{3}, \\
f^{1}:=\frac{1}{2}\left(e^{1}-\sqrt{-1} e^{2}\right), & f^{2}:=\frac{1}{2}\left(e^{1}+\sqrt{-1} e^{2}\right), & f^{3}:=e^{3}
\end{array}
$$

For $a>0$, let the non-zero Christoffel symbols be given by:

$$
\begin{array}{lll}
\Gamma\left(f_{1}, f_{3}, f^{1}\right)=a, & \Gamma\left(f_{3}, f_{1}, f^{1}\right)=a, & \Gamma\left(f_{2}, f_{3}, f^{2}\right)=a \\
\Gamma\left(f_{3}, f_{2}, f^{2}\right)=a, & \Gamma\left(f_{1}, f_{2}, f^{3}\right)=a, & \Gamma\left(f_{2}, f_{1}, f^{3}\right)=a \\
\Gamma\left(f_{3}, f_{3}, f^{3}\right)=\frac{a^{2}+1}{a} . &
\end{array}
$$

We may then compute:

$$
\begin{aligned}
\rho_{1} & =\left(\begin{array}{ccc}
0 & 3 a^{2}+1 & 0 \\
3 a^{2}+1 & 0 & 0 \\
0 & 0 & 3 a^{2}+\frac{1}{a^{2}}+4
\end{array}\right), \\
\rho_{2} & =\left(\begin{array}{ccc}
0 & 2 a^{2} & 0 \\
2 a^{2} & 0 & 0 \\
0 & 0 & 3 a^{2}+\frac{1}{a^{2}}+2
\end{array}\right), \quad \rho=\left(\begin{array}{ccc}
0 & a^{2}+1 & 0 \\
a^{2}+1 & 0 & 0 \\
0 & 0 & 2
\end{array}\right) .
\end{aligned}
$$

Relative to the underlying real basis $\left\{e_{1}, e_{2}, e_{3}\right\}$ we have:

$$
\begin{aligned}
\rho_{1} & =\left(\begin{array}{ccc}
3 a^{2}+1 & 0 & 0 \\
0 & 3 a^{2}+1 & 0 \\
0 & 0 & 3 a^{2}+\frac{1}{a^{2}}+4
\end{array}\right), \\
\rho_{2} & =\left(\begin{array}{ccc}
2 a^{2} & 0 & 0 \\
0 & 2 a^{2} & 0 \\
0 & 0 & 3 a^{2}+\frac{1}{a^{2}}+2
\end{array}\right), \quad \rho=\left(\begin{array}{ccc}
a^{2}+1 & 0 & 0 \\
0 & a^{2}+1 & 0 \\
0 & 0 & 2
\end{array}\right) .
\end{aligned}
$$

We take a basis $\left\{e_{1, \mu}, e_{2, \mu}, e_{3, \mu}\right\}$ for $\mathbb{R}^{3 \ell}$ where $1 \leq \mu \leq \ell$. Let the non-zero Christoffel symbols be given by:
$\Gamma\left(f_{1, \mu}, f_{3, \mu}, f^{1, \mu}\right)=a_{\mu}, \quad \Gamma\left(f_{3, \mu}, f_{1, \mu}, f^{1, \mu}\right)=a_{\mu}, \Gamma\left(f_{2, \mu}, f_{3, \mu}, f^{2, \mu}\right)=a_{\mu}$,
$\Gamma\left(f_{3, \mu}, f_{2, \mu}, f^{2, \mu}\right)=a_{\mu}, \quad \Gamma\left(f_{1, \mu}, f_{2, \mu}, f^{3, \mu}\right)=a_{\mu}, \Gamma\left(f_{2, \mu}, f_{1, \mu}, f^{3, \mu}\right)=a_{\mu}$,
$\Gamma\left(f_{3, \mu}, f_{3, \mu}, f^{3, \mu}\right)=\frac{1+a_{\mu}^{2}}{a_{\mu}}, \quad \Gamma\left(f_{1, \mu}, f_{1, \mu}, f^{1, \mu+1}\right)=1, \Gamma\left(f_{2, \mu}, f_{2, \mu}, f^{2, \mu+1}\right)=1$.

Here we let $\mu$ be defined modulo $\ell$ so $f_{i, \ell+1}=f_{i, 1}$. This defines corresponding real Christoffel symbols. Let $\vec{\lambda} \in \mathbb{C}^{\ell}$. We assume $\left|\lambda_{\mu}\right|=1$ so $\lambda_{\mu}^{-1}=\bar{\lambda}_{\mu}$. Let

$$
\begin{array}{lll}
T\left(f_{1, \mu}\right)=\lambda_{\mu} f_{1, \mu}, & T\left(f_{2, \mu}\right)=\bar{\lambda}_{\mu} f_{2, \mu}, & T\left(f_{3, \mu}\right)=f_{3, \mu} \\
T\left(f^{1, \mu}\right)=\bar{\lambda}_{\mu} f^{1, \mu}, & T\left(f^{2, \mu}\right)=\lambda_{\mu} f^{2, \mu}, & T\left(f^{3, \mu}\right)=f^{3, \mu}
\end{array}
$$

To ensure that $T \Gamma=\Gamma$, we must have $\lambda_{\mu}^{2}=\lambda_{\mu+1}$. Setting $\lambda_{\ell+1}=\lambda_{1}$ then leads to the relation $\lambda_{1}^{2^{\ell}}=\lambda_{1}$ and thus the cyclic group $\mathbb{Z}_{2^{\ell}-1}$ is a subgroup of $G_{\Gamma}^{+}$; elements of arbitrarily large order can be obtained as $\ell \rightarrow \infty$. We complete the proof by showing there are no elements of infinite order in $G_{\Gamma}^{+}$. If we diagonalize $\rho_{2}$ relative to $\rho$, then the resulting eigenspaces must be preserved by any element $T \in G_{\Gamma}^{+}$. The decomposition is given by

$$
\left\{\left(f_{1, \mu}, \frac{3 a_{\mu}^{2}+1}{2 a_{\mu}^{2}}\right),\left(f_{2, \mu}, \frac{3 a_{\mu}^{2}+1}{2 a_{\mu}^{2}}\right),\left(f_{3, \mu}, \frac{6 a_{\mu}^{4}+a_{\mu}^{2}+4 a_{\mu}^{2}}{2 a_{\mu}^{2}}\right)\right\} .
$$

For suitable choice of the $a_{\mu}$, the eigenvalues

$$
\left\{\frac{3 a_{\mu}^{2}+1}{2 a_{\mu}^{2}}, \frac{6 a_{\mu}^{4}+a_{\mu}^{2}+4 a_{\mu}^{2}}{2 a_{\mu}^{2}}\right\}
$$

will all be distinct. Thus $T$ preserves the spaces $\operatorname{Span}\left\{f_{1, \mu}, f_{2, \mu}\right\}$ and $\operatorname{Span}\left\{f_{3, \mu}\right\}$ individually. Since $\Gamma\left(f_{3, \mu}, f_{3, \mu}, f^{3, \mu}\right) \neq 0$, we have $T f_{3, \mu}=1$. Let $T_{\mu}$ be the restriction of $T$ to $\operatorname{Span}\left\{f_{1, \mu}, f_{2, \mu}\right\}$. Since $T_{\mu}^{2} \in \operatorname{SO}(2)$, we have $T_{\mu}^{2} f_{1, \mu}=\lambda_{\mu}$ and $T_{\mu}^{2} f_{2, \mu}=\bar{\lambda}_{\mu}$ for some $\lambda_{\mu} \in S^{1}$. It now follows that $T_{\mu}^{2}$ has order at most $2^{\ell}-1$ so there are no elements of infinite order in $G_{\Gamma}^{+}$.

## 4. The action of $\operatorname{GL}(m, \mathbb{R})$ on $\mathcal{Z}(p, q)$ and on $\mathcal{W}(p, q)$

Let $\rho_{0}$ be a symmetric bilinear form of signature $(p, q)$. Let

$$
\mathrm{SO}\left(\rho_{0}\right):=\left\{T \in \mathrm{GL}^{+}(m, \mathbb{R}): T^{*} \rho_{0}=\rho_{0}\right\}
$$

The following is a quite general remark.
Lemma 4.1. Let $\mathcal{O}$ be a $\mathrm{GL}^{+}(m, \mathbb{R})$ invariant subset of $\mathcal{Z}(p, q)$ or of $\mathcal{W}(p, q)$. If action of $\mathrm{SO}\left(\rho_{0}\right)$ on $\mathcal{O}$ is proper, then the action of $\mathrm{GL}^{+}(m, \mathbb{R})$ on $\mathcal{O}$ is proper.
Proof. Assume that the action of $\mathcal{S O}\left(\rho_{0}\right)$ on $\mathcal{O}$ is proper. Suppose given $\Gamma_{n} \in \mathcal{O}$ and $T_{n} \in \mathrm{GL}^{+}(m, \mathbb{R})$ which satisfy

$$
\Gamma_{n} \rightarrow \Gamma \in \mathcal{O} \text { and } \tilde{\Gamma}_{n}:=T_{n} \Gamma_{n} \rightarrow \tilde{\Gamma} \in \mathcal{O}
$$

We must extract a convergent sequence of the $\left\{T_{n}\right\}$. Make a change of basis to suppose that $\rho_{0}=\operatorname{diag}(-1, \ldots,-1,+1, \ldots,+1)$ relative to the standard basis $\mathcal{B}$ for $\mathbb{R}$. Choose $S \in \mathrm{GL}^{+}(m, \mathbb{R})$ so that $S \rho_{\tilde{\Gamma}}=\rho_{0}$. Then $S T_{n} \Gamma_{n} \rightarrow S \tilde{\Gamma}$. Extracting a convergent subsequence from $S T_{n}$ is equivalent to extracting a convergent subsequence from $T_{n}$. Thus we may assume without loss of generality that $\rho_{\tilde{\Gamma}}=\rho_{0}$. Since $T_{n} \Gamma_{n} \rightarrow \tilde{\Gamma}, \rho_{T_{n} \Gamma_{n}} \rightarrow \rho_{0}$. We may apply the

