# A POLAR REPRESENTATION OF A REGULARITY OF A DUAL QUATERNIONIC FUNCTION IN CLIFFORD ANALYSIS 

Ji Eun Kim and Kwang Ho Shon


#### Abstract

The paper gives the regularity of dual quaternionic functions and the dual Cauchy-Riemann system in dual quaternions. Also, the paper researches the polar representation and properties of a dual quaternionic function and their regular quaternionic functions.


## 1. Introduction

A dual number $z$ is consisted of real numbers $x$ and $y$ associated with a real unit 1 and the dual unit $\varepsilon$, where $\varepsilon^{2}=0$. A dual number is denoted in the form $z=x+\varepsilon y$. Thus, the dual numbers are elements of the two dimensional real algebra

$$
D=R[\varepsilon]=\left\{z=x+\varepsilon y \mid x, y \in \mathbb{R}, \varepsilon^{2}=0, \varepsilon \neq 0\right\}
$$

generated by 1 and $\varepsilon$ (see [17]).
The algebra of dual numbers has been studied by Clifford [1] and its applications to mechanics are due to Study [20]. Dual algebra has been often used for the field of displacement analysis, kinematic synthesis and dynamic analysis of spatial mechanisms. Dual numbers can be represented as follows ([3]):

1. Gaussian representation: $z=x+\varepsilon y$,
2. Polar representation: $z=r(1+\varepsilon \phi)$,
3. Exponential representation: $z=r \exp (\varepsilon \phi)$, where $r=x(x \neq 0), \phi=\frac{y}{x}$ and $\exp (\varepsilon \phi)=1+\varepsilon \phi$.

The dual number has a geometrical property which is investigated detail in $[4,17]$.

Clifford [1] also has studied the following algebra
$\mathbb{H}=\left\{p=z_{1}+z_{2} j \mid z_{1}=x_{0}+x_{1} i, z_{2}=x_{2}+x_{3} i, x_{r} \in \mathbb{R}(r=0,1,2,3)\right\}$

Received March 8, 2016; Revised June 14, 2016.
2010 Mathematics Subject Classification. 32W50, 32A99, 30G35, 11E88.
Key words and phrases. quaternion, dual number, polar representation, regularity, Clifford analysis.

This work was supported by the Dongguk University Research Fund of 2017.
called the set of quaternions. Here imaginary basis elements $i, j$ and $k$ satisfy the following conditions:

$$
i^{2}=j^{2}=k^{2}=-1, i j=-j i=k, j k=-k j=i, k i=-i k=j .
$$

For two quaternions $p=z_{1}+z_{2} j$ and $q=w_{1}+w_{2} j$, they are given the rules of the addition and multiplication as follows, respectively,

$$
p+q=\left(z_{1}+w_{1}\right)+\left(z_{2}+w_{2}\right) j
$$

and

$$
p q=\left(z_{1} w_{1}-z_{2} \overline{w_{2}}\right)+\left(z_{1} w_{2}+z_{2} \overline{w_{1}}\right) j
$$

where $\overline{w_{k}}=y_{k 0}-y_{k 1} i$ for $w_{k}=y_{k 0}+y_{k 1} i, y_{k j} \in \mathbb{R}, k=1,2, j=0,1$. Kajiwara et al. $[5,6]$ applied the theory on a closed densely defined operator and a priori estimate for the adjoint operator in a Hilbert space and brconvex domains. We $[9,10,11,12]$ researched corresponding Cauchy-Riemann systems and properties of functions with values in special quaternions such as reduced quaternions, split quaternions and dual split quaternions. We [13, 14, 15] investigated a regular functions defined by the differential operators of special quaternion number systems. Porter [19] gave an explicit solution to the linear equation in the quaternions $\mathbb{H}$.

This paper gives expressions of the differential operators and the exponential functions in dual quaternions. The paper researches the polar representation of dual quaternionic functions by using a dual Cauchy-Riemann system and their regularity of that functions in dual quaternions.

## 2. Preliminaries

For $p=x_{0}+x_{1} i+x_{2} j+x_{3} k \in \mathbb{H}$, we denote by $S c(p)$ the scalar part, and by $\operatorname{Vec}(p)$ the spatial vector part:

$$
p=S c(p)+V e c(p),
$$

where $S c(p)=x_{0}$ and $\operatorname{Vec}(p)=x_{1} i+x_{2} j+x_{3} k$ with $x_{r} \in \mathbb{R}(r=0,1,2,3)$. Then for $p, q \in \mathbb{H}$, we have

$$
\begin{aligned}
p+q= & S c(p)+S c(q)+V e c(p)+V e c(q), \\
p q= & S c(p) S c(q)-V e c(p) \cdot V e c(q) \\
& +S c(p) V e c(q)+V e c(p) S c(q)+V e c(p) \times V e c(q),
\end{aligned}
$$

where $S c(q)=y_{0}, \operatorname{Vec}(q)=y_{1} i+y_{2} j+y_{3} k$ with $y_{r} \in \mathbb{R}(r=0,1,2,3)$, the symbol • is a usual inner product,

$$
V e c(p) \cdot V e c(q)=x_{1} y_{1}+x_{2} y_{2}+x_{3} y_{3},
$$

and the symbol $\times$ is a usual outer product,

$$
\operatorname{Vec}(p) \times \operatorname{Vec}(q)=\left(x_{2} y_{3}-x_{3} y_{2}\right) i-\left(x_{1} y_{3}-x_{3} y_{1}\right) j+\left(x_{1} y_{2}-x_{2} y_{1}\right) k .
$$

