

TRAVELING WAVE SOLUTIONS FOR HIGHER DIMENSIONAL NONLINEAR EVOLUTION EQUATIONS USING THE $(\frac{G}{G})$ - EXPANSION METHOD

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ABSTRACT. In the present paper, we construct the traveling wave solutions involving parameters of nonlinear evolution equations in the mathematical physics via the (3+1)- dimensional potential- YTSE equation, the (3+1)- dimensional generalized shallow water equation, the (3+1)- dimensional Kadomtsev- Petviashvili equation, the (3+1)- dimensional modified KdV- Zakharov- Kuznetsev equation and the (3+1)- dimensional Jimbo-Miwa equation by using a simple method which is called the $(\frac{G}{G})$ - expansion method, where $G = G(\xi)$ satisfies a second order linear ordinary differential equation. When the parameters are taken special values, the solitary waves are derived from the travelling waves. The travelling wave solutions are expressed by hyperbolic, trigonometric and rational functions.

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1. Introduction

In recent years, the exact solutions of nonlinear PDEs have been investigated by many authors(see for example [1-47]) who are interested in nonlinear physical phenomena. Many powerful methods have been presented by those authors such as the inverse scattering transform [3], the Backlund transform [14,15], the generalized Riccati equation [17,28], the Jacobi elliptic function expansion [7,12,25,27,29,33,36], the extended tanh- function method [1,8,34,35,44], the F-expansion method [2,18-20], the exp-function expansion method [6,9,30,42,43],

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the sub- ODE method [13,21], the extended sinh-cosh and sine-cosine methods [22], the complex hyperbolic function method [37], the truncated Painleve expansion [40] and so on.

In the present paper, we shall use a simple method which is called the $(\frac{G}{G})$ -expansion method [5,24,38,39,46,47]. This method is firstly proposed by the Chinese Mathematicians Wang et al [24] for which the traveling wave solutions of nonlinear equations are obtained. The main idea of this method is that the traveling wave solutions of nonlinear equations can be expressed by a polynomial in $(\frac{G}{G})$, where $G = G(\xi)$ satisfies the second order linear ordinary differential equation $G''(\xi) + \lambda G'(\xi) + \mu G(\xi) = 0$, where $\xi = x + y + z - Vt$ while λ, μ and V are constants. The degree of this polynomial can be determined by considering the homogeneous balance between the highest order derivatives and the nonlinear terms appearing in the given nonlinear equations. The coefficients of this polynomial can be obtained by solving a set of algebraic equations resulted from the process of using the proposed method. This new method will play an important role in expressing the traveling wave solutions for nonlinear evolution equations via the (3+1)- dimensional potential- YTSF equation, the (3+1)- dimensional generalized shallow water equation, the (3+1)- dimensional Kadomtsev- Petviashvili equation, the (3+1)- dimensional modified KdV- Zakharov- Kuznetsev equation and the (3+1)- dimensional Jimbo-Miwa equation in terms of hyperbolic, trigonometric and rational functions.

2. Description of the $(\frac{G}{G})$ -expansion method

Suppose we have the following nonlinear partial differential equation:

$$P(u, u_t, u_x, u_y, u_z, u_{tt}, u_{xt}, u_{xx}, u_{xy}, u_{yy}, u_{yt}, u_{zz}, u_{zt}, u_{zx}, u_{zy}, \dots) = 0, \quad (2.1)$$

where $u = u(x, y, z, t)$ is an unknown function, P is a polynomial in $u(x, y, z, t)$ and its partial derivatives in which the highest order derivatives and the nonlinear terms are involved. In the following we give the main steps [24] of the $(\frac{G}{G})$ -expansion method:

Step 1. The traveling wave variable

$$u(x, y, z, t) = u(\xi), \quad \xi = x + y + z - Vt, \quad (2.2)$$

where V is a constant, permits us reducing Eq. (2.1) to an ODE for $u = u(\xi)$ in the form

$$P(u, u', u'', u''', \dots) = 0. \quad (2.3)$$

Step 2. Suppose the solution of Eq(2.3) can be expressed by a polynomial in $(\frac{G}{G})$ as follows:

$$u(\xi) = \sum_{i=0}^n \alpha_i \left(\frac{G'}{G}\right)^i, \quad (2.4)$$

where $G = G(\xi)$ satisfies the following second order linear ordinary differential equation:

$$G'' + \lambda G' + \mu G = 0, \tag{2.5}$$

where α_i, V, λ and μ are constants to be determined later provided $\alpha_n \neq 0$. The positive integer "n" can be determined by considering the homogeneous balance between the highest order derivatives and the nonlinear terms appearing in Eq (2.3).

Step 3. Substituting (2.4) into (2.3) and using Eq (2.5), collecting all terms with the same power of $(\frac{G}{G})$ together and then equating each coefficient of the resulted polynomial to zero, yields a set of algebraic equations for α_i, V, λ and μ .

Step 4. Since the general solution of Eq (2.5) has been well known for us, then substituting α_i, V and the general solution of Eq (2.5) into (2.4) we have more traveling wave solutions of the nonlinear partial differential equation (2.1).

3. Some applications

In this section, we apply the $(\frac{G}{G})$ - expansion method to construct the traveling wave solutions for the (3+1)- dimensional potential- YTSF equation, the (3+1)- dimensional generalized shallow water equation, the (3+1)- dimensional Kadomtsev- Petviashvili equation, the (3+1)- dimensional modified KdV- Zakharov- Kuznetsev equation and the (3+1)- dimensional Jimbo-Miwa equation which are very important nonlinear evolution equations in the mathematical physics and have been paid attention by many researchers.

3.1. Example 1. The (3+1)- dimensional potential- YTSF equation.

We start with the (3+1)- dimensional potential- YTSF equation [4,32] in the form

$$-4u_{xt} + u_{xxxx} + 4u_x u_{xz} + 2u_{xx} u_z + 3u_{yy} = 0. \tag{3.1}$$

The (3+1)- dimensional evolution equation (3.1) was recently derived by Yu et al [32]. This equation was called the potential- YTSF equation and it was developed by using the strong symmetry. Let us now solve the system (3.1) by $(\frac{G}{G})$ - expansion method. To this end, we see that the following traveling wave variable

$$u(x, y, z, t) = u(\xi), \quad \xi = x + y + z - Vt, \tag{3.2}$$

where V is a constant, permits us converting (3.1) into the following ODE:

$$C - 4Vu' + u''' + 3u'^2 + 3u' = 0, \tag{3.3}$$

where C is an integration constant. Suppose that the solutions of ODE (3.3) can be expressed by a polynomial in $(\frac{G}{G})$ as follows:

$$u(\xi) = \sum_{i=0}^n \alpha_i \left(\frac{G}{G}\right)^i, \tag{3.4}$$

where α_i are arbitrary constants, while $G(\xi)$ satisfies the second order linear ODE (2.5).

Considering the homogeneous balance between the highest order derivative and the nonlinear term in (3.3), we deduce that $n = 1$. Thus, we get

$$u(\xi) = \alpha_1 \left(\frac{G'}{G} \right) + \alpha_0, \quad (3.5)$$

where α_0 and α_1 are constants provided $\alpha_1 \neq 0$. Substituting (3.5) into (3.3) along with (2.5), collecting all terms with the same powers of $\left(\frac{G'}{G} \right)$ and setting them to zero. Consequently, we have the following system of algebraic equations:

$$\begin{aligned} \alpha_1 \lambda [4V - (\lambda^2 + 8\mu) + 6\alpha_1 \mu - 3] &= 0, \\ \alpha_1 [4V - (7\lambda^2 + 8\mu) + 3\alpha_1 \lambda^2 + 6\alpha_1 \mu - 3] &= 0, \\ -12\alpha_1 \lambda + 6\alpha_1^2 \lambda &= 0, \\ -6\alpha_1 + 3\alpha_1^2 &= 0, \\ C + 4\alpha_1 \mu V - \mu \alpha_1 (\lambda^2 + 2\mu) + 3\alpha_1^2 \mu^2 - 3\alpha_1 \mu &= 0, \end{aligned}$$

which can be solved, to get

$$\alpha_1 = 2, \quad V = \frac{1}{4}(\lambda^2 - 4\mu + 3), \quad C = 0. \quad (3.6)$$

