

**POSITIVE RADIAL SOLUTIONS FOR A CLASS OF
ELLIPTIC SYSTEMS CONCENTRATING ON SPHERES
WITH POTENTIAL DECAY**

PAULO CESAR CARRIÃO, NARCISO HORTA LISBOA,
AND OLIMPIO HIROSHI MIYAGAKI

ABSTRACT. We deal with the existence of positive radial solutions concentrating on spheres for the following class of elliptic system

$$(S) \quad \begin{cases} -\varepsilon^2 \Delta u + V_1(x)u = K(x)Q_u(u, v) & \text{in } \mathbb{R}^N, \\ -\varepsilon^2 \Delta v + V_2(x)v = K(x)Q_v(u, v) & \text{in } \mathbb{R}^N, \\ u, v \in W^{1,2}(\mathbb{R}^N), u, v > 0 & \text{in } \mathbb{R}^N, \end{cases}$$

where ε is a small positive parameter; $V_1, V_2 \in C^0(\mathbb{R}^N, [0, \infty))$ and $K \in C^0(\mathbb{R}^N, (0, \infty))$ are radially symmetric potentials; Q is a $(p+1)$ -homogeneous function and p is subcritical, that is, $1 < p < 2^* - 1$, where $2^* = 2N/(N-2)$ is the critical Sobolev exponent for $N \geq 3$.

1. Introduction

This work has been motivated by some papers appeared in recent years concerning the Schrödinger equation

$$(NLS) \quad i\hbar \frac{\partial \psi}{\partial t} + \frac{\hbar^2}{2} \Delta \psi - V(x)\psi + K(x)|\psi|^{p-1}\psi = 0, \quad x \in \mathbb{R}^N,$$

where \hbar denotes the Plank constant, i is the imaginary unit and $p \in (1, \frac{N+2}{N-2})$. This equation appears in many fields of physic, in particular, when we describe Bose-Einstein condensates (see [30] and [34]) and the propagation of light in some nonlinear optical material (see [35]).

For application or motivation, we can cite also, for instance, [32, 33] where are studied the evolution of two orthogonal pulse envelope in birefringent optical fibers, see also [29]. System of type (S) is also important for industrial applications in fiber communications systems [27, 28]. Finally we would to recall that system of type (S) can describe other physical phenomena, such

Received March 18, 2012.

2010 *Mathematics Subject Classification.* 35J50, 35B06, 35A15, 35B25.

Key words and phrases. Schrödinger operator, radial solution, variational method, singular perturbation.

The second author was partially supported by Fapemig/Brazil.

The third author was partially supported by CNPq/Brazil and INCTMat/Brazil.

as Kerr-like photorefractive media in optics, (cf. [1, 20, 21, 22]), Hartree-Fock theory for double condensate [25]. See [31] and [37] for more applications in physical and chemical phenomenas.

Here we are concerned with the existence of standing waves (semiclassical states) of the nonlinear Schrödinger equations for small ε , that is, solutions of the form $\psi(x, t) = \exp(-iEt/\varepsilon)u(x)$. Notice that after a simple rescaling and putting $V(x) - E = V(x)$, ψ satisfies (NLS) if and only if u solve the elliptic equation

$$(NLS)_\varepsilon \quad -\varepsilon^2 \Delta u + V(x)u = K(x)u^p, \quad u > 0, \quad x \in \mathbb{R}^N.$$

The most characteristic feature of $(NLS)_\varepsilon$ is that its solution u_ε concentrate as $\varepsilon \rightarrow 0$. When this concentration set is a single point (resp. finite points), these solutions are called, in the literature, spike solution (resp. multi-bump solutions). When the potential $V > 0$, beginning from the pioneering paper by Floer and Weinstein [24], a great number of work has been devoted to study spike or multi-bump solutions for $(NLS)_\varepsilon$ (see [5] and references there in). Studying in this case ($V > 0$), Ambrosetti-Malchiodi-Ni in [6] constructed solutions concentrating on spheres for $(NLS)_\varepsilon$. Ambrosetti-Ruiz in [9] extended this result to the case of decaying potentials. See also [4], [7], [10], [12], [13], [15] and [23]. In the critical frequency, that means $\inf_{\mathbb{R}^N} V(x) = 0$, spike solutions have been constructed in [16], [17], [18] and [19], which concentrate on the zero of the potential V as $\varepsilon \rightarrow 0$. In those papers also are constructed “small” solutions concentrating on spheres near zeroes of the potentials. On the other hand, Alves [2] and Alves-Soares [3] studied, by using the Mountain Pass Theorem due to Ambrosetti-Rabinowitz [8], the elliptic system (S) , when V_1 and V_2 are globally lower bounded away from zero. The authors showed that the solution $(u_\varepsilon, v_\varepsilon)$ concentrates around local minima of the potentials V_1 and V_2 .

Motivated by the above papers, we are going to construct solutions concentrating on spheres for a class of the elliptic system with decaying potentials, where V_1, V_2 and K are radially symmetric potentials satisfying:

(V) $V_1, V_2 \in C^0(\mathbb{R}^N, [0, \infty))$ are such that

$$\liminf_{|x| \rightarrow \infty} |x|^2 V(x) \equiv 4\lambda > 0,$$

where $V(x) = \min\{V_1(x), V_2(x)\}$ and the zero set of V , $Z = \{x \in \mathbb{R}^N : V(x) = 0\}$ is non-empty;

(K) $K \in C^0(\mathbb{R}^N, (0, \infty))$ is limited.

The function $Q \in C^1([0, +\infty) \times [0, +\infty), \mathbb{R})$ is a homogeneous function of degree $p + 1$, with $1 < p < \frac{N+2}{N-2}$, $N \geq 3$ and verify:

(Q₁) There exists $C > 0$ such that

$$\begin{cases} |Q_u(u, v)| \leq C(|u|^p + |v|^p), \forall u, v \geq 0, \\ |Q_v(u, v)| \leq C(|u|^p + |v|^p), \forall u, v \geq 0; \end{cases}$$

(Q₂) There exist $\eta_1, \eta_2 > 0$ such that

$$\eta_1(|u|^{p+1} + |v|^{p+1}) \leq Q(u, v) \leq \eta_2(|u|^{p+1} + |v|^{p+1}) \quad \forall u, v > 0;$$

(Q₃) $Q_u(0, 1), Q_v(1, 0) > 0$;

(Q₄) $Q(u, v) > 0 \quad \forall u, v > 0$;

(Q₅) $Q_u(u, v), Q_v(u, v) \geq 0 \quad \forall u, v \geq 0$.

Remark 1. (a) Since Q is a C^1 homogeneous function of degree $p + 1$, then $(p + 1)Q(u, v) = uQ_u(u, v) + vQ_v(u, v)$ and ∇Q is a homogeneous function of degree p .

(b) Note that the right hand side of (Q₂) can be obtained from (Q₁), (a) and the Young inequality.

(c) These kind of hypotheses were introduced for instance in [2] and [36].

(d) Our prototype of Q is $Q(u, v) = (au + bv)^{p+1}$, $u, v \geq 0$ and $a, b > 0$.

Our main result is the following.

Theorem 1. *Suppose that (Q₁)-(Q₅), (V) and (K) hold. Let $A \subset Z$ be an isolated compact subset of Z such that $0 \notin A$ and $V_1 \equiv V_2$ in A . Then for ε sufficiently small, (S) has a solution $(u_\varepsilon, v_\varepsilon) \in W^{1,2}(\mathbb{R}^N) \times W^{1,2}(\mathbb{R}^N)$, u_ε and v_ε radially symmetric functions, such that*

$$(1) \quad \lim_{\varepsilon \rightarrow 0} \|u_\varepsilon\|_{L^\infty(\mathbb{R}^N)} = \lim_{\varepsilon \rightarrow 0} \|v_\varepsilon\|_{L^\infty(\mathbb{R}^N)} = 0$$

and

$$(2) \quad \liminf_{\varepsilon \rightarrow 0} \varepsilon^{-2/(p-1)} \|u_\varepsilon + v_\varepsilon\|_{L^\infty(\mathbb{R}^N)} > 0.$$

Moreover, for each $\delta > 0$, there are constants $C, c > 0$ such that

$$(3) \quad u_\varepsilon(x), v_\varepsilon(x) \leq C \exp(-c/\varepsilon) [1 + (|x|/2R_0)^{\omega_\varepsilon}] \quad \forall x \in \mathbb{R}^N \setminus A^{4\delta},$$

where $\omega_\varepsilon \equiv -\frac{(N-2) + \sqrt{(N-2)^2 + 4\lambda/\varepsilon^2}}{2}$, $A^d \equiv \{x \in \mathbb{R}^N \mid d(x, A) \leq d\}$ and R_0 is a positive constant given by (V).

The proof of Theorem 1 is made adapting closely arguments used in [16] and [17], more exactly, the minimization techniques with two constraints in order to construct the spike solutions concentrating on sphere near of the zeros of V_1 and V_2 . Actually, one of the constraints represents a type of the penalization of the nonlinearity. The proof of the decay estimate of the solution is slightly different those made in [16] and [17]. Here, in our case, we use some ideas in [11], as well as, those in [16] and [17], combining Moser iterations, classical elliptic estimates and comparison principle we obtain the decay estimate of the solutions desired.

2. Proof of Theorem 1

First of all by a scaling we see that system (S) is equivalent to

$$(\tilde{S}) \quad \begin{cases} -\Delta u + V_1(\varepsilon x)u = K(\varepsilon x)Q_u(u, v) \text{ in } \mathbb{R}^N, \\ -\Delta v + V_2(\varepsilon x)v = K(\varepsilon x)Q_v(u, v) \text{ in } \mathbb{R}^N, \\ u, v \in W^{1,2}(\mathbb{R}^N), u, v > 0 \text{ in } \mathbb{R}^N. \end{cases}$$

Let A be the isolated compact subset of Z as assumed in the theorem. We choose $\delta > 0$ such that $0 \notin A^{\delta}$, and $A^{8\delta} \cap (Z \setminus A) = \emptyset$, where $A^\delta \equiv \{x \in \mathbb{R}^N \mid d(x, A) \leq \delta\}$. We define $A_\varepsilon^\delta \equiv \{x \in \mathbb{R}^N \mid \varepsilon x \in A^\delta\}$. Let $C_{0,rad}^\infty(\mathbb{R}^N)$ be the class of radially symmetric functions in $C_0^\infty(\mathbb{R}^N)$, where $C_0^\infty(\mathbb{R}^N)$ is the set of functions on $C^\infty(\mathbb{R}^N)$ with compact support. Let E_ε the completion of $C_{0,rad}^\infty(\mathbb{R}^N) \times C_{0,rad}^\infty(\mathbb{R}^N)$ with respect to the norm

$$\|(u, v)\|_\varepsilon = \left(\int_{\mathbb{R}^N} [|\nabla u|^2 + |\nabla v|^2 + V_1(\varepsilon x)u^2 + V_2(\varepsilon x)v^2] dx \right)^{1/2}.$$

We observe that $E_\varepsilon = E_{V_1,\varepsilon} \times E_{V_2,\varepsilon}$, where $E_{V_i,\varepsilon}$ is the completion of $C_{0,rad}^\infty(\mathbb{R}^N)$ with the norm $\|u\|_{V_i,\varepsilon} = \left(\int_{\mathbb{R}^N} [|\nabla u|^2 + V_i(\varepsilon x)u^2] dx \right)^{1/2}$, $i = 1, 2$. Thus, $\|(u, v)\|_\varepsilon^2 = \|u\|_{V_1,\varepsilon}^2 + \|v\|_{V_2,\varepsilon}^2$.

We fix a constant γ with $\gamma(p-1)/(p+1) > 2$. We define a function χ_ε by

$$\chi_\varepsilon(x) = \begin{cases} \varepsilon^{-(N-1)-3(p+1)/(p-1)} & \text{if } |x| \leq R_0/\varepsilon, x \notin A_\varepsilon^{4\delta}, \\ (|x|/\varepsilon)^\gamma & \text{if } |x| \geq R_0/\varepsilon, \\ 0 & \text{if } x \in A_\varepsilon^{4\delta}, \end{cases}$$

where $R_0 \geq 1$ is fixed so that $V(x) > 0$ for $|x| \geq R_0$ and $Z^{8\delta} \subset B(0, R_0)$.

Now we consider the following minimization problem

$$(4) \quad M_\varepsilon = \inf \left\{ \|(u, v)\|_\varepsilon^2 \mid \int_{\mathbb{R}^N} K(\varepsilon x)Q(u, v) dx = 1, \int_{\mathbb{R}^N} \chi_\varepsilon(x)Q(u, v) dx \leq 1, (u, v) \in E_\varepsilon \right\}.$$

First, using the same type of arguments developed in [16], we have the following lemma.

Lemma 2. $\lim_{\varepsilon \rightarrow 0} \varepsilon^{(N-1)(p-1)/(p+1)} M_\varepsilon = 0$.