Gram-Schmidt process to the standard basis $\mathcal{B}$ construct a basis $\mathcal{B}_{n}$ for $\mathbb{R}^{m}$ which is an orthonormal basis for $\rho_{T_{n} \Gamma_{n}}$; since $\rho_{T_{n} \Gamma_{n}} \rightarrow \rho_{\tilde{\Gamma}}$, the Gram-Schmidt process does not fail, i.e., we are not trying to normalize a null vector at some stage. Thus the Gram-Schmidt process yields a sequence $S_{n} \in \mathrm{GL}^{+}(m, \mathbb{R})$ so that $S_{n} \rightarrow$ id and so $S_{n} \rho_{T_{n} \Gamma_{n}}=\rho_{S_{n} T_{n} \Gamma_{n}}=\rho_{0}$. Again, extracting a convergent subsequence from $S_{n} T_{n}$ is equivalent to extracting a convergent sequence from $T_{n}$ and hence we may assume without loss of generality that $\rho_{T_{n} \Gamma_{n}}=\rho_{0}$ for $n$ sufficiently large. We have $\Gamma_{n}=T_{n}^{-1} \tilde{\Gamma}_{n} \rightarrow \Gamma$. Extracting a convergent subsequence from $\left\{T_{n}\right\}$ is equivalent to extracting a convergent subsequence from $\left\{T_{n}^{-1}\right\}$. Thus we may interchange the roles of $\left\{\Gamma_{n}, \Gamma\right\}$ and $\left\{\tilde{\Gamma}_{n}, \tilde{\Gamma}\right\}$ and apply the argument given above to assume without loss of generality that $\rho_{\Gamma}=\rho_{0}$ and $\rho_{\Gamma_{n}}=\rho_{0}$ as well. But since $\rho_{0}=\rho_{T_{n} \Gamma_{n}}=T_{n} \rho_{\Gamma_{n}}=T_{n} \rho_{0}, T_{n} \in \operatorname{SO}\left(\rho_{0}\right)$. By hypothesis, as desired, we can extract a convergent sequence.

The proof of Theorem 1.8(1). Suppose that $(p, q) \in\{(m, 0),(0, m)\}$. Then $\rho_{0}$ is definite and hence $\mathrm{SO}\left(\rho_{0}\right)$ is compact. Thus any sequence of elements $T_{n}$ in $\mathrm{SO}\left(\rho_{0}\right)$ has a convergent subsequence so the action of $\mathrm{SO}\left(\rho_{0}\right)$ on $\mathcal{W}(p, q)$ or on $\mathcal{Z}(p, q)$ is proper. Thus by Lemma 4.1, the same is true of the action by $\mathrm{GL}^{+}(m, \mathbb{R})$.

The proof of Theorem 1.8(2). Let $\Gamma_{n} \in \mathcal{W}(p, q)$ with $\Gamma_{n} \rightarrow \Gamma \in \mathcal{W}(p, q)$ and with $G_{\Gamma_{n}}^{+} \neq\{\mathrm{id}\}$. We wish to show $G_{\Gamma}^{+} \neq\{\mathrm{id}\}$. This will show that the Christoffel symbols with non-trivial isotropy subgroup form a closed set and correspondingly that the Christoffel symbols with trivial isotropy subgroup form an open set.

Choose id $\neq T_{n} \in G_{\Gamma_{n}}^{+}$. Since the action of $\mathrm{GL}^{+}(m, \mathbb{R})$ is proper, we can choose a convergent subsequence $T_{n_{k}} \rightarrow T$. We must guard against the possibility that $T=\mathrm{id}$. Let exp be the exponential map from the Lie algebra $\mathfrak{s o}\left(\rho_{0}\right)$ to $\mathrm{SO}\left(\rho_{0}\right)$. Put a Euclidean metric on $\mathfrak{s o}\left(\rho_{0}\right)$. There exists $\varepsilon>0$ so that exp is a diffeomorphism from the open ball $B_{3 \varepsilon}(0)$ of radius $3 \varepsilon$ in $\mathfrak{s o}\left(\rho_{0}\right)$ to a neighborhood of id in $\operatorname{SO}\left(\rho_{0}\right)$. Let $\mathcal{O}_{k}:=\exp \left(B_{k \varepsilon}(0)\right)$ for $k=1,2$. If $T \in \mathcal{O}_{1}$ with $T \neq \mathrm{id}$, then $T=\exp _{P}(\xi)$ for $0<\|\xi\|<\varepsilon$. Choose $k(T) \in \mathbb{N}$ so that $\varepsilon<k(T)\|\xi\|<2 \varepsilon$. We then have $T^{k(T)} \in \mathcal{O}_{2}-\mathcal{O}_{1}$. We return to our sequence id $\neq T_{n} \in G_{\Gamma_{n}} \subset \mathrm{SO}\left(\rho_{0}\right)$. By replacing $T_{n_{k}}$ by $T_{n_{k}}^{k\left(T_{n}\right)}$, we can assume additionally that $T_{n_{k}} \in \mathcal{O}_{1}^{c}$ and thus $T \in \mathcal{O}_{\varepsilon}^{c}$ so $T \neq \mathrm{id}$. Since by continuity $T \Gamma=\Gamma$, we conclude as desired that $G_{\Gamma}^{+} \neq \mathrm{id}$.

The proof of Theorem 1.8(3). These Assertions follow from Theorem 1.5 and from Assertions (1,2).

## 5. The proof Theorem 1.9

Let $(p, q)$ be given with $m=p+q \geq 3, p \geq 1$, and $q \geq 1$. We must show that there exists $\Gamma \in \mathcal{Z}(p, q)$ so that $G_{\Gamma}^{+}$is non-compact. It then follows that the action of $\mathrm{GL}^{+}(2, \mathbb{R})$ on $\mathcal{Z}(p, q)$ is not proper.

Suppose first $m=3$. We may work in the torsion free setting. We consider the structure of Theorem 1.12(1) and set $\Gamma_{12}{ }^{3}=\Gamma_{21}{ }^{3}=a, \Gamma_{13}{ }^{1}=\Gamma_{31}{ }^{1}=b$, $\Gamma_{23}{ }^{2}=\Gamma_{32}{ }^{2}=c, \Gamma_{33}{ }^{3}=d$. We may then compute that

$$
\rho_{\Gamma}=\left(\begin{array}{ccc}
0 & a d & 0 \\
a d & 0 & 0 \\
0 & 0 & -b^{2}+b d+c(-c+d)
\end{array}\right)
$$

By adjusting the parameters $\{a, b, c, d\}$ suitably, we can obtain either signature $(1,2)$ or signature $(2,1)$. Let $T_{\alpha} e_{1}=\alpha e_{1}, T_{\alpha} e_{2}=\alpha^{-1} e_{2}, T_{\alpha} e_{3}=e_{3}$. This gives a Lie group isomorphic to $\operatorname{SO}(1,1)$. We verify that $T_{\alpha} \Gamma=\Gamma$ for any $\alpha$. The sequence $T_{n}$ obtained by taking $\alpha=n$ then satisfies $\left\|T_{n} e_{1}\right\| \rightarrow \infty$. Consequently no subsequence of this sequence converges in $\mathrm{GL}^{+}(2, \mathbb{R})$. Therefore, $G_{\Gamma}^{+}$is non compact and the action of $\mathrm{GL}^{+}(2, \mathbb{R})$ is not proper.

If $m>3$, extend the structure considered above by adding the (possibly) non-zero Christoffel symbols $\Gamma_{u v}{ }^{3}=\Gamma_{v u}{ }^{3}=\varepsilon_{u v}$ for $4 \leq u \leq v \leq m$ where $\varepsilon_{u v}$ are to be determined. Recall from Equation (2) that $\rho_{j k}=\Gamma_{i n}{ }^{i} \Gamma_{j k}{ }^{n}-\Gamma_{j n}{ }^{i} \Gamma_{i k}{ }^{n}$. The new Christoffel symbols involving $\varepsilon_{u v}$ make no contribution to $\rho_{j k}$ for indices $1 \leq j \leq k \leq 3$ and only contribute to $\Gamma_{33}{ }^{3} \Gamma_{j k}{ }^{3}$ for $4 \leq j, k \leq m$. Thus $\rho_{\Gamma, a b}=d \varepsilon_{a b}$ for $4 \leq a \leq b \leq m$. So by adding in these terms, we can obtain any indefinite signature in dimension at least 3 .

