

Design of Tunable Flat-top Bandpass Filter Based on Two Long-period Fiber Gratings and Core Mode Blocker

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We propose a tunable flat-top bandpass filter to pass light in a customized wavelength band by using long-period fiber gratings (LPFG) structure. The LPFG structure is composed of a core mode blocker in between two LPFGs. The bandpass spectrum of the proposed structure is obtained in overlapped wavelength band of two LPFGs operating on the same modes. To analyze the properties, we introduce a mathematical matrix model for the structure. We theoretically demonstrate flexibility of the flat-top bandpass filter with various bandwidths.

Keywords : Flat-top bandpass, Long-period fiber gratings, Core mode blocker, Matrix model
OCIS codes : (350.2770) Gratings; (350.2460) Filters, interference; (230.2285) Fiber devices and optical amplifiers; (070.1170) Analog optical signal processing

I. INTRODUCTION

In long-period fiber gratings (LPFG), a fundamental core mode and the multiple cladding modes are coupled, which all propagate in the same (forward) direction. Due to unique features of low insertion loss, low back-reflection and excellent polarization insensitivity, LPFGs have attracted great interest in the optical communications and sensor applications [1, 2]. Some researchers have studied the manufacture of the core mode blocker [3, 4], band pass filter using the core mode blocker [5-7], and tunable bandpass filter using coil heater [8, 9].

The core mode blocker is a device to block off the propagated light in the core, and pass the propagated light in the cladding. Some researchers have developed bandpass filters using LPFGs, and some of them have studied the topics based on a core mode blocker as shown in Fig. 1. In Fig. 1, the LPFG₁ is used to interact between core and cladding at resonant wavelength, and light existing in the

core is extinguished by the core mode blocker. The light in the resonant wavelength is propagated through the cladding to the second LPFG. The resonant light existing in the cladding by the LPFG₂ is transferred back into the core. The central points of the research are to study the physical phenomenon [3, 4] and the manufacture of the bandpass filter [5-7, 9]. However the design of the flat-top bandpass filter [11] with specific customized wavelength band using an accurate analysis model was not proposed yet.

In the paper, we present the matrix model to analyze the

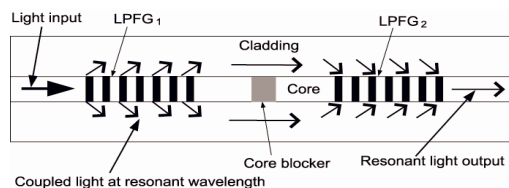


FIG. 1. The LPFG structure with core mode blocker.

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LPFG structures with the core mode blocker. The proposed matrix model is a useful mathematical tool to describe the properties of the structure. We analyze the physical properties of the LPFG structure with a core mode blocker by using the proposed model. By using specially coupled feature of only the same modes in overlapped wavelength region, we demonstrate the design of the flat-top bandpass filter with the customized bandwidth.

II. THE PROPOSED MATHEMATICAL MATRIX MODEL

The structure in Fig. 1 is modeled as the FIR (Finite Impulse Response) filter block diagram as shown in Fig. 2(a). Figure 2(b) is the signal flow graph to analyze both LPFG₁ and LPFG₂, which is called the multiport lattice filter model [10-12]. Here, D_1 and D_2 are the delay of the diagonal matrix, CB is matrix for core mode blocker, E_{co} and $E_{cl}^{(i)}$ are the E-fields (electric fields) of the fundamental core mode and the i th cladding mode, respectively. And β_{co} and $\beta_{cl}^{(i)}$ are the propagation constant of the fundamental core mode and the i th cladding mode for the LPFGs, L_1 and L_2 are the length of LPFG₁ and LPFG₂, and d_1 , d_2 , and d_b are the length for D_1 , D_2 , and CB .