Proof. Let $x_0 \in A$. Then, for any $a > 0$, there exists $b > 0$ such that $V_1(x), V_2(x) \in [0, a]$ for $|x - x_0| \leq b$. Without loss of generality, we can assume $|x_0| = 1$ so that $S_\varepsilon^\delta \subset A_\varepsilon^\delta$, where S is the unit sphere in \mathbb{R}^N . Then, using change of variables (polar coordinates) and setting $u(r + 1/\varepsilon) = \bar{u}(r)$, $v(r + 1/\varepsilon) = \bar{v}(r)$, we obtain that

$$M_\varepsilon \leq C_0 \frac{\int_{S_\varepsilon^\delta} [|\nabla u(x)|^2 + |\nabla v(x)|^2 + a((u(x))^2 + (v(x))^2)] dx}{\left(\int_{S_\varepsilon^\delta} Q(u(x), v(x)) dx \right)^{2/(p+1)}}$$

$$\leq C\varepsilon^{-(N-1)(p-1)/(p+1)} \times \frac{\int_{-\delta/\varepsilon}^{\delta/\varepsilon} [(\bar{u}'(r))^2 + (\bar{v}'(r))^2 + a((\bar{u}(r))^2 + (\bar{v}(r))^2)] dr}{\left(\int_{-\delta/\varepsilon}^{\delta/\varepsilon} Q(\bar{u}(r), \bar{v}(r)) dr\right)^{2/(p+1)}}$$

where C_0 and C are positive constants independent of ε . Here was used that $\chi_\varepsilon(x) = 0, \forall x \in S_\varepsilon^\delta$; V_1 and V_2 are radially symmetric, and $V_1(\varepsilon x), V_2(\varepsilon x) < a, \forall x \in S_\varepsilon^\delta$. Setting $\bar{u}(r) = u(\sqrt{ar})$ and $\bar{v}(r) = v(\sqrt{ar})$, where $a > 0$ is arbitrary, we obtain,

$$\lim_{\varepsilon \rightarrow 0} \varepsilon^{(N-1)(p-1)/(p+1)} M_\varepsilon \leq C a^{(p+3)/2(p+1)} \inf_{u, v \in C_0^\infty(-\infty, \infty)} \frac{\int_{-\infty}^{\infty} [(u')^2 + (v')^2 + u^2 + v^2] dr}{\left(\int_{-\infty}^{\infty} Q(u, v) dr\right)^{2/(p+1)}}$$

Then, since a is arbitrary and the last infimum is bounded, the lemma follows. □

Lemma 3. For sufficiently small $\varepsilon > 0, M_\varepsilon$ is achieved at $(\bar{u}_\varepsilon, \bar{v}_\varepsilon) \in E_\varepsilon$ which satisfies for some $\alpha_\varepsilon > 0 \geq \beta_\varepsilon$,

$$(S_{\alpha_\varepsilon, \beta_\varepsilon}) \begin{cases} -\Delta \bar{u}_\varepsilon + V_1(\varepsilon x) \bar{u}_\varepsilon = \alpha_\varepsilon K(\varepsilon x) Q_u(\bar{u}_\varepsilon, \bar{v}_\varepsilon) + \beta_\varepsilon \chi_\varepsilon(x) Q_u(\bar{u}_\varepsilon, \bar{v}_\varepsilon) & \text{in } \mathbb{R}^N, \\ -\Delta \bar{v}_\varepsilon + V_2(\varepsilon x) \bar{v}_\varepsilon = \alpha_\varepsilon K(\varepsilon x) Q_v(\bar{u}_\varepsilon, \bar{v}_\varepsilon) + \beta_\varepsilon \chi_\varepsilon(x) Q_v(\bar{u}_\varepsilon, \bar{v}_\varepsilon) & \text{in } \mathbb{R}^N, \\ \bar{u}_\varepsilon \geq 0, \bar{v}_\varepsilon \geq 0 & \text{in } \mathbb{R}^N. \end{cases}$$

Proof. Let $\{(\bar{u}_\varepsilon^j, \bar{v}_\varepsilon^j)\}_j \subset E_\varepsilon$ be a minimizing sequence for M_ε . We can assume $\{(\bar{u}_\varepsilon^j, \bar{v}_\varepsilon^j)\}_j \subset C_{0,rad}^\infty(\mathbb{R}^N) \times C_{0,rad}^\infty(\mathbb{R}^N)$, since $C_{0,rad}^\infty(\mathbb{R}^N) \times C_{0,rad}^\infty(\mathbb{R}^N)$ is dense in E_ε . We take $R_j > 0$ such that $\text{supp}(\bar{u}_\varepsilon^j) \subset B(0, R_j)$ and $\text{supp}(\bar{v}_\varepsilon^j) \subset B(0, R_j), j \geq 1$. For a fixed $\varepsilon > 0$, we can assume that $R_0/\varepsilon < R_1 < R_2 < \dots$ and $\lim_{m \rightarrow \infty} R_m = \infty$. We define

$$E_\varepsilon^m \equiv E_\varepsilon \cap \left(W_0^{1,2}(B(0, R_m)) \times W_0^{1,2}(B(0, R_m)) \right).$$

We consider a restricted minimization problem

$$(5) \quad M_\varepsilon^m = \inf \left\{ \|(u, v)\|_\varepsilon^2 \mid \int_{\mathbb{R}^N} K(\varepsilon x) Q(u, v) dx = 1, \int_{\mathbb{R}^N} \chi_\varepsilon(x) Q(u, v) dx \leq 1, (u, v) \in E_\varepsilon^m \right\}.$$

Now, we will prove that there exists a non-negative minimizer $(u_\varepsilon^m, v_\varepsilon^m)$ of M_ε^m such that $M_\varepsilon \leq M_\varepsilon^m$ and $\lim_{m \rightarrow \infty} M_\varepsilon^m = M_\varepsilon$. Indeed, let $\{(u_\varepsilon^k, v_\varepsilon^k)\}_k$ be a minimizing sequence for M_ε^m . Then it follows that $\{(u_\varepsilon^k, v_\varepsilon^k)\}_k$ is bounded. Since E_ε^m is reflexive, there exists $(u_\varepsilon^m, v_\varepsilon^m) \in E_\varepsilon^m$ such that $\{(u_\varepsilon^k, v_\varepsilon^k)\}_k$ is weakly convergent to $(u_\varepsilon^m, v_\varepsilon^m)$, up to subsequence. Thus, $u_\varepsilon^k \rightharpoonup u_\varepsilon^m$ weakly in $E_{V_1, \varepsilon}$ and $v_\varepsilon^k \rightharpoonup v_\varepsilon^m$ weakly in $E_{V_2, \varepsilon}$ as $k \rightarrow \infty$. Since $E_{V_i, \varepsilon} \cap W_0^{1,2}B((0, R_m))$

is compactly imbedded in $L^{p+1}(B(0, R_m))$, with $i = 1, 2$ and $2 < p + 1 < 2^*$, from (Q_2) we have

$$(6) \quad \int_{B(0, R_m)} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m)dx = \lim_{k \rightarrow \infty} \int_{B(0, R_m)} K(\varepsilon x)Q(u_\varepsilon^k, v_\varepsilon^k)dx = 1,$$

and

$$\int_{B(0, R_m)} \chi_\varepsilon Q(u_\varepsilon^m, v_\varepsilon^m)dx = \lim_{k \rightarrow \infty} \int_{B(0, R_m)} \chi_\varepsilon Q(u_\varepsilon^k, v_\varepsilon^k)dx \leq 1.$$

Since $\{(u_\varepsilon^k, v_\varepsilon^k)\}_k$ is weakly convergent to $(u_\varepsilon^m, v_\varepsilon^m)$, we have

$$\|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2 \leq \liminf_{k \rightarrow \infty} \|(u_\varepsilon^k, v_\varepsilon^k)\|_\varepsilon^2 = M_\varepsilon^m \leq \|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2.$$

Thus, $(u_\varepsilon^m, v_\varepsilon^m)$ is a minimizer for M_ε^m . Since $|\nabla |u_\varepsilon^m|| = |\nabla u_\varepsilon^m|$ and $|\nabla |v_\varepsilon^m|| = |\nabla v_\varepsilon^m|$ we see that $\|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2 = \||u_\varepsilon^m|, |v_\varepsilon^m|\|_\varepsilon^2$. Then, there exists a non-negative minimizer $(u_\varepsilon^m, v_\varepsilon^m)$ of M_ε^m . Now, we observe that for any $j \geq 1$,

$$\lim_{k \rightarrow \infty} \|(u_\varepsilon^k, v_\varepsilon^k)\|_\varepsilon^2 \leq \|(u_\varepsilon^j, v_\varepsilon^j)\|_\varepsilon^2.$$

In fact, for any $j \leq k$, $B(0, R_j) \subset B(0, R_k)$. Thus,

$$W_0^{1,2}(B(0, R_j)) \subset W_0^{1,2}(B(0, R_k)).$$

Consequently, $E_\varepsilon^j \subset E_\varepsilon^k$. This implies that $M_\varepsilon^j = \|(u_\varepsilon^j, v_\varepsilon^j)\|_\varepsilon^2 \geq \|(u_\varepsilon^k, v_\varepsilon^k)\|_\varepsilon^2 = M_\varepsilon^k$. We note that

$$M_\varepsilon \leq \lim_{j \rightarrow \infty} M_\varepsilon^j = \lim_{j \rightarrow \infty} \|(u_\varepsilon^j, v_\varepsilon^j)\|_\varepsilon^2 \leq \lim_{j \rightarrow \infty} \|(\bar{u}_\varepsilon^j, \bar{v}_\varepsilon^j)\|_\varepsilon^2 = M_\varepsilon.$$

Therefore, $M_\varepsilon^m \rightarrow M_\varepsilon$ as $m \rightarrow \infty$. Thus $\{(u_\varepsilon^m, v_\varepsilon^m)\}_m$ is a minimizing sequence for M_ε .

Since $(u_\varepsilon^m, v_\varepsilon^m)$ is a minimizer for M_ε^m , there exist Lagrange multipliers $\alpha_\varepsilon^m, \beta_\varepsilon^m \in \mathbb{R}$ such that $(u_\varepsilon^m, v_\varepsilon^m)$ satisfies the system $(S_{\alpha_\varepsilon^m, \beta_\varepsilon^m})$ in $B(0, R_m)$. Taking a subsequence if necessary, we can assume that for some $(\bar{u}_\varepsilon, \bar{v}_\varepsilon) \in E_\varepsilon$, $\{(u_\varepsilon^m, v_\varepsilon^m)\}_m$ converges weakly to $(\bar{u}_\varepsilon, \bar{v}_\varepsilon)$ in E_ε as $m \rightarrow \infty$. Since

$$\int_{\mathbb{R}^N} \chi_\varepsilon Q(u_\varepsilon^m, v_\varepsilon^m)dx \leq 1,$$

it follows that for any $R \geq \frac{R_0}{\varepsilon}$,

$$\int_{\mathbb{R}^N \setminus B(0, R)} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m)dx \leq C(\varepsilon/R)^\gamma$$

for some $C > 0$. By the Dominated Convergence Theorem of Lebesgue, we obtain $\int_{B(0, R)} K(\varepsilon x)Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon)dx \geq 1 - C(\varepsilon/R)^\gamma$. This implies that

$$\int_{\mathbb{R}^N} K(\varepsilon x)Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon)dx = \lim_{R \rightarrow \infty} \int_{B(0, R)} K(\varepsilon x)Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon)dx \geq 1.$$

We claim that

$$(7) \quad \int_{\mathbb{R}^N} K(\varepsilon x)Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon)dx = 1.$$

In fact, arguing by contradiction, we assume that $\int_{\mathbb{R}^N} K(\varepsilon x)Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon)dx > 1$. Then there exists $\bar{R} > 0$ so that $\int_{B(0, \bar{R})} K(\varepsilon x)Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon)dx > 1$. Hence, we get $\lim_{m \rightarrow \infty} \int_{B(0, \bar{R})} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m)dx = \int_{B(0, \bar{R})} K(\varepsilon x)Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon)dx > 1$. Thus, there exists $m_0 \in \mathbb{N}$ such that $\int_{B(0, \bar{R})} K(\varepsilon x)Q(u_\varepsilon^{m_0}, v_\varepsilon^{m_0})dx > 1$. But this is impossible, since $\int_{\mathbb{R}^N} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m)dx = 1$ for all $m \in \mathbb{N}$.

Since $\int_{B(0, T)} \chi_\varepsilon Q(u_\varepsilon^m, v_\varepsilon^m)dx \leq 1$ for each $T > 0$ we get, again using the Dominated Convergence Theorem of Lebesgue, that

$$\int_{B(0, T)} \chi_\varepsilon Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon)dx \leq 1$$

for each $T > 0$. Consequently,

$$(8) \quad \int_{\mathbb{R}^N} \chi_\varepsilon Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon)dx \leq 1.$$

Since $\|(\bar{u}_\varepsilon, \bar{v}_\varepsilon)\|_\varepsilon^2 \leq \liminf_{m \rightarrow \infty} \|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2 = M_\varepsilon$, we infer that $(\bar{u}_\varepsilon, \bar{v}_\varepsilon)$ is a minimizer of M_ε .