## 6. The two dimensional setting. The proof of Theorem 1.11

Let $\Gamma_{2}$ be the structure of Definition 1.10. Let $\operatorname{SO}(2)=\left\{T_{\theta}: 0 \leq \theta<2 \pi\right\}$ and $\operatorname{SO}(1,1)=\left\{\tilde{T}_{a}: a \neq 0\right\}$ where

$$
T_{\theta}:=\left(\begin{array}{rr}
\cos (\theta) & \sin (\theta)  \tag{7}\\
-\sin (\theta) & \cos (\theta)
\end{array}\right) \text { and } \tilde{T}_{a}:=\left(\begin{array}{ll}
a & 0 \\
0 & a^{-1}
\end{array}\right)
$$

Suppose $\Gamma \in \mathcal{Z}(1,1)$ so the structure group is $\mathrm{SO}(1,1)$. Then

$$
\begin{equation*}
\tilde{T}_{a}^{*} \Gamma_{i j}^{k}=a^{\varepsilon} \text { for } \varepsilon= \pm 1+ \pm 1+ \pm 1 \in\{ \pm 1, \pm 3\} \tag{8}
\end{equation*}
$$

Suppose $\tilde{T}_{a}^{*} \Gamma=\Gamma$. Choose $\Gamma_{i j}{ }^{k} \neq 0$. Then $\tilde{T}_{a}^{*} \Gamma_{i j}{ }^{k}=\Gamma_{i j}{ }^{k}$ implies $a=1$ and $\tilde{T}_{a}=\mathrm{id}$. Suppose $\Gamma_{n} \rightarrow \Gamma$ and $g_{n} \Gamma_{n} \rightarrow \tilde{\Gamma}$. We apply the argument of Lemma 4.1 to see that we may suppose

$$
\rho_{s, \Gamma}=\left(\begin{array}{cc}
0 & 1 \\
1 & 0
\end{array}\right) \text { and } g_{n}=\tilde{T}_{a_{n}} \in \mathrm{SO}(1,1)
$$

We use Equation (8) to see that the $\alpha_{n}$ must converge. Assertion 1 now follows.
Assume $\rho_{\Gamma}$ is definite so the structure group is $\mathrm{SO}(2)$. Assume $G_{\Gamma}$ is non trivial. Let id $\neq T_{\theta}^{*} \Gamma \in G_{\Gamma}^{+}$. We complexify. Let $\left\{e_{1}, e_{2}\right\}$ be the standard basis for $\mathbb{R}^{2}$. Set $z:=e_{1}+\sqrt{-1} e_{2}$. Then $\{z, \bar{z}\}$ is a $\mathbb{C}$ basis for $\mathbb{R}^{2} \otimes_{\mathbb{R}}$ $\mathbb{C}$. We extend $\Gamma$ to be complex linear to define the corresponding complex Christoffel symbols $\left\{\Gamma_{z z^{z}}, \Gamma_{z \bar{z}}{ }^{z}, \Gamma_{\bar{z} z}{ }^{z}\right\}$. Since the underlying structure is real, the remaining symbols are determined:

$$
\bar{\Gamma}_{z z}{ }^{z}=\Gamma_{\bar{z} \bar{z}} \bar{z}, \quad \bar{\Gamma}_{z \bar{z}}^{z}=\Gamma_{z \bar{z}} \bar{z}^{z}, \quad \bar{\Gamma}_{\bar{z} \bar{z}}^{z}=\Gamma_{z z} \bar{z} .
$$

Let $\alpha:=e^{\sqrt{-1} \theta}$. Since $T_{\theta} e_{1}=\cos (\theta) e_{1}-\sin (\theta) e_{2}$ and $T_{\theta} e_{2}=\sin (\theta) e_{1}+\cos (\theta) e_{1}$,

$$
\begin{aligned}
T_{\theta} z & =(\cos (\theta)+\sqrt{-1} \sin (\theta)) e_{1}+(-\sin (\theta)+\sqrt{-1} \cos (\theta)) e_{2} \\
& =(\cos (\theta)+\sqrt{-1} \sin (\theta))\left(e_{1}+\sqrt{-1} e_{2}\right)=\alpha z .
\end{aligned}
$$

Since $T_{\theta} z=\alpha z$, we have dually that $T_{\theta} z^{*}=\bar{\alpha} z^{*}$. Consequently when we raise indices, we have:

$$
\left(T_{\theta}^{*} \Gamma\right)_{z z}{ }^{z}=\alpha \alpha \bar{\alpha} \Gamma_{z z}{ }^{z}, \quad\left(T_{\theta}^{*} \Gamma\right)_{z \bar{z}}^{z}=\alpha \bar{\alpha} \bar{\alpha} \Gamma_{z \bar{z}}{ }^{z}, \quad\left(T_{\theta}^{*} \Gamma\right)_{\bar{z} \bar{z}}{ }^{z}=\bar{\alpha} \bar{\alpha} \bar{\alpha} \Gamma_{\bar{z} \bar{z}}{ }^{z} .
$$

We have assumed that $T_{\theta} \neq \mathrm{id}$. If $\Gamma_{z z}{ }^{z} \neq 0$, then $\alpha=1$. Similarly, if $\Gamma_{z \bar{z}}{ }^{z} \neq 0$, then $\bar{\alpha}=1$. So the only possibility left is that $\Gamma_{\bar{z} \bar{z}} z^{z} \neq 0$ in which case $\alpha^{3}=\mathrm{id}$ and $\theta=\frac{2 \pi}{3}$ or $\theta=\frac{4 \pi}{3}$. This implies $G_{\Gamma}=\mathbb{Z}_{3}$. By making a coordinate rotation and then rescaling, we may assume $\Gamma_{\bar{z} \bar{z}} z=2 \sqrt{2}$. We have:

$$
\begin{aligned}
0 & =\Gamma_{z z}^{z}=\left\{\Gamma_{11}^{1}+2 \Gamma_{12}^{2}-\Gamma_{22}^{1}\right\}+\sqrt{-1}\left\{-\Gamma_{11}^{2}+2 \Gamma_{12}^{1}+\Gamma_{22}^{2}\right\}, \\
0 & =\Gamma_{z \bar{z}}^{z}=\left\{\Gamma_{11}^{1}+\Gamma_{22}^{1}\right\} \quad+\sqrt{-1}\left\{-\Gamma_{11}^{2}-\Gamma_{22}^{2}\right\}, \\
2 \sqrt{2} & =\Gamma_{\bar{z} \bar{z}}{ }^{z}=\left\{\Gamma_{11}^{1}-2 \Gamma_{12}^{2}-\Gamma_{22}^{1}\right\}+\sqrt{-1}\left\{-\Gamma_{11}^{2}-2 \Gamma_{12}^{1}+\Gamma_{22}^{2}\right\} .
\end{aligned}
$$

We solve these equations to obtain the structure $\Gamma_{2}$. The remainder of the argument is similar to that given in the indefinite setting and is therefore omitted.

Remark 6.1. We normalized the coordinates so that $\Gamma_{\bar{z} \bar{z}}{ }^{z}=2 \sqrt{2}$. Subsequently, in the proof of Theorem $1.12(3)$, we shall deal with the full orbit $\mathrm{GL}^{+}(2, \mathbb{R}) \Gamma_{2}$ and will not adopt this normalization.

## 7. The proof of Theorem 1.12

Let $\mathcal{M}=\left(\mathbb{R}^{3}, \Gamma\right) \in \mathcal{Z}(p, q)$ for $p+q=3$. We suppose id $\neq T \in G_{\Gamma}^{+}$. In Section 7.1, we will show there is an axis of rotation $\xi$ which is not a null vector so $T \xi= \pm \xi$. It then follows that $T$ preserves $\xi^{\perp}$ so the problem becomes, in a certain sense, 2-dimensional. In Section 7.2, we show that if $T$ is an element of order at least 4 , then $\operatorname{dim}\left\{G_{\Gamma}^{+}\right\} \geq 1$; this focuses attention on the elements of order 2 and order 3 and in Section 7.3, we establish a technical result for elements of order 2 and 3. We use these results in Sections 7.4-7.7 to complete the proof of Theorem 1.12.

### 7.1. The axis of rotation

We begin our study with the following result:
Lemma 7.1. Let $\mathcal{M}=\left(\mathbb{R}^{3}, \Gamma\right) \in \mathcal{Z}(p, q)$ where $p+q=3$. If $T \in G_{\Gamma}^{+}$, then there exists a non-null vector $\xi$ so $T \xi=\xi$.

Proof. Let $\rho:=\rho_{s, \Gamma}$ provide a non-degenerate symmetric bilinear inner-product on $\mathbb{R}^{3}$. Let $T \in G_{\Gamma}^{+} \subset \mathrm{SO}(\rho)$. The characteristic polynomial of $T$ is a cubic polynomial. Since every cubic polynomial has a real root, there exists $e_{1} \neq 0$ so $T e_{1}=a e_{1}$ for $a \neq 0$. Our first task is to show that we can choose an eigenvector
which is not a null vector. Suppose, to the contrary, that $\rho\left(e_{1}, e_{1}\right)=0$. Choose $e_{3}$ so that $\rho\left(e_{1}, e_{3}\right)=1$. By subtracting an appropriate multiple of $e_{1}$, we can assume $\rho\left(e_{3}, e_{3}\right)=0$. Choose $e_{2} \in \operatorname{Span}\left\{e_{1}, e_{3}\right\}^{\perp}$. Normalize $e_{2}$ so $\rho\left(e_{2}, e_{2}\right)=$ $\pm 1$. Express

$$
T e_{1}=a e_{1}, \quad T e_{2}=t_{21} e_{1}+t_{22} e_{2}+t_{23} e_{3}, \quad T e_{3}=t_{31} e_{1}+t_{32} e_{2}+t_{33} e_{3}
$$

As $\rho\left(T e_{1}, T e_{2}\right)=\rho\left(e_{1}, e_{2}\right)=0, t_{23}=0$. As $\rho\left(T e_{1}, T e_{3}\right)=\rho\left(e_{1}, e_{3}\right)=1$, $t_{33}=a^{-1}$. Consequently,

$$
T e_{1}=a e_{1}, \quad T e_{2}=t_{21} e_{1}+t_{22} e_{2}, \quad T e_{3}=t_{31} e_{1}+t_{32} e_{2}+a^{-1} e_{3}
$$

We have $\operatorname{det}(T)=t_{22}=1$. Thus $t_{22}=1$. This shows that the eigenvalues of $T$ are $\left\{1, a, a^{-1}\right\}$. Consequently, we could have chosen $e_{1}$ in the first instance so $T e_{1}=1$ and we may therefore assume $a=1$ and consequently $a$ is the only eigenvalue of $T$. This implies that

$$
T e_{1}=e_{1}, \quad T e_{2}=t_{21} e_{1}+e_{2}, \quad T e_{3}=t_{31} e_{1}+t_{32} e_{1}+e_{3}
$$

