

## CHARACTERISTIC FUNCTION OF SYMMETRIC DAMEK-RICCI SPACE

SINHWI KIM AND JEONGHYEONG PARK

**ABSTRACT.** A Damek-Ricci space is a typical locally harmonic manifold, which is a generalization of the rank one symmetric space of the non-compact type. In this paper, we determine explicitly the characteristic function of a Damek-Ricci space by calculating the determinant of a Jacobi tensor.

### 1. Introduction

A Riemannian manifold  $(M, g)$  is locally harmonic at  $p \in M$  if every volume density  $\sqrt{\det(g_{ij})}$  is a function of the Riemannian distance from  $p$  on some neighborhood of  $p$  in  $M$ . There are several equivalent definitions for locally harmonic manifolds ([2], p. 156). One of them is the following:

**Theorem 1.** *A Riemannian manifold  $M = (M, g)$  is locally harmonic at  $p \in M$  if and only if the equality*

$$\Delta \Omega_p = f_p(\Omega_p) \quad \left( \Omega_p = \frac{1}{2} r_p^2 \right)$$

*holds for a certain smooth function  $f_p$  on  $[0, \varepsilon(p))$ , where  $\varepsilon(p)$  is the injectivity radius at  $p \in M$  and  $r_p$  is the Riemannian distance from  $p$ .*

It is known that the function  $f_p$  in Theorem 1 does not depend on the choice of  $p \in M$  ([2], Proposition 6.16). The function  $f = f_p$  ( $p \in M$ ) is called the characteristic function of a harmonic manifold  $M = (M, g)$ . The characteristic function plays an important role in the geometry of harmonic manifolds and there are many applications such as [5, 8, 9].

The characteristic function has been determined previously for all the rank one symmetric spaces except for the Cayley projective plane and the Cayley

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hyperbolic plane [8, 9]. Very recently, this function of the Cayley projective plane and the Cayley hyperbolic plane has just been determined explicitly [6].

Damek-Ricci spaces are important noncompact harmonic manifolds. Nonsymmetric Damek-Ricci spaces are counterexamples of the Lichnerowicz conjecture: “every locally harmonic manifold is a locally symmetric space.” and there are many nonsymmetric Damek-Ricci spaces, in addition to symmetric Damek-Ricci spaces. Complex hyperbolic space  $\mathbb{CH}^n(-1)$ , quaternion hyperbolic space  $\mathbb{QH}^n(-1)$ , and Cayley hyperbolic plane  $\mathbb{CH}^2(-1)$  are all symmetric Damek-Ricci spaces.

Damek-Ricci spaces are harmonic, so they have the characteristic function. In [1], the characteristic function of a Damek-Ricci space was discovered but it is not the explicit form. In this paper, we can get the explicit form of the characteristic function of a Damek-Ricci space by using another method. In this article, we shall prove the following theorem by using the Jacobi tensor.

**Theorem 2.** *Let  $S$  be a symmetric Damek-Ricci space. Then, the characteristic function as a harmonic manifold is given by*

$$(1.1) \quad \Delta\Omega = 1 + \sqrt{\frac{\Omega}{2}}((n+m)\coth(\sqrt{\frac{\Omega}{2}}) + m\tanh(\sqrt{\frac{\Omega}{2}})).$$

We aimed our paper to be self-contained as much as possible. The authors thank to the referee for the kind suggestions.

## 2. Preliminaries

In this section, we prepare a brief review on the geometry of the Damek-Ricci space.

A Damek-Ricci space is a one-dimensional extension of a generalized Heisenberg group and a Lie group with the Lie algebra of Iwasawa type. It is a solvable Lie group with a left invariant metric, and is a Riemannian homogeneous Hadamard manifold which is harmonic (see [1] for details). Concretely, the definition of a Damek-Ricci space is the following [1, 4].

**Definition.** Let  $\mathfrak{s}$  be a Lie algebra with inner product  $\langle \cdot, \cdot \rangle$  satisfying

$$(2.1) \quad \mathfrak{s} = \mathfrak{o} \oplus \mathfrak{h} \oplus \mathfrak{a}$$

orthogonally where  $\mathfrak{a}$  is a one-dimensional subspace of  $\mathfrak{s}$ ,  $\mathfrak{o} \oplus \mathfrak{h} = [\mathfrak{s}, \mathfrak{s}]$ , and the linear maps

$$(2.2) \quad \begin{aligned} J : \mathfrak{h} &\rightarrow \text{End}(\mathfrak{o}), \quad J_X := J(X), \\ J_X^2 &= -Id_{\mathfrak{o}}, \quad \forall X \in \mathfrak{h} \end{aligned}$$

are given. The simply connected Lie group  $S$  with the Lie algebra  $\mathfrak{s}$  is called a *Damek-Ricci space*.