The norm for a quaternion is

$$
|p|^{2}:=p p^{*}=S c(p)^{2}+V e c(p) \cdot \operatorname{Vec}(p),
$$

where $p^{*}=S c(p)-V e c(p)$, and the inverse of $p$ is

$$
p^{-1}=\frac{p^{*}}{|p|^{2}}
$$

For a unit quaternion, $|p|=1$, it is given by:

$$
p=(\cos (\theta / 2), \mathbf{n} \sin (\theta / 2))
$$

and

$$
\begin{gathered}
x_{0}=\cos (\theta / 2), x_{1}=n_{1} \sin (\theta / 2) \\
x_{2}=n_{2} \sin (\theta / 2), x_{3}=n_{3} \sin (\theta / 2)
\end{gathered}
$$

where an angle $\theta$ and axis $\mathbf{n}=\left(n_{1}, n_{2}, n_{3}\right)$ of rotation (see [7]).
We consider the following form:

$$
\mathbb{D}_{q}=\left\{Z=p_{1}+\varepsilon p_{2} \mid p_{r} \in \mathbb{H}, \varepsilon^{2}=0, r=1,2\right\} \cong \mathbb{H}^{2} \cong \mathbb{R}^{8},
$$

where $\varepsilon$ is the dual unit that commutes with $i, j$ and $k$. The dual quaternion $Z=p_{1}+\varepsilon p_{2} \in \mathbb{D}_{q}$ is also written as a linear combination of a scalar, denoted by $S c(Z)$, and a spatial vector, denoted by $\operatorname{Vec}(Z)$ (see [7, 8]):

$$
Z=S c(Z)+V e c(Z)=S c\left(p_{1}\right)+V e c\left(p_{1}\right)+\varepsilon\left\{S c\left(p_{2}\right)+V e c\left(p_{2}\right)\right\}
$$

where

$$
S c(Z)=S c\left(p_{1}\right)+\varepsilon S c\left(p_{2}\right), V e c(Z)=V e c\left(p_{1}\right)+\varepsilon V e c\left(p_{2}\right)
$$

with $p_{1}, p_{2} \in \mathbb{H}$.
For two elements $Z$ and $W=S c(W)+V e c(W)=S c\left(q_{1}\right)+V e c\left(q_{1}\right)+$ $\varepsilon\left\{S c\left(q_{2}\right)+V e c\left(q_{2}\right)\right\}$ of $\mathbb{D}_{q}$, we give the addition and the multiplication on $\mathbb{D}_{q}$ as follows:

$$
Z+W=S c(Z)+S c(W)+\varepsilon\{V e c(Z)+V e c(W)\}
$$

and

$$
\begin{aligned}
Z W= & S c(Z) S c(W)-V e c(Z) \cdot V e c(W)+S c(Z) V e c(W) \\
& +S c(W) \operatorname{Vec}(Z)+V e c(Z) \times V e c(W),
\end{aligned}
$$

where
$V e c(Z) \cdot V e c(W)=V e c\left(p_{1}\right) \cdot V e c\left(q_{1}\right)+\varepsilon\left\{V e c\left(p_{1}\right) \cdot V e c\left(q_{2}\right)+V e c\left(p_{2}\right) \cdot V e c\left(q_{1}\right)\right\}$
and
$V e c(Z) \times V e c(W)=V e c\left(p_{1}\right) \times V e c\left(q_{1}\right)+\varepsilon\left\{V e c\left(p_{1}\right) \times V e c\left(q_{2}\right)+V e c\left(p_{2}\right) \times V e c\left(q_{1}\right)\right\}$.
We give the complex conjugate element of $\mathbb{D}_{q}$ :

$$
Z^{*}=S c\left(p_{1}\right)-V e c\left(p_{1}\right)+\varepsilon\left\{S c\left(p_{2}\right)-V e c\left(p_{2}\right)\right\}
$$

It is also written as

$$
Z^{*}=S c(Z)-V e c(Z)
$$

and the modulus of $Z$, denoted by $|Z|$, is described by
$|Z|^{2}:=S c(Z) S c\left(Z^{*}\right)+V e c(Z) \cdot V e c\left(Z^{*}\right)=\left\{S c\left(p_{1}\right)\right\}^{2}+V e c\left(p_{1}\right) \cdot V e c\left(p_{1}\right)=\left|p_{1}\right|^{2}$.