If $t_{21}=0$, then we may take $\xi=e_{2}$ to establish the desired result. Thus $t_{21}=b \neq 0$. Because $\rho\left(T e_{2}, T e_{3}\right)=0$ and $\rho\left(T e_{3}, T e_{3}\right)=0$,

$$
T e_{1}=e_{1}, \quad T e_{2}=b e_{1}+e_{2}, \quad T e_{3}=\frac{1}{2} b^{2} \rho\left(e_{2}, e_{2}\right) e_{1}-b e_{2}+e_{3} \text { for } b \neq 0
$$

Note that $T^{*} \Gamma$ is a polynomial in $b$. If we replace $T$ by $T^{n}$, we replace $b$ by $n b$. Thus if $\Gamma_{i j}{ }^{k} \neq 0$, all the coefficients of $b^{k}$ must vanish in $\left(T^{*} \Gamma\right)_{i j}{ }^{k}$. We linearize the problem and work modulo terms which are quadratic and of higher order in $b$ and concentrate on the relations provided by the linear terms. Let $\left\{e^{1}, e^{2}, e^{3}\right\}$ be the dual basis. Expand:

$$
\begin{array}{lll}
T e_{1} \equiv e_{1}, & T e_{2} \equiv b e_{1}+e_{2}, & T e_{3} \equiv-b e_{2}+e_{3}, \\
T e^{1} \equiv e^{1}-b e^{2}, & T e^{2} \equiv e^{2}+b e^{3}, & T e^{3} \equiv e^{3}, \\
T^{*} \Gamma_{12}{ }^{3} \equiv b \Gamma_{11}{ }^{3}+\Gamma_{12}{ }^{3}, & T^{*} \Gamma_{13}{ }^{3} \equiv-b \Gamma_{12}{ }^{3}+\Gamma_{13}{ }^{3} \\
T^{*} \Gamma_{23}{ }^{3} \equiv b \Gamma_{13}{ }^{3}-b \Gamma_{22}{ }^{3}+\Gamma_{23}{ }^{3}, & T^{*} \Gamma_{33}{ }^{3} \equiv-2 b \Gamma_{23}{ }^{3}+\Gamma_{33}{ }^{3} .
\end{array}
$$

We set the terms involving $b$ to zero to see:

$$
\Gamma_{11}^{3}=0, \quad \Gamma_{12}^{3}=0, \quad \Gamma_{13}^{3}=c_{1}, \quad \Gamma_{22}^{3}=c_{1}, \quad \Gamma_{23}^{3}=0 .
$$

We continue the expansion

$$
\begin{array}{ll}
T^{*} \Gamma_{12}^{2} \equiv b \Gamma_{11}^{2}+b \Gamma_{12}^{3}+\Gamma_{12}^{2}, & T^{*} \Gamma_{13}^{2} \equiv b \Gamma_{13}^{3}-b \Gamma_{12}^{2}+\Gamma_{13}^{2} \\
T^{*} \Gamma_{22}^{2} \equiv 2 b \Gamma_{12}^{2}+b \Gamma_{22}^{3}+\Gamma_{22}^{2}, & T^{*} \Gamma_{23}^{2} \equiv b \Gamma_{13}^{2}-b \Gamma_{22}^{2}+b \Gamma_{23}^{3}+\Gamma_{23}^{2}, \\
T^{*} \Gamma_{33}^{2} \equiv-2 b \Gamma_{23}^{2}+b \Gamma_{33}{ }^{3}+\Gamma_{33}{ }^{2} .
\end{array}
$$

We set the terms involving $b$ to zero. We use the previous relations to $c_{1}=0$ and

$$
\begin{array}{lllll}
\Gamma_{11}^{2}=0, & \Gamma_{12}^{2}=0, & \Gamma_{22}^{2}=c_{2}, & \Gamma_{13}^{2}=c_{2}, & \Gamma_{23}^{2}=c_{3}, \\
\Gamma_{11}{ }^{3}=0, & \Gamma_{12}{ }^{3}=0, & \Gamma_{13}^{3}=0, & \Gamma_{22}^{3}=0, & \Gamma_{23}^{3}=0 \\
\Gamma_{33}^{3}=2 c_{3} . & & &
\end{array}
$$

We continue the computation:

$$
\begin{array}{ll}
T^{*} \Gamma_{12}^{1} \equiv b \Gamma_{11}{ }^{1}+\Gamma_{12}{ }^{1}, & \\
T^{*} \Gamma_{13}{ }^{1} \equiv-b \Gamma_{12}^{1}-b \Gamma_{13}{ }^{2}+\Gamma_{13}{ }^{1}, & T^{*} \Gamma_{22}{ }^{1} \equiv 2 b \Gamma_{12}^{1}-b \Gamma_{22}^{2}+\Gamma_{22}{ }^{1} \\
T^{*} \Gamma_{23}{ }^{1} \equiv b \Gamma_{13}{ }^{1}-b \Gamma_{22}{ }^{1}-b \Gamma_{23}{ }^{2}+\Gamma_{23}{ }^{1}, & T^{*} \Gamma_{33}{ }^{1} \equiv-2 b \Gamma_{23}{ }^{1}-b \Gamma_{33}{ }^{2}+\Gamma_{33}{ }^{1}
\end{array}
$$

We set the terms involving $b$ to zero and use the previous relations to see $c_{2}=0$ and obtain:

$$
\begin{array}{lllll}
\Gamma_{11}^{1}=0, & \Gamma_{12}^{1}=0, & \Gamma_{22}^{1}=c_{5}, & \Gamma_{13}{ }^{1}=c_{3}+c_{5}, & \Gamma_{23}^{1}=c_{4}, \\
\Gamma_{11}^{2}=0, & \Gamma_{12}^{2}=0, & \Gamma_{22}^{2}=0, & \Gamma_{13}{ }^{2}=0, & \Gamma_{23}{ }^{2}=c_{3}, \\
\Gamma_{11}^{3}=0, & \Gamma_{12}^{3}=0, & \Gamma_{22}^{3}=0, & \Gamma_{13}{ }^{3}=0, & \Gamma_{23}{ }^{3}=0, \\
\Gamma_{33}{ }^{1}=c_{6}, & \Gamma_{33}^{2}=-2 c_{4}, & \Gamma_{33}^{3}=2 c_{3} . & &
\end{array}
$$

By Equation (2), $\rho_{j k}=\Gamma_{i n}{ }^{i} \Gamma_{j k}{ }^{n}-\Gamma_{j n}{ }^{i} \Gamma_{i k}{ }^{n}$. Consequently

$$
\rho_{k 1}=\rho_{1 k}=\Gamma_{i n}{ }^{i} \Gamma_{1 k}{ }^{n}-\Gamma_{1 n}{ }^{i} \Gamma_{i k}{ }^{n}=\delta_{k, 3} \Gamma_{i 1}{ }^{i} \Gamma_{13}{ }^{1}-\Gamma_{13}{ }^{1} \Gamma_{1 k}{ }^{3}=0-0=0 .
$$

This shows that $\rho$ is singular which is false.
Thus we can choose an eigenvector which is not a null vector. Choose $\xi$ so $T \xi=a \xi$ and $\rho(\xi, \xi) \neq 0$. Since $\rho(\xi, \xi)=\rho(T \xi, T \xi)=a^{2} \rho(\xi, \xi)$, we conclude $a^{2}=1$. Suppose that $a \neq 1$ so $a=-1$. Let $V=\xi^{\perp}$. Then $\rho_{V}:=\left.\rho\right|_{V}$ is non-degenerate. Furthermore, $T \xi^{\perp}=\xi^{\perp}$. Let $T_{V}:=\left.T\right|_{V}$. Since $T \xi=-\xi$ and $\operatorname{det}(T)=1, \operatorname{det}\left(T_{V}\right)=-1$. The characteristic polynomial $T_{V}$ takes the form $\lambda^{2}+c_{1} \lambda+\operatorname{det}\left(T_{V}\right)=0$. Since $\operatorname{det}\left(T_{V}\right)=-1$, there are 2 real eigenvalues of $T_{V}$ and $T_{V}$ is diagonal. Thus we can choose a basis $\{\eta, \sigma\}$ for $V$ so $T \eta=\lambda_{1} \eta$ and $T \sigma=\lambda_{2} \sigma$ where $\lambda_{1} \lambda_{2}=-1$. Since $\rho(\eta \sigma)=\rho(T \eta, T \sigma)=\lambda_{1} \lambda_{2} \rho(\eta, \sigma)=$ $-\rho(\eta, \sigma)$, we have $\rho(\eta, \sigma)=0$. Since $\rho_{V}$ is non-degenerate, neither $\eta$ nor $\sigma$ is null. Thus $\lambda_{1}^{2}=\lambda_{2}^{2}=1$. Since $\lambda_{1} \lambda_{2}=-1$, there must exist a +1 eigenvector of $T$ which is not null.

### 7.2. Elements of order at least 4

The following subgroups of $\mathrm{GL}^{+}(3, \mathbb{R})$ will play a central role in what follows. We generalize Equation (7) to the 3 dimensional setting to define:

$$
\begin{aligned}
& \mathrm{SO}(2):=\left\{T_{\theta}:=\left(\begin{array}{rrr}
\cos (\theta) & \sin (\theta) & 0 \\
-\sin (\theta) & \cos (\theta) & 0 \\
0 & 0 & 1
\end{array}\right) \text { for } 0 \leq \theta \leq 2 \pi\right\}, \\
& \mathrm{SO}(1,1):=\left\{\tilde{T}_{a}:=\left(\begin{array}{lll}
a & 0 & 0 \\
0 & a^{-1} & 0 \\
0 & 0 & 1
\end{array}\right) \text { for } a \in \mathbb{R}-\{0\}\right\}
\end{aligned}
$$

Lemma 7.2. Let $T \in G_{\Gamma}^{+}$for $\Gamma \in \mathcal{Z}(p, q)$ with $p+q=3$. If $T$ has order at least 4, then $\operatorname{dim}\left\{G_{\Gamma}^{+}\right\}=1$ and after making a suitable choice of basis we have either that $\mathrm{SO}(2) \subset G_{\Gamma}^{+}$or that $\mathrm{SO}(1,1) \subset G_{\Gamma}^{+}$.

