

EINSTEIN'S CONNECTION IN 5-DIMENSIONAL *ES*-MANIFOLD

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ABSTRACT. The manifold $*g-ESX_n$ is a generalized n -dimensional Riemannian manifold on which the differential geometric structure is imposed by the unified field tensor $*g^{\lambda\nu}$ through the *ES*-connection which is both Einstein and semi-symmetric. The purpose of the present paper is to prove a necessary and sufficient condition for a unique Einstein's connection to exist in 5-dimensional $*g-ESX_5$ and to display a surveyable tensorial representation of 5-dimensional Einstein's connection in terms of the unified field tensor, employing the powerful recurrence relations in the first class.

1. Preliminaries

This paper is a direct continuation of our previous paper [1], which will be denoted by I in the present paper. All considerations in this paper are based on the results and symbolism of I . Whenever necessary, they will be quoted in the present paper. In this section, we introduce a brief collection of basic concepts, notations, and results of I , which are frequently used in the present paper([2],[3],[4]).

(a) n -dimensional $*g$ -unified field theory

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Let X_n be an n -dimensional generalized Riemannian manifold referred to a real coordinate system x^ν , which obeys the coordinate transformations $x^\nu \rightarrow x^{\nu'}$ for which

$$(1.1) \quad \det\left(\frac{\partial x'}{\partial x}\right) \neq 0$$

In $n - g - UFT$ the manifold X_n is endowed with a real nonsymmetric tensor $g_{\lambda\mu}$, which may be decomposed into its symmetric part $h_{\lambda\mu}$ and skew-symmetric part $k_{\lambda\mu}$:

$$(1.2) \quad g_{\lambda\mu} = h_{\lambda\mu} + k_{\lambda\mu}$$

where

$$(1.3) \quad \mathfrak{g} = \det(g_{\lambda\mu}) \neq 0, \quad \mathfrak{h} = \det(h_{\lambda\mu}) \neq 0, \quad \mathfrak{k} = \det(k_{\lambda\mu})$$

In $n - *g - UFT$ the algebraic structure on X_n is imposed by the basic real tensor $*g^{\lambda\nu}$ defined by

$$(1.4) \quad g_{\lambda\mu} *g^{\lambda\nu} = g_{\mu\lambda} *g^{\nu\lambda} = \delta_\mu^\nu$$

It may be also decomposed into its symmetric part $*h^{\lambda\nu}$ and skew-symmetric part $*k^{\lambda\nu}$:

$$(1.5) \quad *g^{\lambda\nu} = *h^{\lambda\nu} + *k^{\lambda\nu}$$

Since $\det(*h^{\lambda\nu}) \neq 0$, we may define a unique tensor $*h_{\lambda\mu}$ by

$$(1.6) \quad *h_{\lambda\mu} *h^{\lambda\nu} = \delta_\mu^\nu$$

In $n - *g - UFT$ we use both $*h^{\lambda\nu}$ and $*h_{\lambda\mu}$ as tensors for raising and/or lowering indices of all tensors in X_n in the usual manner. We then have

$$(1.7) \quad *k_{\lambda\mu} = *k^{\rho\sigma} *h_{\lambda\rho} *h_{\mu\sigma}, \quad *g_{\lambda\mu} = *g^{\rho\sigma} *h_{\lambda\rho} *h_{\mu\sigma}$$

so that

$$(1.8) \quad *g_{\lambda\mu} = *h_{\lambda\mu} + *k_{\lambda\mu}$$

The differential geometric structure on X_n is imposed by the tensor $*g^{\lambda\nu}$ by means of a connection $\Gamma_{\lambda}^{\nu}{}_{\mu}$ defined by a system of equations

$$(1.9) \quad D_\omega *g^{\lambda\nu} = -2S_{\omega\alpha}{}^\nu *g^{\lambda\alpha}$$

where D_ω denotes the symbol of the covariant derivative with respect to $\Gamma_{\lambda}^{\nu}{}_{\mu}$ and $S_{\lambda\mu}{}^\nu$ is the torsion tensor of $\Gamma_{\lambda}^{\nu}{}_{\mu}$. Under certain conditions the system (1.9) admits a unique solutions $\Gamma_{\lambda}^{\nu}{}_{\mu}$.

