

ON THE EXISTENCE OF POSITIVE SOLUTION FOR A CLASS OF NONLINEAR ELLIPTIC SYSTEM WITH MULTIPLE PARAMETERS AND SINGULAR WEIGHTS

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ABSTRACT. This study concerns the existence of positive solution for the following nonlinear system

$$\begin{cases} -\operatorname{div}(|x|^{-ap} |\nabla u|^{p-2} \nabla u) = |x|^{-(a+1)p+c_1} (\alpha_1 f(v) + \beta_1 h(u)), & x \in \Omega, \\ -\operatorname{div}(|x|^{-bq} |\nabla v|^{q-2} \nabla v) = |x|^{-(b+1)q+c_2} (\alpha_2 g(u) + \beta_2 k(v)), & x \in \Omega, \\ u = v = 0, & x \in \partial\Omega, \end{cases}$$

where Ω is a bounded smooth domain of \mathbb{R}^N with $0 \in \Omega$, $1 < p, q < N$, $0 \leq a < \frac{N-p}{p}$, $0 \leq b < \frac{N-q}{q}$ and $c_1, c_2, \alpha_1, \alpha_2, \beta_1, \beta_2$ are positive parameters. Here $f, g, h, k : [0, \infty) \rightarrow [0, \infty)$ are nondecreasing continuous functions and

$$\lim_{s \rightarrow \infty} \frac{f(Ag(s)^{\frac{1}{q-1}})}{s^{p-1}} = 0$$

for every $A > 0$.

We discuss the existence of positive solution when f, g, h and k satisfy certain additional conditions. We use the method of sub-super solutions to establish our results.

1. Introduction

The paper deal with the existence of positive solution for the nonlinear system

$$(1) \quad \begin{cases} -\operatorname{div}(|x|^{-ap} |\nabla u|^{p-2} \nabla u) = |x|^{-(a+1)p+c_1} (\alpha_1 f(v) + \beta_1 h(u)), & x \in \Omega, \\ -\operatorname{div}(|x|^{-bq} |\nabla v|^{q-2} \nabla v) = |x|^{-(b+1)q+c_2} (\alpha_2 g(u) + \beta_2 k(v)), & x \in \Omega, \\ u = v = 0, & x \in \partial\Omega, \end{cases}$$

where Ω is a bounded smooth domain of \mathbb{R}^N with $0 \in \Omega$, $1 < p, q < N$, $0 \leq a < \frac{N-p}{p}$, $0 \leq b < \frac{N-q}{q}$ and $c_1, c_2, \alpha_1, \alpha_2, \beta_1, \beta_2$ are positive parameters. Here $f, g, h, k : [0, \infty) \rightarrow [0, \infty)$ are nondecreasing continuous functions.

Elliptic problems involving more general operator, such as the degenerate quasilinear elliptic operator given by $-\operatorname{div}(|x|^{-ap} |\nabla u|^{p-2} \nabla u)$, were motivated

Received May 21, 2011.

2010 *Mathematics Subject Classification.* 35J55, 35J65.

Key words and phrases. singular weights, nonlinear elliptic system, multiple parameters.

by the following Caffarelli, Kohn and Nirenberg's inequality (see [10], [25]). The study of this type of problem is motivated by its various applications, for example, in fluid mechanics, in newtonian fluids, in flow through porous media and in glaciology (see [8, 13]). On the other hand, quasilinear elliptic systems has an extensive practical background. It can be used to describe the multiplicate chemical reaction catalyzed by the catalyst grains under constant or variant temperature, it can be used in the theory of quasiregular and quasiconformal mappings in Riemannian manifolds with boundary (see [17, 23]) and can be a simple model of tubular chemical reaction, more naturally, it can be a correspondence of the stable station of dynamical system determined by the reaction-diffusion system, see Ladde and Lakshmikantham et al. [21]. More naturally, it can be the populations of two competing species [15]. So, the study of positive solutions of elliptic systems has more practical meanings. We refer to [1], [2], [9], [18] for additional results on elliptic problems.