Since every element of the set $\{\varepsilon p \mid p \in \mathbb{H}\}$ has no inverse, the inverse of a dual quaternion is given by

$$
Z^{-1}=\frac{Z^{\dagger}}{\left|p_{1}\right|^{2}} \in \mathbb{D}_{q} \quad\left(p_{1} \neq 0\right)
$$

where

$$
Z^{\dagger}=\left(\left|p_{1}\right|^{2}-\varepsilon p_{2} p_{1}^{*}\right) p_{1}^{-1}
$$

where $p_{1}^{-1}=\frac{p_{1}^{*}}{\left|p_{1}\right|^{2}}$, called the dual conjugate of $Z$ with $Z Z^{\dagger}=Z^{\dagger} Z=p_{1} p_{1}^{*}=$ $\left|p_{1}\right|^{2}$.

Plucker [18] gave screw coordinates so that we can rewrite dual quaternions in a form of the spherical linear interpolation. Screw parameters have the form $(\theta, d, \mathbf{I}, \mathbf{m})$, where
$\begin{cases}\theta & \text { is the angle of rotation, } \\ d & \text { is the translation along the axis, } \\ \mathbf{I} & \text { is the vector line direction, } \\ \mathbf{m}=\mathbf{p} \times \mathbf{I} & \text { is the line moment with } \mathbf{p} \text { is a point on a given line. }\end{cases}$

From the above components, Daniilidis [2] converted a unit dual quaternion to screw coordinates as follows:

$$
\begin{gather*}
S c\left(p_{1}\right)=\cos (\theta / 2), V e c\left(p_{1}\right)=\mathbf{I} \sin (\theta / 2), S c\left(p_{2}\right)=-\frac{d}{2} \sin (\theta / 2)  \tag{2.1}\\
V e c\left(p_{2}\right)=\mathbf{I} \frac{d}{2} \cos (\theta / 2)+\mathbf{m} \sin (\theta / 2) \tag{2.2}
\end{gather*}
$$

Referring [2], we can write the following representation of a unit dual quaternion

$$
Z=\cos \left(\frac{\theta+\varepsilon d}{2}\right)+(\mathbf{I}+\varepsilon \mathbf{m}) \sin \left(\frac{\theta+\varepsilon d}{2}\right)=\cos (\phi)+\mathbf{v} \sin (\phi)
$$

where $\mathbf{v}=\mathbf{I}+\varepsilon \mathbf{m}$ and $\phi=\frac{\theta+\varepsilon d}{2}$. By the properties of trigonometric functions, we have the representation of a unit dual quaternion

$$
\begin{aligned}
Z= & \rho \cos (\phi)+\mathbf{v} \rho \sin (\phi) \\
= & \cos \left(\frac{\theta}{2}\right) \cos \left(\frac{\varepsilon d}{2}\right)-\sin \left(\frac{\theta}{2}\right) \sin \left(\frac{\varepsilon d}{2}\right) \\
& +\mathbf{v}\left\{\sin \left(\frac{\theta}{2}\right) \cos \left(\frac{\varepsilon d}{2}\right)+\cos \left(\frac{\theta}{2}\right) \sin \left(\frac{\varepsilon d}{2}\right)\right\} .
\end{aligned}
$$

From the representation of a Taylor series, since $\cos \left(\frac{\varepsilon d}{2}\right)=1$ and $\sin \left(\frac{\varepsilon d}{2}\right)=\frac{\varepsilon d}{2}$, we have

$$
Z=\cos \left(\frac{\theta}{2}\right)-\sin \left(\frac{\theta}{2}\right) \frac{\varepsilon d}{2}+\mathbf{v}\left\{\sin \left(\frac{\theta}{2}\right)+\cos \left(\frac{\theta}{2}\right) \frac{\varepsilon d}{2}\right\}
$$

where $\rho=|Z|^{2}$. From the equations (2.1), we have

$$
Z=S c\left(p_{1}\right)+\varepsilon S c\left(p_{2}\right)+\mathbf{v}\left\{-\frac{2}{d} S c\left(p_{2}\right)+\varepsilon S c\left(p_{1}\right) \frac{d}{2}\right\}:=p+\mathbf{v} q
$$

where

$$
p=S c\left(p_{1}\right)+\varepsilon S c\left(p_{2}\right), q=-\frac{2}{d} S c\left(p_{2}\right)+\varepsilon S c\left(p_{1}\right) \frac{d}{2} .
$$

Since we have

$$
\mathbf{v}^{2}=\mathbf{I}^{2}=-1,
$$

we obtain a corresponding Euler's formula for a unit dual quaternion:

$$
\exp (\mathbf{v} \phi)=\sum_{n=0}^{\infty} \frac{1}{n!}(\mathbf{v} \phi)^{n}=\cos (\phi)+\mathbf{v} \sin (\phi)
$$

Proposition 2.1. For any unit dual quaternion, we have

1. $\exp \left(\mathbf{v} \phi_{1}\right) \exp \left(\mathbf{v} \phi_{2}\right)=\exp \left(\mathbf{v}\left(\phi_{1}+\phi_{2}\right)\right)$,
2. $\frac{\exp \left(\mathbf{v} \phi_{1}\right)}{\exp \left(\mathbf{v} \phi_{2}\right)}=\exp \left(\mathbf{v}\left(\phi_{1}-\phi_{2}\right)\right)$.