Now, we will prove that in system $(S_{\alpha_\varepsilon^m, \beta_\varepsilon^m})$, $\alpha_\varepsilon^m > 0 \geq \beta_\varepsilon^m$. In fact, using same ideas in [14], we take $\xi_0, \xi_1 \in C_0^\infty(\mathbb{R}^N)$ non-negative radially symmetric functions with $\text{supp}(\xi_0) \subset \text{int}(A_\varepsilon^{4\delta})$ and $\text{supp}(\xi_1) \subset \{x \in \mathbb{R}^N \mid |x| < d(0, A_\varepsilon^{4\delta})\}$. Define

$$D(s, t) \equiv \int_{B(0, R_m)} K(\varepsilon x)Q((1 + t\xi_0 - s\xi_1)(u_\varepsilon^m, v_\varepsilon^m))dx.$$

The function D is continuously differentiable in a neighborhood of $(0, 0)$. We note that $D(0, 0) = 1$ and $\frac{\partial}{\partial t}D(0, 0) = (p+1) \int_{B(0, R_m)} K(\varepsilon x)\xi_0 Q(u_\varepsilon^m, v_\varepsilon^m) dx > 0$. By the implicit function theorem, for small $\tau > 0$ there exists $t \in C^1(-\tau, \tau)$ such that

$$t(0) = 0 \text{ and } D(s, t(s)) = 1 \text{ for all } s \in (-\tau, \tau).$$

Hence

$$(9) \quad (p+1) \int_{B(0, R_m)} K(\varepsilon x)(t'(0)\xi_0 - \xi_1)Q(u_\varepsilon^m, v_\varepsilon^m) dx = 0.$$

Moreover, using the definition of χ_ε and the fact that $\chi_\varepsilon \xi_0 \equiv 0$ in $B(0, R_m)$, we obtain

$$(10) \quad \begin{aligned} & \frac{d}{ds} \Big|_{s=0} \int_{B(0, R_m)} \chi_\varepsilon Q((1 + t(s)\xi_0 - s\xi_1)(u_\varepsilon^m, v_\varepsilon^m)) dx \\ &= -(p+1)\varepsilon^{-(N-1)-3(p+1)/(p-1)} \int_{\text{supp}(\xi_1)} \xi_1 Q(u_\varepsilon^m, v_\varepsilon^m) dx < 0. \end{aligned}$$

This implies that there exists $c > 0$ such that for any $s \in (0, c)$,

$$\int_{B(0, R_m)} \chi_\varepsilon Q((1 + t(s)\xi_0 - s\xi_1)(u_\varepsilon^m, v_\varepsilon^m)) dx < 1.$$

Since $(u_\varepsilon^m, v_\varepsilon^m)$ is a minimizer for M_ε^m , we have

$$\begin{aligned} (11) \quad 0 &\leq \frac{d}{ds} \Big|_{s=0} \int_{B(0, R_m)} [|\nabla((1 + t(s)\xi_0 - s\xi_1)u_\varepsilon^m)|^2 + |\nabla((1 + t(s)\xi_0 - s\xi_1)v_\varepsilon^m)|^2 \\ &\quad + (1 + t(s)\xi_0 - s\xi_1)^2 (V_1(\varepsilon x)(u_\varepsilon^m)^2 + V_2(\varepsilon x)(v_\varepsilon^m)^2)] dx \\ &= 2 \int_{B(0, R_m)} [\nabla u_\varepsilon^m \cdot \nabla((t'(0)\xi_0 - \xi_1)u_\varepsilon^m) + \nabla v_\varepsilon^m \cdot \nabla((t'(0)\xi_0 - \xi_1)v_\varepsilon^m) \\ &\quad + (t'(0)\xi_0 - \xi_1)(V_1(\varepsilon x)(u_\varepsilon^m)^2 + V_2(\varepsilon x)(v_\varepsilon^m)^2)] dx. \end{aligned}$$

Using $(t'(0)\xi_0 - \xi_1)(u_\varepsilon^m, v_\varepsilon^m)$ as test function in $(S_{\alpha_\varepsilon^m, \beta_\varepsilon^m})$, the homogeneity of Q , the definition of χ_ε and (9), we deduce that

$$\begin{aligned} 0 &\leq \int_{B(0, R_m)} [\nabla u_\varepsilon^m \cdot \nabla((t'(0)\xi_0 - \xi_1)u_\varepsilon^m) + \nabla v_\varepsilon^m \cdot \nabla((t'(0)\xi_0 - \xi_1)v_\varepsilon^m) \\ &\quad + (t'(0)\xi_0 - \xi_1)(V_1(\varepsilon x)(u_\varepsilon^m)^2 + V_2(\varepsilon x)(v_\varepsilon^m)^2)] dx \\ &= (p + 1)\alpha_\varepsilon^m \int_{B(0, R_m)} (t'(0)\xi_0 - \xi_1)K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m) dx \\ &\quad + (p + 1)\beta_\varepsilon^m \int_{B(0, R_m)} \chi_\varepsilon(t'(0)\xi_0 - \xi_1)Q(u_\varepsilon^m, v_\varepsilon^m) dx \\ &= -(p + 1)\beta_\varepsilon^m \varepsilon^{-(N-1)-3(p+1)/(p-1)} \int_{\text{supp}(\xi_1)} \xi_1 Q(u_\varepsilon^m, v_\varepsilon^m) dx. \end{aligned}$$

By (10) and (11) we conclude that $\beta_\varepsilon^m \leq 0$.

Now, taking $(u_\varepsilon^m, v_\varepsilon^m)$ as test function in $(S_{\alpha_\varepsilon^m, \beta_\varepsilon^m})$ and using (6) we obtain

$$(12) \quad \|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2 = (p + 1)\alpha_\varepsilon^m + (p + 1)\beta_\varepsilon^m \int_{B(0, R_m)} \chi_\varepsilon Q(u_\varepsilon^m, v_\varepsilon^m) dx.$$

This implies that $\alpha_\varepsilon^m > 0$.

Now we will show that $(\bar{u}_\varepsilon, \bar{v}_\varepsilon)$ satisfies the system $(S_{\alpha_\varepsilon, \beta_\varepsilon})$. We claim that $\{\alpha_\varepsilon^m\}_m$ is bounded for small $\varepsilon > 0$. Indeed, arguing by contradiction assume, without loss of generality, that $\lim_{m \rightarrow \infty} \alpha_\varepsilon^m = \infty$. For any $\sigma > 0$, choose a function $\phi_\sigma \in C_0^\infty(\text{int}(A_\varepsilon^{4\delta}))$ such that $0 \leq \phi_\sigma \leq 1$, $\phi_\sigma(x) = 1$ for $d(x, \partial A_\varepsilon^{4\delta}) \geq \sigma$, and $|\nabla \phi_\sigma| \leq 2/\sigma$. Using $\phi_\sigma(u_\varepsilon^m, v_\varepsilon^m)$ as test function in $(S_{\alpha_\varepsilon^m, \beta_\varepsilon^m})$ and that $\chi_\varepsilon \phi_\sigma \equiv 0$, we obtain

$$\begin{aligned} &\int_{\mathbb{R}^N} [|\nabla u_\varepsilon^m|^2 \phi_\sigma + \nabla u_\varepsilon^m \cdot \nabla \phi_\sigma u_\varepsilon^m + |\nabla v_\varepsilon^m|^2 \phi_\sigma + \nabla v_\varepsilon^m \cdot \nabla \phi_\sigma v_\varepsilon^m \\ &\quad + \phi_\sigma (V_1(\varepsilon x)(u_\varepsilon^m)^2 + V_2(\varepsilon x)(v_\varepsilon^m)^2)] dx \end{aligned}$$

$$(13) \quad = (p + 1)\alpha_\varepsilon^m \int_{\mathbb{R}^N} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m)\phi_\sigma dx.$$

From $\inf_{x \in \text{supp}(|\nabla\phi_\sigma|)} V(\varepsilon x) > 0$ and the properties of ϕ_σ , we have

$$(14) \quad \begin{aligned} & \int_{\mathbb{R}^N} [|\nabla u_\varepsilon^m|^2 \phi_\sigma + \nabla u_\varepsilon^m \cdot \nabla \phi_\sigma u_\varepsilon^m + |\nabla v_\varepsilon^m|^2 \phi_\sigma + \nabla v_\varepsilon^m \cdot \nabla \phi_\sigma v_\varepsilon^m \\ & + \phi_\sigma (V_1(\varepsilon x) (u_\varepsilon^m)^2 + V_2(\varepsilon x) (v_\varepsilon^m)^2)] dx \\ & \leq C \|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2 \end{aligned}$$

for some $C > 0$, independent of m . By (13), (14) and the fact that

$$\{\|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2\}_m$$

is a bounded sequence, we see that for some $C > 0$, independent of m ,

$$\int_{\mathbb{R}^N} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m)\phi_\sigma dx \leq C/\alpha_\varepsilon^m.$$

Thus,

$$(15) \quad \lim_{m \rightarrow \infty} \int_{\{x \in A_\varepsilon^{4\delta} \mid d(x, \partial A_\varepsilon^{4\delta}) \geq \sigma\}} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m) dx = 0.$$

From the condition $\int_{\mathbb{R}^N} \chi_\varepsilon Q(u_\varepsilon^m, v_\varepsilon^m) dx \leq 1$ and from the definition of χ_ε , we have

$$(16) \quad \int_{\mathbb{R}^N \setminus B(0, R_0/\varepsilon)} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m) dx \leq C (\varepsilon/R_0)^\gamma$$

and

$$(17) \quad \int_{B(0, R_0/\varepsilon) \setminus A_\varepsilon^{4\delta}} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m) dx \leq C\varepsilon^{(N-1)+3(p+1)/(p-1)}$$

for some positive constant C . Now, using $\int_{\mathbb{R}^N} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m) dx = 1$, (15), (16) and (17) we infer that

$$\begin{aligned} & \liminf_{m \rightarrow \infty} \int_{\{x \in A_\varepsilon^{4\delta} \mid d(x, \partial A_\varepsilon^{4\delta}) \leq \sigma\}} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m) dx \\ & \geq 1 - C\varepsilon^{(N-1)+3(p+1)/(p-1)} - C (\varepsilon/R_0)^\gamma > 0 \end{aligned}$$

for small $\varepsilon > 0$ and for each $\sigma > 0$. Then for each $\sigma > 0$ there is a sequence $\{x_m\}_m$ in $A_\varepsilon^{4\delta}$ such that $\lim_{m \rightarrow \infty} d(x_m, \partial A_\varepsilon^{4\delta}) = 0$ and $Q(u_\varepsilon^m(x_m), v_\varepsilon^m(x_m)) = 1$. Since $A_\varepsilon^{4\delta}$ is a compact subset of \mathbb{R}^N , we see that $\lim_{m \rightarrow \infty} x_m = x_0 \in A_\varepsilon^{4\delta}$, up to subsequence. This implies that $x_0 \in \partial A_\varepsilon^{4\delta}$ and $\lim_{m \rightarrow \infty} |x_m| = |x_0| = r_0 > 0$ so that for each $\sigma > 0$

$$(18) \quad \liminf_{m \rightarrow \infty} \int_{D_{r_0}^\sigma} K(\varepsilon x)Q(u_\varepsilon^m, v_\varepsilon^m) dx > 0,$$

where $D_{r_0}^\sigma$ is defined by $D_{r_0}^\sigma \equiv \{x \in \mathbb{R}^N \mid r_0 - \sigma \leq |x| \leq r_0 + \sigma\}$. To reach a contradiction of (18), we will prove the following statements:

$$(19) \quad \int_{D_{r_0}^\sigma} [((u_\varepsilon^m - 1)_+)^2 + ((v_\varepsilon^m - 1)_+)^2] dx \leq C\sigma^{2/N} \|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2$$

for m large and some positive constant C , independent of σ ;

$$(20) \quad \int_{D_{r_0}^\sigma} [|\nabla(u_\varepsilon^m - 1)_+|^2 + |\nabla(v_\varepsilon^m - 1)_+|^2 + V_1(\varepsilon x)((u_\varepsilon^m - 1)_+)^2 + V_2(\varepsilon x)((v_\varepsilon^m - 1)_+)^2] dx \leq \|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2,$$

$$(21) \quad \int_{D_{r_0}^\sigma} K(\varepsilon x)Q((u_\varepsilon^m - 1)_+, (v_\varepsilon^m - 1)_+) dx \leq C\sigma^{s(p+1)/N}$$

for some $s \in (0, 1)$ and $C > 0$, C independent of σ ; and

$$(22) \quad \int_{D_{r_0}^\sigma} [(u_\varepsilon^m)^{p+1} + (v_\varepsilon^m)^{p+1}] dx \leq C_0 \int_{D_{r_0}^\sigma} K(\varepsilon x)Q((u_\varepsilon^m - 1)_+, (v_\varepsilon^m - 1)_+) dx + C_0\sigma$$

for some positive constant C_0 .

To prove the assertion (19) note that, by the Poincaré inequality, there is a positive constant C_1 so that

$$\begin{aligned} \int_{D_{r_0}^\sigma} ((u_\varepsilon^m - 1)_+)^2 dx &\leq (|D_{r_0}^\sigma|/\omega_N)^{2/N} \int_{D_{r_0}^\sigma} |\nabla(u_\varepsilon^m - 1)_+|^2 dx \\ &\leq C_1\sigma^{2/N} \int_{D_{r_0}^\sigma} |\nabla u_\varepsilon^m|^2 dx \\ &\leq C_1\sigma^{2/N} \|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2. \end{aligned}$$

Similarly $\int_{D_{r_0}^\sigma} ((v_\varepsilon^m - 1)_+)^2 dx \leq C_2\sigma^{2/N} \|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2$ for some constant $C_2 > 0$ and the inequality (19) follows.

The verification of (20) is immediate.

For the statement (21), we use the interpolation inequality, Sobolev inequality, (19) and (20) to find

$$\begin{aligned} &\int_{D_{r_0}^\sigma} ((u_\varepsilon^m - 1)_+)^{p+1} dx \\ &\leq C_0 \left(\int_{D_{r_0}^\sigma} ((u_\varepsilon^m - 1)_+)^2 dx \right)^{s_1(p+1)/2} \times \left(\int_{D_{r_0}^\sigma} |\nabla(u_\varepsilon^m - 1)_+|^2 dx \right)^{(1-s_1)(p+1)/2} \\ &\leq C_1 (\sigma^{2/N} \|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2)^{s_1(p+1)/2} \times (\|(u_\varepsilon^m, v_\varepsilon^m)\|_\varepsilon^2)^{(1-s_1)(p+1)/2} \\ &\leq C_2 \sigma^{s_1(p+1)/N} \end{aligned}$$

for $s_1 \in (0, 1)$ and for some constants $C_0, C_1, C_2 > 0$, independent of σ . Similarly, we get $\int_{D_{r_0}^\sigma} ((v_\varepsilon^m - 1)_+)^{p+1} dx \leq C_3 \sigma^{s_2(p+1)/N}$ for some constants $C_3 > 0$ and $s_2 \in (0, 1)$. Using this information and (Q_2) , (21) follows.