Proof. Choose a unit vector $e_{3}$ so $T e_{3}=e_{3}$. Let $V=e_{3}^{\perp}$ and let $\rho_{V}:=\left.\rho\right|_{V}$; $\rho_{V}$ is non-degenerate. Furthermore, $T$ preserves $V$. Let $T_{V}:=\left.T\right|_{V}$. Since $\operatorname{det}(T)=1$ and $T e_{3}=e_{3}, T_{V} \in \mathrm{SO}\left(\rho_{V}\right)$. We apply the same argument as that used to establish Lemma 1.11.
Case 1. Suppose $\rho_{V}$ is indefinite. Choose a hyperbolic basis $\left\{e^{1}, e^{2}\right\}$ for $V$ so $\rho_{V}=e^{1} \otimes e^{2}+e^{2} \otimes e_{1}$. Since $T_{V} \in \operatorname{SO}\left(\rho_{V}\right)$, there exists $a$ so $T=T_{a}$ takes the form $T e_{1}=a e_{1}$ and $T e_{2}=a^{-1} e_{2}$. Since $T$ has order at least $4, a \neq 1$. We compute

$$
T \Gamma_{i j}^{k}=a^{\epsilon(i j k)} \Gamma_{i j}^{k} \text { for } \epsilon(i j k)=\delta_{1 i}-\delta_{2 i}+\delta_{1 j}-\delta_{1 k}+\delta_{2 k} .
$$

Since $\epsilon(i j k) \in(0, \pm 1, \pm 2, \pm 3)$ and $a \neq \pm 1$, we conclude $\Gamma_{i j}{ }^{k}=0$ for $\epsilon(i j k) \neq 0$. But this implies that $T_{b} \in G_{\Gamma}^{+}$for any $b$ and hence $\mathrm{SO}(1,1) \subset G_{\Gamma}^{+}$.
Case 2. Suppose that $\rho_{V}$ is indefinite. We complexify and set

$$
\begin{array}{lll}
f_{1}:=e_{1}+\sqrt{-1} e_{2}, & f_{2}:=e_{1}-\sqrt{-1} e_{2}, & f_{3}:=e_{3}, \\
f^{1}:=\frac{1}{2}\left(e^{1}-\sqrt{-1} e^{2}\right), & f^{2}:=\frac{1}{2}\left(e^{1}+\sqrt{-1} e^{2}\right), & f^{3}:=e^{3}
\end{array}
$$

Since $T \in \mathrm{SO}(2)$, the complex eigenvalues of $T$ are $\left\{\alpha, \bar{\alpha}=\alpha^{-1}\right\}$. Since we can diagonalize $T$ over $\mathbb{C}$, we may choose the notation so $T f_{1}=\alpha f_{1}$ and $T f_{2}=\alpha^{-1} f_{2}$. The analysis of Case 1 pertains and thus $T \Gamma_{i j}{ }^{k}=\alpha^{\epsilon(i j k)} \Gamma_{i j}{ }^{k}$. By assumption, $\alpha \neq 1, \alpha^{2} \neq 1$, and $\alpha^{3} \neq 1$. Thus once again $\Gamma_{i j}{ }^{k}=0$ if $\epsilon(i j k) \neq 0$ and we may conclude $\mathrm{SO}(2) \subset G_{\Gamma}^{+}$.

### 7.3. Elements of order 2 and of order 3

Lemma 7.2 focuses attention on the elements of order 2 and of order 3. The following is a useful technical result.
Lemma 7.3. $T_{i} \in G_{\Gamma}^{+}$for $i=1,2$. If $T_{1}, T_{2}$, and $T_{1} T_{2}$ have order 3 , then either $T_{1} T_{2}^{2}=\mathrm{id}$ or $T_{1} T_{2}^{2}$ has order 2 .
Proof. Suppose $T \in \operatorname{GL}(3, \mathbb{R})$ can be written in the form

$$
T=\left(\begin{array}{rrr}
\cos (\theta) & \sin (\theta) & 0  \tag{9}\\
-\sin (\theta) & \cos (\theta) & 0 \\
0 & 0 & 1
\end{array}\right)
$$

Then $\operatorname{Tr}(T)=2 \cos (\theta)+1$ so $\operatorname{Tr}(T)$ determines $T$ up to conjugacy in this setting:
(1) $T$ has order $1 \Leftrightarrow \cos (\theta)=+1 \Leftrightarrow \operatorname{Tr}(T)=+3$.
(2) $T$ has order $2 \Leftrightarrow \cos (\theta)=-1 \Leftrightarrow \operatorname{Tr}(T)=-1$.
(3) $T$ has order $3 \Leftrightarrow \cos (\theta)=-\frac{1}{2} \Leftrightarrow \operatorname{Tr}(T)=0$.

Since $T_{1} \in G_{\Gamma}^{+}$has order 3, we may choose a non-null vector so $T e_{3}= \pm e_{3}$. Since $T^{3}=\mathrm{id}, T e_{3}=e_{3}$. Consequently, $T_{1}$ has the form given in Equation (9). Let $P_{1}:=e_{3}^{\perp}$ be the rotation plane of $T_{1}$ and, similarly, let $P_{2}$ be the rotation plane of $T_{2}$. Since $\operatorname{dim}\left\{P_{1} \cap P_{2}\right\} \geq \operatorname{dim}\left\{P_{1}\right\}+\operatorname{dim}\left\{P_{2}\right\}-3=1$, we can choose
a unit vector $e_{1} \in P_{1} \cap P_{2}$. Let $\left\{e_{2}, f_{2}\right\}$ be unit vectors so $T_{1} e_{1}=-\frac{1}{2} e_{1}+\frac{\sqrt{3}}{2} e_{2}$ and $T_{2} e_{1}=-\frac{1}{2} e_{1}+\frac{\sqrt{3}}{2} f_{2}$.
Case 1. Assume $\rho$ is definite. Decompose $f_{2}=x e_{2}+y e_{3}$ where $x^{2}+y^{2}=1$. Let $f_{3}:=-y e_{2}+x e_{3}$ be a unit vector which spans the rotation axis of $T_{2}$. We then have:

$$
\begin{array}{rll}
f_{1}=e_{1}, & f_{2}=x e_{2}+y e_{3}, & f_{3}=-y e_{2}+x e_{3}, \\
e_{1}=f_{1}, & e_{2}=x f_{2}-y f_{3}, & e_{3}=y f_{2}+x f_{3}, \\
T_{1} e_{1}=-\frac{1}{2} e_{1}+\frac{\sqrt{3}}{2} e_{2}, & T_{1} e_{2}=-\frac{\sqrt{3}}{2} e_{1}-\frac{1}{2} e_{2}, & T_{1} e_{3}=e_{3}, \\
T_{2} e_{1}=-\frac{1}{2} e_{1}+\frac{\sqrt{3}}{2} f_{2}, & T_{2} f_{2}=-\frac{\sqrt{3}}{2} e_{1}-\frac{1}{2} f_{2}, & T_{2} f_{3}=f_{3}, \\
T_{1} T_{2} e_{1}=T_{1}\left\{-\frac{1}{2} e_{1}+\frac{\sqrt{3}}{2}\left(x e_{2}+y e_{3}\right)\right\}=\frac{1}{4}(1-3 x) e_{1}+\star e_{2}+\star e_{3}, \\
T_{1} T_{2} e_{2}=T_{1} T_{2}\left(x f_{2}-y f_{3}\right)=T_{1}\left(-\frac{x \sqrt{3}}{2} e_{1}-\frac{x}{2} f_{2}-y f_{3}\right) \\
& =T_{1}\left(-\frac{x \sqrt{3}}{2} e_{1}+\left(-\frac{x^{2}}{2}+y^{2}\right) e_{2}+\star e_{3}\right) \\
& =\star e_{1}+\frac{1}{4}\left(-3 x+x^{2}-2 y^{2}\right) e_{2}+\star e_{3},
\end{array} \begin{array}{rlrl} 
\\
T_{1} T_{2} e_{3} & =T_{1} T_{2}\left(y f_{2}+x f_{3}\right)=T_{1}\left(\star e_{1}-\frac{1}{2} y f_{2}+x f_{3}\right) \\
& =T_{1}\left(\star e_{1}+\star e_{2}+\left(-\frac{1}{2} y^{2}+x^{2}\right) e_{3}\right) & =\star e_{1}+\star e_{2}+\frac{1}{4}\left(4 x^{2}-2 y^{2}\right) e_{3}, & \\
\operatorname{Tr}\left(T_{1} T_{2}\right) & =\frac{1}{4}\left(1-3 x-3 x+x^{2}+4 x^{2}-4 y^{2}\right)=\frac{3}{4}\left(-1-2 x+3 x^{2}\right) .
\end{array}
$$

