

A Design of Mid-wave Infrared Integral Catadioptric Optical System with Wide FOV

Lin yao Yu, Hong guang Jia*, Qun Wei, Hu hai Jiang, Tian yi Zhang, and Chao Wang

Changchun Institute of Optics, Fine Mechanics and Physics, CAS, Changchun 130033, China

(Received January 22, 2013 : revised March 11, 2013 : accepted March 21, 2013)

In order to deduce the difficulty of fixing the Ritchey-Chretien (R-C) dual reflective optical system and enhance the stability of the secondary mirror, a compact integral structure is presented here composed of two transmitting and two reflective aspheric surfaces. The four surfaces were manufactured from a single germanium lens and integrated together. The two reflective surfaces formed by coating the inner reflecting films were assembled in one lens. It makes the installation of the two mirrors easier and the structure of the secondary mirror more stable. A design of mid-wave infrared (MWIR) compact imaging system is presented with a spectral range chosen as 3.7-4.8 μm . The effective focal length is $f=90$ mm. The field of view (FOV) for the lens is 4.88° . It has good imaging capability with Modulation Transfer Function (MTF) of all field of view more than 0.55 close to the diffraction limitation. Outdoor experiments were carried out and it is shown that the integral catadioptric optical system performs well on imaging.

Keywords : Optical design, R-C systems, Infrared catadioptric, Aspherical

OCIS codes : (080.2740) Geometric optical design; (080.4035) Mirror system design; (120.4570) Optical design of instruments

I. INTRODUCTION

In recent years, the miniaturization and lightweight are paid more and more attention in a variety of optical system designs for aircraft and missiles [1]. Despite the small FOV, Ritchey-Chretien (R-C) dual reflective optical system has become the main choice for the domestic and foreign researchers because of its smaller, lighter structure and larger caliber [2-4]. However the assembly and the support mounting of the secondary mirror have been the main difficulties for Ritchey-Chretien dual reflective optical system: (1) the tolerance of R-C system is very sensitive which makes the assembly and adjustment difficult; (2) in order to reduce the obscuration of light caused by the brackets holding the secondary mirror, the size of the brackets usually is limited which leads to the instability of the secondary mirror [5]. The brackets will also be deformed when the missiles are accelerating forward. It decreases the image quality seriously, and reduces the accuracy of the guidance [4, 5].

An integral catadioptric structure is adopted in this paper. The major and secondary mirrors were manufactured from a single lens so that there is no need to assemble and

adjust them. The supporting structure of the secondary mirror will be removed which will not only increase the light irradiating the primary mirror, but also enhance the stability of imaging. The accuracy of the assembly and adjustment is converted to the process of manufacturing and now with the advanced single point diamond turning machine the machining accuracy can meet the requirements sufficiently. An optical design with the wavelength 3.7~4.8 μm was presented and 100% cold aperture efficiency was implemented. The imaging quality was close to the diffractive limitation.

II. R-C REFLECTIVE SYSTEM DESIGN

In order to design the integral catadioptric system, an initial R-C dual-reflective system ought to be designed first. Then the integral catadioptric can be achieved by adding material that transmits mid-infrared rays between the two mirrors. The final system can be gained after the optimization using Zemax or Code V. The structure of classical R-C reflective system is shown in Fig. 1. Changing the surface of primary mirror into hyperboloid can not only correct

*Corresponding author: jiahg@ciomp.ac.cn

Color versions of one or more of the figures in this paper are available online.

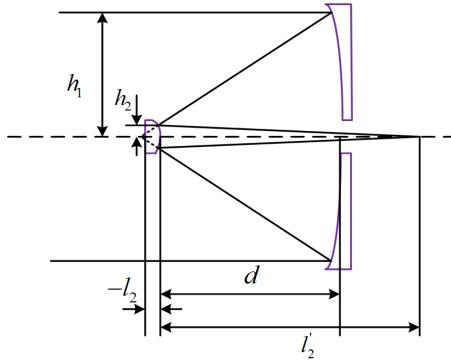


FIG. 1. R-C optical system configuration.

the spherical and coma aberration, but also enlarge the field of the optical system.