In Fig. 2(a), the E-fields coming into and out from the structure can be written to be

$$\begin{bmatrix} E_{co}(out) \\ E_{cl}^{(1)}(out) \\ \vdots \\ E_{cl}^{(N)}(out) \end{bmatrix} = Q \begin{bmatrix} E_{co}(in) \\ E_{cl}^{(1)}(in) \\ \vdots \\ E_{cl}^{(N)}(in) \end{bmatrix}, \quad Q \equiv M_2 D_2 C B D_1 M_1 \quad (1)$$

where $D_k = \text{diag}(e^{-j\beta_{co}d_k}, e^{-j\beta_{cl}^{(1)}d_k}, \dots, e^{-j\beta_{cl}^{(N)}d_k})$, ($k = 1, 2$) are

diagonal matrices and $CB = \text{diag}(0, e^{-j\beta_{cl}^{(1)}d_b}, \dots, e^{-j\beta_{cl}^{(N)}d_b})$ is a diagonal matrix with zero determinant.

By assuming the LPFG_k, ($k = 1, 2$) are uniform long period fiber gratings, we can easily get the matrix form of the $(N+1) \times (N+1)$ complex matrix M_k , ($k = 1, 2$) for the LPFG_k with codirectional interactions since the LPFG_k has the multiport lattice filter structure in [10] as follows:

$$M_k \equiv D_G G^{(N)} G^{(N-1)} \dots G^{(1)}, \quad (2)$$

where $G^{(p)}$ is $(N+1) \times (N+1)$ identity matrix except for the four entries adopted from $F^{(p)}$ in Fig. 2(b). The relation between $G^{(p)}$ and $F^{(p)}$ are as follows:

$$\begin{aligned} G_{1,1}^{(p)} &= F_{1,1}^{(p)}, & G_{1,p+1}^{(p)} &= F_{1,2}^{(p)}, & G_{p+1,1}^{(p)} &= F_{2,1}^{(p)} = -(F_{1,2}^{(p)})^*, \\ G_{p+1,p+1}^{(p)} &= F_{2,2}^{(p)} = (F_{1,1}^{(p)})^*. \end{aligned}$$

(Here, $B_{k,l}$ denotes the (k, l) element of the matrix B .) The matrix D_G represents the phase shift of the LPFG₂, $D_G = \text{diag}(e^{-j\beta_{co}L_k}, e^{-j\beta_{cl}^{(1)}L_k}, \dots, e^{-j\beta_{cl}^{(N)}L_k})$ and

$$F^{(p)} \equiv \begin{bmatrix} (\cos(S_p L_k) - j \frac{\delta_p}{S_p} \sin(S_p L_k)) e^{j\delta_p L_k} & -j \frac{\kappa_p}{S_p} \sin(S_p L_k) e^{j\delta_p L_k} \\ -j \frac{\kappa_p}{S_p} \sin(S_p L_k) e^{-j\delta_p L_k} & (\cos(S_p L_k) + j \frac{\delta_p}{S_p} \sin(S_p L_k)) e^{-j\delta_p L_k} \end{bmatrix}, \quad (3)$$

and $S_p \equiv \sqrt{(K_p)^2 + (\delta_p)^2}$, L_k is the length of the LPFG_k, δ_p is the detuning factor for the p th cladding mode, and κ_p is the coupling coefficient between the fundamental core mode and the p th cladding mode [10-13].

Note that M_k in (2) can be regarded as a multiport lattice section. Because of the decoupling property of modes, the matrices $G^{(p)}$ in M_k commute with each other so that the