Since $T_{1} T_{2}$ has order $3, \operatorname{Tr}\left(T_{1} T_{2}\right)=0$. The equation $-1-2 x+3 x^{2}=0$ then implies $x=1$ or $x=-\frac{1}{3}$. Since $T_{2}$ has order $3, T_{2}^{2}$ also has order 3 . Introduce similar notation $\left\{\tilde{x}, \tilde{y}, \tilde{f}_{2}, \tilde{f}_{3}\right\}$ for $T_{2}^{2}$ to expand $\operatorname{Tr}\left(T_{1} T_{2}^{2}\right)=\frac{3}{4}\left(-1-2 \tilde{x}+3 \tilde{x}^{2}\right)$. Because $T_{2}^{2} e_{1}=-\frac{1}{2} e_{1}-\frac{\sqrt{3}}{2} f_{2}$, we see that $\tilde{f}_{2}=-f_{2}$ and thus $\tilde{x}=-x$. We complete proof if $\rho$ is definite by computing:

$$
\operatorname{Tr}\left(T_{1} T_{2}^{2}\right)=\frac{3}{4}\left(-1+2 x+3 x^{2}\right)=\left\{\begin{array}{rrr}
3 & \text { if } & x=1 \\
-1 & \text { if } & x=-\frac{1}{3}
\end{array}\right\} .
$$

Case 2. Assume $\rho$ is indefinite. In the hyperbolic setting, we have $x^{2}-y^{2}=1$, but the same argument pertains. We compute:

$$
\begin{array}{lll}
T_{1} e_{1}=-\frac{1}{2} e_{1}+\frac{\sqrt{3}}{2} e_{2}, & T_{1} e_{2}=-\frac{\sqrt{3}}{2} e_{1}-\frac{1}{2} e_{2}, & T_{1} e_{3}=e_{3}, \\
T_{2} e_{1}=-\frac{1}{2} e_{1}+\frac{\sqrt{3}}{2} f_{2}, & T_{2} f_{2}=-\frac{\sqrt{3}}{2} e_{1}-\frac{1}{2} f_{2}, & T_{2} f_{3}=f_{3}, \\
f_{1}=e_{1}, & f_{2}=x e_{2}+y e_{3}, & f_{3}=y e_{2}+x e_{3}, \\
e_{1}=f_{1}, & e_{2}=x f_{2}-y f_{3}, & e_{3}=-y f_{2}+x f_{3} . \\
T_{1} T_{2} e_{1}=T_{1}\left\{\left(-\frac{1}{2} e_{1}+\frac{\sqrt{3}}{2}\left(x e_{2}+y e_{3}\right)\right\}=\frac{1}{4}(1-3 x) e_{1}+\star e_{2}+\star e_{3},\right. \\
T_{1} T_{2} e_{2}=T_{1} T_{2}\left(x f_{2}-y f_{3}\right)=T_{1}\left(-\frac{x \sqrt{3}}{2} f_{1}-\frac{x}{2} f_{2}-y f_{3}\right) \\
& =T_{1}\left(-\frac{x \sqrt{3}}{2} e_{1}+\left(-\frac{x^{2}}{2}-y^{2}\right) e_{2}+\star e_{3}\right) & \\
& =\star e_{1}+\left(-\frac{3 x}{4}+\frac{x^{2}}{4}+\frac{1}{2} y^{2}\right) e_{2}+\star e_{3}, &
\end{array}
$$

$$
\begin{aligned}
T_{1} T_{2} e_{3} & =T_{1} T_{2}\left(-y f_{2}+x f_{3}\right)=T_{1}\left(\star f_{1}+\frac{1}{2} y f_{2}+x f_{3}\right) \\
& =T_{1}\left(\star e_{1}+\star e_{2}+\left(\frac{1}{2} y^{2}+x^{2}\right) e_{3}\right) \\
& =\star e_{1}+\star e_{2}+\left(x^{2}+\frac{1}{2} y^{2}\right) e_{3} \\
\operatorname{Tr}\left(T_{1} T_{2}\right) & =\frac{1}{4}\left(1-3 x-3 x+x^{2}+2 y^{2}+4 x+2 y^{2}\right)=\frac{3}{4}\left(-1-2 x+3 x^{2}\right)
\end{aligned}
$$

The remainder of the argument is the same as in the definite setting. Assertion 2 now follows.

Let $\nu(T)$ be the order of $T \in \mathrm{GL}(2, \mathbb{R})$. In what follows we list the (possibly) non-zero Christoffel symbols up to the symmetry $\Gamma_{i j}{ }^{k}=\Gamma_{j i}{ }^{k}$. Let

$$
S_{j} e_{i}=\left\{\begin{align*}
e_{i} & \text { if } i=j  \tag{10}\\
-e_{i} & \text { if } i \neq j
\end{align*}\right\}
$$

Let $s_{3}$ be the symmetric group of all permutations on 3 elements; $s_{3}$ is a nonAbelian group of order 6 . Let $a_{4}$ be the alternating group of permutations of 4 elements; $a_{4}$ is a non-Abelian group of order 12. We assume there exists $T \in G_{\Gamma}^{+}$of order at least 3 as otherwise Assertion 4 holds of Theorem 1.12 holds. By Assertion 1 of Lemma 7.1, we may choose a non-null vector $e_{3}$ so that $T e_{3}= \pm e_{3}$ and $\rho\left(e_{3}, e_{3}\right)= \pm 1$. Let $V:=e_{3}^{\perp}$, let $\rho_{V}:=\left.\rho\right|_{V}$, and let $T_{V}:=\left.T\right|_{V}$. Then $\rho_{V}$ is non-degenerate and $T_{V} \in \operatorname{SO}\left(\left.\rho\right|_{V}\right)$. We divide the proof of Theorem 1.12 into 4 cases.

### 7.4. The proof of Theorem 1.12(1)

Assume $\rho_{V}$ is indefinite. Choose a hyperbolic basis for $V$ so

$$
\rho_{V}\left(e_{1}, e_{1}\right)=\rho_{V}\left(e_{2}, e_{2}\right)=0 \text { and } \rho_{V}\left(e_{1}, e_{2}\right)=1
$$

Suppose first $T e_{3}=e_{3}$. Since $\operatorname{det}\left(T_{V}\right)=+1, T_{V} e_{1}=\alpha e_{1}$ and $T_{V} e_{2}=\alpha^{-1} e_{2}$. Since $\nu(T) \geq 3, \alpha \neq \pm 1$. Since $T^{n}$ preserves $\Gamma$, the only possible non-zero Christoffel symbols are $\left\{\Gamma_{13}{ }^{1}, \Gamma_{23}{ }^{2}, \Gamma_{12}{ }^{3}, \Gamma_{33}{ }^{3}\right\}$. Consequently, $\Gamma$ has the form given in Assertion 1 and the Ricci tensor is as given; to ensure $\rho$ is nondegenerate, $(a, b, c, d)$ satisfy the given constraints. We adopt the notation of Equation (10) to define $S_{2}$. Then $S_{2}^{*} \Gamma=\Gamma$ implies $d=0$ which is false. Thus, in particular, $S_{2} \notin G_{\Gamma}^{+}$so $G_{\Gamma}^{+} \neq \mathrm{SO}(\rho)$.

Suppose $S \in G_{\Gamma}^{+}-\mathrm{SO}(1,1)$. Since any two distinct connected 1-dimensional subgroups generate $\mathrm{SO}(\rho)$ and since $G_{\Gamma}^{+} \neq \mathrm{SO}(\rho), S$ must normalize $\mathrm{SO}(1,1)$ and in particular preserves $V$ and $e_{3}$. Since $S \notin \mathrm{SO}(1,1), S e_{3}=-e_{3}$. But $S^{*} \Gamma_{33}{ }^{3}=-\Gamma_{33}{ }^{3}$ and hence $d=0$. This is not possible. Thus $G_{\Gamma}^{+}=\mathrm{SO}(1,1)$.

Next suppose $T e_{3}=-e_{3}$. Since $T_{V} \in \mathrm{O}\left(\rho_{V}\right)$ and $\operatorname{det}\left(T_{V}\right)=-1, T e_{1}=\alpha e_{2}$ and $T e_{2}=\alpha^{-1} e_{1}$ for some $\alpha \neq 0$. We then have $T^{2}=$ id which is false as we assumed that $\nu(T) \geq 3$. This completes the analysis of the case when $\rho_{V}$ is indefinite.

### 7.5. The proof of Theorem $1.12(2)$

Assume that $\rho_{V}$ is definite and that $T$ has order at least 4. If $T e_{3}=-e_{3}$, then $T_{V} \in \mathrm{O}(2)-\mathrm{SO}(2)$ and $T^{2}=\mathrm{id}$ which is false. Thus $T e_{3}=e_{3}$ and $T_{V} \in \mathrm{SO}(2)$. Let $\left\{e_{1}, e_{2}\right\}$ be an orthonormal basis for $V$. Since $T_{V} \in \mathrm{SO}(V)$, $T_{V}$ is a rotation through an angle $\theta$ on $e_{3}^{\perp}$. We use the argument used to prove Theorem 1.11. Set $f_{1}=e_{1}+\sqrt{-1} e_{2}, f_{2}=e_{1}-\sqrt{-1} e_{2}$, and $f_{3}=e_{3}$. We then have $T f_{1}=e^{\sqrt{-1} \theta} f_{1}, T f_{2}=e^{-\sqrt{-1} \theta} f_{2}$, and $T f_{3}=f_{3}$. Let $\tilde{\Gamma}_{i j}^{k}$ be the complex Christoffel symbols relative to the basis $\left\{f_{1}, f_{2}, f_{3}\right\}$. We have:

$$
T^{*} \tilde{\Gamma}_{i j}^{k}=e^{\sigma_{i j}{ }^{k} \sqrt{-1} \theta} \tilde{\Gamma}_{i j}^{k} \text { for } \sigma_{i j}^{k}:=\left\{\delta_{1 i}-\delta_{2 i}\right\}+\left\{\delta_{1 j}-\delta_{2 j}\right\}+\left\{\delta_{2 k}-\delta_{1 k}\right\}
$$

Consequently, $\sigma_{i j}{ }^{k} \in\{-3,-2,-1,0,1,2,3\}$. Since $T_{V}$ does not have order 1, 2 or $3, e^{\sigma_{i j}{ }^{k} \sqrt{-1} \theta} \neq 1$ for $\sigma_{i j}^{k} \neq 0$ and thus $\tilde{\Gamma}_{i j}{ }^{k}=0$ for $\sigma_{i j}{ }^{k} \neq 0$. The possible elements with $\sigma_{i j}{ }^{k}=0$ are $\left\{\tilde{\Gamma}_{12}{ }^{3}, \tilde{\Gamma}_{13}{ }^{1}, \tilde{\Gamma}_{23}{ }^{2}, \tilde{\Gamma}_{33}{ }^{3}\right\}$. After disentangling the notation, we conclude that $\Gamma=\Gamma_{\mathrm{SO}(2)}(a, b, c, d)$ has the form where the (potentially) non-zero entries are given in Assertion 2 of Theorem 1.12 and, consequently, $\mathrm{SO}(2) \subset \Gamma_{\mathrm{SO}(2)}(a, b, c, d)$. One then computes the Ricci tensor and obtains the required conditions on $(a, b, c, d)$.