The meanings of the symbols are demonstrated in Fig. 1. The distance of the target is infinite $l_1 = \infty$, the incidence angle $u_1 = 0$; α is the obscuration ratio of the secondary mirror; β is the magnification of the secondary mirror, $\beta = l_2/l_1$; the focal length of the R-C system is the focal length of primary mirror multiplying the magnification of the secondary mirror; e_1^2 is the quadratic coefficient of the primary mirror; e_2^2 is the quadratic coefficient of the secondary mirror. From the Gaussian optics formula the curvature of the secondary mirror can be derived [6].

$$R_2 = \frac{\alpha\beta}{1+\beta} R_1 \quad (1)$$

There is no chromatic aberration in the R-C reflective system. When the monochromatic aberration is taken into consideration, the quadratic coefficient of the two mirrors can be derived in the aplanatic condition. The five monochromatic aberrations are spherical aberration, coma, astigmatism, field curvature, distortion. S_I , S_{II} denotes the Seidel coefficient of spherical aberration, coma respectively which can be expressed in equation (2) and (3) [7]:

$$S_I = \sum_{i=1}^m h_i P_i + \sum_{i=1}^m h_i^4 K_i \quad (2)$$

$$S_{II} = \sum_{i=1}^m y_i P_i - J \sum_{i=1}^m W_i + \sum_{i=1}^m h_i^3 y_i K_i \quad (3)$$

Where m is the total number of the surfaces. In R-C reflective system m equals 2, and

$$P = \left(\frac{u' - u}{\frac{1}{n'} - \frac{1}{n}} \right)^2 \left(\frac{u'}{n'} - \frac{u}{n} \right) \quad (4)$$

$$W = \left(\frac{u' - u}{\frac{1}{n'} - \frac{1}{n}} \right) \left(\frac{u'}{n'} - \frac{u}{n} \right) \quad (5)$$

$$K = -\frac{e^2}{R^3} (n' - n) \quad (6)$$

In the above equations n and n' are the refractive index of object space and image space, respectively. For R-C reflective system, the refractive index can be concluded as $n_1 = n_2 = 1$, $n'_1 = n'_2 = -1$. R is the vertex curvature of the mirror. For computational advantage, we define the height of the light irradiating at the primary mirror as $h_1 = 1$, the focal length of the system as $f = 1$ and the angle of incidence light as $u_1 = 0$. Then we can obtain, $u_1 = u_2 = \beta$, $u'_2 = 1$, $J = 1$, $R_1 = 2/\alpha$, $R_2 = 2\alpha/(\beta+1)$. The equation (4)~(6) can be written as follows

$$P_1 = -\frac{\beta^3}{4}, P_2 = \frac{(1-\beta)^2(1+\beta)}{4} \quad (7)$$

$$W_1 = \frac{\beta^2}{2}, W_2 = \frac{1-\beta^2}{2} \quad (8)$$

$$K_1 = \frac{e_1^2}{4} \beta^3, K = -\frac{e_2^2 (1+\beta)^3}{4 \alpha^3} \quad (9)$$

Then the spherical aberration, coma can be written as follows:

$$S_I = \left[\frac{\alpha(\beta-1)^2(\beta+1)}{4} - \frac{\alpha(\beta+1)^3}{4} e_2^2 \right] - \frac{\beta^3}{4} (1-e_1^2) \quad (10)$$

$$S_{II} = \frac{1-\alpha}{\alpha} \left[\frac{\alpha(\beta+1)^3}{4\beta} e_2^2 - \frac{\alpha(\beta-1)^2(\beta+1)}{4} \right] - \frac{1}{2} \quad (11)$$