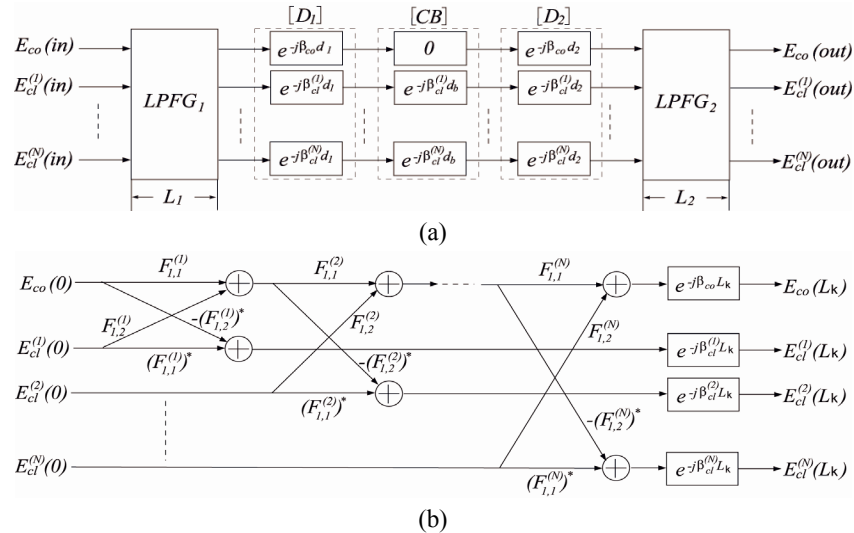


FIG. 2. (a) An equivalent block diagram for Fig. 1. (b) Multiport lattice filter model.

2×2 subsections may be interchanged freely. If LPFG₁ and LPFG₂ are extended as piecewise uniform LPFGs, we can also easily get the M_k using the multiport lattice filter model [10]. Now, the overall transmission coefficient $t \equiv E_{\omega}(out)/E_{\omega}(in)$ with $E_{cl}^{(i)}(in) = 0$, ($1 \leq i \leq N$) is easily seen to be $t = Q_{1,1}$ which is the (1,1)-element of Q .

As shown in Fig. 2(a), after the coupled light on the LPFG₁ is absorbed at CB, and then in the LPFG₂, the propagated light in the cladding is coupled with the same modes from cladding to core, due to the orthogonality of modes as shown in Fig. 2(b) [10]. Because the overall transmission is calculated as the complex multiplying of two bandpass filters, we cannot get the spectrum in the core if two filters have entirely different bandpass regions.

III. EXAMPLES

We analytically calculated the transmission spectrum curves for the concatenated LPFGs with a core mode blocker in the following examples by computing the transmission coefficient $t = Q_{1,1}$ in (1). The parameters of the fiber used in these examples are as follows:

$n_{co} = 1.449$, $n_{cl} = 1.444$, $n_{air} = 1$, $r_{co} = 4.5 \mu\text{m}$, $r_{cl} = 62.5 \mu\text{m}$ where n_{cl} is the refractive index of the cladding, n_{air} is the refractive index of air, r_{co} is the radius of the core, and r_{cl} is the radius of the cladding.