Finally, to obtain (22), we note that

$$\begin{aligned} \int_{D_{r_0}^\sigma} (u_\varepsilon^m)^{p+1} dx &\leq \int_{D_{r_0}^\sigma \cap \{u_\varepsilon^m \geq 1\}} ((u_\varepsilon^m - 1)_+ + 1)^{p+1} dx + |D_{r_0}^\sigma| \\ &\leq 2^p \int_{D_{r_0}^\sigma} ((u_\varepsilon^m - 1)_+)^{p+1} dx + (2^p + 1) |D_{r_0}^\sigma|. \end{aligned}$$

Also, $\int_{D_{r_0}^\sigma} (v_\varepsilon^m)^{p+1} dx \leq 2^p \int_{D_{r_0}^\sigma} ((v_\varepsilon^m - 1)_+)^{p+1} dx + (2^p + 1) |D_{r_0}^\sigma|$. Therefore

$$\begin{aligned} (23) \quad &\int_{D_{r_0}^\sigma} [(u_\varepsilon^m)^{p+1} + (v_\varepsilon^m)^{p+1}] dx \\ &\leq 2^p \int_{D_{r_0}^\sigma} [((u_\varepsilon^m - 1)_+)^{p+1} + ((v_\varepsilon^m - 1)_+)^{p+1}] dx + 2(2^p + 1) |D_{r_0}^\sigma|. \end{aligned}$$

Using (23), (Q_2) and the fact that $|D_{r_0}^\sigma| \leq C\sigma$ for all small $\sigma > 0$ and for some positive constant C , we obtain (22). From (Q_2) , (21) and (22) it follows that

$$\begin{aligned} \int_{D_{r_0}^\sigma} K(\varepsilon x) Q(u_\varepsilon^m, v_\varepsilon^m) dx &\leq C_0 \eta_2 \int_{D_{r_0}^\sigma} [(u_\varepsilon^m)^{p+1} + (v_\varepsilon^m)^{p+1}] dx \\ &\leq C_1 \int_{D_{r_0}^\sigma} K(\varepsilon x) Q((u_\varepsilon^m - 1)_+, (v_\varepsilon^m - 1)_+) dx + C_1 \sigma \\ &\leq C(\sigma^{s(p+1)/N} + \sigma) \end{aligned}$$

for some $s \in (0, 1)$ and for some constants $C_0, C_1, C > 0$, independent of σ and m . Therefore,

$$\liminf_{m \rightarrow \infty} \int_{D_{r_0}^\sigma} K(\varepsilon x) Q(u_\varepsilon^m, v_\varepsilon^m) dx \leq C(\sigma^{s(p+1)/N} + \sigma)$$

for all $\sigma > 0$ small. But this contradicts (18), given the arbitrariness of $\sigma > 0$. Thus, we conclude that $\{\alpha_\varepsilon^m\}_m$ is bounded. This implies that $\lim_{m \rightarrow \infty} \alpha_\varepsilon^m = \alpha_\varepsilon \geq 0$, up to subsequence. Using (12) and the fact that

$$0 \leq \int_{\mathbb{R}^N} \chi_\varepsilon(x) Q(u_\varepsilon^m, v_\varepsilon^m) dx \leq 1$$

for all $m \in \mathbb{N}$, we get $\lim_{m \rightarrow \infty} \beta_\varepsilon^m = \beta_\varepsilon \leq 0$. Since $(u_\varepsilon^m, v_\varepsilon^m)$ is solution of $(S_{\alpha_\varepsilon^m, \beta_\varepsilon^m})$ we have that

$$\begin{aligned} &\int_{\mathbb{R}^N} [\nabla u_\varepsilon^m \cdot \nabla \varphi + \nabla v_\varepsilon^m \cdot \nabla \psi + V_1(\varepsilon x) u_\varepsilon^m \varphi + V_2(\varepsilon x) v_\varepsilon^m \psi] dx \\ &= \alpha_\varepsilon^m \int_{\mathbb{R}^N} K(\varepsilon x) [\varphi Q_u(u_\varepsilon^m, v_\varepsilon^m) + \psi Q_v(u_\varepsilon^m, v_\varepsilon^m)] dx \end{aligned}$$

$$(24) \quad + \beta_\varepsilon^m \int_{\mathbb{R}^N} \chi_\varepsilon [\varphi Q_u(u_\varepsilon^m, v_\varepsilon^m) + \psi Q_v(u_\varepsilon^m, v_\varepsilon^m)] dx$$

for any $\varphi, \psi \in C_{0,rad}^\infty(\mathbb{R}^N)$. Finally, taking the limit in (24) as $m \rightarrow \infty$, we see that

$$\begin{aligned} & \int_{\mathbb{R}^N} [\nabla \bar{u}_\varepsilon \cdot \nabla \varphi + \nabla \bar{v}_\varepsilon \cdot \nabla \psi + V_1(\varepsilon x) \varphi \bar{u}_\varepsilon + V_2(\varepsilon x) \psi \bar{v}_\varepsilon] dx \\ &= \alpha_\varepsilon \int_{\mathbb{R}^N} K(\varepsilon x) [\varphi Q_u(\bar{u}_\varepsilon, \bar{v}_\varepsilon) + \psi Q_v(\bar{u}_\varepsilon, \bar{v}_\varepsilon)] dx \\ & \quad + \beta_\varepsilon \int_{\mathbb{R}^N} \chi_\varepsilon [\varphi Q_u(\bar{u}_\varepsilon, \bar{v}_\varepsilon) + \psi Q_v(\bar{u}_\varepsilon, \bar{v}_\varepsilon)] dx \end{aligned}$$

for any $\varphi, \psi \in C_{0,rad}^\infty(\mathbb{R}^N)$. Therefore, $(\bar{u}_\varepsilon, \bar{v}_\varepsilon)$ satisfies $(S_{\alpha_\varepsilon, \beta_\varepsilon})$. From (7), $\beta_\varepsilon \leq 0$, the homogeneity of Q and the fact that $(\bar{u}_\varepsilon, \bar{v}_\varepsilon)$ is solution of $(S_{\alpha_\varepsilon, \beta_\varepsilon})$, we conclude that $\|(\bar{u}_\varepsilon, \bar{v}_\varepsilon)\|_\varepsilon^2 \leq (p+1)\alpha_\varepsilon$ and therefore $\alpha_\varepsilon > 0$. This completes the proof of lemma. \square

Claim: For ε small,

$$(25) \quad \int_{\mathbb{R}^N} \chi_\varepsilon Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx < 1.$$

This claim is one of the crucial steps of our work. We will postpone its proof for while. If this is the case, for any $\varphi, \psi \in C_{0,rad}^\infty(\mathbb{R}^N)$, we define

$$\varphi_s \equiv (\bar{u}_\varepsilon + s\varphi) \left(\int_{\mathbb{R}^N} K(\varepsilon x) Q(\bar{u}_\varepsilon + s\varphi, \bar{v}_\varepsilon + s\psi) dx \right)^{-1/(p+1)}$$

and

$$\psi_s \equiv (\bar{v}_\varepsilon + s\psi) \left(\int_{\mathbb{R}^N} K(\varepsilon x) Q(\bar{u}_\varepsilon + s\varphi, \bar{v}_\varepsilon + s\psi) dx \right)^{-1/(p+1)}.$$

From (7) we conclude that $(\varphi_0, \psi_0) = (\bar{u}_\varepsilon, \bar{v}_\varepsilon)$. Since Q is homogeneous of degree $p+1$, we obtain $\int_{\mathbb{R}^N} K(\varepsilon x) Q(\varphi_s, \psi_s) dx = 1$. Also, by (25), $\int_{\mathbb{R}^N} \chi_\varepsilon Q(\varphi_s, \psi_s) dx < 1$ for small $|s|$.

Thus,

$$\begin{aligned} 0 &= \frac{d}{ds} \|(\varphi_s, \psi_s)\|_\varepsilon^2 \Big|_{s=0} \\ &= -2M_\varepsilon/(p+1) \int_{\mathbb{R}^N} K(\varepsilon x) [\varphi Q_u(\bar{u}_\varepsilon, \bar{v}_\varepsilon) + \psi Q_v(\bar{u}_\varepsilon, \bar{v}_\varepsilon)] dx \\ & \quad + 2 \int_{\mathbb{R}^N} [\nabla \bar{u}_\varepsilon \cdot \nabla \varphi + \nabla \bar{v}_\varepsilon \cdot \nabla \psi + V_1(\varepsilon x) \bar{u}_\varepsilon \varphi + V_2(\varepsilon x) \bar{v}_\varepsilon \psi] dx. \end{aligned}$$

This implies that $(\bar{u}_\varepsilon, \bar{v}_\varepsilon)$ satisfies the system $(S_{M_\varepsilon/(p+1)}, 0)$. Then, as the functions Q_u and Q_v are homogeneous of degree p , we deduce that $(\tilde{u}_\varepsilon, \tilde{v}_\varepsilon)$, where $\tilde{u}_\varepsilon = (M_\varepsilon/(p+1))^{1/(p-1)} \bar{u}_\varepsilon$ and $\tilde{v}_\varepsilon = (M_\varepsilon/(p+1))^{1/(p-1)} \bar{v}_\varepsilon$, is a solution of (\tilde{S}) .

Lemma 4. $\lim_{\varepsilon \rightarrow 0} \varepsilon^{(N-1)(p-1)/(p+1)} \alpha_\varepsilon = 0$.

Proof. Assume by contradiction, taking a subsequence if necessary, that

$$\lim_{\varepsilon \rightarrow 0} \varepsilon^{(N-1)(p-1)/(p+1)} \alpha_\varepsilon = \alpha \in (0, \infty].$$

For any $\sigma > 0$, we choose $\phi_\sigma \in C_0^\infty(\text{int}(A_\varepsilon^{4\delta}))$ satisfying $0 \leq \phi_\sigma \leq 1$, $\phi_\sigma(x) = 1$ for $d(x, \partial A_\varepsilon^{4\delta}) \geq \sigma$, and $|\nabla \phi_\sigma| \leq 2/\sigma$. Using $\phi_\sigma(\bar{u}_\varepsilon, \bar{v}_\varepsilon)$ as test function in the system $(S_{\alpha_\varepsilon, \beta_\varepsilon})$ and the fact that $\chi_\varepsilon \phi_\sigma \equiv 0$, we have

$$\begin{aligned} & \int_{\mathbb{R}^N} [|\nabla \bar{u}_\varepsilon|^2 \phi_\sigma + \nabla \bar{u}_\varepsilon \cdot \nabla \phi_\sigma \bar{u}_\varepsilon + |\nabla \bar{v}_\varepsilon|^2 \phi_\sigma + \nabla \bar{v}_\varepsilon \cdot \nabla \phi_\sigma \bar{v}_\varepsilon \\ & + \phi_\sigma (V_1(\varepsilon x)(\bar{u}_\varepsilon)^2 + V_2(\varepsilon x)(\bar{v}_\varepsilon)^2)] dx \\ (26) \quad & = (p+1)\alpha_\varepsilon \int_{\mathbb{R}^N} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) \phi_\sigma dx. \end{aligned}$$

From $\inf_{x \in \text{supp}(|\nabla \phi_\sigma|)} V(\varepsilon x) > 0$ and the properties of ϕ_σ , we obtain

$$\begin{aligned} & \int_{\mathbb{R}^N} [|\nabla \bar{u}_\varepsilon|^2 \phi_\sigma + \nabla \bar{u}_\varepsilon \cdot \nabla \phi_\sigma \bar{u}_\varepsilon + |\nabla \bar{v}_\varepsilon|^2 \phi_\sigma + \nabla \bar{v}_\varepsilon \cdot \nabla \phi_\sigma \bar{v}_\varepsilon \\ & + \phi_\sigma (V_1(\varepsilon x)(\bar{u}_\varepsilon)^2 + V_2(\varepsilon x)(\bar{v}_\varepsilon)^2)] dx \\ (27) \quad & \leq C \|(\bar{u}_\varepsilon, \bar{v}_\varepsilon)\|_\varepsilon^2 \end{aligned}$$

for some $C > 0$, independent of $\varepsilon > 0$. From (26) and (27) it follows that $\int_{\mathbb{R}^N} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) \phi_\sigma dx \leq C \|(\bar{u}_\varepsilon, \bar{v}_\varepsilon)\|_\varepsilon^2 / \alpha_\varepsilon$ for some positive constant C , independent of $\varepsilon > 0$. By Lemma 2, for each $\sigma > 0$,

$$\lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^N} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) \phi_\sigma dx = 0.$$

Then

$$(28) \quad \lim_{\varepsilon \rightarrow 0} \int_{\{x \in A_\varepsilon^{4\delta} \mid d(x, \partial A_\varepsilon^{4\delta}) \geq \sigma\}} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx = 0.$$

From (8) and the definition of χ_ε , we get

$$(29) \quad \int_{\mathbb{R}^N \setminus B(0, R_0/\varepsilon)} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx \leq C (\varepsilon/R_0)^\gamma$$

and

$$(30) \quad \int_{B(0, R_0/\varepsilon) \setminus A_\varepsilon^{4\delta}} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx \leq C \varepsilon^{(N-1)+3(p+1)/(p-1)}$$

for some $C > 0$. From (7), (28), (29) and (30), we see that for each $\sigma > 0$,

$$(31) \quad \liminf_{\varepsilon \rightarrow 0} \int_{\{x \in A_\varepsilon^{4\delta} \mid d(x, \partial A_\varepsilon^{4\delta}) \leq \sigma\}} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx > 0.$$

From (31), for every $\sigma > 0$, there exists a sequence $\{x_m\}_m$ in $A_\varepsilon^{4\delta}$ such that $\lim_{m \rightarrow \infty} d(x_m, \partial A_\varepsilon^{4\delta}) = 0$. Therefore, there exists some $x_0 \in \partial A^{4\delta}$, with

$\lim_{m \rightarrow \infty} x_m = \frac{x_0}{\varepsilon}$, and $\omega > 0$ such that for any $\sigma > 0$,

$$(32) \quad \liminf_{\varepsilon \rightarrow 0} \int_{\{x \in \mathbb{R}^N \mid ||x| - |x_0|/\varepsilon| \leq \sigma\}} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx \geq \omega,$$

otherwise we would have $\lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^N} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx = 0$. But this is impossible because it contradicts (7).