From the first order design and the initial structure parameters as α and β , equation (10) and (11) can be evaluated. The conic constants correcting for $S_I = 0$, $S_{II} = 0$ are given as follows:

$$e_1^2 = 1 + \frac{2\alpha}{(1-\alpha)\beta^2} \quad (12)$$

$$e_2^2 = \frac{\frac{2\beta}{1-\alpha} + (1+\beta)(1-\beta)^2}{(1+\beta)^3} \quad (13)$$

The initial structure of the R-C reflective system can be calculated by equation (1), (12), (13) for a given focal length of primary mirror, the obscuration ratio and the magnification of secondary mirror.

III. OPTICAL SYSTEM DESIGN AND ANALYSIS

3.1. Optical Design Requirements

The working brand of the optical system is 3.7~4.8 μm and the center wavelength is 4.2 μm . The medium wave detector is J-T refrigeration type, 320×256 array plane and the pixel size is 30 μm . The efficiency focal length of the system is 90 mm. F number is 1.86.

The field of view (FOV) can be calculated from equation (14) and (15)

$$\tan(FOV_1/2) = \frac{a}{f} \quad (14)$$

$$\tan(FOV_2/2) = \frac{b}{f} \quad (15)$$

Where a and b are the half length and half breadth of the focal plane in the detector, f is the effective focal length of optical system. In this optical system a is 4.8 mm, b is 3.84 mm and f is 90 mm. Then the FOV_2 of the MIR system is 4.88°. In this design we take the FOV_2 as the

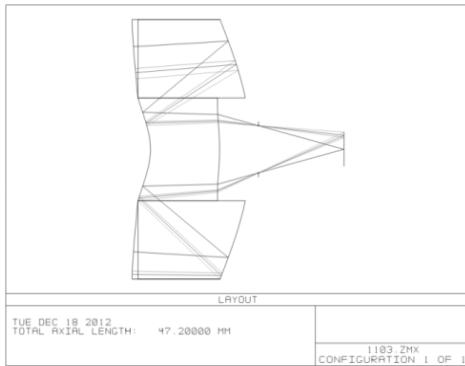


FIG. 2. Optical configuration of the catadioptric system.

full field of view.

The center obscuration is less than 40%. In order to resist the stray light irradiation, the cold aperture efficiency should be 100%. The full field energy in a pixel should be more than 80%.

3.2. Optical Design Results

After optimizing the system for several times using Zemax-EE, an integral catadioptric system with a total length of 47 mm was realized as shown in Fig. 2. The system is made of four parts and the structure of the major part is shown in Fig. 3. In the front surface there are the incidence surface and the secondary mirror. In the rear surface there are the emergent surface and the primary mirror. The lens data is shown in Table 1 from ZEMAX lens data format.

The first surface of the lens is a concave surface which can effectively reduce the angle when light irradiates at the primary mirror. So this surface can not only decrease the off axis aberration of the system, especially spherical aberration and coma, but also enlarge the FOV of the optical system.

The fourth surface of the lens is convex. It can provide certain positive focal power which will decrease the power caused by the two mirrors. In this situation the redistribution of the power can reduce the aberration of the system. How to resist stray light is an important problem needed to be concerned in the reflective system [8]. In the design, the

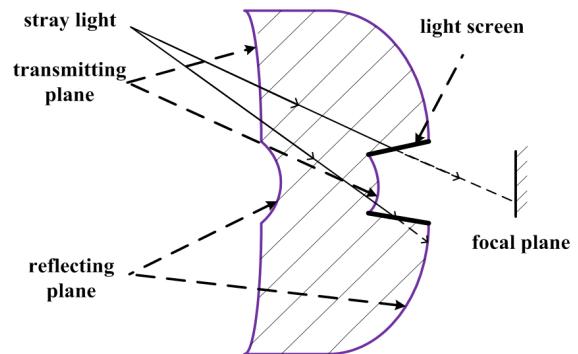


FIG. 3. Structure of the catadioptric system.