We complete the proof of Assertion 2 by showing $\mathrm{SO}(2)=G_{\Gamma}^{+}$. Adopt the notation of Equation (10) to define $S_{2}$. If $S_{2}^{*} \Gamma=\Gamma$, then $a=0$ which is false. Thus $G_{\Gamma}^{+} \neq \mathrm{SO}(\rho)$. Suppose there exists $S \in G_{\Gamma}^{+}-\mathrm{SO}(2)$. As any 2 distinct connected 1-dimensional Lie subgroups of $\mathrm{SO}(\rho)$ generate $\mathrm{SO}(\rho)$ and as $G_{\Gamma}^{+} \neq \mathrm{SO}(\rho), S$ must normalize $\mathrm{SO}(2)$ so, in particular, $S$ preserves $V$ and, consequently, $S e_{3}= \pm e_{3}$. If $S e_{3}=e_{3}$, then $S \in \mathrm{SO}(2)$ which is false. Thus $S e_{3}=-e_{3}$. Since $\operatorname{det}\left(S_{V}\right)=-1, S$ fixes a vector of $V$. Choose the basis so $S e_{2}=e_{2}$. It then follows $S e_{1}=-e_{1}$ so $S=S_{2} \in G_{\Gamma}^{+}$which contradicts the argument we have just given. Thus $G_{\Gamma}^{+}=\mathrm{SO}(2)$ and Assertion 2 holds.

### 7.6. The proof of Theorem $1.12(3)$

If there exists an element $T \in G_{\Gamma}^{+}$of order at least 4 , then the analysis given above to examine Assertion 1 or Assertion 2 pertains and $G_{\Gamma}^{+}=\mathrm{SO}(1,1)$ or $G_{\Gamma}^{+}=\mathrm{SO}(2)$. Thus we conclude that the order of any element $T \in G_{\Gamma}^{+}$ is at most 3. Furthermore, if $T$ has order 3 and if $T e_{3}= \pm e_{3}$, then $\rho_{V}$ is definite. Finally, since Assertion 4 does not hold, there exists an element of order 3. Fix such an element. In the proof of Assertion 3, we ignored the case where $\sigma_{i j}{ }^{k}= \pm 3$. When we include these cases, we must allow $\tilde{\Gamma}_{11}{ }^{2}$ and the complex conjugate $\tilde{\Gamma}_{22}{ }^{1}$. This shows $\Gamma=e \Gamma_{2}+\Gamma_{\mathrm{SO}}(a, b, c, d)$ is as described in Assertion 2. If $e=0$, then the analysis of Assertion 2 pertains and $\Gamma_{G}^{+}=\mathrm{SO}(2)$. Thus we may assume $e \neq 0$.

In Assertion 3, $G_{\Gamma}^{+}$can be bigger than $\mathbb{Z}_{3}$ in certain instances and we must examine these exceptional structures. We suppose there exists $S \in G_{\Gamma}^{+}-$ $\left\{1, T, T^{2}\right\}$. We wish to show $S$ can be chosen so $S$ has order 2 . Suppose to the contrary that $S$ has order 3 . If $T S$ has order 2 , we have an element of order 2 .

So we assume $T S$ has order 3. But the Assertion 2 of Lemma 7.3 shows that $T S^{2}=\mathrm{id}$ or $T S^{2}$ has order 2. Since $S \notin\left\{1, T, T^{2}\right\}, T S^{2} \neq \mathrm{id}$.

We use the argument used to prove the second assertion of Theorem 1.11. We had the exceptional structure $\Gamma_{2}$. We made a coordinate rotation to ensure $\Gamma_{z z}{ }^{\bar{z}}=1$. If instead, we assume that $\Gamma_{z z}{ }^{\bar{z}}=4 e+4 \sqrt{-1} f$, then we obtain a slightly more general form

$$
\begin{array}{llllll}
\Gamma_{11}^{1}=e, & \Gamma_{11}^{2}=-f, & \Gamma_{11}^{3}=a, & \Gamma_{12}^{1}=-f, & \Gamma_{12}^{2}=-e, & \Gamma_{12}^{3}=0, \\
\Gamma_{13}{ }^{1}=b, & \Gamma_{13}^{2}=c, & \Gamma_{13}^{3}=0, & \Gamma_{22}{ }^{1}=-e, & \Gamma_{22}^{2}=f, & \Gamma_{22}^{3}=a, \\
\Gamma_{23}{ }^{1}=-c, & \Gamma_{23}^{2}=b, & \Gamma_{23}^{3}=0, & \Gamma_{33}{ }^{1}=0, & \Gamma_{33}^{2}=0, & \Gamma_{33}^{3}=d .
\end{array}
$$

Let $S$ be an element of order 2 in $G_{\Gamma}^{+}$. Let $N(S,-1)$ be the -1 eigenspace of $S$ and let $V=e_{3}^{\perp}$. Since $\operatorname{dim}\{N(S,-1)\}+\operatorname{dim}\{V\}-3=1, N(S,-1)$ intersects $V$. If $N(S,-1)=V$, then $S$ commutes with $T$ so $T S$ is an element of order 6 which is impossible. Thus $N(S,-1) \cap V$ is 1 -dimensional. We choose the basis for $V$ so $N(S,-1) \cap V=e_{2} \cdot \mathbb{R}$; this implies $f=0$. We have $N\left(T S T^{-1},-1\right) \cap V=T e_{2} \cdot \mathbb{R}$ and $N\left(T^{2} S T^{-2},-1\right) \cap V=T^{2} e_{2} \cdot \mathbb{R}$. Thus $\left\{S, T S T^{-1}, T^{2} S T^{-2}\right\}$ are 3 distinct elements of order 2 in $G_{\Gamma}^{+}$. We have already ruled out the case $S e_{3}=e_{3}$. There are 3 remaining cases:
Case 1. Suppose $S e_{3}=-e_{3}$. This means $S \in \mathrm{O}(2)-\mathrm{SO}(2)$. So we can choose the basis so $S e_{1}=e_{1}$ and $S e_{2}=-e_{2}$ ). This implies $f=0, a=0, b=0$, and $d=0$. We obtain $\rho=\operatorname{diag}\left(-2 e^{2},-2 e^{2}, 2 c^{2}\right)$. In particular, $\rho$ is indefinite. We can renormalize the coordinates so $e=1$ and $c=1$ to obtain the structure of Assertion 1. We have $S T S^{-1}=T^{-1} ;\left\{\mathrm{id}, T, T^{2}, S, S T, S T^{2}\right\}$ is a non-Abelian group of order 3 and hence isomorphic to $s_{3}$. We note that $\left\{T, T^{2}\right\}$ are the elements of order 3 in $s_{3}$ and $\left\{S, T S T^{-1}, T^{2} S T^{-2}\right\}$ are the elements of order 2 in $s_{3}$.

Suppose that $\tilde{S}$ is an element of order 2 in $G_{\Gamma}^{+}$which does not belong to $s_{3}$. Since we have established the signature is indefinite, the argument we will give in Case 2 below shows that $\tilde{S} e_{3}=-e_{3}$ as well. Since $S \tilde{S} e_{3}=e_{3}, S \tilde{S}$ belongs to $\mathrm{SO}(2) \cap G_{\Gamma}^{+}=\left\{\mathrm{id}, T, T^{2}\right\}$ and $\tilde{S} \in s_{3}$. Thus the elements of order 2 in $G_{\Gamma}^{+}$ are given by $\left\{S, T S T^{-1}, T^{2} S T^{-2}\right\}$. Suppose $\tilde{T}$ is an element of $G_{\Gamma}^{+}$which is not in $s_{3}$. Thus $\tilde{T}$ has order 3. Since $\tilde{T} T$ is not in $s_{3}, \tilde{T} T$ has order 3 as well. Lemma 7.3 then implies $\tilde{T} T^{2}$ has order 1 or order 2 and hence belongs to $s_{3}$. This implies $\tilde{T}$ belongs to $s_{3}$. This contradiction then shows, as desired,that $G_{\Gamma}^{+}=s_{3}$.
Case 2. Suppose $S e_{3} \neq \pm e_{3}$ and $\rho$ is indefinite. We rescale $e_{1}$ to assume $e=1$. We rescale $e_{2}$ and $e_{3}$ to assume

$$
\rho\left(e_{1}, e_{1}\right)=\rho\left(e_{2}, e_{2}\right)=-\rho\left(e_{3}, e_{3}\right) \neq 0 \text { and } \rho\left(e_{i}, e_{j}\right)=0 \text { for } i \neq j
$$

We find $\{x, y\}$ so $x^{2}-y^{2}=1$ with $y \neq 0$. It then follows $x \neq 0$. Express:

$$
\begin{array}{ll}
S e_{1}=x e_{1}+y e_{3}, & S e_{2}=-e_{2}, \\
S e^{1}=x e^{1}-y e^{3}, & S e^{2}=-e^{2},
\end{array} \quad S e^{3}=y e^{1}-x e_{3}, x e^{3} .
$$

