# FRACTIONAL GAGLIARDO-NIRENBERG INEQUALITY 

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#### Abstract

A fractional Gagliardo-Nirenberg inequality is established. A sharp fractional Sobolev inequality is discussed as a direct consequence.


## 1. Introduction

Fractional differential equations have been of increasing importance for the past decades due to their diverse applications in science and engineering. Some of such areas are fluid flow, solute transport or dynamical processes in self-similar and porous structures, diffusive transport, material viscoelastic theory, electromagnetic theory, optics and signal processing, bio-sciences, economics, geology, and astrophysics. Fractional integral inequalities certainly provide fundamental tools for the study of fractional differential equations and optimizing problems including fractional Laplacian.

In this paper, we establish a fractional version of Gagliardo-Nirenberg inequality(Theorem 2.1). We make an effort to deliver a simple and direct proof in order to present a useful upper bound of the constant based on the effective upper bound of the constant in Hardy-LittlewoodSobolev inequality. In fact, when we set up the fractional GagliardoNirenberg inequalities, we found Wang's preprint [9] in which fractional Gagliardo-Nirenberg inequalities have been built via Littlewood-Paley decomposition. We present a totally different way of proof - we exhibit a self-contained simple and elegant proof, and provide a useful upper bound for the constant which was not accessible in Wang's proof. This upper bound plays an important role in actual applications [5]. A long

[^0]standing problem of the optimality of the constants in the (classical) Gagliardo-Nirenberg inequalities is still open. Using this line of the proof, however, we could establish a sharp form of a fractional Sobolev inequality (Corollary 2.2).

Among some equivalent definitions of the fractional Laplacian, we employ it as

$$
\sqrt{-\Delta}^{s} \phi:=\mathcal{F}^{-1}\left(|\cdot|^{s} \mathcal{F}(\phi)\right)
$$

where $\hat{u}=\mathcal{F}(u)$ represents the Fourier transform of $u$ on $\mathbb{R}^{n}$ defined by

$$
\hat{f}(\xi)=\mathcal{F}(f)(\xi)=\int_{\mathbb{R}^{n}} f(x) e^{-i x \cdot \xi} d x
$$

for $f \in L^{1}\left(\mathbb{R}^{n}\right) \cap L^{2}\left(\mathbb{R}^{n}\right)$.

## 2. Fractional Gagliardo-Nirenberg inequality

Gagliardo-Nirenberg inequality for fractional Laplacian is presented, and a sharp form of the fractional Sobolev inequality is obtained as a corollary.

Theorem 2.1. Let $m, q, \theta \in \mathbb{R} \backslash\{0\}$ with $q \neq m \theta>0,0<s<n$, $1<p<\frac{n}{s}$ and $1<\frac{r}{q-m \theta}$. Then the inequality

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}|u(x)|^{q} d x \leq C_{0}\left(\int_{\mathbb{R}^{n}}\left|\sqrt{-\Delta}^{s} u(x)\right|^{p} d x\right)^{\frac{m \theta}{p}}\left(\int_{\mathbb{R}^{n}}|u(x)|^{r} d x\right)^{\frac{q-m \theta}{r}} \tag{2.1}
\end{equation*}
$$

holds for the indices with the relation

$$
\begin{equation*}
m \theta\left(\frac{1}{p}-\frac{s}{n}\right)+\frac{q-m \theta}{r}=1 \tag{2.2}
\end{equation*}
$$

The sharp constant $C_{0}$ satisfies
$C_{0}^{\frac{1}{m \theta}} \leq \frac{2^{1-s}}{\pi^{\frac{s}{2}}} \frac{\left[\Gamma\left(\frac{n}{2}+1\right)\right]^{\frac{s}{n}} \Gamma\left(\frac{n-s}{2}\right)}{\Gamma\left(\frac{s}{2}\right) \Gamma\left(\frac{n}{2}\right)} \frac{1-\frac{1}{p}+\frac{s}{n}}{s p}\left(\left[\frac{1-\frac{s}{n}}{1-\frac{1}{p}}\right]^{1-\frac{s}{n}}+\left[\frac{1-\frac{s}{n}}{\frac{1}{p}-\frac{s}{n}}\right]^{1-\frac{s}{n}}\right)$.
In particular, when $m=q$, we have a fractional version of GagliardoNirenberg inequality:

$$
\begin{equation*}
\left(\int_{\mathbb{R}^{n}}|u(x)|^{q} d x\right)^{\frac{1}{q}} \leq C_{0}^{\frac{1}{q}}\left(\int_{\mathbb{R}^{n}}\left|\sqrt{-\Delta}^{s} u(x)\right|^{p} d x\right)^{\frac{\theta}{p}}\left(\int_{\mathbb{R}^{n}}|u(x)|^{r} d x\right)^{\frac{1-\theta}{r}} \tag{2.3}
\end{equation*}
$$

where the indices satisfy the relation

$$
\begin{equation*}
\theta\left(\frac{1}{p}-\frac{s}{n}\right)+\frac{1-\theta}{r}=\frac{1}{q} \tag{2.4}
\end{equation*}
$$