For the regular case, that is, when $a = b = 0$, $c_1 = p$ and $c_2 = q$, the quasilinear elliptic equation has been studied by several authors (see [3, 4, 12]). See [5, 14] where the authors discussed the system (1) when $p = q = 2$, $\alpha_1 = \alpha_2$, $\beta_1 = \beta_2 = 0$, f, g are increasing and $f, g \geq 0$. In [19], the authors extended the study of [14], to the case when no sign conditions on $f(0)$ or $g(0)$ were required and in [20] they extend this study to the case when $p = q > 1$. Here we focus on further extending the study in [3] for the quasilinear elliptic systems involving singularity. Due to this singularity in the weights, the extensions are challenging and nontrivial. Our approach is based on the method of sub-super solutions, see [11, 16]. Several methods have been used to treat quasilinear equations and systems. In the scalar case, weak solutions can be obtained through variational methods which provide critical points of the corresponding energy functional, an approach which is also fruitful in the case of potential systems, i.e., the nonlinearities on the right hand side are the gradient of a C^1 -functional [7]. However, due to the loss of the variational structure, the treatment of nonvariational systems like (1) is more complicated and is based mostly on topological methods [6].

2. Preliminaries

In this paper, we denote $W_0^{1,p}(\Omega, \|x\|^{-ap})$, the completion of $C_0^\infty(\Omega)$, with respect to the norm $\|u\| = (\int_\Omega \|x\|^{-ap} |\nabla u|^p dx)^{\frac{1}{p}}$. To precisely state our existence result we consider the eigenvalue problem

$$(2) \quad \begin{cases} -div(|x|^{-sr} |\nabla \phi|^{r-2} \nabla \phi) = \lambda |x|^{-(s+1)r+t} |\phi|^{r-2} \phi, & x \in \Omega, \\ \phi = 0, & x \in \partial\Omega. \end{cases}$$

For $r = p$, $s = a$ and $t = c_1$, let $\phi_{1,p}$ be the eigenfunction corresponding to the first eigenvalue $\lambda_{1,p}$ of (2) such that $\phi_{1,p}(x) > 0$ in Ω , and $\|\phi_{1,p}\|_\infty = 1$ and for $r = q$, $s = b$ and $t = c_2$, let $\phi_{1,q}$ be the eigenfunction corresponding to the first eigenvalue $\lambda_{1,q}$ of (2) such that $\phi_{1,q}(x) > 0$ in Ω , and $\|\phi_{1,q}\|_\infty = 1$

(see [22, 24]). It can be shown that $\frac{\partial \phi_{1,r}}{\partial n} < 0$ on $\partial\Omega$ for $r = p, q$. Here n is the outward normal. This result is well known and hence, depending on Ω , there exist positive constants $m, \delta, \sigma_p, \sigma_q$ such that

$$(3) \quad \lambda_{1,r} |x|^{-(s+1)r+t} \phi_{1,r}^r - |x|^{-sr} |\nabla \phi_{1,r}|^r \leq -m, \quad x \in \bar{\Omega}_\delta,$$

$$(4) \quad \phi_{1,r} \geq \sigma_r, \quad x \in \Omega_0 = \Omega \setminus \bar{\Omega}_\delta,$$

with $r = p, q; s = a, b; t = c_1, c_2$ and $\bar{\Omega}_\delta = \{x \in \Omega \mid d(x, \partial\Omega) \leq \delta\}$ (see [22]). We will also consider the unique solution $(\zeta_p(x), \zeta_q(x)) \in W_0^{1,p}(\Omega, \|x\|^{-ap}) \times W_0^{1,q}(\Omega, \|x\|^{-bq})$ for the system

$$\begin{cases} -\operatorname{div}(|x|^{-ap} |\nabla \zeta_p|^{p-2} \nabla \zeta_p) = |x|^{-(a+1)p+c_1}, & x \in \Omega, \\ -\operatorname{div}(|x|^{-bq} |\nabla \zeta_q|^{q-2} \nabla \zeta_q) = |x|^{-(b+1)q+c_2}, & x \in \Omega, \\ u = v = 0, & x \in \partial\Omega, \end{cases}$$

to discuss our existence result. It is known that $\zeta_r(x) > 0$ in Ω and $\frac{\partial \zeta_r(x)}{\partial n} < 0$ on $\partial\Omega$, for $r = p, q$ (see [22]).

3. Existence results

In this section, we shall establish our existence result via the method of sub and supersolutions. A pair of nonnegative functions $(\psi_1, \psi_2), (z_1, z_2)$ are called a subsolution and supersolution of (1) if they satisfy $(\psi_1, \psi_2) = (0, 0) = (z_1, z_2)$ on $\partial\Omega$ and