It has been shown in [5] that if the system (1.9) admits $\Gamma_{\lambda}^{\nu}{}_{\mu}$, it must be of the form

$$(1.10) \quad \Gamma_{\lambda}^{\nu}{}_{\mu} = * \left\{ \begin{array}{c} \nu \\ \lambda\mu \end{array} \right\} + U^{\nu}{}_{\lambda\mu} + S_{\lambda\mu}{}^{\nu}.$$

where

$$(1.11) \quad U_{\nu\lambda\mu} = S_{(\lambda\mu)\nu}^{100} + 2 S_{\nu(\lambda\mu)}^{(10)0}$$

(b) Some notations and results

The following quantities are frequently used in our further considerations:

$$(1.12) \quad *g = \det(*g_{\lambda\mu}), \quad *h = \det(*h_{\lambda\mu}), \quad *k = \det(*k_{\lambda\mu})$$

$$(1.13) \quad *g = \frac{*g}{*h}, \quad *k = \frac{*k}{*h}.$$

$$(1.14) \quad K_p = *k_{[\alpha_1}{}^{\alpha_1} *k_{\alpha_2}{}^{\alpha_2} \dots *k_{\alpha_p]}{}^{\alpha_p}, \quad (p = 0, 1, 2, \dots).$$

$$(1.15) \quad {}^{(0)}*k_{\lambda}{}^{\nu} = \delta_{\lambda}^{\nu}, \quad {}^{(p)}*k_{\lambda}{}^{\nu} = *k_{\lambda}{}^{\alpha} {}^{(p-1)}*k_{\alpha}{}^{\nu} \quad (p = 1, 2, \dots).$$

$$(1.16) \quad K_{\omega\mu\nu} = \nabla_{\nu} *k_{\omega\mu} + \nabla_{\omega} *k_{\nu\mu} + \nabla_{\mu} *k_{\omega\nu}$$

where ∇_{ω} is the symbolic vector of the covariant derivative with respect to the christoffel symbols $* \left\{ \begin{array}{c} \nu \\ \lambda\mu \end{array} \right\}$ defined by $*h_{\lambda\mu}$ in the usual way.

In X_n it was proved in [5] that

$$(1.17) \quad K_0 = 1, \quad K_n = *k \text{ if } n \text{ is even, and } K_n = 0 \text{ if } n \text{ is odd.}$$

$$(1.18) \quad *g = 1 + K_2 + \dots + K_{n-\sigma}.$$

$$(1.19) \quad \sum_{s=0}^{n-\sigma} K_s {}^{(n-s)}*k_{\lambda}{}^{\nu} = 0 \quad (p = 0, 1, 2, \dots).$$

We also use the following useful abbreviations, denoting an arbitrary tensor $T_{\omega\mu\nu}$ skew-symmetric in the first two indices by T :

$$(1.20) \quad \overset{pqr}{T} = \overset{pqr}{T}_{\omega\mu\lambda} = T_{\alpha\beta\gamma} {}^{(p)}*k_{\omega}{}^{\alpha} {}^{(q)}*k_{\mu}{}^{\beta} {}^{(r)}*k_{\lambda}{}^{\gamma}$$

and for an arbitrary tensor T_{\dots} for $p = 1, 2, 3, \dots$:

$$(1.21) \quad {}^{(p)}T_{\dots}^{\nu\dots} = {}^{(p-1)} * k_{\alpha}^{\nu} T_{\dots}^{\alpha\dots}.$$

On the other hand, it has shown in [6] that the tensor $S_{\lambda\mu}{}^{\nu}$ satisfies

$$(1.22) \quad S = B - 3 \overset{(110)}{S}$$

where

$$(1.23) \quad 2B_{\omega\mu\nu} = K_{\omega\mu\nu} + 3K_{\alpha[\mu\beta} * k_{\omega]}^{\alpha} k_{\nu}^{\beta}$$

In our subsequent chapter, we start with the relation (1.22) to solve the system (1.9). Furthermore, for the first class, the nonholonomic solution of (1.22) may be given by

$$(1.24) \quad M S_{xyz} = B_{xyz}$$

or equivalently

$$(1.25) \quad 4M S_{xyz} = (2 + \underset{z}{M} \underset{x}{M} + \underset{z}{M} \underset{y}{M}) K_{xyz} + \underset{z}{M} (\underset{x}{M} + \underset{z}{M}) K_{zxy} + \underset{z}{M} (\underset{y}{M} + \underset{z}{M}) K_{yzx}$$

where

$$(1.26) \quad M = 1 + \underset{xyz}{M} \underset{x}{M} \underset{y}{M} + \underset{y}{M} \underset{z}{M} \underset{x}{M} + \underset{z}{M} \underset{x}{M} \underset{y}{M}$$

Therefore, in virtue of (1.24), we see that a necessary and sufficient condition for the system (1.9) to have a unique solution in the first class is

$$(1.27) \quad M \neq 0 \text{ for all } x, y, z$$

DEFINITION 1.1. A connection $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$ is said to be *semi-symmetric* if its torsion tensor $S_{\lambda\mu}{}^{\nu}$ is of the form

$$(1.28) \quad S_{\lambda\mu}{}^{\nu} = 2\delta_{[\lambda}^{\nu} X_{\mu]}.$$

for an arbitrary non-null vector X_{μ} .