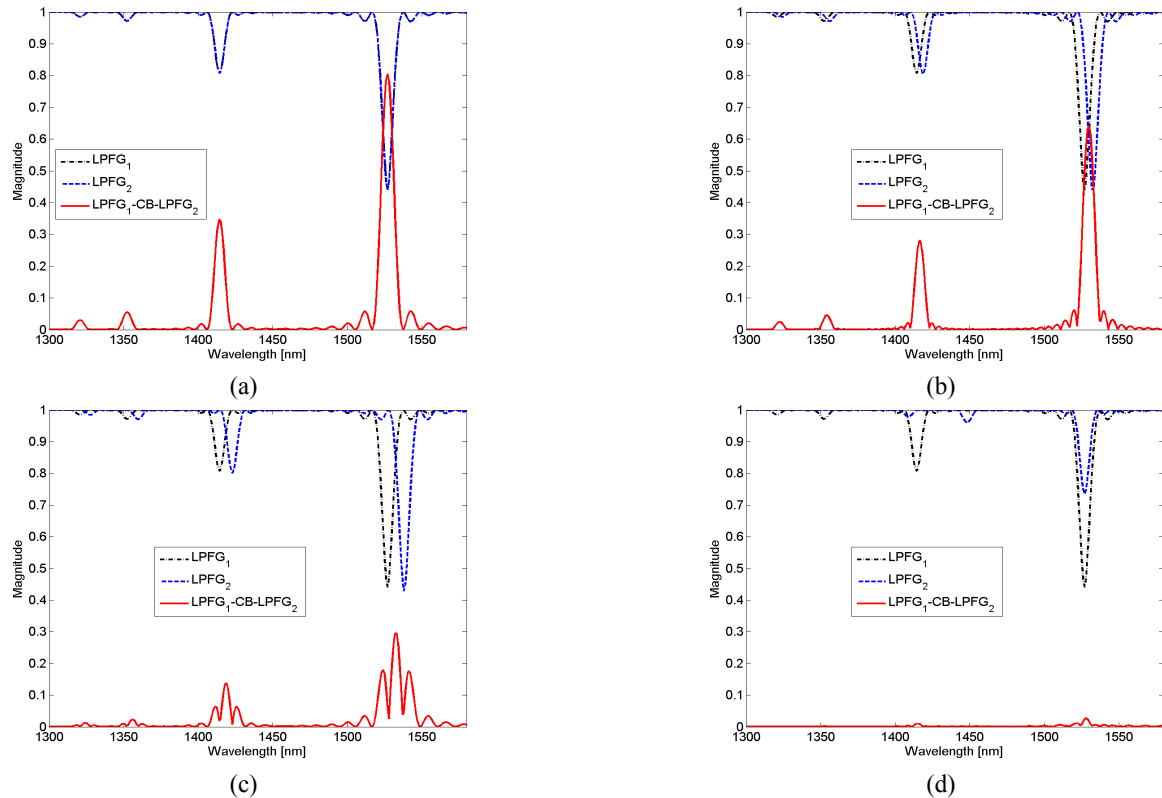


FIG. 3. $\Lambda_1 = 441.24 \mu\text{m}$ is fixed and Λ_2 is changed ((a) $\Lambda_2 = 441.24 \mu\text{m}$, (b) $\Lambda_2 = 443.44 \mu\text{m}$ (c) $\Lambda_2 = 445.65 \mu\text{m}$ and (d) $\Lambda_2 = 501.28 \mu\text{m}$).

3.1. Example 1 for Uniform LPFG with Both Core Mode Blocker and Delay

Figure 3 shows the transmission for the modes LP_{0i} , ($i = 1, \dots, 4$) in the wavelength range between 1300 nm and 1580 nm, where LPFG₁ and LPFG₂ have uniform gratings. Their lengths are $L_1 = 100\Lambda_1$ and $L_2 = 100\Lambda_2$, respectively, the grating period for LPFG₁ is $\Lambda_1 = 441.24 \mu\text{m}$, Λ_2 for LPFG₂ is varied ($441.24 \mu\text{m}$ for Fig. 3(a), $443.44 \mu\text{m}$ for Fig. 3(b), $445.65 \mu\text{m}$ for Fig. 3(c), and $501.28 \mu\text{m}$ for Fig. 3(d) and their induced index changes are $\Delta n_1 = n_2 = 0.00011$, respectively. The length of the D₁ and the D₂ are $d_1 = 400 \Lambda_1$ and $d_2 = 400\Lambda_1$, respectively. The length of the CB is assumed to be $d_b = 0$.

To analyze the properties of the concatenated LPFGs with the core mode blocker, we have calculated the transmission spectra along the change of the Λ_2 as shown in Fig 3. Figures 3(a)-3(c) show the transmission along the overlapped region of the spectrum band of the same modes. We can only get the bandpass spectra by coupling the same modes in the overlapped region. Figure 3(d) shows the transmission for the coupling with the different modes (LP_{04} of LPFG₁ and LP_{03} for LPFG₂). From the result, we know that the bandpass filter cannot be obtained from coupling the different modes.