We fix $\sigma > 0$ and choose a radially symmetric function $\psi_\sigma \in C_0^\infty$ so that

$$\psi_\sigma(x) = \begin{cases} 0 & \text{if } ||x| - |x_0|/\varepsilon| \geq 2\sigma, \\ 1 & \text{if } ||x| - |x_0|/\varepsilon| \leq \sigma, \end{cases}$$

$0 \leq \psi_\sigma \leq 1$ and $|\nabla \psi_\sigma| \leq 3/\sigma$. From (32) it follows that

$$(33) \quad \liminf_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^N} K(\varepsilon x) Q(\psi_\sigma \bar{u}_\varepsilon, \psi_\sigma \bar{v}_\varepsilon) dx \geq \omega.$$

Now, we claim that

$$(34) \quad \lim_{\varepsilon \rightarrow 0} \varepsilon^{(N-1)(p-1)/(p+1)} \|(\psi_\sigma \bar{u}_\varepsilon, \psi_\sigma \bar{v}_\varepsilon)\|_\varepsilon^2 = 0.$$

Indeed, by the Cauchy-Schwarz inequality, the boundedness of the gradient of ψ_σ , and by the fact that $\alpha_0 = \inf_{x \in \text{supp}(\psi_\sigma)} V(\varepsilon x) > 0$, we have

$$\begin{aligned} & \|(\psi_\sigma \bar{u}_\varepsilon, \psi_\sigma \bar{v}_\varepsilon)\|_\varepsilon^2 \\ & \leq C_1 \int_{\text{supp}(\psi_\sigma)} [|\nabla \bar{u}_\varepsilon|^2 + |\nabla \bar{v}_\varepsilon|^2 + (\bar{u}_\varepsilon)^2 + (\bar{v}_\varepsilon)^2 + V_1(\varepsilon x)(\bar{u}_\varepsilon)^2 \\ & \quad + V_2(\varepsilon x)(\bar{v}_\varepsilon)^2] dx \\ & \leq C_1 \int_{\text{supp}(\psi_\sigma)} [|\nabla \bar{u}_\varepsilon|^2 + |\nabla \bar{v}_\varepsilon|^2 + \frac{1}{\alpha_0} V_1(\varepsilon x)(\bar{u}_\varepsilon)^2 + \frac{1}{\alpha_0} V_2(\varepsilon x)(\bar{v}_\varepsilon)^2 \\ & \quad + V_1(\varepsilon x)(\bar{u}_\varepsilon)^2 + V_2(\varepsilon x)(\bar{v}_\varepsilon)^2] dx \\ & \leq C_2 \|(\bar{u}_\varepsilon, \bar{v}_\varepsilon)\|_\varepsilon^2 = C_2 M_\varepsilon \end{aligned}$$

for some positive constants C_1 and C_2 , independent of $\varepsilon > 0$. By Lemma 2, (34) follows.

On other hand, putting $D_\varepsilon \equiv \{x \in \mathbb{R}^N \mid |x_0|/\varepsilon - 2\sigma \leq |x| \leq |x_0|/\varepsilon + 2\sigma\}$, we see that

$$\begin{aligned} & \liminf_{\varepsilon \rightarrow 0} \varepsilon^{(N-1)(p-1)/(p+1)} \|(\psi_\sigma \bar{u}_\varepsilon, \psi_\sigma \bar{v}_\varepsilon)\|_\varepsilon^2 \\ & \geq \liminf_{\varepsilon \rightarrow 0} \left(\left[\int_{\mathbb{R}^N} K(\varepsilon x) Q(\psi_\sigma \bar{u}_\varepsilon, \psi_\sigma \bar{v}_\varepsilon) dx \right]^{2/(p+1)} \right) \\ & \quad \times \liminf_{\varepsilon \rightarrow 0} \left(\varepsilon^{(N-1)(p-1)/(p+1)} \inf_{u, v \in C_0^1(D_\varepsilon)} \frac{\|(u, v)\|_\varepsilon^2}{\left[\int_{\mathbb{R}^N} K(\varepsilon x) Q(u, v) dx \right]^{2/(p+1)}} \right) \\ & \geq C \omega^{2/(p+1)} \liminf_{\varepsilon \rightarrow 0} ((|x_0| - 2\sigma\varepsilon)^{N-1} (|x_0| + 2\sigma\varepsilon)^{-2(N-1)/(p+1)}) J_\sigma \\ (35) \quad & = C \omega^{2/(p+1)} |x_0|^{(N-1)(p-1)/(p+1)} J_\sigma > 0 \end{aligned}$$

for some $C > 0$, where

$$J_\sigma \equiv \inf_{g_0, g_1 \in C_0^1(-2\sigma, 2\sigma)} \frac{\int_{-2\sigma}^{2\sigma} [(g'_0(s))^2 + (g'_1(s))^2 + \alpha_0((g_0(s))^2 + (g_1(s))^2)] ds}{[\int_{-2\sigma}^{2\sigma} |g_0(s)|^{p+1} ds]^{2/(p+1)} + [\int_{-2\sigma}^{2\sigma} |g_1(s)|^{p+1} ds]^{2/(p+1)}}.$$

From (34) and (35) we have a contradiction. So, to conclude the proof of the lemma, we will have a verification for (35).

Using (Q_2) , change of variables and setting $g_0(s + |x_0|/\varepsilon) = \bar{g}_0(s)$, $g_1(s + |x_0|/\varepsilon) = \bar{g}_1(s)$, we deduce that

$$\begin{aligned} & \frac{\|(u, v)\|_\varepsilon^2}{[\int_{\mathbb{R}^N} K(\varepsilon x) Q(u, v) dx]^{2/(p+1)}} \\ & \geq C\varepsilon^{-(N-1)(p-1)/(p+1)} (|x_0| - 2\sigma\varepsilon)^{N-1} (|x_0| + 2\sigma\varepsilon)^{-2(N-1)/(p+1)} \\ & \quad \times \frac{\int_{-2\sigma}^{2\sigma} [(\bar{g}'_0(s))^2 + (\bar{g}'_1(s))^2 + \alpha_0((\bar{g}_0(s))^2 + (\bar{g}_1(s))^2)] ds}{\left(\int_{-2\sigma}^{2\sigma} |\bar{g}_0(s)|^{p+1} ds\right)^{2/(p+1)} + \left(\int_{-2\sigma}^{2\sigma} |\bar{g}_1(s)|^{p+1} ds\right)^{2/(p+1)}} \\ & \geq C\varepsilon^{-(N-1)(p-1)/(p+1)} (|x_0| - 2\sigma\varepsilon)^{N-1} (|x_0| + 2\sigma\varepsilon)^{-2(N-1)/(p+1)} J_\sigma \end{aligned}$$

for some positive constant C . Then

$$\begin{aligned} & \varepsilon^{(N-1)(p-1)/(p+1)} \inf_{u, v \in C_0^1(D_\varepsilon)} \frac{\|(u, v)\|_\varepsilon^2}{[\int_{\mathbb{R}^N} K(\varepsilon x) Q(u, v) dx]^{2/(p+1)}} \\ (36) \quad & \geq C(|x_0| - 2\sigma\varepsilon)^{N-1} (|x_0| + 2\sigma\varepsilon)^{-2(N-1)/(p+1)} J_\sigma. \end{aligned}$$

Combining (33) and (36) we obtain (35). The proof of the lemma is complete. \square

Lemma 5. *If $(\bar{u}_\varepsilon, \bar{v}_\varepsilon)$ and α_ε are as above, then*

$$\lim_{\varepsilon \rightarrow 0} \|(\alpha_\varepsilon)^{1/(p-1)} \bar{u}_\varepsilon\|_{L^\infty(\mathbb{R}^N)} = \lim_{\varepsilon \rightarrow 0} \|(\alpha_\varepsilon)^{1/(p-1)} \bar{v}_\varepsilon\|_{L^\infty(\mathbb{R}^N)} = 0.$$

Proof. Let $w_\varepsilon = (\alpha_\varepsilon)^{1/(p-1)} (\bar{u}_\varepsilon + \bar{v}_\varepsilon)$. By (Q_1) , (Q_5) and the fact that $(\bar{u}_\varepsilon, \bar{v}_\varepsilon)$ is solution $(S_{\alpha_\varepsilon, \beta_\varepsilon})$ it follows that

$$(37) \quad -\Delta w_\varepsilon + V(\varepsilon x) w_\varepsilon \leq CK(\varepsilon x) (w_\varepsilon)^p \text{ in } \mathbb{R}^N$$

for some positive constant C .

Now, we claim that

$$(38) \quad \lim_{\varepsilon \rightarrow 0} \|w_\varepsilon\|_{L^\infty(\{x \in \mathbb{R}^N \mid |y|/\varepsilon - 1 \leq |x| \leq |y|/\varepsilon + 1\})} = 0$$

for all $y \in \mathbb{R}^N \setminus \{0\}$ and

$$(39) \quad \lim_{\varepsilon \rightarrow 0} \|w_\varepsilon\|_{L^\infty(B(0, r_0/\varepsilon))} = 0$$

for some constant $r_0 > 0$.

Proof of (38): First of all we will show that

$$(40) \quad \lim_{\varepsilon \rightarrow 0} \int_{B(y/\varepsilon, 2)} (\alpha_\varepsilon)^{(p+1)/(p-1)} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx = 0,$$

all $y \in \mathbb{R}^N \setminus \{0\}$. Suppose that

$$\liminf_{\varepsilon \rightarrow 0} \int_{B(y/\varepsilon, 2)} (\alpha_\varepsilon)^{(p+1)/(p-1)} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx > 0$$

for some $y \in \mathbb{R}^N \setminus \{0\}$. As K , \bar{u}_ε and \bar{v}_ε are radially symmetric functions, it follows that

$$\begin{aligned} & (\varepsilon/|y|)^{N-1} \int_{\{x \in \mathbb{R}^N \mid |y|/\varepsilon - 2 \leq |x| \leq |y|/\varepsilon + 2\}} (\alpha_\varepsilon)^{(p+1)/(p-1)} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx \\ & \geq C \int_{B(y/\varepsilon, 2)} (\alpha_\varepsilon)^{(p+1)/(p-1)} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx. \end{aligned}$$

This implies that

$$\liminf_{\varepsilon \rightarrow 0} (\varepsilon/|y|)^{N-1} \int_{\{x \in \mathbb{R}^N \mid |y|/\varepsilon - 2 \leq |x| \leq |y|/\varepsilon + 2\}} (\alpha_\varepsilon)^{(p+1)/(p-1)} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx > 0.$$

In view of (7) and Lemma 4, we have a contradiction. Similarly,

$$\limsup_{\varepsilon \rightarrow 0} \int_{B(y/\varepsilon, 2)} (\alpha_\varepsilon)^{(p+1)/(p-1)} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx = 0, \forall y \in \mathbb{R}^N \setminus \{0\}$$

and the proof of (40) is complete.

Note that

$$\int_{B(y/\varepsilon, 2)} (w_\varepsilon)^{p+1} dx \leq C \int_{B(y/\varepsilon, 2)} (\alpha_\varepsilon)^{(p+1)/(p-1)} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx$$

for some $C > 0$. From this and (40) we see that

$$(41) \quad \lim_{\varepsilon \rightarrow 0} \int_{B(y/\varepsilon, 2)} (w_\varepsilon)^{p+1} dx = 0$$

for all $y \in \mathbb{R}^N \setminus \{0\}$. Now we fix $\varepsilon > 0$. Using (37) and the fact that w_ε is radially symmetric we deduce, by the Moser iteration argument (see Theorem 9.20 in [26]), that

$$\|w_\varepsilon\|_{L^\infty(\{x \in \mathbb{R}^N \mid |y|/\varepsilon - 1 \leq |x| \leq |y|/\varepsilon + 1\})} \leq C \left(\int_{B(y/\varepsilon, 2)} (w_\varepsilon)^{p+1} dx \right)^{1/(p+1)}$$

for some positive constant C , independent of $\varepsilon > 0$. Using this and (41) we obtain (38).