Proof. From the corresponding Euler's formula for a dual quaternion, we have

$$
\begin{aligned}
\exp \left(\mathbf{v} \phi_{1}\right) \exp \left(\mathbf{v} \phi_{2}\right) & =\left\{\cos \left(\phi_{1}\right)+\mathbf{v} \sin \left(\phi_{1}\right)\right\}\left\{\cos \left(\phi_{2}\right)+\mathbf{v} \sin \left(\phi_{2}\right)\right\} \\
& =\cos \left(\phi_{1}+\phi_{2}\right)+\mathbf{v} \sin \left(\phi_{1}+\phi_{2}\right) \\
& =\exp \left(\mathbf{v}\left(\phi_{1}+\phi_{2}\right)\right)
\end{aligned}
$$

and

$$
\begin{aligned}
\frac{\exp \left(\mathbf{v} \phi_{1}\right)}{\exp \left(\mathbf{v} \phi_{2}\right)} & =\left\{\cos \left(\phi_{1}\right)+\mathbf{v} \sin \left(\phi_{1}\right)\right\}\left\{\cos \left(\phi_{2}\right)-\mathbf{v} \sin \left(\phi_{2}\right)\right\} \\
& =\cos \left(\phi_{1}-\phi_{2}\right)+\mathbf{v} \sin \left(\phi_{1}-\phi_{2}\right) \\
& =\exp \left(\mathbf{v}\left(\phi_{1}-\phi_{2}\right)\right) .
\end{aligned}
$$

Therefore, we obtain the results.
Proposition 2.2. Let $Z=\cos (\phi)+\mathbf{v} \sin (\phi)$ be a unit dual quaternion. Then we have

$$
\begin{equation*}
Z^{n}=(\cos (\phi)+\mathbf{v} \sin (\phi))^{n}=\cos (n \phi)+\mathbf{v} \sin (n \phi) \tag{2.3}
\end{equation*}
$$

for all integer $n$.
Proof. From the induction for integers $n$, the equation (2.3) is obtained.

## 3. Hyperholomorphic function in dual quaternions

Let $\Omega$ be a bounded open set in $\mathbb{H}^{2}$. A function $F$ is given by

$$
F: \Omega \rightarrow \mathbb{D}_{q} ; F(Z)=f_{1}\left(p_{1}, p_{2}\right)+\varepsilon f_{2}\left(p_{1}, p_{2}\right)
$$

where

$$
\begin{aligned}
& f_{1}=g_{1}\left(z_{1}, z_{2}, w_{1}, w_{2}\right)+g_{2}\left(z_{1}, z_{2}, w_{1}, w_{2}\right) j \quad \text { and } \\
& f_{2}=h_{1}\left(z_{1}, z_{2}, w_{1}, w_{2}\right)+h_{2}\left(z_{1}, z_{2}, w_{1}, w_{2}\right) j
\end{aligned}
$$

are quaternionic functions, $g_{r}$ and $h_{r}(r=1,2)$ are complex-valued functions.

Definition. A function $F$ is said to be hyperholomorphic on $\Omega=D \cap L$, where $D$ is an open subset of $\mathbb{D}_{q}$ and $L=\mathbb{H} \backslash\{0\}+\varepsilon \mathbb{H}$, with values in $\mathbb{D}_{q}$ if the limit (3.4)
$\frac{d F(Z)}{d Z}:=\lim _{\zeta \rightarrow 0}\{F(Z+\zeta)-F(Z)\} \zeta^{-1}=\lim _{\zeta \rightarrow 0} \frac{\{F(Z+\zeta)-F(Z)\} \zeta^{*}}{\eta_{1} \eta_{1}^{*}} \quad\left(\eta_{1} \neq 0\right)$ exists, where $\zeta=\eta_{1}+\varepsilon \eta_{2} \rightarrow 0$ means $\eta_{1} \rightarrow 0$ and $\eta_{2} \rightarrow 0$ which are referred by [16].
Theorem 3.1. A function $F$ is hyperholomorphic on $\Omega$ with values in $\mathbb{D}_{q}$ if and only if the following conditions are held:

$$
\left\{\begin{array}{l}
\lim _{\substack{\eta_{1} \rightarrow 0 \\
\eta_{2} \rightarrow 0}}\{F(Z+\zeta)-F(Z)\} \eta_{1}^{-1} \quad \text { exists and }  \tag{3.5}\\
\lim _{\substack{\eta_{1} \rightarrow 0, \eta_{2} \rightarrow 0}}\left\{f_{1}\left(p_{1}+\eta_{1}, p_{2}+\eta_{2}\right)-f_{1}\left(p_{1}, p_{2}\right)\right\} \eta_{2}^{-1}=0
\end{array}\right.
$$