3.2. Example 2 for Flat-top Bandpass Filters

We can design a flat-top bandpass filter with the customized bandwidth using the property of coupling with the same modes and the overlapped band. We have synthesized

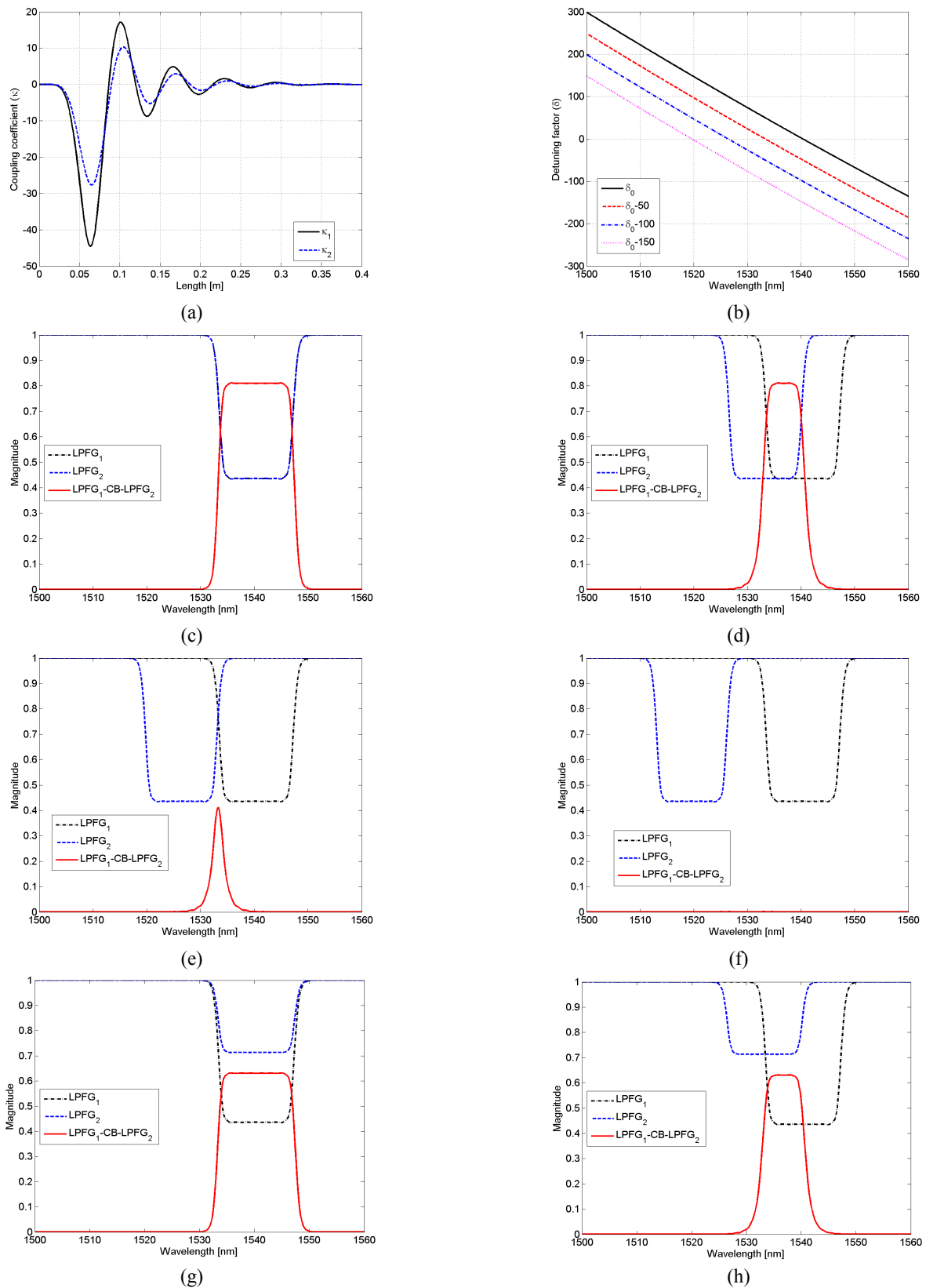


FIG. 4. (a) Coupling coefficients. (b) Detuning factors. (c) The κ_1 and the δ_0 are utilized for the LPFG₁ and the LPFG₂. (d) LPFG₁ (κ_1 and δ_0 is used.) and LPFG₂ (κ_1 and $\delta_0=50$ is used.). (e) LPFG₁ (κ_1 and δ_0 is used.) and LPFG₂ (κ_1 and $\delta_0=100$ is used.). (f) LPFG₁ (κ_1 and δ_0 is used.) and LPFG₂ (κ_1 and $\delta_0=150$ is used.). (g) LPFG₁ (κ_1 and δ_0 is used.) and LPFG₂ (κ_2 and δ_0 is used.). (h) LPFG₁ (κ_1 and δ_0 is used.) and LPFG₂ (κ_2 and $\delta_0=50$ is used.).

the tunable flat-top filter by considering a single cladding mode as shown in Fig. 4. We have utilized the coupling coefficients κ and the detuning factors δ , respectively, as shown in the Figs. 4(a) and 4(b). The coupling coefficients in Fig. 4(a) can be found by the Gel'fand-Levitan-Marchenko coupled equations [14, 15].

The κ_1 in Fig. 4(a) is used for the LPFG₁ and LPFG₂ of Figs. 4(c)-4(f). In Fig. 4(b), the δ_0 is used for the LPFG₁ and LPFG₂ is used the δ_0 for Fig. 4(c), the δ_0 -50 for Fig. 4(d), the δ_0 -100 for Fig. 4(e), and the δ_0 -150 for Fig. 4(f). We have assumed the delay is $d_1 = d_2 = 0$ and the core mode blocker is $CB = \text{diag}(0, 1, 1, \dots, 1)$, ideally. The 3dB bandwidth of the LPFG₁ and LPFG₂ is 13.9 nm. In the case, we can control the bandwidth of the flat-top bandpass filter in the wavelength range less than 13.9 nm $\pm \varepsilon$ (where ε is a small variation by calculating the equation (1)). The 3dB bandwidth of the obtained flat-top filter in Figs. 4(c) and 4(d) are 14.2 nm ($\varepsilon = 0.3$ nm) and 8 nm, respectively.