Proof of (39): From (8), the definition of χ_ε and the fact that $0 \notin A^{4\delta}$, it follows that there is a constant $r_0 > 0$ such that

$$(42) \quad \int_{B(0, 2r_0/\varepsilon)} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx \leq C \varepsilon^{(N-1)+3(p+1)/(p-1)}$$

for small $\varepsilon > 0$ and for some $C > 0$. By (42) and Lemma 4, we have

$$(43) \quad \int_{B(0,2r_0/\varepsilon)} (\alpha_\varepsilon)^{(p+1)/(p-1)} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx \leq C\varepsilon^{3(p+1)/(p-1)}$$

for small $\varepsilon > 0$ and for some $C > 0$. As (Q_2) and (43) are satisfied, we see that

$$(44) \quad \int_{B(0,2r_0/\varepsilon)} (w_\varepsilon)^{p+1} dx \leq C\varepsilon^{3(p+1)/(p-1)}$$

for small $\varepsilon > 0$ and for some $C > 0$. Then using Theorem 9.20 in [26] and (44) we conclude that

$$\|w_\varepsilon\|_{L^\infty(B(0,r_0/\varepsilon))} \leq C\varepsilon^{N/(p+1)}\varepsilon^{3/(p-1)}$$

for some positive constant C , independent of $\varepsilon > 0$. This shows (39). From (38) and (39) the lemma follows. \square

In the next lemma, we also will use the arguments developed by Byeon and Wang in [18] adapted to our case.

Lemma 6. $\liminf_{\varepsilon \rightarrow 0} \varepsilon^{-2} \alpha_\varepsilon > 0$.

Proof. On the contrary, we assume for a subsequence, still denoted by ε , that $\varepsilon^{-2} \alpha_\varepsilon \rightarrow 0$ as $\varepsilon \rightarrow 0$. Let ϕ be a cut-off function such that $\phi(x) = 1$ for $x \in A_\varepsilon^{4\delta}$, $\phi(x) = 0$ for $x \notin A_\varepsilon^{5\delta}$, $0 \leq \phi \leq 1$ and $|\nabla\phi| \leq c\varepsilon$, $c > 0$. Then, it follows that

$$(45) \quad \lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^N} [|\nabla(\phi\bar{u}_\varepsilon)|^2 + |\nabla(\phi\bar{v}_\varepsilon)|^2] dx = 0.$$

In fact, since $V(\varepsilon x) \geq \delta_0 > 0$ for all $x \in \text{supp}(|\nabla\phi|)$, we see that

$$\begin{aligned} & \int_{\mathbb{R}^N} [|\nabla(\phi\bar{u}_\varepsilon)|^2 + |\nabla(\phi\bar{v}_\varepsilon)|^2] dx \\ & \leq 2(c\varepsilon)^2 \delta_0^{-1} \int_{\text{supp}(|\nabla\phi|)} [V_1(\varepsilon x)(\bar{u}_\varepsilon)^2 + V_2(\varepsilon x)(\bar{v}_\varepsilon)^2] dx \\ & \quad + 2 \int_{\mathbb{R}^N} [|\nabla\bar{u}_\varepsilon|^2 + |\nabla\bar{v}_\varepsilon|^2] dx \\ & \leq C \int_{\mathbb{R}^N} [|\nabla\bar{u}_\varepsilon|^2 + |\nabla\bar{v}_\varepsilon|^2 + V_1(\varepsilon x)(\bar{u}_\varepsilon)^2 + V_2(\varepsilon x)(\bar{v}_\varepsilon)^2] dx = C \|(\bar{u}_\varepsilon, \bar{v}_\varepsilon)\|_\varepsilon^2 \end{aligned}$$

for some positive constant C , independent of small $\varepsilon > 0$. This and Lemma 3 imply

$$\int_{\mathbb{R}^N} [|\nabla(\phi\bar{u}_\varepsilon)|^2 + |\nabla(\phi\bar{v}_\varepsilon)|^2] dx \leq C_1 \alpha_\varepsilon \leq C_1 \varepsilon^{-2} \alpha_\varepsilon$$

for some $C_1 > 0$ independent of small $\varepsilon > 0$. Hence, we get (45).

Now, using change of variables, (Q_2) , the Hölder inequality and Sobolev imbedding results, we see that

$$\int_{A_\varepsilon^{4\delta}} K(\varepsilon x) Q(\bar{u}_\varepsilon(x), \bar{v}_\varepsilon(x)) dx$$

$$\begin{aligned} &\leq C |A^{4\delta}|^{(2^*-(p+1))/2^*} \varepsilon^{-N} \left(\int_{A^{4\delta}} (\phi(y/\varepsilon) \bar{u}_\varepsilon(y/\varepsilon))^{2^*} dy \right)^{(p+1)/2^*} \\ &\quad + C |A^{4\delta}|^{(2^*-(p+1))/2^*} \varepsilon^{-N} \left(\int_{A^{4\delta}} (\phi(y/\varepsilon) \bar{v}_\varepsilon(y/\varepsilon))^{2^*} dy \right)^{(p+1)/2^*} \\ &\leq C |A^{4\delta}|^{(2^*-(p+1))/2} \varepsilon^{N(p-1)/2} \left(\int_{\mathbb{R}^N} |\nabla(\phi \bar{u}_\varepsilon)|^2 dx \right)^{(p+1)/2} \\ &\quad + C |A^{4\delta}|^{(2^*-(p+1))/2} \varepsilon^{N(p-1)/2} \left(\int_{\mathbb{R}^N} |\nabla(\phi \bar{v}_\varepsilon)|^2 dx \right)^{(p+1)/2} \end{aligned}$$

for some $C > 0$ independent of ε . From (45) it follows that

$$(46) \quad \lim_{\varepsilon \rightarrow 0} \int_{A_\varepsilon^{4\delta}} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx = 0.$$

From (8) and the definition of χ_ε , we conclude that

$$(47) \quad \lim_{\varepsilon \rightarrow 0} \int_{B(0, R_0/\varepsilon) \setminus A_\varepsilon^{4\delta}} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx = 0$$

and

$$(48) \quad \lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^N \setminus B(0, R_0/\varepsilon)} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx = 0.$$

As a consequence, from (46), (47) and (48) we have

$$\lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^N} K(\varepsilon x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx = 0.$$

But this is a contradiction with (7). The proof of lemma is complete. □

Completion of the proof for Theorem 1. To complete the proof of Theorem 1, we use arguments developed in [11], [16], [17] and [18]. We define $U_\varepsilon \equiv (C\alpha_\varepsilon)^{1/(p-1)} (\bar{u}_\varepsilon + \bar{v}_\varepsilon)$, where $C > 0$ was obtained in (37). Note that

$$(49) \quad -\Delta U_\varepsilon + V(\varepsilon x) U_\varepsilon \leq K(\varepsilon x) (U_\varepsilon)^p \text{ in } \mathbb{R}^N.$$

By Lemma 5,

$$(50) \quad \lim_{\varepsilon \rightarrow 0} \|U_\varepsilon\|_{L^\infty(\mathbb{R}^N)} = 0.$$

Let

$$2c = \inf \{V(x) \mid x \in B(0, 3R_0) \setminus Z^\delta\} > 0.$$

So, we conclude that

$$(51) \quad U_\varepsilon(x) \leq \exp(-cd(x, \partial(B(0, 3R_0/\varepsilon) \setminus Z_\varepsilon^\delta)))$$

for all $x \in B(0, 3R_0/\varepsilon) \setminus Z_\varepsilon^\delta$ and for some $c > 0$. Indeed, from (49) and (50) it follows that

$$(52) \quad \Delta U_\varepsilon - cU_\varepsilon \geq 0 \text{ in } B(0, 3R_0/\varepsilon) \setminus Z_\varepsilon^\delta$$

for small $\varepsilon > 0$. Putting $F_\varepsilon(x) = \exp(-\sqrt{cd}(x, \partial(B(0, 3R_0/\varepsilon) \setminus Z_\varepsilon^\delta)))$ we deduce that

$$(53) \quad \Delta F_\varepsilon - cF_\varepsilon < 0 \text{ in } B(0, 3R_0/\varepsilon) \setminus Z_\varepsilon^\delta.$$

By (50), (52), (53) and the comparison principle we obtain (51). Using (51) we have that

$$(54) \quad U_\varepsilon(x) \leq \exp(-c\delta/\varepsilon) \text{ in } Z_\varepsilon^{3\delta} \setminus Z_\varepsilon^{2\delta}$$

for some constant $c > 0$.

For a connected component Λ of $\text{int}(Z^{4\delta} \setminus A^{4\delta})$, we consider the first eigenvalue problem on Λ ,

$$(55) \quad \begin{cases} -\Delta\phi = \lambda_1\phi & \text{in } \Lambda, \\ \phi = 0 & \text{on } \partial\Lambda. \end{cases}$$

We can assume that $\max_{x \in \Lambda \cap \partial Z^{3\delta}} \phi(x) \geq 1$. Now, we claim that

$$(56) \quad U_\varepsilon(x) \leq C \exp(-c\delta/\varepsilon) \text{ in } \Lambda_\varepsilon \cap Z_\varepsilon^{3\delta},$$

where $\Lambda_\varepsilon = \{x \in \mathbb{R}^N \mid \varepsilon x \in \Lambda\}$. To justify the assertion (56), we define $\phi_\varepsilon(x) \equiv \exp(-c\delta/\varepsilon)\phi(\varepsilon x)$. Using (Q₂) and (8), we see that

$$(57) \quad \int_{\mathbb{R}^N} \chi_\varepsilon(U_\varepsilon)^{p+1} dx \leq C(\alpha_\varepsilon)^{(p+1)/(p-1)}$$

for some positive constant C . Now, let $z \in Z_\varepsilon^{3\delta} \setminus A_\varepsilon^{3\delta}$. Then, by Lemma 4, (57) and the definition of χ_ε , we conclude that

$$(58) \quad \int_{B(z, \delta/\varepsilon)} (U_\varepsilon)^{p+1} dx \leq C\varepsilon^{3(p+1)/(p-1)}$$

for small $\varepsilon > 0$. From Theorem 9.20 in [26] and (58) it follows that

$$\sup_{B(z, \delta/2\varepsilon)} U_\varepsilon \leq C\varepsilon^{3/(p-1)}$$

for small $\varepsilon > 0$. Thus,

$$(59) \quad U_\varepsilon \leq C\varepsilon^{3/(p-1)} \text{ in } Z_\varepsilon^{3\delta} \setminus A_\varepsilon^{3\delta}.$$

From (49) and (59) we have

$$(60) \quad \Delta U_\varepsilon + C_1\varepsilon^3 U_\varepsilon \geq 0 \text{ in } \Lambda_\varepsilon \cap Z_\varepsilon^{3\delta}$$

for some positive constant C_1 . Since ϕ satisfies (55), we deduce that, for small $\varepsilon > 0$,

$$(61) \quad \Delta\phi_\varepsilon + C_1\varepsilon^3\phi_\varepsilon \leq 0 \text{ in } \Lambda_\varepsilon \cap Z_\varepsilon^{3\delta}.$$

From (54) and the fact that $\phi(x) \geq 1$ for $x \in \Lambda \cap \partial Z^{3\delta}$, we conclude that $(U_\varepsilon - \phi_\varepsilon)_+ = 0$ on $\Lambda_\varepsilon \cap (Z_\varepsilon^{3\delta} \setminus Z_\varepsilon^{2\delta})$. From (60) and (61) we see that

$$(62) \quad -\Delta(U_\varepsilon - \phi_\varepsilon) \leq C_1\varepsilon^3(U_\varepsilon - \phi_\varepsilon) \text{ in } \Lambda_\varepsilon \cap Z_\varepsilon^{3\delta}.$$

As in [11], using $(U_\varepsilon - \phi_\varepsilon)_+$ as a test function in (62) and the Poincaré inequality, we obtain

$$\begin{aligned}
 (63) \quad & \int_{\Lambda_\varepsilon \cap Z_\varepsilon^{3\delta}} |\nabla (U_\varepsilon - \phi_\varepsilon)_+|^2 dx \\
 & \leq C_1 \varepsilon^3 \int_{\Lambda_\varepsilon \cap Z_\varepsilon^{3\delta}} ((U_\varepsilon - \phi_\varepsilon)_+)^2 dx \\
 & \leq C_1 \varepsilon^3 (|\Lambda_\varepsilon \cap Z_\varepsilon^{3\delta}| / \omega_N)^{2/N} \int_{\Lambda_\varepsilon \cap Z_\varepsilon^{3\delta}} |\nabla (U_\varepsilon - \phi_\varepsilon)_+|^2 dx \\
 & \leq C \varepsilon \int_{\Lambda_\varepsilon \cap Z_\varepsilon^{3\delta}} |\nabla (U_\varepsilon - \phi_\varepsilon)_+|^2 dx
 \end{aligned}$$

for some $C > 0$. From (63) it follows that $(U_\varepsilon - \phi_\varepsilon)_+ = 0$ in $\Lambda_\varepsilon \cap Z_\varepsilon^{3\delta}$ for small $\varepsilon > 0$. This shows (56). From (51) and (56), we deduce that for some $C, c > 0$,

$$(64) \quad \|U_\varepsilon\|_{L^\infty(B(0, 3R_0/\varepsilon - \delta/\varepsilon) \setminus A_\varepsilon^{4\delta})} \leq C \exp(-c\delta/\varepsilon).$$

Our next goal is to prove that

$$(65) \quad U_\varepsilon(x) \leq C (\varepsilon / |x|)^{\gamma/(p+1)}$$

for all $x \in \mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)$, where $C > 0$ is a constant independent of y .