Proof. Since the dual part of a dual quaternion has no inverse elements, we use the dual conjugation of $Z$ as follows:

$$
\begin{aligned}
& \lim _{\zeta \rightarrow 0}\{F(Z+\zeta)-F(Z)\} \zeta^{-1}=\lim _{\substack{\eta_{1} \rightarrow 0, \eta_{2} \rightarrow 0}} \frac{\{F(Z+\zeta)-F(Z)\}\left(\eta_{1}^{*}-\varepsilon \eta_{2}^{\dagger}\right)}{\eta_{1} \eta_{1}^{*}} \\
= & \lim _{\substack{\eta_{1} \rightarrow 0, \eta_{2} \rightarrow 0}}\{F(Z+\zeta)-F(Z)\} \eta_{1}^{-1} \\
& -\varepsilon \lim _{\substack{\eta_{1} \rightarrow 0, \eta_{2} \rightarrow 0}}\left\{f_{1}\left(p_{1}+\eta_{1}, p_{2}+\eta_{2}\right)-f_{1}\left(p_{1}, p_{2}\right)\right\} \eta_{2}^{-1}\left(\eta_{2} \eta_{1}^{-1}\right)^{2} .
\end{aligned}
$$

For the existence of the above limit, the limit

$$
\lim _{\substack{\eta_{1} \rightarrow 0, \eta_{2} \rightarrow 0}}\left\{f_{1}\left(p_{1}+\eta_{1}, p_{2}+\eta_{2}\right)-f_{1}\left(p_{1}, p_{2}\right)\right\} \eta_{2}^{-1}
$$

has to be independent to $\left(\eta_{2} \eta_{1}^{-1}\right)^{2}$. Thus, we obtain the following equation:

$$
\lim _{\substack{\eta_{1} \rightarrow 0, \eta_{2} \rightarrow 0}}\left\{f_{1}\left(p_{1}+\eta_{1}, p_{2}+\eta_{2}\right)-f_{1}\left(p_{1}, p_{2}\right)\right\} \eta_{2}^{-1}=0
$$

Conversely, if the conditions (3.5) are satisfied for the function $F$, then the limit

$$
\lim _{\zeta \rightarrow 0}\{F(Z+\zeta)-F(Z)\} \zeta^{-1}
$$

exists. From the definition of a hyperholomorphic function in $\mathbb{D}_{q}$, the function $F$ is hyperholomorphic.

We give the left differential operators in $\mathbb{D}_{q}$.

$$
D_{1}:=\frac{\partial}{\partial z_{1}}-j \frac{\partial}{\partial \overline{z_{2}}} \quad \text { and } \quad D_{1}^{*}=\frac{\partial}{\partial \overline{z_{1}}}+j \frac{\partial}{\partial \overline{z_{2}}}
$$

where $\frac{\partial}{\partial z_{r}}$ and $\frac{\partial}{\partial \bar{z}_{r}}(r=1,2)$ are usual complex differential operators and $j$ is an imaginary basis element in $\mathbb{H}$.

Remark 3.2. From the representation of differential operators in $\mathbb{D}_{q}$, we have

$$
\begin{aligned}
F D_{1} & =\left\{g_{1}+g_{2} j+\varepsilon\left(h_{1}+h_{2} j\right)\right\}\left(\frac{\partial}{\partial z_{1}}-j \frac{\partial}{\partial \overline{z_{2}}}\right) \\
& =\left\{\frac{\partial g_{1}}{\partial z_{1}}+\frac{\partial g_{2}}{\partial \overline{z_{2}}}+\left(\frac{\partial g_{2}}{\partial \overline{z_{1}}}-\frac{\partial g_{1}}{\partial z_{2}}\right) j\right\}+\varepsilon\left\{\frac{\partial h_{1}}{\partial z_{1}}+\frac{\partial h_{2}}{\partial \overline{z_{2}}}+\left(\frac{\partial h_{2}}{\partial \overline{z_{1}}}-\frac{\partial h_{1}}{\partial z_{2}}\right) j\right\}
\end{aligned}
$$

and

$$
\begin{aligned}
F D_{1}^{*} & =\left\{g_{1}+g_{2} j+\varepsilon\left(h_{1}+h_{2} j\right)\right\}\left(\frac{\partial}{\partial \overline{z_{1}}}+j \frac{\partial}{\partial \overline{z_{2}}}\right) \\
& =\left\{\frac{\partial g_{1}}{\partial \overline{z_{1}}}-\frac{\partial g_{2}}{\partial \overline{z_{2}}}+\left(\frac{\partial g_{2}}{\partial z_{1}}+\frac{\partial g_{1}}{\partial z_{2}}\right) j\right\}+\varepsilon\left\{-\frac{\partial h_{2}}{\partial \overline{z_{2}}}+\frac{\partial h_{1}}{\partial \overline{z_{1}}}+\left(\frac{\partial h_{1}}{\partial z_{2}}+\frac{\partial h_{2}}{\partial z_{1}}\right) j\right\} .
\end{aligned}
$$