From the results, we can observe that the bandwidth of the bandpass is changed along the overlapped band region of the flat-top band-rejection of both LPFG₁ and LPFG₂. If we can make the ideal band-rejection filter, we can synthesize a sharp bandpass filter with the customized bandwidth as well as the bandpass filters with a tunable bandwidth. Especially, if the bandwidth of LPFG₁ and LPFG₂ aren't overlapped as shown in Fig. 4(f), we cannot synthesize the desired filter.

We also demonstrated the transmission along the overlapped region, when the coupling coefficients are different as shown in Figs. 4(g) and 4(h). The coupling coefficients κ_1 for the LPFG₁ and κ_2 the LPFG₂ are utilized as shown in Fig. 4(a). The detuning factors δ for Figs. 4(g) and 4(h) are the same ones as for Figs. 4(c) and 4(d). We have also used the delay and the core mode blocker as above examples.

From the Figs 4(g) and 4(h), we have obtained similar results to Figs. 4(c) and 4(d) although the LPFG₁ and LPFG₂ have different magnitude levels. The 3 dB bandwidth of the LPFG₂ is 14.1 nm. The 3 dB bandwidth of the obtained flat-top filter in Figs. 4(c) and 4(d) are 14.2 nm ($\varepsilon = 0.1$ nm) and 8 nm, respectively. The results show that the magnitude level of the flat-top bandpass filter is dependent on the magnitudes of the LPFG₁ and LPFG₂.

IV. CONCLUSION

The analytical model for the LPFG with the core mode blocker is proposed. We have also analyzed and described the physical phenomenon of the structure by using the proposed model. We have proposed a design method of the flat-top bandpass filter with customized bandwidth. We have also demonstrated, through computer simulations, the bandwidth of the flat-top filters controlled by tuning the overlapped band region.

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