Substituting (3.6) into (3.5) yields

$$u(\xi) = 2 \left(\frac{G'}{G} \right) + \alpha_0, \quad (3.7)$$

where

$$\xi = x + y + z - \frac{t}{4}(\lambda^2 - 4\mu + 3). \quad (3.8)$$

Solving Eq. (2.5), we deduce for $\lambda^2 - 4\mu > 0$ that

$$\frac{G'}{G} = \frac{1}{2} \sqrt{\lambda^2 - 4\mu} \left(\frac{A \cosh(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi) + B \sinh(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi)}{A \sinh(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi) + B \cosh(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi)} \right) - \frac{\lambda}{2}, \quad (3.9)$$

where A and B are arbitrary constants.

Substituting (3.9) into (3.7) we deduce the following three types of traveling wave solutions:

Case 1. If $\lambda^2 - 4\mu > 0$, then we have

$$u(\xi) = \sqrt{\lambda^2 - 4\mu} \left(\frac{A \cosh(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi) + B \sinh(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi)}{A \sinh(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi) + B \cosh(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi)} \right) + \alpha_0 - \lambda, \quad (3.10)$$

Case 2 . If $\lambda^2 - 4\mu < 0$, then we have

$$u(\xi) = \sqrt{4\mu - \lambda^2} \left(\frac{-A \sin(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi) + B \cos(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi)}{A \cos(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi) + B \sin(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi)} \right) + \alpha_0 - \lambda, \tag{3.11}$$

Case 3. If $\lambda^2 - 4\mu = 0$, then we have

$$u(\xi) = 2 \left(\frac{B}{A + B \xi} \right) + \alpha_0 - \lambda. \tag{3.12}$$

In particular, if $A = 0, B \neq 0, \lambda > 0, \mu = 0$, then we deduce from (3.10) that

$$u(\xi) = \lambda \tanh\left(\frac{\lambda}{2}\xi\right) + \alpha_0 - \lambda, \tag{3.13}$$

which represents the solitary wave solutions of the (3+1)- dimensional potential-YTSF equation (3.1) .

3.2. Example 2. The (3+1)- dimensional generalized shallow water equation. In this subsection, we study the following (3+1)- dimensional generalized shallow water equation [16]:

$$u_{xxxx} - 3u_{xx}u_y - 3u_xu_{xy} + u_{yt} - u_{xz} = 0. \tag{3.14}$$

Let us now solve Eq. (3.14) by the proposed method. To this end, we see that the traveling wave variable (3.2) permits us converting (3.14) into the following ODE:

$$C + u''' - 3u'^2 - (V + 1)u' = 0, \tag{3.15}$$

where C is a constant of integration. Considering the homogeneous balance between highest order derivative and nonlinear term in (3.15), we get

$$u(\xi) = \alpha_1 \left(\frac{G'}{G} \right) + \alpha_0. \tag{3.16}$$

Substituting (3.16) into (3.15) along with (2.5), collecting all terms with the same powers of $\left(\frac{G'}{G}\right)$ and setting them to zero. Consequently, we have the following system of algebraic equations:

$$\begin{aligned} \alpha_1 \lambda [V - (\lambda^2 + 8\mu) - 6\alpha_1 \mu + 1] &= 0, \\ \alpha_1 [V - (7\lambda^2 + 8\mu) - 3\alpha_1 (\lambda^2 + 2\mu) + 1] &= 0, \\ 12\alpha_1 \lambda + 6\alpha_1^2 \lambda &= 0, \\ 6\alpha_1 + 3\alpha_1^2 &= 0, \\ C + \alpha_1 \mu V - \mu \alpha_1 (\lambda^2 + 2\mu) - 3\alpha_1^2 \mu^2 + \alpha_1 \mu &= 0, \end{aligned}$$

which can be solved, to get

$$\alpha_1 = -2, \quad V = \lambda^2 - 4\mu - 1, \quad C = 0, \tag{3.17}$$

Substituting (3.17) into (3.16) yields

$$u(\xi) = -2 \left(\frac{G}{G} \right) + \alpha_0, \quad (3.18)$$

where $\xi = x + y + z - t(\lambda^2 - 4\mu - 1)$. On substituting (3.9) into (3.18), we deduce the following three types of traveling wave solutions:

Case 1 . If $\lambda^2 - 4\mu > 0$, then we have

$$u(\xi) = -\sqrt{\lambda^2 - 4\mu} \left(\frac{A \cosh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi) + B \sinh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi)}{A \sinh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi) + B \cosh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi)} \right) + \alpha_0 + \lambda, \quad (3.19)$$

Case 2 . If $\lambda^2 - 4\mu < 0$, then we have

$$u(\xi) = -\sqrt{(4\mu - \lambda^2)} \left(\frac{-A \sin(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi) + B \cos(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi)}{A \cos(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi) + B \sin(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi)} \right) + \alpha_0 + \lambda, \quad (3.20)$$

Case 3 . If $\lambda^2 - 4\mu = 0$, then we have

$$u(\xi) = -2 \left(\frac{B}{A + B \xi} \right) + \alpha_0 + \lambda, \quad (3.21)$$

In particular, if $A = 0, B \neq 0, \lambda > 0, \mu = 0$, then we deduce from (3.19) that

$$u(\xi) = -\lambda \tanh\left(\frac{\lambda}{2}\xi\right) + \alpha_0 + \lambda, \quad (3.22)$$

which represents the solitary wave solutions of the (3+1)- dimensional generalized shallow water equation (3.14).

3.3. Example 3. The (3+1)- dimensional Kadomtsev- Petviashvili equation. In this subsection, we consider the (3+1)- dimensional Kadomtsev- Petviashvili equation [11,42,45] in the form:

$$u_{xt} + 6(u_x)^2 + 6uu_{xx} - u_{xxxx} - u_{yy} - u_{zz} = 0, \quad (3.23)$$

which describes the dynamics of solitons and nonlinear wave in plasmas and superfluids. Let us now solve this equation by the proposed method. To this end, we see that the traveling wave variable (3.2) permits us converting (3.23) into the following ODE :

$$-Vu + 3u^2 - u'' - 2u = 0, \quad (3.24)$$

where the constants of integrations are assumed to be zero. Considering the homogeneous balance between highest order derivative and nonlinear term in (3.24), we get

$$u(\xi) = \alpha_2 \left(\frac{G}{G} \right)^2 + \alpha_1 \left(\frac{G}{G} \right) + \alpha_0, \quad (3.25)$$

where α_i ($i = 0, 1, 2$) are constants provided $\alpha_2 \neq 0$. Substituting (3.25) into (3.24) along with (2.5), collecting all terms with the same powers of $\left(\frac{G'}{G}\right)$ and setting them to zero. Consequently, we have the following system of algebraic equations:

$$\begin{aligned} -\alpha_1 V + 6\alpha_0 \alpha_1 - (6\alpha_2 \lambda \mu + 2\alpha_1 \mu + \alpha_1 \lambda^2) - 2\alpha_1 &= 0, \\ -\alpha_2 V + 3(\alpha_1^2 + 2\alpha_0 \alpha_2) - (8\alpha_2 \mu + 3\alpha_1 \lambda + 4\alpha_2 \lambda^2) - 2\alpha_2 &= 0, \\ 6\alpha_1 \alpha_2 - (2\alpha_1 + 10\alpha_2 \lambda) &= 0, \\ -6\alpha_2 + 3\alpha_2^2 &= 0, \\ -\alpha_0 V - 2\alpha_2 \mu^2 + 3\alpha_0^2 - \alpha_1 \lambda \mu - 2\alpha_0 &= 0, \end{aligned}$$

which can be solved, to get

$$\alpha_2 = 2, \quad \alpha_1 = 2\lambda, \quad V = 6\alpha_0 - \lambda^2 - 8\mu - 2, \tag{3.26}$$

Substituting (3.26) into (3.25) yields

$$u(\xi) = 2 \left(\frac{G'}{G}\right)^2 + 2\lambda \left(\frac{G'}{G}\right) + \alpha_0, \tag{3.27}$$

where $\xi = x + y + z - t(6\alpha_0 - \lambda^2 - 8\mu - 2)$. On substituting (3.9) into (3.27), we deduce the following three types of traveling wave solutions:

Case 1 . If $\lambda^2 - 4\mu > 0$, then we have $u(\xi)$

$$= \frac{1}{2}(\lambda^2 - 4\mu) \left(\frac{A \cosh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi) + B \sinh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi)}{A \sinh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi) + B \cosh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi)} \right)^2 + \alpha_0 - \frac{\lambda^2}{2}, \tag{3.28}$$

Case 2 . If $\lambda^2 - 4\mu < 0$, then we have $u(\xi)$

$$= \frac{1}{2}(4\mu - \lambda^2) \left(\frac{-A \sin(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi) + B \cos(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi)}{A \cos(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi) + B \sin(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi)} \right)^2 + \alpha_0 - \frac{\lambda^2}{2}, \tag{3.29}$$

Case 3 . If $\lambda^2 - 4\mu = 0$, then we have

$$u(\xi) = 2 \left(\frac{B}{A + B \xi} \right)^2 + \alpha_0 - \frac{\lambda^2}{2}, \tag{3.30}$$

In particular, if $A = 0, B \neq 0, \lambda > 0, \mu = 0$, then we deduce from (3.28) that

$$u(\xi) = -\frac{\lambda^2}{2} \operatorname{sech}^2\left(\frac{\lambda}{2}\xi\right) + \alpha_0, \tag{3.31}$$

which represents the solitary wave solutions of the (3+1)- dimensional Kadomtsev-Petviashvili equation(3.23).

3.4. Example 4. The (3+1)- dimensional modified KdV- Zakharov-Kuznetsev equation. In this subsection, we consider the (3+1)- dimensional modified KdV- Zakharov- Kuznetsev equation [27] in the form:

$$u_t + \alpha u^2 u_x + u_{xxx} + u_{xyy} + u_{xzz} = 0, \quad (3.32)$$

where α is a nonzero constant. Xu [27] has discussed Eq. (3.32) using an elliptic equation method and found new types of elliptic function solutions. Let us now solve this equation by the proposed method. To this end, we see that the traveling wave variable (3.2) permits us converting (3.32) into the following ODE

$$C - Vu + \frac{1}{3}\alpha u^3 + 3u'' = 0, \quad (3.33)$$

where C is a constant of integration. Considering the homogeneous balance between highest order derivative and nonlinear term in (3.33), we get

$$u(\xi) = \alpha_1 \left(\frac{G}{G'} \right) + \alpha_0. \quad (3.34)$$

Substituting (3.34) into (3.33) along with (2.5), collecting all terms with the same powers of $\left(\frac{G}{G'} \right)$ and setting them to zero. Consequently, we have the following system of algebraic equations:

$$\begin{aligned} -\alpha_1 V + \alpha \alpha_1 \alpha_0^2 + 3\alpha_1 (\lambda^2 + 2\mu) &= 0, \\ \alpha \alpha_0 \alpha_1^2 + 9\alpha_1 \lambda &= 0, \\ \frac{1}{3} \alpha \alpha_1^3 + 6\alpha_1 &= 0, \\ C - \alpha_0 V + \frac{1}{3} \alpha \alpha_0^3 + 3\mu \lambda \alpha_1 &= 0, \end{aligned}$$

which can be solved, to get

$$\alpha_1 = \pm \frac{6i}{\sqrt{2\alpha}}, \quad \alpha_0 = \pm \frac{3i\lambda}{\sqrt{2\alpha}}, \quad V = -\frac{3}{2}\lambda^2 + 6\mu, \quad C = 0, \quad (3.35)$$

Substituting (3.35) into (3.34) yields

$$u(\xi) = \pm \frac{6i}{\sqrt{2\alpha}} \left(\frac{G}{G'} \right) \pm \frac{3i\lambda}{\sqrt{2\alpha}}, \quad (3.36)$$

where $\xi = x + y + z + \frac{3}{2}t(\lambda^2 - 4\mu)$. On substituting (3.9) into (3.36), we deduce the following three types of traveling wave solutions:

Case 1 . If $\lambda^2 - 4\mu > 0$, then we have

$$u(\xi) = \pm 3i \sqrt{\frac{(\lambda^2 - 4\mu)}{2\alpha}} \left(\frac{A \cosh\left(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi\right) + B \sinh\left(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi\right)}{A \sinh\left(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi\right) + B \cosh\left(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi\right)} \right), \quad (3.37)$$

Case 2 . If $\lambda^2 - 4\mu < 0$, then we have

$$\begin{aligned}
 &u(\xi) \\
 &= \pm 3i \sqrt{\frac{(4\mu - \lambda^2)}{2\alpha}} \left(\frac{-A \sin(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi) + B \cos(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi)}{A \cos(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi) + B \sin(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi)} \right),
 \end{aligned} \tag{3.38}$$

Case 3 . If $\lambda^2 - 4\mu = 0$, then we have

$$u(\xi) = \pm \frac{6i}{\sqrt{2\alpha}} \left(\frac{B}{A + B \xi} \right), \tag{3.39}$$

In particular, if $A = 0, B \neq 0, \lambda > 0, \mu = 0$, then we deduce from (3.37) that

$$u(\xi) = \pm \frac{3 i \lambda}{\sqrt{2\alpha}} \tanh\left(\frac{\lambda}{2}\xi\right), \tag{3.40}$$

which represents the solitary wave solutions of the (3+1)- dimensional modified KdV- Zakharov- Kuznetsev equation(3.32).

3.5. Example 5. The (3+1)- dimensional Jimbo-Miwa equation. In this subsection, we consider the(3+1)- dimensional Jimbo-Miwa equation [10,23] in the form:

$$u_{xxxxy} + 3u_y u_{xx} + 3u_x u_{xy} + 2u_{yt} - 3u_{xz} = 0. \tag{3.41}$$

This equation is firstly investigated by Jimbo-Miwa and its certain soliton solutions are obtained in [10] . Then its studied on many manifolds by several authors regarding its solutions, symmetries and integrability properties. Wazwaz [25,26] successfully studied one- soliton solution to equation (3.41) by means of the tanh-coth method. He also employed the Hirota’s bilinear method to this equation and confirmed that it is completely integrable and it admits multiple - soliton solutions of any order. Let us now solve this equation by the proposed method. To this end, we see that the traveling wave variable (3.2) permits us converting (3.41) into the following ODE :

$$C + u''' + 3u'^2 - (2V + 3)u' = 0, \tag{3.42}$$

where C is a constant of integration. Considering the homogeneous balance between highest order derivative and nonlinear term in (3.42), we get