Definition. Let $\Omega$ be a bounded open set in $\mathbb{H}^{2}$. Then a function $F$ is said to be hyperholomorphic on $\Omega$ with values in $\mathbb{D}_{q}$ if $F$ satisfies the following conditions:
(i) each component of $F, f_{1}$ and $f_{2}$, is a continuously differentiable function and
(ii) $F$ satisfies the following equations

$$
\begin{cases}F D_{1}^{*}=0 & \text { and }  \tag{3.6}\\ \frac{\partial f_{1}}{\partial y_{r}}=0 & (r=0,1,2,3)\end{cases}
$$

In detail, the equations (3.6) is equivalent to the following system

$$
\frac{\partial g_{1}}{\partial \overline{z_{1}}}=\frac{\partial g_{2}}{\partial \overline{z_{2}}}, \frac{\partial g_{2}}{\partial z_{1}}=-\frac{\partial g_{1}}{\partial z_{2}}, \frac{\partial h_{1}}{\partial \overline{z_{1}}}=\frac{\partial h_{2}}{\partial \overline{z_{2}}}, \frac{\partial h_{2}}{\partial z_{1}}=-\frac{\partial h_{1}}{\partial z_{2}},
$$

called the corresponding Cauchy-Riemann system on $\mathbb{D}_{q}$.
Let $\Omega$ be an open set in $\mathbb{H}^{2}$. A function can be written by

$$
\begin{aligned}
& F: \Omega \rightarrow \mathbb{D}_{q} \\
& F(Z)=F(p(\rho, \phi), q(\rho, \phi))=M(p(\rho, \phi), q(\rho, \phi))+\mathbf{v} N(p(\rho, \phi), q(\rho, \phi)),
\end{aligned}
$$

where

$$
M=S c\left(f_{1}\right)+\varepsilon S c\left(f_{2}\right) \text { and } \quad N=-\frac{2}{d} S c\left(f_{2}\right)+\varepsilon S c\left(f_{1}\right) \frac{d}{2}
$$

are dual-valued functions and $S c\left(f_{1}\right)$ and $S c\left(f_{2}\right)$ are real-valued functions.
Theorem 3.3. Let $\Omega$ be a bounded open set in $\mathbb{H}^{2}$. If a function $F=M+\mathbf{v} N$ is hyperholomorphic on $\Omega$ with values in $\mathbb{D}_{q}$, then the following equations hold:

$$
\begin{equation*}
\frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{1}\right)}=\frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{2}\right)} \quad \text { and } \quad \frac{4}{d^{2}} \frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{1}\right)}=\frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{2}\right)} \tag{3.7}
\end{equation*}
$$

Proof. From the definition of the hyperholomorphic function $F$ on $\Omega$ with values in $\mathbb{D}_{q}$ and Definition 3, if the following limit

$$
\lim _{\substack{\lambda_{1} \rightarrow 0, \lambda_{2} \rightarrow 0}} \frac{S c\left(f_{1}\right)+\varepsilon S c\left(f_{2}\right)+\mathbf{v}\left\{\frac{-2}{d} S c\left(f_{2}\right)+\varepsilon \frac{d}{2} S c\left(f_{1}\right)\right\}}{\lambda_{1}+\varepsilon \lambda_{2}+\mathbf{v}\left(\frac{-2}{d} \lambda_{2}+\varepsilon \frac{d}{2} \lambda_{1}\right)}
$$

exists, then the function $F$ is hyperholomorphic, where $\lambda_{1}=S c\left(\eta_{1}\right)$ and $\lambda_{2}=$ $S c\left(\eta_{2}\right)$. By the definition of the existence of the limit and calculating of the complex conjugation of a dual quaternion, we have

$$
\begin{aligned}
& \frac{\partial S c\left(f_{1}\right)\left(1+\mathbf{I} \varepsilon \frac{d}{2}\right)+\partial S c\left(f_{2}\right)\left(\varepsilon-\frac{2}{d} \mathbf{v}\right)}{\partial S c\left(p_{1}\right)\left(1+\mathbf{I} \varepsilon \frac{d}{2}\right)} \\
= & \frac{\partial S c\left(f_{1}\right)\left(1+\mathbf{I} \varepsilon \frac{d}{2}\right)+\partial S c\left(f_{2}\right)\left(\varepsilon-\frac{2}{d} \mathbf{v}\right)}{\partial S c\left(p_{2}\right)\left(\varepsilon-\frac{2}{d} \mathbf{v}\right)}
\end{aligned}
$$