Let $\Delta_{i j}{ }^{k}:=\left(S_{*} \Gamma\right)_{i j}{ }^{k}-\Gamma_{i j}{ }^{k}$. We have $0=\Delta_{22}{ }^{2}=-2 f$. Thus $f=0$. We then have

$$
0=\Delta_{11}^{2}=-2 c x y, \quad 0=\Delta_{22}^{1}=1-x-a y, \quad 0=\Delta_{12}^{2}=1-x+b y
$$

Since $x y \neq 0$, we have $c=0, b=-a$, and $x=1-a y$. We impose these relations and compute:

$$
0=\Delta_{23}^{2}=2 a+\left(1-a^{2}\right) y
$$

This implies $a^{2} \neq 1$. We obtain therefore

$$
y=\frac{2 a}{a^{2}-1} \quad \text { and } \quad x=\frac{1+a^{2}}{1-a^{2}} .
$$

We verify $x^{2}-y^{2}=1$. Since $y \neq 0, a \neq 0$. The relations $\Delta_{11}{ }^{1}=0$ and $\Delta_{33}{ }^{3}=0$ imply

$$
3+3 a^{2}+2 a^{4}+2 a d=0, \quad 10 a^{3}+6 a^{5}+d+3 a^{4} d=0
$$

We eliminate $d$ to see:

$$
d=-\left(3+3 a^{2}+2 a^{4}\right)(2 a)^{-1}=-\left(10 a^{3}+6 a^{5}\right)\left(1+3 a^{4}\right)^{-1} .
$$

After cross multiplying and simplifying, we obtain $3\left(-1+a^{2}\right)^{2}\left(1+3 a^{2}+2 a^{4}\right)=$ 0 . This implies $a= \pm 1$ which is not permitted. Thus, as claimed earlier, Case 2 is impossible.
Case 3. Suppose $S e_{3} \neq \pm e_{3}$ and $\rho$ is definite. We may assume $e=1$ and $f=0$.

$$
\rho\left(e_{1}, e_{1}\right)=\rho\left(e_{2}, e_{2}\right)=\rho\left(e_{3}, e_{3}\right) \neq 0 \text { and } \rho\left(e_{i}, e_{j}\right)=0 \text { for } i \neq j .
$$

We find $\{x, y\}$ with $x^{2}+y^{2}=1$ with $y \neq 0$ so

$$
\begin{aligned}
& S e_{1}=x e_{1}+y e_{3}, \quad S e_{2}=-e_{2}, \quad S e_{3}=-y e_{1}-x e_{3}, \\
& S e^{1}=x e^{1}+y e^{3}, \quad S e^{2}=-e^{2}, \quad S e^{3}=y e^{1}-x e^{3}
\end{aligned}
$$

We compute

$$
0=\Delta_{12}{ }^{3}=c y^{2}, \quad 0=\Delta_{22}{ }^{1}=1-x+a y, \quad 0=\Delta_{12}{ }^{2}=1-x+b y
$$

We obtain $c=0, a=b$, and $x=1+b y$. We impose these relations and compute:

$$
\begin{aligned}
& 0=\Delta_{23}^{2}=-2 a-y-a^{2} y \text { so } \\
& x=\left(1-a^{2}\right)\left(1+a^{2}\right)^{-1} \text { and } y=-2 a\left(1+a^{2}\right)^{-1} .
\end{aligned}
$$

We note $x^{2}+y^{2}=1$. Furthermore, since $y \neq 0, a \neq 0$. The relations $0=\Delta_{11}{ }^{1}$ and $0=\Delta_{13}{ }^{1}$ imply

$$
0=3-3 a^{2}+2 a^{4}+2 a d \text { and } 0=1-3 a^{2}+4 a^{4}+a d-a^{3} d
$$

If $a= \pm 1$, the second equation is inconsistent and thus $a \neq \pm 1$ so $x \neq 0$. We set $d=-\frac{\left(3-3 a^{2}+2 a^{4}\right)}{2 a}$ and substitute this into the second relation to see
$2 a-4 a^{3}=0$. Since $a \neq 0, a= \pm \frac{1}{\sqrt{2}}$. We solve to see $b= \pm \frac{1}{\sqrt{2}}$ and $d=\mp \sqrt{2}$. This gives rise to two possibilities:

$$
\begin{array}{llllll}
a=b=\frac{1}{\sqrt{2}}, & c=0, & d=-\sqrt{2}, & e=1, & f=0, & x=\frac{1}{3},
\end{array} \quad y=-\frac{2 \sqrt{2}}{3}, ~\left(\frac{1}{\sqrt{2}}, \quad c=0, \quad d=\sqrt{2}, \quad e=1, \quad f=0, \quad x=\frac{1}{3}, \quad y=\frac{2 \sqrt{2}}{3} .\right.
$$

We remark that $x=\frac{1}{3}$ corresponds to $\tilde{x}=\frac{1}{3}$ in the analysis of Case 1 in the proof of Lemma 7.3; this is, of course, not an accident that this value surfaces again. If we consider $e_{2} \rightarrow-e_{2}$ and $e_{3} \rightarrow-e_{3}$, we simply interchange these two solutions. So there is really only one solution. This gives rise to the exceptional case given in Assertion (3b).

Let $\tilde{S} \in G_{\Gamma}^{+}$have order 2. Then there exists $\xi:=x e_{1}+y e_{2}$ for some $x^{2}+y^{2}=1$ so $\tilde{S} \xi=-\xi$. But then $\rho\left(\nabla_{\xi} \xi, \xi\right)=0$. Expanding this out yields the relation $-3 x y^{2}+x^{3}=0$. Since $x^{2}+y^{2}=1$, either $x=0$ and $\xi= \pm e_{2}$ or $x= \pm \frac{\sqrt{3}}{2}$ and $y= \pm \frac{1}{2}$. Thus the line through $\xi$ is a rotation of $\pm \frac{2 \pi}{3}$ from the line through $e_{2}$ and the argument above shows $\tilde{S}$ is unique. This shows that $\left\{S, T S T^{-1}, T^{2} S T^{-2}\right\}$ are the 3 elements of order 2 and conjugation by $T$ permutes them cyclically. Let $\tilde{T}$ be another element of order 3. By Assertion 2 of Lemma 7.1, if $\tilde{T}$ is another element of order 3 , then either $\tilde{T} T$ has order 2 or $\tilde{T} T^{2}=$ id or $\tilde{T} T^{2}$ has order 2. But in any event, this implies $\tilde{T}$ belongs to the subgroup generated by $T$ and $S$.

### 7.7. The proof Theorem $1.12(4)$

Suppose every element of $G_{\Gamma}^{+}$is of order 2. Then $A B A B=$ id implies $A B A^{-1}=B$ since every element is idempotent. Thus $G_{\Gamma}^{+}$is Abelian. We can simultaneously diagonalize the elements of $G_{\Gamma}^{+}$. Since we are dealing with $\mathbb{R}^{3}$, either $G_{\Gamma}^{+}=\mathbb{Z}_{2} \oplus \mathbb{Z}_{2}$ or $G_{\Gamma}^{+}=\mathbb{Z}_{2}$. We suppose the former possibility pertains. Let $S_{1}$ and $S_{2}$ generate $\mathbb{Z}_{2} \oplus \mathbb{Z}_{2}$. These idempotent matrices commute and thus can be simultaneously diagonalized. Since $\operatorname{det}\left(S_{i}\right)=+1$, each $S_{i}$ has two -1 eigenvalues and one +1 eigenvalue and $H=\left\{\mathrm{id}, T_{1}, T_{2}, T_{3}\right\}$. If $\Gamma$ is invariant under this action, then $\Gamma_{i j}{ }^{k}$ must contain each index exactly once. We compute:

$$
\Gamma_{12}^{3}=a^{3}, \quad \Gamma_{13}^{2}=a^{3}, \quad \Gamma_{23}^{1}=a^{1}, \quad \rho=-2 \operatorname{diag}\left(a^{2} a^{3}, a^{1} a^{3}, a^{1} a^{2}\right) .
$$

If we rescale and let $\tilde{e}_{i}=\mu_{i} e_{i}$, then

$$
\left(a^{1}, a^{2}, a^{3}\right) \rightarrow\left(\mu_{2} \mu_{3} \mu_{1}^{-1} a^{1}, \mu_{1} \mu_{3} \mu_{2}^{-1} a^{2}, \mu_{1} \mu_{2} \mu_{3}^{-1} a^{3}\right)
$$

We can certainly rescale to assume $a^{1}=1$. To preserve this normalization, we require $\mu_{1}=\mu_{2} \mu_{3}$. We then have $\tilde{a}^{2}=\mu_{2} \mu_{3} \mu_{3} \mu_{2}^{-1} a^{2}=\mu_{3}^{2} a^{-2}$ and $\tilde{a}^{3}=$ $\mu_{2} \mu_{3} \mu_{2} \mu_{3}^{-1}=\mu_{2}^{2} a^{3}$. Consequently we can obtain $a^{2}= \pm 1$ and $a^{3}= \pm 1$. So the possibilities become $\vec{a} \in\{(1,-1,-1),(1,-1,1),(1,1,-1),(1,1,1)\}$. By permuting the elements we can get $\vec{a} \in\{(1,-1,-1),(1,1,-1),(1,1,1)\}$. By replacing $e_{i} \rightarrow-e_{i}$, we can get $\vec{a} \in\{(1,-1,-1),(1,1,1)\}$ as claimed. The first possibility is discussed in Case 4a. The second case contains the permutation
$e_{1} \rightarrow e_{2} \rightarrow e_{3} \rightarrow e_{1}$ and is discussed in Assertion 3 b ; the structure group is $a_{4}$ and is not $\mathbb{Z}_{2} \oplus \mathbb{Z}_{2}$.
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