$$\begin{aligned} & \int_{\Omega} |x|^{-ap} |\nabla \psi_1|^{p-2} |\nabla \psi_1| \cdot \nabla w \, dx \\ & \leq \int_{\Omega} |x|^{-(a+1)p+c_1} (\alpha_1 f(\psi_2) + \beta_1 h(\psi_1)) w \, dx, \\ & \int_{\Omega} |x|^{-bq} |\nabla \psi_2|^{q-2} |\nabla \psi_2| \cdot \nabla w \, dx \\ & \leq \int_{\Omega} |x|^{-(b+1)q+c_2} (\alpha_2 g(\psi_1) + \beta_2 k(\psi_2)) w \, dx, \\ & \int_{\Omega} |x|^{-ap} |\nabla z_1|^{p-2} |\nabla z_1| \cdot \nabla w \, dx \\ & \geq \int_{\Omega} |x|^{-(a+1)p+c_1} (\alpha_1 f(z_2) + \beta_1 h(z_1)) w \, dx, \\ & \int_{\Omega} |x|^{-bq} |\nabla z_2|^{q-2} |\nabla z_2| \cdot \nabla w \, dx \\ & \geq \int_{\Omega} |x|^{-(b+1)q+c_2} (\alpha_2 g(z_1) + \beta_2 k(z_2)) w \, dx \end{aligned}$$

for all $w \in W = \{w \in C_0^\infty(\Omega) \mid w \geq 0, x \in \Omega\}$. Then the following result holds:

Lemma 3.1 (See [22]). *Suppose there exist sub and super-solutions (ψ_1, ψ_2) and (z_1, z_2) respectively of (1) such that $(\psi_1, \psi_2) \leq (z_1, z_2)$. Then (1) has a solution (u, v) such that $(u, v) \in [(\psi_1, \psi_2), (z_1, z_2)]$.*

We make the following assumptions:

(H1) $f, g, h, k : [0, \infty) \rightarrow [0, \infty)$ are C^1 nondecreasing functions such that

$$\lim_{s \rightarrow \infty} f(s) = \lim_{s \rightarrow \infty} g(s) = \lim_{s \rightarrow \infty} h(s) = \lim_{s \rightarrow \infty} k(s) = \infty.$$

(H2) For all $A > 0$,

$$\lim_{s \rightarrow \infty} \frac{f(Ag(s)^{\frac{1}{q-1}})}{s^{p-1}} = 0.$$

(H3)

$$\lim_{s \rightarrow \infty} \frac{h(s)}{s^{p-1}} = \lim_{s \rightarrow \infty} \frac{k(s)}{s^{q-1}} = 0.$$

Now we are ready to state our existence result.

Theorem 3.2. *Assume (H1)-(H3) hold. Then there exists a positive large solution of system (1) when $\alpha_1 + \beta_1$ and $\alpha_2 + \beta_2$ are large.*

Proof. Since f, g, h, k are continuous and nondecreasing, we have $f(x), g(x), h(x), k(x) \geq -k_0$ for all $x \geq 0$ and for some $k_0 > 0$. Choose $r > 0$ such that

$$r \leq \min\{|x|^{-(a+1)p+c_1}, |x|^{-(b+1)q+c_2}\},$$

in $\bar{\Omega}_\delta$. We shall verify that

$$\begin{aligned} & (\psi_1, \psi_2) \\ &= \left(\left[\frac{(\alpha_1 + \beta_1)k_0r}{m} \right]^{\frac{1}{p-1}} \left(\frac{p-1}{p} \right) \phi_{1,p}^{\frac{p}{p-1}}, \left[\frac{(\alpha_2 + \beta_2)k_0r}{m} \right]^{\frac{1}{q-1}} \left(\frac{q-1}{q} \right) \phi_{1,q}^{\frac{q}{q-1}} \right), \end{aligned}$$

is a sub-solution of (1). Let $w \in W$. Then a calculation shows that

$$\begin{aligned} & \int_{\Omega} |x|^{-ap} |\nabla \psi_1|^{p-2} \nabla \psi_1 \nabla w \, dx \\ &= \left(\frac{(\alpha_1 + \beta_1)k_0r}{m} \right) \int_{\Omega} |x|^{-ap} \phi_{1,p} |\nabla \phi_{1,p}|^{p-2} \nabla \phi_{1,p} \nabla w \, dx \\ &= \left(\frac{(\alpha_1 + \beta_1)k_0r}{m} \right) \int_{\Omega} |x|^{-ap} |\nabla \phi_{1,p}|^{p-2} \nabla \phi_{1,p} [\nabla(\phi_{1,p}w) - |\nabla \phi_{1,p}|^p w] \, dx \\ &= \left(\frac{(\alpha_1 + \beta_1)k_0r}{m} \right) \int_{\Omega} [\lambda_{1,p} |x|^{-(a+1)p+c_1} \phi_{1,p}^p - |x|^{-ap} |\nabla \phi_{1,p}|^p] w \, dx. \end{aligned}$$