Let $y \in \mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)$. From (8), the definition of χ_ε and of the fact that \bar{u}_ε and \bar{v}_ε are radially symmetric functions, we have

$$\begin{aligned}
 & \int_{B(y, 2)} (\alpha_\varepsilon)^{(p+1)/(p-1)} Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx \\
 & \leq \frac{C}{|y|^{N-1}} \int_{\{x \in \mathbb{R}^N \mid |y| - 2 \leq |x| \leq |y| + 2\}} (\alpha_\varepsilon)^{(p+1)/(p-1)} Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx \\
 (66) \quad & \leq C (\varepsilon/R_0)^{N-1} 2^\gamma (\varepsilon/|y|)^\gamma (\alpha_\varepsilon)^{(p+1)/(p-1)}
 \end{aligned}$$

for some constant $C > 0$. Thus, from (Q_2) , (66) and Lemma 4 it follows that

$$(67) \quad \int_{B(y, 2)} (U_\varepsilon)^{p+1} dx \leq C (\varepsilon/|y|)^\gamma$$

for sufficiently small $\varepsilon > 0$ and for some positive constant C . Then, from (67) and Theorem 9.20 in [26], we have that

$$\sup_{B(y, 1)} U_\varepsilon \leq C_0 (\varepsilon/|y|)^{\gamma/(p+1)} \leq C_1 (\varepsilon/|x|)^{\gamma/(p+1)}$$

for some constants $C_0, C_1 > 0$, for small $\varepsilon > 0$ and for any $x \in B(y, 1)$. Hence, (65) follows. We define

$$\omega_\varepsilon \equiv -\frac{(N-2) + \sqrt{(N-2)^2 + 4\lambda/\varepsilon^2}}{2}.$$

Consequently, $(\omega_\varepsilon)^2 + (N - 2)\omega_\varepsilon - \frac{\lambda}{\varepsilon^2} = 0$. Then, setting $\Psi_\varepsilon(x) = |x|^{\omega_\varepsilon}$, we conclude from condition (V) that

$$(68) \quad \begin{aligned} -\Delta\Psi_\varepsilon(x) + V(\varepsilon x)\Psi_\varepsilon(x) &\geq (2\lambda/\varepsilon^2 - (\omega_\varepsilon)^2 - (N - 2)\omega_\varepsilon)\gamma^{\omega_\varepsilon-2} \\ &= \frac{\lambda}{\varepsilon^2|x|^2}\Psi_\varepsilon(x) \text{ for } |x| \geq R_0/\varepsilon. \end{aligned}$$

Using (65), (68) and the fact that $\gamma(p - 1)/(p + 1) > 2$, we have

$$(69) \quad -\Delta\Psi_\varepsilon + V(\varepsilon x)\Psi_\varepsilon \geq K(\varepsilon x)(U_\varepsilon)^{p-1}\Psi_\varepsilon$$

for all $x \in \mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)$ and small $\varepsilon > 0$. From (64), we deduce that for some $C, c > 0$,

$$(70) \quad U_\varepsilon \leq C \exp(-c/\varepsilon) \text{ in } \partial B(0, 2R_0/\varepsilon).$$

Let $\tilde{\Psi}_\varepsilon(x) = C \exp(-c/\varepsilon)(\frac{2R_0}{\varepsilon})^{-\omega_\varepsilon}\Psi_\varepsilon(x)$. We claim that

$$(71) \quad U_\varepsilon(x) \leq C \exp(-c/\varepsilon)(2R_0/\varepsilon)^{-\omega_\varepsilon}\tilde{\Psi}_\varepsilon(x)$$

for all $x \in \mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)$ and some constants $C, c > 0$. In fact, as a consequence of (70), $(U_\varepsilon - \tilde{\Psi}_\varepsilon)_+ = 0$ on $\partial B(0, 2R_0/\varepsilon)$. From (69), we note that

$$(72) \quad -\Delta\tilde{\Psi}_\varepsilon + V(\varepsilon x)\tilde{\Psi}_\varepsilon \geq K(\varepsilon x)(U_\varepsilon)^{p-1}\tilde{\Psi}_\varepsilon$$

for all $x \in \mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)$. As in [11], using (49) and (72) we see that

$$(73) \quad -\Delta(U_\varepsilon - \tilde{\Psi}_\varepsilon) + V(\varepsilon x)(U_\varepsilon - \tilde{\Psi}_\varepsilon) \leq (U_\varepsilon)^{p-1}K(\varepsilon x)(U_\varepsilon - \tilde{\Psi}_\varepsilon)$$

for all $x \in \mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)$. Multiplying both sides of (73) by $(U_\varepsilon - \tilde{\Psi}_\varepsilon)_+$ and integrating by parts, we obtain

$$(74) \quad \begin{aligned} &\int_{\mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)} [|\nabla(U_\varepsilon - \tilde{\Psi}_\varepsilon)_+|^2 + V(\varepsilon x)((U_\varepsilon - \tilde{\Psi}_\varepsilon)_+)^2] dx \\ &\leq \int_{\mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)} (U_\varepsilon)^{p-1}K(\varepsilon x)((U_\varepsilon - \tilde{\Psi}_\varepsilon)_+)^2 dx. \end{aligned}$$

Using (V), (65) and the fact that $\gamma(p - 1)/(p + 1) - 2 > 0$, we deduce that, for some constants $C_0, C_1 > 0$,

$$(75) \quad \begin{aligned} K(\varepsilon x)(U_\varepsilon(x))^{p-1} &\leq C_0\varepsilon^{\gamma(p-1)/(p+1)} \frac{1}{|x|^{\gamma(p-1)/(p+1)-2}} \frac{4\lambda}{|x|^2} \\ &\leq C_1\varepsilon^{\gamma(p-1)/(p+1)}V(\varepsilon x) \end{aligned}$$

for all $x \in \mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)$ and small $\varepsilon > 0$. (74) and (75) imply

$$\begin{aligned} &\int_{\mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)} [|\nabla(U_\varepsilon - \tilde{\Psi}_\varepsilon)_+|^2 + V(\varepsilon x)((U_\varepsilon - \tilde{\Psi}_\varepsilon)_+)^2] dx \\ &\leq C_1\varepsilon^{\gamma(p-1)/(p+1)} \int_{\mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)} V(\varepsilon x)((U_\varepsilon - \tilde{\Psi}_\varepsilon)_+)^2 dx. \end{aligned}$$

This implies that, for sufficiently small $\varepsilon > 0$, $(U_\varepsilon - \tilde{\Psi}_\varepsilon)_+ = 0$ in $\mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)$ and the proof of (71) is over.

Verification of (25). Indeed, from (Q_2) , (64) and Lemma 6, we infer that

$$\begin{aligned}
 Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) &\leq 2\eta_2(\bar{u}_\varepsilon + \bar{v}_\varepsilon)^{p+1} \\
 &= C_0(\alpha_\varepsilon)^{-(p+1)/(p-1)}(U_\varepsilon)^{p+1} \\
 (76) \quad &\leq C\varepsilon^{-2(p+1)/(p-1)} \exp(-c_1/\varepsilon) \text{ in } B(0, 3R_0/\varepsilon - \delta/\varepsilon) \setminus A_\varepsilon^{4\delta}
 \end{aligned}$$

for small $\varepsilon > 0$ and for some constants $c_1, C_0, C > 0$. Thus, using the definition of χ_ε and (76) we see that

$$(77) \quad \int_{B(0, R_0/\varepsilon) \setminus A_\varepsilon^{4\delta}} \chi_\varepsilon(x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx \leq C_3 \varepsilon^{-(2N-1)-5(p+1)/(p-1)} \exp(-c_1/\varepsilon)$$

and

$$(78) \quad \int_{B(0, 2R_0/\varepsilon) \setminus B(0, R_0/\varepsilon)} \chi_\varepsilon(x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx \leq C_4 \varepsilon^{-(2(p+1)/(p-1)+2\gamma+N)} \exp(-c_1/\varepsilon)$$

for some constants $C_3, C_4 > 0$, independent of ε . Moreover, from (Q_2) , Lemma 6 and (71) it follows that, for some constants $C_5, c_2 > 0$,

$$(79) \quad Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) \leq C_5 \varepsilon^{-2(p+1)/(p-1)} \exp(-c_2/\varepsilon) (2R_0)^{-(p+1)\omega_\varepsilon} \varepsilon^{(p+1)\omega_\varepsilon} |x|^{(p+1)\omega_\varepsilon}$$

for all $x \in \mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)$ and small $\varepsilon > 0$. Then, combining (79) with the definition of χ_ε , we have

$$\begin{aligned}
 &\int_{\mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)} \chi_\varepsilon(x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx \\
 &\leq C_5 \exp(-c_2/\varepsilon) (2R_0)^{-(p+1)\omega_\varepsilon} \\
 &\quad \times \varepsilon^{-2(p+1)/(p-1)+(p+1)\omega_\varepsilon-\gamma} \int_{\mathbb{R}^N \setminus B(0, 2R_0/\varepsilon)} |x|^{\gamma+(p+1)\omega_\varepsilon} dx \\
 (80) \quad &= C_6 \frac{1}{-\gamma - (p+1)\omega_\varepsilon - N} (2R_0)^{\gamma+N} \varepsilon^{-2(p+1)/(p-1)-2\gamma-N} \exp(-c_2/\varepsilon)
 \end{aligned}$$

for some constant $C_6 > 0$. From (77), (78), (80) and of the fact that $\chi_\varepsilon \equiv 0$ in $A_\varepsilon^{4\delta}$, we deduce that

$$\lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^N} \chi_\varepsilon(x) Q(\bar{u}_\varepsilon, \bar{v}_\varepsilon) dx = 0.$$

This proves (25).

As a consequence of (25) we have $\beta_\varepsilon = 0$. Using (7), the homogeneity of Q , Lemma 3 and $(\bar{u}_\varepsilon, \bar{v}_\varepsilon)$ as test function in $(S_{\alpha_\varepsilon}, 0)$, we obtain $\frac{M_\varepsilon}{p+1} = \alpha_\varepsilon$. This implies that $(u_\varepsilon, v_\varepsilon)$, where $u_\varepsilon(x) = (\alpha_\varepsilon)^{1/(p-1)} \bar{u}_\varepsilon(\varepsilon^{-1}x)$ and $v_\varepsilon(x) = (\alpha_\varepsilon)^{1/(p-1)} \bar{v}_\varepsilon(\varepsilon^{-1}x)$, satisfies (S) . Note that of (64) and (71), we have

$$(81) \quad u_\varepsilon(x), v_\varepsilon(x) \leq C \exp(-c\delta/\varepsilon) \forall x \in B(0, 2R_0) \setminus A^{4\delta}$$

and

$$(82) \quad u_\varepsilon(x), v_\varepsilon(x) \leq C \exp(-c/\varepsilon)(|x|/2R_0)^{\omega_\varepsilon} \quad \forall x \in \mathbb{R}^N \setminus B(0, 2R_0).$$

The property (1) is proved in Lemma 5. We now show the property (2), i.e.,

$$\liminf_{\varepsilon \rightarrow 0} \varepsilon^{-2/(p-1)} \|u_\varepsilon + v_\varepsilon\|_{L^\infty(\mathbb{R}^N)} > 0.$$

We define $W_\varepsilon \equiv \varepsilon^{-2/(p-1)}(u_\varepsilon + v_\varepsilon)$. Then, it suffices to show that

$$\liminf_{\varepsilon \rightarrow 0} \|W_\varepsilon\|_{L^\infty(\mathbb{R}^N)} > 0.$$

From (Q_1) and of the fact that $(u_\varepsilon, v_\varepsilon)$ is solution of (S) we see that, for some $C > 0$,

$$(83) \quad -\Delta W_\varepsilon + \frac{1}{\varepsilon^2} V(x)W_\varepsilon \leq CK(x)(W_\varepsilon)^p \text{ in } \mathbb{R}^N.$$

Multiplying both sides of (83) by W_ε and integrating by parts, we obtain

$$(84) \quad \int_{\mathbb{R}^N} [|\nabla W_\varepsilon|^2 + \frac{1}{\varepsilon^2} V(x)(W_\varepsilon)^2] dx \leq C \int_{\mathbb{R}^N} (W_\varepsilon)^{p+1} dx \\ \leq C \|W_\varepsilon\|_{L^\infty(\mathbb{R}^N)}^{p-1} \int_{A^{5\delta}} (W_\varepsilon)^2 dx + C \|W_\varepsilon\|_{L^\infty(\mathbb{R}^N)}^{(p-1)/2} \int_{\mathbb{R}^N \setminus A^{5\delta}} (W_\varepsilon)^{(p-1)/2} (W_\varepsilon)^2 dx.$$

Now, to conclude our proof once more, we will use the arguments developed by Byeon-Wang in [17] and [18]. We take $\varphi \in C_0^\infty(\text{int}(A^{5\delta}))$ such that $\varphi(x) = 1$ for $x \in A^{4\delta}$. As $\inf_{x \in \text{supp}(\varphi) \setminus A^{4\delta}} V(x) > 0$ and $\inf_{x \in A^{5\delta} \setminus A^{4\delta}} V(x) > 0$, it follows, by definition of φ and by the Poincaré inequality, that

$$(85) \quad \int_{A^{5\delta}} (W_\varepsilon)^2 dx \leq 2 \int_{A^{5\delta}} [(\varphi W_\varepsilon)^2 + (1 - \varphi)^2 (W_\varepsilon)^2] dx \\ \leq C_0 \int_{\text{supp}(\varphi)} |\nabla(\varphi W_\varepsilon)|^2 dx + C_1 \frac{1}{\varepsilon^2} \int_{A^{5\delta} \setminus A^{4\delta}} V(x)(W_\varepsilon)^2 dx \\ \leq C_2 \frac{1}{\varepsilon^2} \int_{\text{supp}(\varphi) \setminus A^{4\delta}} V(x)(W_\varepsilon)^2 dx + 2C_0 \int_{\mathbb{R}^N} |\nabla W_\varepsilon|^2 dx \\ + C_1 \frac{1}{\varepsilon^2} \int_{A^{5\delta} \setminus A^{4\delta}} V(x)(W_\varepsilon)^2 dx \\ \leq C \int_{\mathbb{R}^N} [|\nabla W_\varepsilon|^2 + \frac{1}{\varepsilon^2} V(x)(W_\varepsilon)^2] dx$$

for some positive constants C_0, C_1, C_2 and C , independent of small $\varepsilon > 0$. On the other hand, using the Hölder inequality and Sobolev embedding results, we get

$$(86) \quad \int_{\mathbb{R}^N \setminus A^{5\delta}} (W_\varepsilon)^{(p-1)/2} (W_\varepsilon)^2 dx \leq C \left(\int_{\mathbb{R}^N \setminus A^{5\delta}} (W_\varepsilon)^{N(p-1)/4} dx \right)^{2/N} \int_{\mathbb{R}^N} |\nabla W_\varepsilon|^2 dx.$$