TABLE 1. Lens data from ZEMAX (unit: thickness/radius in mm.)

| Suface | Type | Radius | Thickness | Glass |
|--------|--------------|----------|-----------|-----------|
| OBJ | Standard | Infinity | Infinity | Air |
| 1 | Even Asphere | -109.67 | 25.6 | GERMANIUM |
| 2 | Even Asphere | -52.13 | -23.2 | Mirror |
| 3 | Even Asphere | -14.26 | 16 | Mirror |
| 4 | Even Asphere | -48.93 | 9 | Air |
| STO | Standard | Infinity | 19.8 | Air |
| IMA | Standard | Infinity | - | Air |

emergent surface is placed between the two mirrors, which can shorten the length of the system and make it more compact. Secondly the circumference formed between the primary mirror and the emergent surface can work as a light screen by blacking. The stray light from other FOV will be restricted, for instance the light irradiating at the primary mirror or directly at the focal plane of the detector as shown in Fig. 3.

Due to the large field of view, a lot of off-axis aberration is induced. The two transmission and two reflective surfaces are all high order aspherical surfaces used to correct the aberration when we finished the optimization at last. The whole lens is made of germanium which is easy for manufacturing. The surface sag is shown as equation (16)

$$z = \frac{cr^2}{1 + \sqrt{1 - (1 + c^2)r^2}} + \sum_{i=1}^6 \alpha_i r^{2i} \quad (16)$$

Where c is the curvature (the reciprocal of the radius), r is the radial coordinate in lens unit, k is the conic constant, and α_i is the coefficient. The even aspheric parameters are shown in Table 2.

The aperture stop is placed at 19.8 mm from the image plane and the diameter of exit pupil is set 10.6 mm to make the cold aperture efficiency reach 100%.

The optical transfer function, spot diagram and energy encircle curve were shown in Figure 4, Figure 5 and Figure 6, wherein the optical transfer function reached 0.55 close to the diffraction limit at the Nyquist frequency (16.67 lines per millimeter) and the geometry of the spot size was less than 35 μm . The full field energy in a pixel was more than 95%. Due to the wide FOV of the system, distortion is needed to be considered as shown in Fig. 7. The maximum distortion is less than 0.5% which can be ignored. The optical system presented an excellent image quality.

TABLE 2. Even aspheric parameters

| Surface | Conic(k) | 4 th Order Term | 6 th Order Term | 8 th Order Term | 10 th Order Term |
|---------|----------|----------------------------|----------------------------|----------------------------|-----------------------------|
| 1 | -19.82 | 2.175e-6 | -9.774e-10 | 2.746e-13 | 2.161e-16 |
| 2 | -4.36 | -2.054e-6 | 2.281e-9 | -1.062e-12 | 4.064e-16 |
| 3 | -3.07 | 7.84e-5 | -4.054e-7 | 1.356e-9 | -2.338e-12 |
| 4 | 0.92 | 4.564e-5 | 1.946e-7 | 4.23e-10 | -1.183e-11 |

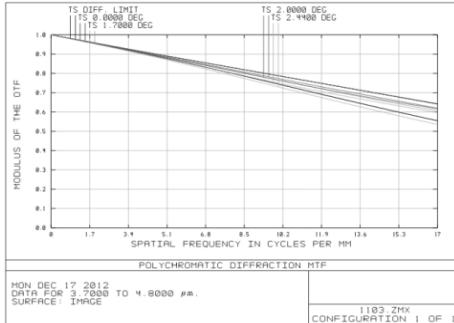


FIG. 4. MTF plot for the catadioptric system.