By arranging the above equation, we have

$$
\left\{\frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{1}\right)}-\frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{2}\right)}\right\}\left(2 \varepsilon-\frac{2}{d} \mathbf{v}\right)+\left\{\frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{2}\right)}-\frac{4}{d^{2}} \frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{1}\right)}\right\}(1+\mathbf{I} \varepsilon d)=0
$$

Therefore, we obtain the equations (3.7).
Theorem 3.4. Let $\Omega$ be a bounded open set in $\mathbb{H}^{2}$ and a function $F=M+\mathbf{v} N$ be hyperholomorphic on $\Omega$ with values in $\mathbb{D}_{q}$. Then the following equations hold:

$$
\begin{equation*}
\frac{\rho}{2}(1+\varepsilon) \frac{\partial M}{\partial \rho}=\frac{\partial N}{\partial \phi} \quad \text { and } \quad \frac{\rho}{2}(1+\varepsilon) \frac{\partial N}{\partial \rho}=-\frac{\partial M}{\partial \phi} \tag{3.8}
\end{equation*}
$$

Proof. From the chain rule of multi variables calculus, we have

$$
\begin{aligned}
& \frac{\partial M}{\partial \rho}=\frac{\partial p}{\partial \rho} \frac{\partial M}{\partial p}+\frac{\partial q}{\partial \rho} \frac{\partial M}{\partial q}, \frac{\partial M}{\partial \phi}=\frac{\partial p}{\partial \phi} \frac{\partial M}{\partial p}+\frac{\partial q}{\partial \phi} \frac{\partial M}{\partial q} \\
& \frac{\partial N}{\partial \rho}=\frac{\partial p}{\partial \rho} \frac{\partial N}{\partial p}+\frac{\partial q}{\partial \rho} \frac{\partial N}{\partial q}, \frac{\partial N}{\partial \phi}=\frac{\partial p}{\partial \phi} \frac{\partial N}{\partial p}+\frac{\partial q}{\partial \phi} \frac{\partial N}{\partial q}
\end{aligned}
$$

Since we have the following equations:

$$
\begin{gathered}
\frac{\partial p}{\partial \rho}=\cos \phi, \frac{\partial q}{\partial \rho}=\sin \phi \\
\frac{\partial p}{\partial \phi}=-\rho(1+\varepsilon) \sin \phi, \frac{\partial q}{\partial \phi}=\rho(1+\varepsilon) \cos \phi
\end{gathered}
$$

we have

$$
\begin{aligned}
\frac{\partial M}{\partial \rho} & =\cos \phi \frac{\partial M}{\partial p}+\sin \phi \frac{\partial M}{\partial q} \\
\frac{\partial M}{\partial \phi} & =-\rho \sin \phi(1+\varepsilon) \frac{\partial M}{\partial p}+\rho \cos \phi(1+\varepsilon) \frac{\partial M}{\partial q} \\
\frac{\partial N}{\partial \rho} & =\cos \phi \frac{\partial N}{\partial p}+\sin \phi \frac{\partial N}{\partial q}
\end{aligned}
$$

$$
\frac{\partial N}{\partial \phi}=-\rho \sin \phi(1+\varepsilon) \frac{\partial N}{\partial p}+\rho \cos \phi(1+\varepsilon) \frac{\partial N}{\partial q}
$$

where

$$
\begin{aligned}
& \frac{\partial M}{\partial p}=\frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{1}\right)}+\varepsilon \frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{1}\right)}+\frac{\partial S c\left(p_{2}\right)}{\partial p}\left\{\frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{2}\right)}+\varepsilon \frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{2}\right)}\right\}, \\
& \frac{\partial M}{\partial q}=\frac{\partial S c\left(p_{1}\right)}{\partial q}\left\{\frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{1}\right)}+\varepsilon \frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{1}\right)}\right\}-\frac{d}{2}\left\{\frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{2}\right)}+\varepsilon \frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{2}\right)}\right\}, \\
& \frac{\partial N}{\partial p}=-\frac{2}{d} \frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{1}\right)}+\varepsilon \frac{d}{2} \frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{1}\right)}+\frac{\partial S c\left(p_{2}\right)}{\partial p}\left\{-\frac{2}{d} \frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{2}\right)}+\varepsilon \frac{d}{2} \frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{2}\right)}\right\}, \\
& \frac{\partial N}{\partial q}=\frac{\partial S c\left(p_{1}\right)}{\partial q}\left\{-\frac{2}{d} \frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{1}\right)}+\varepsilon \frac{d}{2} \frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{1}\right)}\right\}-\frac{d}{2}\left\{-\frac{2}{d} \frac{\partial S c\left(f_{2}\right)}{\partial S c\left(p_{2}\right)}+\varepsilon \frac{d}{2} \frac{\partial S c\left(f_{1}\right)}{\partial S c\left(p_{2}\right)}\right\} .
\end{aligned}
$$