We will use the notations  $n := \dim \mathfrak{o}$ ,  $m := \dim \mathfrak{h}$ , and the Jacobi operator  $R_v(w) := R(w, v)v$  for all tangent vectors  $v, w \in TS$  in [1] for a Damek-Ricci

space  $S$ . For a symmetric Damek-Ricci space, the Jacobi operator has constant eigenvalues with multiplicities [1]:

**Theorem 3.** *Let  $S$  be a symmetric Damek-Ricci space. Let  $V + Y + sA$  be a unit vector in  $\mathfrak{s}$  where  $V \in \mathfrak{o}, Y \in \mathfrak{h}, sA \in \mathfrak{a}$ . The eigenvalues and multiplicities of  $R_{V+Y+sA}$  are  $0, 1; -1/4, n; -1, m$ .*

### 3. Proof of Theorem 2

This chapter uses the same method as [6]. First, we denote by  $\gamma = \gamma(t)$  the normal geodesic in  $(S, g)$  through the identity  $e = \gamma(0)$  with the initial direction  $\gamma'(0) = y_0$ . By Theorem 3, there is an orthonormal basis  $\{y_0, y_1, \dots, y_n, y_{\bar{1}}, \dots, y_{\bar{m}}\}$  of  $\mathfrak{s}$  such that

$$(3.1) \quad \begin{aligned} R_{y_0}(y_j) &= -\frac{1}{4}y_j, \\ R_{y_0}(y_{\bar{k}}) &= -y_{\bar{k}} \end{aligned}$$

for  $1 \leq j \leq n$  and  $1 \leq k \leq m$  when identifying  $\mathfrak{s}$  and  $T_e S$ . Then there is a parallel frame field  $\{y_1(t), \dots, y_n(t), y_{\bar{1}}(t), \dots, y_{\bar{m}}(t)\}$  along  $\gamma$  such that

$$(3.2) \quad \begin{aligned} y_j(0) &= y_j (1 \leq j \leq n), \\ y_{\bar{k}}(0) &= y_{\bar{k}} (1 \leq k \leq m). \end{aligned}$$

Since  $S$  is locally symmetric, by (3.1) and (3.2),

$$(3.3) \quad \begin{aligned} R_{\gamma'(t)}(y_j(t)) &= -\frac{1}{4}y_j(t), \\ R_{\gamma'(t)}(y_{\bar{k}}(t)) &= -y_{\bar{k}}(t) \end{aligned}$$

for  $1 \leq j \leq n$  and  $1 \leq k \leq m$ .

Now, let  $Y_i(t)$  ( $1 \leq i \leq n$ ) and  $Y_{\bar{l}}(t)$  ( $1 \leq l \leq m$ ) be the Jacobi vector fields along  $\gamma$  satisfying the following conditions

$$(3.4) \quad \begin{aligned} Y_i(0) &= 0, \quad Y_{\bar{l}}(0) = 0, \\ Y'_i(0) &= (\nabla_{\gamma'} Y_i)(0) = y_i, \quad Y'_{\bar{l}}(0) = (\nabla_{\gamma'} Y_{\bar{l}})(0) = y_{\bar{l}} \end{aligned}$$

for  $1 \leq i \leq n$  and  $1 \leq l \leq m$ . We set as follows along  $\gamma$ :

$$(3.5) \quad \begin{aligned} Y_i(t) &= \sum_{j=1}^n a_{ji}(t)y_j(t) + \sum_{k=1}^m a_{\bar{k}i}(t)y_{\bar{k}}(t), \\ Y_{\bar{l}}(t) &= \sum_{j=1}^n a_{j\bar{l}}(t)y_j(t) + \sum_{k=1}^m a_{\bar{k}\bar{l}}(t)y_{\bar{k}}(t) \end{aligned}$$

for  $1 \leq i \leq n$  and  $1 \leq l \leq m$ .

Since  $Y_i(t)$  ( $1 \leq i \leq n$ ) and  $Y_{\bar{l}}(t)$  ( $1 \leq l \leq m$ ) are Jacobi vector fields along the geodesic, from (3.5), taking account of (3.3), we have the following system

of differential equations along  $\gamma$ :

$$(3.6) \quad \begin{aligned} a''_{ji} - \frac{1}{4}a_{ji} &= 0, \\ a''_{\bar{k}i} &= a''_{j\bar{l}} = 0, \\ a''_{\bar{k}\bar{l}} - a_{\bar{k}\bar{l}} &= 0 \end{aligned}$$

for  $1 \leq i, j \leq n$  and  $1 \leq k, l \leq m$ .

Solving (3.6) under the initial conditions (3.4), we have

$$(3.7) \quad \begin{aligned} a_{ji}(t) &= 2\delta_{ji} \sinh\left(\frac{t}{2}\right), \\ a_{\bar{k}i}(t) &= a_{j\bar{l}}(t) = 0, \\ a_{\bar{k}\bar{l}}(t) &= \delta_{\bar{k}\bar{l}} \sinh t \end{aligned}$$

for  $1 \leq i, j \leq n$  and  $1 \leq k, l \leq m$ .