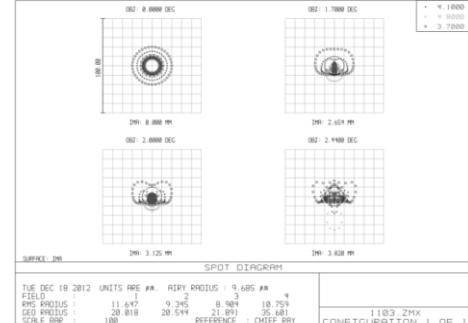


FIG. 5. Spot diagram for the catadioptric system.

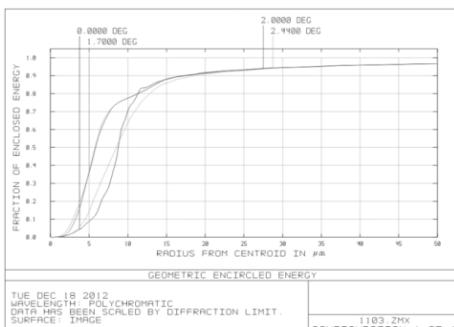


FIG. 6. Energy encircle Curve for the catadioptric system.

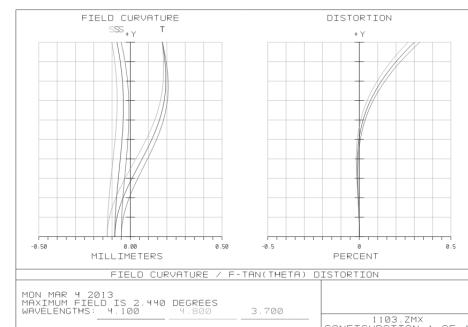


FIG. 7. Field curvature and distortion curve for the catadioptric system.

TABLE 3. Optical tolerances for manufacturing

| Item | Tolerances |
|----------------------------|-----------------|
| Surface 1 | ± 5 fringe |
| Surface 2 | ± 3 fringe |
| Surface 3 | ± 3 fringe |
| Surface 4 | ± 5 fringe |
| Thickness (2-3) | ± 0.01 mm |
| Thickness (other surfaces) | ± 0.02 mm |
| Tilt | ± 30 arcsec |

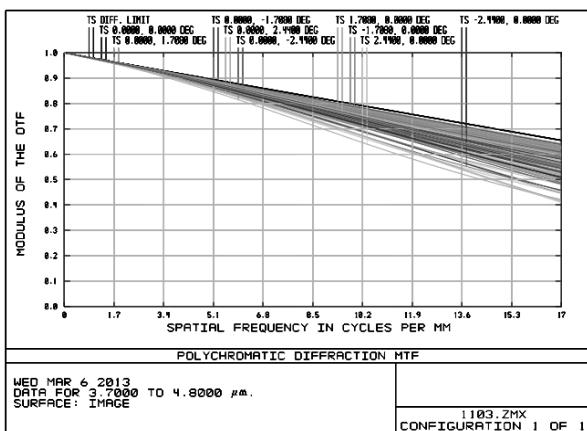


FIG. 8. MTF plot after Mont Carlo calculations.

3.3. The Optical Tolerances for Manufacturing

Just the manufacturing errors concerned, the optical tolerances are shown in Table 3. The thickness between the two mirrors is found to be very sensitive by changing the thickness manually and the tolerance is set ± 0.01 mm. For the other thickness, since they are insensitive, the tolerance is set ± 0.02 mm. Since all the manufacturing was completed in the advanced single point diamond turning machine which can provide very high precision at a scale of 10nm, we can conclude that the tolerances are loose and reasonable. With Monte Carlo calculations the MTF will be changed 0.18 for 17lp/mm at most in the Zemax program as shown in Fig. 8.