A connection which is both semi-symmetric and Einstein is called an ES connection. An n -dimensional generalized Riemannian manifold X_n , on which the differential geometric structure is imposed by $*g^{\lambda\nu}$ by means of an ES connection, is called an n -dimensional $*g-ES$ manifold. We denote this manifold by $*g-ESX_n$ in our further considerations.

In $*g-ESX_5$, the following theorems were proved in I .

THEOREM 1.2. *The basic scalars in $*g-ESX_5$ may be given by*

$$(1.29) \quad \begin{aligned} M_1 &= -M_2 = \sqrt{-L - K} \neq 0 \\ M_3 &= -M_4 = \sqrt{L - K} \neq 0, \quad M_5 = 0 \end{aligned}$$

where

$$(1.30) \quad K = \frac{K_2}{2}, \quad L = \sqrt{\left(\frac{K_2}{2}\right)^2 - K_4}$$

THEOREM 1.3. *The main recurrence relation in the first class is*

$$(1.31) \quad {}^{(p+5)*}k_{\lambda}{}^{\nu} = -K_2 {}^{(p+3)*}k_{\lambda}{}^{\nu} - K_4 {}^{(p+1)*}k_{\lambda}{}^{\nu}, \quad (p = 0, 1, 2, \dots)$$

THEOREM 1.4. *The basic scalars M_x satisfy*

$$(1.32) \quad M_1 + M_2 = M_3 + M_4 = 0$$

$$(1.33) \quad M_1 M_5 = M_2 M_5 = M_3 M_5 = M_4 M_5 = 0$$

$$(1.34) \quad M_1^2 M_3^2 = M_1^2 M_4^2 = M_2^2 M_3^2 = M_2^2 M_4^2 = K_4$$

$$(1.35) \quad M_1^2 + M_3^2 = M_1^2 + M_4^2 = M_2^2 + M_3^2 = M_2^2 + M_4^2 = -K_2$$

In virtue of the above theorem, we have

THEOREM 1.5. *In the first class, the following identities hold for all values of x and y when $x \neq y$*

$$(1.36) \quad M_x^{(4} M_y^{1)} = -M_x^{(3} M_y^{2)} - K_2 M_x^{(2} M_y^{1)}$$

$$(1.37) \quad M_x^{(4} M_y^{3)} = K_4 M_x^{(2} M_y^{1)}$$

$$(1.38) \quad M^4_{x \ y} M^4 = K_4^2 M^2_{x \ y} M^2 + K_2 M^3_{x \ y} M^3 + 2K_4 M^3_{x \ y} M^1$$

$$(1.39) \quad 2M^4_{x \ y} M^2 = -M^3_{x \ y} M^3 - K_2 M^2_{x \ y} M^2 + K_4 M M_{x \ y}$$

THEOREM 1.6. (*Recurrence relations in the first class*) If $T_{\omega\mu\nu}$ is a tensor skew-symmetric in the first two indices, then the following recurrence relations hold in the first class of $5 - *g - ESX_5$:

$$(1.40) \quad T^{(41)r} = - T^{(32)r} - K_2 T^{(21)r}$$

$$(1.41) \quad T^{(43)r} = K_4 T^{(21)r}$$

$$(1.42) \quad T^{44r} = K_4 T^{22r} + K_2 T^{33r} + 2K_4 T^{(31)r}$$

$$(1.43) \quad 2 T^{(42)r} = - T^{33r} - K_2 T^{22r} + K_4 T^{11r}$$

2. Einstein's connection $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$ in the first class

In this section, we shall derive surveyable representations of $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$ in terms of $*g^{\lambda\nu}$, employing the recurrence relations.

In the following theorem, we shall prove two relations in X_n . These relations will be used in our subsequent theorem when we are concerned with the solution of (1.9).