In view of (81) and (82) we deduce that, for some constants $C_1, C_2 > 0$,

$$(87) \quad W_\varepsilon(x) \leq C_1 \varepsilon^{-2/(p-1)} \exp(-c\delta/\varepsilon) \quad \forall x \in B(0, 2R_0) \setminus A^{4\delta}$$

and

$$(88) \quad W_\varepsilon(x) \leq C_2 \varepsilon^{-2/(p-1)} \exp(-c/\varepsilon) (2R_0)^{-\omega_\varepsilon} |x|^{\omega_\varepsilon} \quad \forall x \in \mathbb{R}^N \setminus B(0, 2R_0).$$

From (87) and (88) we have, for some constants $c_1, c_2, C_3, C_4, C_5 > 0$, that

$$(89) \quad \begin{aligned} & \int_{\mathbb{R}^N \setminus A^{5\delta}} (W_\varepsilon)^{N(p-1)/4} dx \\ & \leq C_3 \varepsilon^{-N/2} \exp(-c_1/\varepsilon) + C_4 \varepsilon^{-N/2} \exp(-c_2/\varepsilon) (2R_0)^{-N(p-1)\omega_\varepsilon/4} \\ & \quad \times \int_{\mathbb{R}^N \setminus B(0, 2R_0)} |x|^{N(p-1)\omega_\varepsilon/4} dx \\ & = C_3 \varepsilon^{-N/2} \exp(-c_1/\varepsilon) + C_5 \varepsilon^{-N/2} \exp(-c_2/\varepsilon) (2R_0)^N \frac{1}{-N(p-1)\omega_\varepsilon/4 - N}. \end{aligned}$$

From (89), $\int_{\mathbb{R}^N \setminus A^{5\delta}} (W_\varepsilon)^{N(p-1)/4} dx \leq 1$ for sufficiently small $\varepsilon > 0$. This and (86) imply

$$(90) \quad \int_{\mathbb{R}^N \setminus A^{5\delta}} (W_\varepsilon)^{(p-1)/2} (W_\varepsilon)^2 dx \leq C \int_{\mathbb{R}^N} [|\nabla W_\varepsilon|^2 + \frac{1}{\varepsilon^2} V(x) (W_\varepsilon)^2] dx$$

for small $\varepsilon > 0$. From (84), (85) and (90) it follows that $\|W_\varepsilon\|_{L^\infty(\mathbb{R}^N)}^{p-1} + \|W_\varepsilon\|_{L^\infty(\mathbb{R}^N)}^{(p-1)/2} \geq C$ for some positive constant C . Then $\|W_\varepsilon\|_{L^\infty(\mathbb{R}^N)} \geq C_1 > 0$, where C_1 is a constant independent of $\varepsilon > 0$. This completes the proof of property (2).

Now, we claim that $u_\varepsilon, v_\varepsilon \in W^{1,2}(\mathbb{R}^N)$. In fact, from (81) and (82) we obtain

$$(91) \quad \begin{aligned} \int_{\mathbb{R}^N} (u_\varepsilon)^2 dx & \leq \int_{A^{4\delta}} (u_\varepsilon)^2 dx + C_0 \exp(-c_0/\varepsilon) \\ & \quad + C_1 \exp(-c_1/\varepsilon) (2R_0)^{-2\omega_\varepsilon} \int_{\mathbb{R}^N \setminus B(0, 2R_0)} |x|^{2\omega_\varepsilon} dx \\ & = \int_{A^{4\delta}} (u_\varepsilon)^2 dx + C_0 \exp(-c_0/\varepsilon) \\ & \quad + C_2 \exp(-c_1/\varepsilon) (2R_0)^N \frac{1}{-2\omega_\varepsilon - N} \end{aligned}$$

for some constants $C_0, C_1, C_2, c_0, c_1 > 0$. Using the Lemma 3 and change of variables, we have

$$\begin{aligned} \varepsilon^2 \int_{\mathbb{R}^N} |\nabla u_\varepsilon(x)|^2 dx & = (\alpha_\varepsilon)^{2/(p-1)} \int_{\mathbb{R}^N} |\nabla \bar{u}_\varepsilon(x/\varepsilon)|^2 dx \\ & = \varepsilon^N (\alpha_\varepsilon)^{2/(p-1)} \int_{\mathbb{R}^N} |\nabla \bar{u}_\varepsilon(y)|^2 dy \end{aligned}$$

$$\begin{aligned}
 &\leq \varepsilon^N (\alpha_\varepsilon)^{2/(p-1)} \|(\bar{u}_\varepsilon, \bar{v}_\varepsilon)\|_\varepsilon^2 \\
 (92) \quad &= (p+1) (\alpha_\varepsilon)^{(p+1)/(p-1)} \varepsilon^N.
 \end{aligned}$$

From (91) and (92) we conclude that $u_\varepsilon \in W^{1,2}(\mathbb{R}^N)$. Similarly, we obtain $v_\varepsilon \in W^{1,2}(\mathbb{R}^N)$. The decay property (3) follows from (81) and (82). \square

Acknowledgement. The second author would like to thank the Department of Mathematic of UFMG for their hospitality during his visit.

References

- [1] N. Akhmediev and A. Ankiewicz, *Partially coherent solitons on a finite background*, Phys. Rev. Lett. **82** (1999), 2661–2664.
- [2] C. O. Alves, *Local mountain pass for a class of elliptic system*, J. Math. Anal. Appl. **335** (2007), no. 1, 135–150.
- [3] C. O. Alves and S. H. M. Soares, *Existence and concentration of positive solutions for a class gradient systems*, Nonlinear Differential Equations Appl. **12** (2005), no. 4, 437–457.
- [4] A. Ambrosetti, V. Felli, and A. Malchiodi, *Ground states of nonlinear Schrödinger equations with potentials vanishing at infinity*, J. Eur. Math. Soc. **7** (2005), no. 1, 117–144.
- [5] A. Ambrosetti and A. Malchiodi, *Perturbation Methods and Semilinear Elliptic Problems on \mathbb{R}^N* , Progr. Math., Birkhäuser 240, Boston, 2006.
- [6] A. Ambrosetti, A. Malchiodi, and W.-M. Ni, *Singularly perturbed elliptic equations with symmetry: existence of solutions concentrating on spheres, I*, Comm. Math. Phys. **235** (2003), no. 3, 427–466.
- [7] A. Ambrosetti, A. Malchiodi, and D. Ruiz, *Bound states of nonlinear Schrödinger equations with potentials vanishing at infinity*, J. Anal. Math. **98** (2006), 317–348.
- [8] A. Ambrosetti and P. H. Rabinowitz, *Dual variational methods in critical point theory and applications*, J. Funct. Anal. **14** (1973), 349–381.
- [9] A. Ambrosetti and D. Ruiz, *Radial solutions concentrating on spheres of nonlinear Schrödinger equations with vanishing potentials*, Proc. Roy. Soc. Edinburgh Sect. A **136** (2006), no. 5, 889–907.
- [10] A. Ambrosetti and Z.-Q. Wang, *Nonlinear Schrödinger equations with vanishing and decaying potentials*, Differential Integral Equations **18** (2005), no. 12, 1321–1332.
- [11] M. Badiale, V. Benci, and S. Rolando, *A nonlinear elliptic equation with singular potential and applications to nonlinear field equations*, J. Eur. Math. Soc. **9** (2007), no. 3, 355–381.
- [12] M. Badiale and T. d’Aprile, *Concentration around a sphere for a singularly perturbed Schrödinger equation*, Nonlinear Anal. **49** (2002), no. 7, Ser. A: Theory Methods, 947–985.
- [13] T. Bartsch and S. Peng, *Semiclassical symmetric Schrödinger equations: Existence of solutions concentrating simultaneously on several spheres*, Z. Angew. Math. Phys. **58** (2007), no. 5, 778–804.
- [14] J. Byeon, *Existence of many nonequivalent nonradial positive solutions of semilinear elliptic equations on three-dimensional annuli*, J. Differential Equations **136** (1997), no. 1, 136–165.
- [15] J. Byeon and L. Jeanjean, *Standing waves for nonlinear Schrödinger equations with a general nonlinearity*, Arch. Ration. Mech. Anal. **185** (2007), no. 2, 185–200.
- [16] J. Byeon and Z.-Q. Wang, *Spherical semiclassical states of a critical frequency for Schrödinger equations with decaying potentials*, J. Eur. Math. Soc. **8** (2006), no. 2, 217–228.

- [17] ———, *Standing waves with a critical frequency for nonlinear Schrödinger equations*, Arch. Ration. Mech. Anal. **165** (2002), no. 4, 295–316.
- [18] ———, *Standing waves for nonlinear Schrödinger equations with singular potentials*, Ann. Inst. H. Poincaré Anal. Non Linéaire **26** (2009), no. 3, 943–958.
- [19] ———, *Standing waves with a critical frequency for nonlinear Schrödinger equations. II*, Calc. Var. Partial Differential Equations **18** (2003), no. 2, 207–219.
- [20] D. N. Christodoulides, T. H. Coskun, M. Mitchell, and M. Segev, *Theory of incoherent self-focusing in biased photorefractive media*, Phys. Rev. Lett. **78** (1997), 646–649.
- [21] R. Cipolatti and W. Zumpichiatti, *On the existence and regularity of ground states for a nonlinear system of coupled Schrödinger equations in \mathbb{R}^N* , Comput. Appl. Math. **18** (1999), no. 1, 15–29.
- [22] ———, *Orbitally stable standing waves for a system of coupled nonlinear Schrödinger equations*, Nonlinear Anal. **42** (2000), no. 3, Ser. A: Theory Methods, 445–461.
- [23] E. N. Dancer and S. Yan, *A new type of concentration solutions for a singularly perturbed elliptic problem*, Trans. Amer. Math. Soc. **359** (2007), no. 4, 1765–1790.
- [24] A. Floer and A. Weinstein, *Nonspreading wave packets for the cubic Schrödinger equation with a bounded potential*, J. Funct. Anal. **69** (1986), no. 3, 397–408.
- [25] B. D. Esry, C. H. Greene, J. P. Burke, Jr., and J. L. Bohn, *Hartree-Fock theory for double condensates*, Phys. Rev. Lett. **78** (1997), 3594–3597.
- [26] D. Gilbarg and N. S. Trudinger, *Elliptic Partial Differential Equations of Second Order*, 2nd ed., Grundlehren Math. 224, Springer-Verlag, Berlin Heidelberg, 1983.
- [27] A. Hasegawa and Y. Kodama, *Solitons in Optical Communications*, Academic Press, San Diego, 1995.
- [28] M. N. Islam, *Ultrafast Fiber Switching Devices and Systems*, Cambridge University Press, New York, 1992.
- [29] I. P. Kaminow, *Polarization in optical fibers*, IEEE J. Quantum Electron. **17** (1981), 15–22.
- [30] E. H. Lieb and R. Seiringer, *Proof of Bose-Einstein condensation for dilute trapped gases*, Phys. Rev. Lett. **88** (2002), 170409.
- [31] L. A. Maia, E. Montefusco, and B. Pellacci, *Positive solutions for a weakly coupled nonlinear Schrödinger system*, J. Differential Equations **229** (2006), no. 2, 743–767.
- [32] C. R. Menyuk, *Nonlinear pulse propagation in birefringence optical fiber*, IEEE J. Quantum Electron. **23** (1987), 174–176.
- [33] ———, *Pulse propagation in an elliptically birefringent Kerr medium*, IEEE J. Quantum Electron. **25** (1989), 2674–2682.
- [34] P. Meystre, *Atom Optics*, Springer-Verlag, New York, 2001.
- [35] D. L. Mills, *Nonlinear Optics*, Springer-Verlag, Berlin, 1998.
- [36] D. C. de Moraes Filho and M. A. S. Souto, *Systems of p -laplacean equations involving homogeneous nonlinearities with critical Sobolev exponent degrees*, Comm. Partial Differential Equations **24** (1999), no. 7-8, 1537–1553.
- [37] A. Pomponio, *Coupled nonlinear Schrödinger systems with potentials*, J. Differential Equations **227** (2006), no. 1, 258–281.

PAULO CESAR CARRIÃO
 DEPARTAMENTO DE MATEMÁTICA
 UNIVERSIDADE FEDERAL DE MINAS GERAIS
 31270-010 BELO HORIZONTE(MG), BRAZIL
 E-mail address: carrion@mat.ufmg.br

NARCISO HORTA LISBOA
DEPARTAMENTO DE CIÊNCIAS EXATAS
UNIVERSIDADE ESTADUAL DE MONTES CLAROS
39401-089 MONTES CLAROS(MG), BRAZIL
E-mail address: narciso.lisboa@unimontes.br

OLIMPIO HIROSHI MIYAGAKI
DEPARTAMENTO DE MATEMÁTICA
UNIVERSIDADE FEDERAL DE JUIZ DE FORA
36036-330 JUIZ DE FORA(MG), BRAZIL
E-mail address: ohmiyagaki@gmail.com