3.4. Experiments and Results Analysis

The integral catadioptric lens was finally manufactured using the advanced single point diamond turning machine and assembled with mechanical parts as shown in Fig. 9 (a). The set up of the optical experiment platform of the MWIR imaging system goes as shown in Fig. 9 (b). Two outdoor experiments were carried out to test the performance of the imaging system at night. In Fig. 10 (a) the targets are some buildings. From the image we can identify the edge of the building clearly. In Fig. 10 (b) an air bus was chosen to image. The configuration was distinct and even the heated air out of the engines can be seen from the

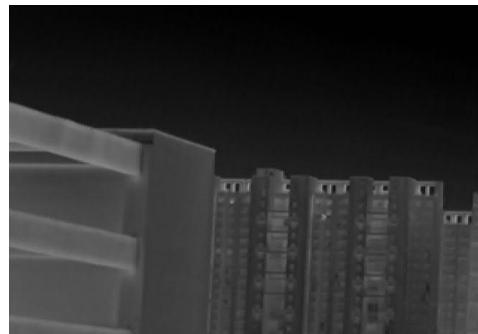


(a)

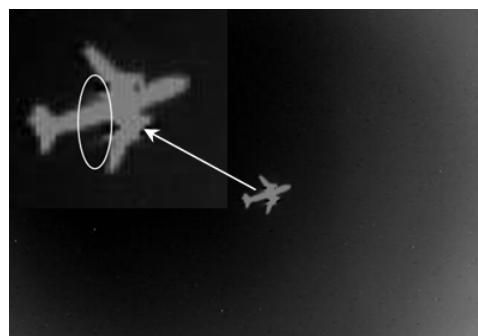


(b)

FIG. 9. The integral catadioptric system: (a)the integral catadioptric lens after manufacuring and asemmbled with mechanical structures; (b)the whole MIR imaging system platform.



(a)



(b)

FIG. 10. Two outdoor experiments carried out at night for different targets: (a) the image of buildings (the distance is about 2 km); (b) the image of an air bus and its amplification.

image. These experiments proved that the integral catadioptric system had an excellent performance for MWIR imaging.

IV. CONCLUSION

This paper presented an MWIR integral catadioptric imaging system. The focal length of the system was 90 mm manufactured from a single germanium lens. The cold aperture efficiency was 100%. This kind of system was easier to assemble and adjust for the secondary mirror is integrated in the lens, providing a very compact structure and making the fixing easier. The circumference formed between the primary mirror and the emergent surface efficiently reduces the stray light and enhances the image quality in the focal plane. With spectral range 3.7-4.8 μm , the system transfer function achieved at 0.55 in 171 p/mm which is close to the diffraction limit. Outdoor imaging experiments were carried out with real integral catadioptric system which provided a good imaging quality.

ACKNOWLEDGMENT

This research is supported by the forecast research of arming under Grant NO. 51301060207.

REFERENCES

- C. Yi and H. Xu, "State of the art and future trend of detectors for infrared imaging seekers," *Infrared and Laser Engineering* **35**, 7-11 (2006).
- R. Blakley, "Cesarrian telescope optical system," *Opt. Eng.* **35**, 3338-3341 (1996).
- J. H. Lee, T. S. Jang, H. S. Yang, and S. W. Rhee, "Optical design of a compact imaging spectrometer for STSAT3," *J. Opt. Soc. Korea* **12**, 262-268 (2008).
- L. Epstein, "Improved geometry for the all reflecting Schmidt telescope," *Appl. Opt.* **12**, 926-928 (1973).
- X. N. Fu, X. J. Cheng, and W. K. Wu, "Design of large aperture mirror support and its stability," *Opt. Precision Eng.* **16**, 179-184 (2008).
- W. D. Joo, M. S. Jung, and S. A. Rodionov, "The general concept of Gaussian beams," *Proc. SPIE* **2778**, 67-68 (1996).
- J. H. Pan, *The Design, Manufacture and Test of the Aspheric Optical Surfaces* (Science Press, Beijing, China, 1994), pp. 10-47.
- J. O. Park, W. K. Jang, S. Kim, H. S. Jang, and S. Lee, "Stray light analysis of high resolution camera for a low-earth-orbit satellite," *J. Opt. Soc. Korea* **15**, 52-55 (2011).