Proof. For convenience, we use the notation $\|u\|_{L^{t}}:=\left(\int_{\mathbb{R}^{n}}|u(x)|^{t} d x\right)^{\frac{1}{t}}$ for any $t \in \mathbb{R} \backslash\{0\}$. First we point out by the standard dilation argument that the index relation (2.2) is necessary. In fact, by replacing $u(\cdot)$ with $u(\delta \cdot)$, we can observe that

$$
\delta^{-n}\|u\|_{L^{q}}^{q} \leq C_{0} \delta^{\left(s-\frac{n}{p}\right) m \theta+\frac{n(m \theta-q)}{r}}\left\|\sqrt{-\Delta}^{s} u\right\|_{L^{p}}^{m \theta}\|u\|_{L^{r}}^{q-m \theta}
$$

for all $\delta>0$, which forces the relation (2.2).
Now, for any $u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ (the Schwartz class), we have

$$
\begin{align*}
\int_{\mathbb{R}^{n}}|u(x)|^{q} d x & =\int_{\mathbb{R}^{n}}|u(x)|^{m \theta}|u(x)|^{q-m \theta} d x \\
& \leq\left\||u|^{m \theta}\right\|_{L^{\bar{p}}}\left\||u|^{q-m \theta}\right\|_{L^{\bar{r}}}, \quad \frac{1}{\bar{p}}+\frac{1}{\bar{r}}=1  \tag{2.5}\\
& =\|u\|_{L^{m \theta \bar{p}}}^{m \theta}\|u\|_{L^{(q-m \theta) \bar{r}}}^{q-m \theta} .
\end{align*}
$$

We set $m \theta \bar{p}:=p_{0}$ and $(q-m \theta) \bar{r}:=r$ to have

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}|u(x)|^{q} d x \leq\|u\|_{L^{p_{0}}}^{m \theta}\|u\|_{L^{r}}^{q-m \theta} \tag{2.6}
\end{equation*}
$$

and $\frac{m \theta}{p_{0}}+\frac{q-m \theta}{r}=1$. Let $\sqrt{-\Delta}^{s} u:=f$, and we have

$$
\begin{aligned}
u(x) & =\frac{1}{2^{s} \pi^{n / 2}} \frac{\Gamma\left(\frac{n-s}{2}\right)}{\Gamma\left(\frac{s}{2}\right)} \int_{\mathbb{R}^{n}} \frac{f(y)}{|x-y|^{n-s}} d y \\
& =\frac{1}{2^{s} \pi^{n / 2}} \frac{\Gamma\left(\frac{n-s}{2}\right)}{\Gamma\left(\frac{s}{2}\right)} \int_{\mathbb{R}^{n}} \frac{\sqrt{-\Delta}^{s} u(y)}{|x-y|^{n-s}} d y
\end{aligned}
$$

Indeed, we may take the Fourier transform on $\sqrt{-\Delta}^{s} u=f$, and take it back to have $u$ after solving for $\widehat{u}$. Therefore Hardy-Littlewood-Sobolev inequality yields

$$
\begin{equation*}
\|u\|_{L^{p_{0}}} \leq \frac{1}{2^{s} \pi^{n / 2}} \frac{\Gamma\left(\frac{n-s}{2}\right)}{\Gamma\left(\frac{s}{2}\right)} C_{1}\left\|\sqrt{-\Delta}^{s} u\right\|_{L^{p}} \tag{2.7}
\end{equation*}
$$

where $C_{1}$ is a positive constant(see the comments after the proof) and $p$ satisfies

$$
\begin{equation*}
\frac{1}{p}+\frac{n-s}{n}=1+\frac{1}{p_{0}} \tag{2.8}
\end{equation*}
$$

This index relation combined with the index relation appeared at (2.6) implies (2.2), and (2.6) together with (2.7) implies (2.1).

The extremals and the best constant $C_{1}$ of Hardy-Littlewood-Sobolev inequality for some special cases are known (see [7] or Section 4.3 in [8]). Thanks to those cases, we have a sharp form of a fractional Sobolev inequality:

Corollary 2.2 (Fractional Sobolev inequality). For $0<s<n$, $1<p<\frac{n}{s}$ and $q=\frac{n p}{n-p s}$, we have

$$
\|u\|_{L^{q}} \leq C_{0}^{1 / q}\left\|\sqrt{-\Delta}^{s} u\right\|_{L^{p}}
$$

The sharp constant for the inequality for the case $p=\frac{2 n}{n+s}$ and $q=\frac{2 n}{n-s}$ is

$$
\left[\frac{\pi^{\frac{n}{2}-s}}{2^{s}} \frac{\Gamma\left(\frac{n}{q}\right)}{\Gamma\left(\frac{n}{p}\right)}\left\{\frac{\Gamma(n)}{\Gamma\left(\frac{n}{2}\right)}\right\}^{s / n}\right]^{1 / q}
$$

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