THEOREM 2.1. *We have*

$$(2.1) \quad B^{(pq)r} = S^{(pq)r} + S^{(p'q')r} + S^{(p'q)r'} + S^{(pq')r'}$$

$$(2.2) \quad 2 B^{(pq)r}{}_{\omega\mu\nu} = K^{(pq)r}{}_{\omega\mu\nu} + K^{r''(pq)}{}_{\nu[\omega\mu]} + \frac{1}{2} (K^{(pq')r'}{}_{\omega\mu\nu} + K^{(p'q)r'}{}_{\omega\mu\nu} + K^{r'p'q}{}_{\nu[\omega\mu]} + K^{r'q'p}{}_{\nu[\omega\mu]})$$

where

$$(2.3) \quad p' = p + 1, \quad q' = q + 1, \quad r' = r + 1, \quad r'' = r + 2$$

Proof. In virtue of (1.22) and (1.20), the first relation (2.1) is obtained as in the following way:

$$(2.4) \quad \begin{aligned} \overset{(pq)r}{B} &= \overset{(pq)r}{B}_{\omega\mu\nu} = \frac{1}{2}B_{\omega\beta\gamma}(\overset{(p)*}{k}_\omega{}^\alpha\overset{(q)*}{k}_\mu{}^\beta + \overset{(q)*}{k}_\omega{}^\alpha\overset{(p)*}{k}_\mu{}^\beta)\overset{(r)*}{k}_\nu{}^\gamma \\ &= \frac{1}{2}(S_{\alpha\beta\gamma} + S_{\epsilon\eta\gamma}{}^*k_\alpha{}^\epsilon{}^*k_\beta{}^\eta + S_{\epsilon\beta\eta}{}^*k_\alpha{}^\epsilon{}^*k_\gamma{}^\eta + S_{\alpha\epsilon\eta}{}^*k_\beta{}^\epsilon{}^*k_\gamma{}^\eta) \times \\ &\quad \times (\overset{(p)*}{k}_\omega{}^\alpha\overset{(q)*}{k}_\mu{}^\beta + \overset{(q)*}{k}_\omega{}^\alpha\overset{(p)*}{k}_\mu{}^\beta)\overset{(r)*}{k}_\nu{}^\gamma \end{aligned}$$

After a lengthy calculation, we note that the right-hand side of the above equation is equal to that of (2.1). Similarly, we verify (2.2) using (1.20) and (1.23). \square

THEOREM 2.2. *A necessary and sufficient condition for the system (1.9) to admit a unique solution $\Gamma_\lambda{}^\nu{}_\mu$ in the first class is that*

$$(2.5) \quad gAB(C^2 - 4K_4D^2) \neq 0$$

where

$$(2.6) \quad \begin{aligned} A &= 1 - K_2 + K_4, & B &= 1 - K_4 \\ C &= 1 - K_2 + 5K_4, & D &= K_2 - 2 \end{aligned}$$

Proof. In virtue of (1.29) and (1.30), the symmetric scalars M_{xyz} defined by (1.26) takes values as in the following 3 cases:

If two of the indices x, y, z are 1, 2 or 3, 4, then

$$(2.7) \quad M_{xyz} = 1 + K + L, \quad 1 + K - L$$

If at least one of x, y, z is 5 and no two take the values 1, 2 nor 3, 4, then

$$(2.8) \quad M_{xyz} = 1 - K + L, \quad 1 - K - L, \quad 1 + \sqrt{K_4}, \quad 1 - \sqrt{K_4}, \quad 1$$

In the remaining cases,

$$(2.9) \quad \begin{aligned} M_{xyz} &= 1 - K - L - 2\sqrt{K_4}, \quad 1 - K + L - 2\sqrt{K_4} \\ &1 - K - L + 2\sqrt{K_4}, \quad 1 - K + L + 2\sqrt{K_4} \end{aligned}$$

It may easily verified that the product of two factors in the right of (2.7) is g , that of five factors in the right of (2.8) is $(1 - K_2 + K_4)(1 - K_4)$, and that of four factors in the right of (2.9) is $(1 - K_2 + 5K_4)^2 - 4K_4(K_2 - 2)^2$. Hence we have proved our assertion (2.5) in virtue of (1.27) and (2.6). \square