$$u(\xi) = \alpha_1 \left(\frac{G}{G} \right) + \alpha_0. \tag{3.43}$$

Substituting (3.43) into (3.42) along with (2.5), collecting all terms with the same powers of $\left(\frac{G}{G} \right)$ and setting them to zero. Consequently, we have the

following system of algebraic equations:

$$\begin{aligned} -\alpha_1\lambda(\lambda^2 + 8\mu) + 6\alpha_1^2\lambda\mu + 2V\lambda\alpha_1 + 3\lambda\alpha_1 &= 0, \\ -\alpha_1(7\lambda^2 + 8\mu) + 3\alpha_1^2(\lambda^2 + 2\mu) + 2V\alpha_1 + 3\alpha_1 &= 0, \\ -12\alpha_1\lambda + 6\alpha_1^2\lambda &= 0, \\ -6\alpha_1 + 3\alpha_1^2 &= 0, \\ C + 2\alpha_1\mu V - \mu\alpha_1(\lambda^2 + 2\mu) + 3\alpha_1^2\mu^2 + 3\alpha_1\mu &= 0, \end{aligned}$$

which can be solved, to get

$$\alpha_1 = 2, \quad V = \frac{1}{2}(\lambda^2 - 4\mu - 3), \quad C = 0, \quad (3.44)$$

Substituting (3.44) into (3.43) yields

$$u(\xi) = 2 \left(\frac{G'}{G} \right) + \alpha_0, \quad (3.45)$$

where $\xi = x + y + z - \frac{1}{2}(\lambda^2 - 4\mu - 3)t$. On substituting (3.9) into (3.45), we deduce the following three types of traveling wave solutions:

Case 1 . If $\lambda^2 - 4\mu > 0$, then we have
 $u(\xi)$

$$= \sqrt{(\lambda^2 - 4\mu)} \left(\frac{A \cosh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi) + B \sinh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi)}{A \sinh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi) + B \cosh(\frac{1}{2}\sqrt{\lambda^2 - 4\mu} \xi)} \right) + \alpha_0 - \lambda, \quad (3.46)$$

Case 2 . If $\lambda^2 - 4\mu < 0$, then we have
 $u(\xi)$

$$= \sqrt{(4\mu - \lambda^2)} \left(\frac{-A \sin(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi) + B \cos(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi)}{A \cos(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi) + B \sin(\frac{1}{2}\sqrt{4\mu - \lambda^2} \xi)} \right) + \alpha_0 - \lambda, \quad (3.47)$$

Case 3 . If $\lambda^2 - 4\mu = 0$, then we have

$$u(\xi) = 2 \left(\frac{B}{A + B \xi} \right) + \alpha_0 - \lambda, \quad (3.48)$$

In particular, if $A = 0, B \neq 0, \lambda > 0, \mu = 0$, then we deduce from (3.46) that

$$u(\xi) = \lambda \tanh\left(\frac{\lambda}{2}\xi\right) + \alpha_0 - \lambda, \quad (3.49)$$

which represents the solitary wave solutions of the the(3+1)- dimensional Jimbo-Miwa equation (3.41).

4. Summary and conclusions

In this paper, we have seen that three types of traveling wave solutions in terms of hyperbolic, trigonometric and rational functions for the (3+1)- dimensional potential- YTSF equation, the (3+1)- dimensional generalized shallow water equation, the (3+1)- dimensional Kadomtsev- Petviashvili equation, the (3+1)- dimensional modified KdV- Zakharov- Kuznetsev equation and the (3+1)- dimensional Jimbo-Miwa equation are successfully found out by using the $(\frac{G}{G})$ - expansion method. This method is more powerful, effective and convenient. The performance of this method is reliable, simple and gives many new solutions. The $(\frac{G}{G})$ - expansion method has more advantages: It is direct and concise. It is also a standard and computerizable method which allows us to solve complicated nonlinear evolution equations in the mathematical physics. We have noted that this method changes the given problems into simple problems which can be solved easily. Also, the models proposed in this article described important applications in physics and engineering.

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