Now, we define  $(n+m) \times (n+m)$ -matrix  $A(t)$  by

$$(3.8) \quad A(t) = \begin{pmatrix} a_{ij}(t) & a_{i\bar{k}}(t) \\ a_{\bar{l}j}(t) & a_{\bar{l}\bar{k}}(t) \end{pmatrix}$$

for  $1 \leq i, j \leq n$  and  $1 \leq k, l \leq m$ . Let  $\Theta_e(q)$  be the volume density function of the geodesic sphere centered at  $e$  through  $q$  for each point  $q$  in a normal neighborhood centered at  $e$ . Then, it is well-known that the following equality

$$(3.9) \quad \Theta_e(\gamma(t)) = \det A(t)$$

holds along the geodesic  $\gamma$  for small  $t$ . From (3.8) with (3.7), we have

$$(3.10) \quad \det A(t) = \left(2 \sinh \frac{1}{2}t\right)^n (\sinh t)^m.$$

Thus, from (3.9) and (3.10), we have

$$(3.11) \quad \ln \Theta_e(\gamma(t)) = n \ln 2 + n \ln \sinh \frac{1}{2}t + m \ln \sinh t.$$

Here, since a Damek-Ricci space  $S$  is a harmonic manifold, the volume density function  $\theta_e$  (and hence, the function  $\Theta_e$ ) is a radial function on a normal neighborhood  $U_e$  centered at the identity  $e$ . Thus,  $\Theta_e$  is determined by its value along the geodesic  $\gamma$ . Thus, from (3.11), we easily see that the function  $\Theta_e$  is given by

$$(3.12) \quad \ln \Theta_e(q) = n \ln 2 + n \ln \sinh \frac{1}{2}t + m \ln \sinh t,$$

where  $q = \gamma(t) \in U_e - \{e\}$  ([3], p. 269).

Now, let  $\phi(t)$  be a smooth function of  $t$  ( $0 < t < \epsilon$ ,  $\epsilon > 0$ ), and consider the function  $f(q)$  on  $U_e$  defined by  $f(q) = \phi(t)$ ,  $t = d(e, q)$ ,  $q \in U_e$ . Then, the following equality holds as in [7] with the sign difference:

$$(3.13) \quad \Delta f = \phi''(t) + \frac{(\Theta_e(\gamma(t)))'}{\Theta_e(\gamma(t))} \phi'(t), \quad q \in \gamma(t),$$

where  $\Delta$  denotes the Laplace-Beltrami operator of  $(S, g)$ . Here, from (3.12), we get

$$\begin{aligned}
 (3.14) \quad \frac{(\Theta_e(\gamma(t)))'}{\Theta_e(\gamma(t))} &= (\ln \Theta_e(\gamma(t)))' \\
 &= \frac{n}{2} \coth \frac{1}{2}t + m \coth t \\
 &= \frac{n+m}{2} \coth \frac{1}{2}t + \frac{m}{2} \tanh \frac{1}{2}t.
 \end{aligned}$$

We here consider the special case where  $\phi(t) = \frac{1}{2}t^2$  ( $t > 0$ ). Then, from (3.13) and (3.14), by direct calculation, we see that

$$(3.15) \quad \Delta\Omega = 1 + \sqrt{\frac{\Omega}{2}} \left\{ (n+m) \coth \sqrt{\frac{\Omega}{2}} + m \tanh \sqrt{\frac{\Omega}{2}} \right\}$$

holds on  $U_e - \{e\}$ . This completes the proof of Theorem 2.

*Remark 1.* From Theorem 2, taking account of the discussion in [6], we may reconfirm the statement that symmetric Damek-Ricci spaces are isometric to the complex hyperbolic space, the quaternionic hyperbolic space or a Cayley hyperbolic plane with the canonical Riemannian metrics, respectively ([1], p. 79).

*Remark 2.* In the proof,  $y_j(t)$  or  $y_{\bar{k}}(t)$  are eigenvectors of the Jacobi operator at each point along the geodesic since they are eigenvectors at the initial point and the Damek-Ricci space  $S$  is locally symmetric. But, if  $S$  is nonsymmetric, usually, neither  $R_{\gamma'(t)}(y_j(t))$  nor  $R_{\gamma'(t)}(y_{\bar{k}}(t))$  are parallel, so we need to calculate  $R_{\gamma'(t)}(y_j(t))$  and  $R_{\gamma'(t)}(y_{\bar{k}}(t))$ . Thus, the following question remains.

**Question.** How can we determine the characteristic function of a nonsymmetric Damek-Ricci space by the Jacobi tensor in the explicit form?

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SINHUI KIM  
DEPARTMENT OF MATHEMATICS  
SUNGKYUNKWAN UNIVERSITY  
SUWON 16419, KOREA  
*E-mail address:* kimsinhui@skku.edu

JEONGHYEONG PARK  
DEPARTMENT OF MATHEMATICS  
SUNGKYUNKWAN UNIVERSITY  
SUWON 16419, KOREA  
*E-mail address:* parkj@skku.edu