THEOREM 2.3. *The system of equations (1.22) in the first class is reduced to the following 25 equations:*

$$(2.10) \left\{ \begin{array}{l} B = S + S^{110} + 2 S^{(10)1} \\ B = S^{(10)1} + S^{(10)1} + S^{(21)1} + S^{(20)2} + S^{112} \\ B = S^{(12)1} + S^{(12)1} + S^{(23)1} + S^{222} + S^{(13)2} \\ B = S^{(20)2} + S^{(20)2} + S^{(31)2} + S^{(30)3} + S^{(21)3} \\ 2 B = 2 S^{(23)1} + 2 S^{(23)1} + 2 K_4 S^{(21)1} + S^{332} + K_4 S^{112} - K_2 S^{222} \\ 2 B = 2 S^{(13)2} + K_4 S^{(13)2} + K_4 S^{112} - K_2 S^{222} - S^{332} - 2 K_2 S^{(21)3} \\ B = S^{(30)3} - K_2 S^{(30)3} - S^{(21)3} + S^{(32)3} + S^{(40)4} + S^{(31)4} \\ B = S^{(21)3} + S^{(21)3} + S^{(32)3} + S^{(31)4} + S^{224} \\ 2 B = 2 S^{(32)3} + 2 K_4 S^{(32)3} + K_4 S^{(21)3} + K_4 S^{114} - K_2 S^{224} + S^{334} \\ B = S^{(40)4} - K_2 S^{(40)4} - K_4 S^{(31)4} + K_2 S^{114} + K_2 K_4 S^{(30)3} + K_2 K_4 S^{(10)3} + K_2 K_4 S^{(30)1} \\ \quad + K_4 S^{(10)1} + K_2 S^{(21)3} + K_2 K_4 S^{(21)1} + K_2 S^{(32)3} + K_4 S^{(32)1} \\ 2 B = 2 S^{(31)4} + 2 K_2 S^{(31)4} + 2 K_2 K_4 S^{(21)3} + 2 K_2 K_4 S^{(21)1} + K_4 S^{114} - K_2 S^{224} - S^{334} \\ B = S^{(10)3} + S^{(10)3} + S^{(21)3} + S^{(20)4} + S^{114} \\ B = S^{(30)1} - K_2 S^{(30)1} - S^{(21)1} + S^{(32)1} + S^{(40)2} + S^{(31)2} \\ B = S^{(20)4} + S^{(20)4} - K_2 S^{(31)4} - K_4 S^{(30)3} - K_2 S^{(30)1} - K_2 S^{(21)3} - K_4 S^{(21)1} \\ B = S^{(40)2} - K_2 S^{(40)2} - K_4 S^{(31)2} - K_4 S^{112} - K_2 S^{(30)3} - K_4 S^{(10)3} - K_2 S^{(21)3} - S^{(32)3} \\ B = S^{110} + S^{110} + 2 S^{220} + S^{(21)1} \\ B = S^{112} + S^{112} + 2 S^{222} + S^{(21)3222} + S^{222} + S^{332} + 2 S^{(32)3} \\ B = (1 + K_2) S^{332} + K_4 S^{222} + 2 K_4 S^{(21)3} + 2 K_4 S^{(31)2} \\ B = S^{224} + S^{224} + S^{334} - 2 K_2 S^{(32)3} - 2 K_4 S^{(32)1} \end{array} \right.$$

$$\begin{cases} B = S^{114} + S^{114} - 2K_2 S^{(21)3} - 2K_4 S^{(21)1} \\ B = (1 + K_2) S^{334} + K_4 S^{224} + 2K_4 S^{(31)4} - 2K_2 K_4 S^{(21)3} - 2K_4^2 S^{(21)1} \\ B = S^{220} + S^{330} + 2 S^{(32)1} \\ B = (1 + K_2) S^{330} + K_4 S^{220} + 2K_4 S^{(31)0} + 2K_4 S^{(21)1} \\ 2 B = 2 S^{(31)0} + K_4 S^{110} - K_2 S^{220} - S^{330} - 2K_2 S^{(21)1} \end{cases}$$

Proof. This assertion follows from (2.1) using (1.31) and (1.40)-(1.43). \square

References

- [1] I.H. Hwang, *A study on the recurrence relations of 5-dimensional ES-manifold*, Korean J. Math. **24** (3) (2016), 319–330.
- [2] D.k. Datta, *Some theorems on symmetric recurrent tensors of the second order*, Tensor (N.S.) **15** (1964), 1105–1136.
- [3] A. Einstein, *The meaning of relativity*, Princeton University Press, 1950.
- [4] R.S. Mishra, *n-dimensional considerations of unified field theory of relativity*, Tensor **9** (1959), 217–225.
- [5] K.T. Chung, *Einstein's connection in terms of $*g^{\lambda\nu}$* , Nuovo cimento Soc. Ital. Fis. B **27** (1963), (X), 1297–1324
- [6] V. Hlavatý, *Geometry of Einstein's unified field theory*, Noordhoop Ltd., 1957

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