Similarly

$$\begin{aligned} & \int_{\Omega} |x|^{-bq} |\nabla \psi_2|^{q-2} \nabla \psi_2 \nabla w \, dx \\ &= \left(\frac{(\alpha_2 + \beta_2)k_0r}{m} \right) \int_{\Omega} [\lambda_{1,q} |x|^{-(b+1)q+c_2} \phi_{1,q}^q - |x|^{-bq} |\nabla \phi_{1,q}|^q] w \, dx. \end{aligned}$$

First we consider the case when $x \in \bar{\Omega}_\delta$. We have $\lambda_{1,p} |x|^{-(a+1)p+c_1} \phi_{1,p}^p - |x|^{-ap} |\nabla \phi_{1,p}|^p \leq -m$ on $\bar{\Omega}_\delta$. Since $\psi_1(x), \psi_2(x) \geq 0$ in Ω , it follows that

$$-k_0r \leq \min\{|x|^{-(a+1)p+c_1} f(\psi_2), |x|^{-(a+1)p+c_1} h(\psi_1)\},$$

in $\bar{\Omega}_\delta$. Hence, we have

$$\begin{aligned} & \left(\frac{(\alpha_1 + \beta_1)k_0r}{m} \right) \int_{\bar{\Omega}_\delta} [\lambda_{1,p} |x|^{-(a+1)p+c_1} \phi_{1,p}^p - |x|^{-ap} |\nabla \phi_{1,p}|^p] w \, dx \\ & \leq -(\alpha_1 + \beta_1)k_0r \int_{\bar{\Omega}_\delta} w \, dx \\ & \leq \int_{\bar{\Omega}_\delta} |x|^{-(a+1)p+c_1} (\alpha_1 f(\psi_2) + \beta_1 h(\psi_1)) w \, dx. \end{aligned}$$

A similar argument shows that

$$\begin{aligned} & \left(\frac{(\alpha_2 + \beta_2)k_0r}{m} \right) \int_{\bar{\Omega}_\delta} [\lambda_{1,q} |x|^{-(b+1)q+c_2} \phi_{1,q}^q - |x|^{-bq} |\nabla \phi_{1,q}|^q] w \, dx \\ & \leq \int_{\bar{\Omega}_\delta} |x|^{-(b+1)q+c_2} (\alpha_2 g(\psi_1) + \beta_2 k(\psi_2)) w \, dx. \end{aligned}$$

On the other hand, on $\Omega \setminus \bar{\Omega}_\delta$, since $\phi_{1,p} \geq \sigma_p, \phi_{1,q} \geq \sigma_q$ for some $0 < \sigma_p, \sigma_q < 1$, if $\alpha_1 + \beta_1$ and $\alpha_2 + \beta_2$ are large, then by (H1) we have

$$f(\psi_2), h(\psi_1), g(\psi_1), k(\psi_2) \geq \frac{k_0r}{m} \max\{\lambda_{1,p}, \lambda_{1,q}\}.$$

Hence

$$\begin{aligned} & \left(\frac{(\alpha_1 + \beta_1)k_0r}{m} \right) \int_{\Omega \setminus \bar{\Omega}_\delta} [\lambda_{1,p} |x|^{-(a+1)p+c_1} \phi_{1,p}^p - |x|^{-ap} |\nabla \phi_{1,p}|^p] w \, dx \\ & \leq \left(\frac{(\alpha_1 + \beta_1)k_0r}{m} \right) \int_{\Omega \setminus \bar{\Omega}_\delta} |x|^{-(a+1)p+c_1} \lambda_{1,p} w \, dx \\ & \leq \int_{\Omega \setminus \bar{\Omega}_\delta} |x|^{-(a+1)p+c_1} (\alpha_1 f(\psi_2) + \beta_1 h(\psi_1)) w \, dx. \end{aligned}$$

Similarly,

$$\begin{aligned} & \left(\frac{(\alpha_2 + \beta_2)k_0r}{m} \right) \int_{\Omega \setminus \bar{\Omega}_\delta} [\lambda_{1,q} |x|^{-(b+1)q+c_2} \phi_{1,q}^q - |x|^{-bq} |\nabla \phi_{1,q}|^q] w \, dx \\ & \leq \int_{\Omega \setminus \bar{\Omega}_\delta} |x|^{-(b+1)q+c_2} (\alpha_2 g(\psi_1) + \beta_2 k(\psi_2)) w \, dx. \end{aligned}$$