From Theorem 3.3, we have the following equations by comparing with the equations (3.7) and the derivative of $S c\left(f_{1}\right)$ and $S c\left(f_{2}\right)$ for $S c\left(p_{1}\right)$ and $S c\left(p_{2}\right)$ :

$$
\frac{\partial M}{\partial q}=\frac{\partial N}{\partial p} \quad \text { and } \quad \frac{\partial M}{\partial p}=\frac{\partial N}{\partial q}
$$

Therefore, the equations (3.8) are obtained.
Example 3.5. Let $F(Z)=Z=\rho \cos \phi+\mathbf{v} \rho \sin \phi$ on $\Omega$ in $\mathbb{H}^{2}$. Then we have

$$
\begin{gathered}
\frac{\partial N}{\partial \phi}=\frac{\partial(\rho \sin \phi)}{\partial \phi}=\frac{1}{2} \rho \cos \phi+\frac{\varepsilon}{2} \rho \cos \phi \\
\frac{\partial M}{\partial \rho}=\cos \phi, \frac{\partial N}{\partial \rho}=\cos \phi
\end{gathered}
$$

and

$$
\frac{\partial M}{\partial \phi}=\frac{\partial(\rho \cos \phi)}{\partial \phi}=-\frac{1}{2} \rho \sin \phi-\frac{\varepsilon}{2} \rho \sin \phi
$$

Therefore, the function $F$ satisfies the equations (3.8).

## References

[1] W. K. Clifford, Preliminary sketch of bi-quaternions, Proc. London Math. Soc. 4 (1873), 381-395.
[2] K. Daniilidis, Hand-eye calibration using dual quaternions, Int. J. Rob. Res. 18 (1999), no. 3, 286-298.
[3] Z. Ercan and S. Yüce, On Properties of the Dual Quaternions, Eur. J. Pure Appl. Math. 4 (2011), no. 2, 142-146.
[4] G. Helzer, Special Relativity with acceleration, Amer. Math. Monthy 107 (2000), no. 3, 215-237.
[5] J. Kajiwara, X. D. Li, and K. H. Shon, Regeneration in complex, quaternion and Clifford analysis, In Finite or Infinite Dimensional Complex Analysis and Applications, pp. 287298, Springer US, New York, USA, 2004.
[6] , Function spaces in complex and Clifford analysis, In Finite or Infinite Dimensional Complex Analysis and Applications, pp. 127-155, Springer US, New York, USA, 2006.
[7] B. Kenwright, A beginners guide to dual-quaternions: what they are, how they work, and how to use them for 3D character hierarchies, wscg.zcu.cz/wscg2012/short/a29-full.pdf.
[8] , Inverse kinematics with dual-quaternions, exponential-maps, and joint limits, Int. J. Adv. Sys. 6 (2013), no. 1-2.
[9] J. E. Kim, S. J. Lim, and K. H. Shon, Regular functions with values in ternary number system on the complex Clifford analysis, Abstr. Appl. Anal. 2013 (2013), Artical ID 136120, 7 pages.
[10] -, Regularity of functions on the reduced quaternion field in Clifford analysis, Abstr. Appl. Anal. 2014 (2014), Artical ID 654798, 8 pages.
[11] J. E. Kim and K. H. Shon, The Regularity of functions on Dual split quaternions in Clifford analysis, Abstr. Appl. Anal. 2014 (2014), Artical ID 369430, 8 pages.
[12] _, Polar Coordinate Expression of Hyperholomorphic Functions on Split Quaternions in Clifford Analysis, Adv. Appl. Clifford Algebr. 25 (2015), no. 4, 915-924.
[13] , Coset of a hypercomplex numbers in Clifford analysis, Bull. Korean Math. Soc. 52 (2015), no. 5, 1721-1728.
[14] , Properties of regular functions with values in bicomplex numbers, Bull. Korean Math. Soc. 53 (2016), no. 2, 507-518.
[15] _, Inverse Mapping Theory on Split Quaternions in Clifford Analysis, To appear in Filomat.
[16] K. Nono, Hyperholomorphic functions of a quaternion variable, Bull. Fukuoka Univ. Ed. 32 (1983), 21-37.
[17] E. Pennestrì and R. Stefanelli, Linear algebra and numerical algorithms using dual numbers, Multibody Syst. Dyn. 18 (2007), no. 3, 323-344.
[18] J. Plucker, On a new geometry of space, Phil. Trans. Roy. Soc. Lond. 155 (1865), 725791.
[19] R. M. Porter, Quaternionic linear and quadratic equations, J. Nat. Geom. 11 (1997), no. 2, 101-106.
[20] E. Study, Geometrie der Dynamen, Leipzig, Germany, 1903.

## Ji Eun Kim

Department of Mathematics
Dongguk University
Gyeonguu-si 38066, Korea
E-mail address: jeunkim@pusan.ac.kr
Kwang Ho Shon
Department of Mathematics
Pusan National University
Busan 46241, Korea
E-mail address: khshon@pusan.ac.kr