Hence

$$\begin{aligned} & \int_{\Omega} |x|^{-ap} |\nabla \psi_1|^{p-2} |\nabla \psi_1| \cdot \nabla w \, dx \\ & \leq \int_{\Omega} |x|^{-(a+1)p+c_1} (\alpha_1 f(\psi_2) + \beta_1 h(\psi_1)g) w \, dx, \\ & \int_{\Omega} |x|^{-bq} |\nabla \psi_2|^{q-2} |\nabla \psi_2| \cdot \nabla w \, dx \\ & \leq \int_{\Omega} |x|^{-(b+1)q+c_2} (\alpha_2 g(\psi_1) + \beta_2 k(\psi_2)) w \, dx, \end{aligned}$$

i.e., (ψ_1, ψ_2) is a sub-solution of (1).

Now, we will prove there exists a M large enough so that

$$(z_1, z_2) = \left(M \zeta_p(x), (\alpha_2 + \beta_2)^{\frac{1}{q-1}} [g(M \|\zeta_p\|_\infty)]^{\frac{1}{q-1}} \zeta_q(x) \right),$$

is a super-solution of (1). A calculation shows that:

$$\begin{aligned} \int_{\Omega} |x|^{-ap} |\nabla z_1|^{p-2} \nabla z_1 \nabla w \, dx &= M^{p-1} \int_{\Omega} |x|^{-ap} |\nabla \zeta_p|^{p-2} \nabla \zeta_p \nabla w \, dx \\ &= M^{p-1} \int_{\Omega} |x|^{-(a+1)p+c_1} w \, dx. \end{aligned}$$

By (H2)-(H3) we can choose M large enough so that

$$\begin{aligned} M^{p-1} &\geq \alpha_1 f \left((\alpha_2 + \beta_2)^{\frac{1}{q-1}} \|\zeta_q\|_\infty g(M \|\zeta_p\|_\infty)^{\frac{1}{q-1}} \right) + \beta_1 h(M \|\zeta_p\|_\infty) \\ &\geq \alpha_1 f \left((\alpha_2 + \beta_2)^{\frac{1}{q-1}} \zeta_q(x) g(M \|\zeta_p\|_\infty)^{\frac{1}{q-1}} \right) + \beta_1 h(M \zeta_p(x)) \\ &= \alpha_1 f(z_2) + \beta_1 h(z_1). \end{aligned}$$

Hence

$$\begin{aligned} &\int_{\Omega} |x|^{-ap} |\nabla z_1|^{p-2} |\nabla z_1| \cdot \nabla w \, dx \\ &\geq \int_{\Omega} |x|^{-(a+1)p+c_1} (\alpha_1 f(z_2) + \beta_1 h(z_1)) w \, dx. \end{aligned}$$

Next, by (H3) for M large enough we have

$$g(M \|\zeta_p\|_\infty) \geq k \left((\alpha_2 + \beta_2)^{\frac{1}{q-1}} g(M \|\zeta_p\|_\infty)^{\frac{1}{q-1}} \|\zeta_q\|_\infty \right).$$

Hence

$$\begin{aligned} (5) \quad &\int_{\Omega} |x|^{-bq} |\nabla z_2|^{q-2} \nabla z_2 \nabla w \, dx \\ &= (\alpha_2 + \beta_2) \int_{\Omega} |x|^{-(b+1)q+c_2} g(M \|\zeta_p\|_\infty) w \, dx \\ &\geq \int_{\Omega} |x|^{-(b+1)q+c_2} [\alpha_2 g(z_1) + \beta_2 g(M \|\zeta_p\|_\infty)] w \, dx \\ &\geq \int_{\Omega} |x|^{-(b+1)q+c_2} \left[\alpha_2 g(z_1) + \beta_2 k \left((\alpha_2 + \beta_2)^{\frac{1}{q-1}} g(M \|\zeta_p\|_\infty)^{\frac{1}{q-1}} \|\zeta_q\|_\infty \right) \right] w \, dx \\ &\geq \int_{\Omega} |x|^{-(b+1)q+c_2} [\alpha_2 g(z_1) + \beta_2 k(z_2)] w \, dx, \end{aligned}$$

i.e., (z_1, z_2) is a super-solution of (1) with $z_i \geq \psi_i$ for M large, $i = 1, 2$. Thus, by [22] there exists a positive solution (u, v) of (1) such that $(\psi, \psi) \leq (u, v) \leq (z_1, z_2)$. This completes the proof of Theorem 3.2. \square

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