

The Pulsed I_d - V_g methodology and Its Application to the Electron Trapping Characterization of High- κ gate Dielectrics

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Abstract—Pulsed current-voltage (I-V) methods are introduced to evaluate the impact of fast transient charge trapping on the performance of high- κ dielectric transistors. Several pulsed I-V measurement configurations and measurement requirements are critically reviewed. Properly configured pulsed I-V measurements are shown to be capable of extracting such device characteristics as trap-free mobility, trap-induced threshold voltage shift (ΔV_t), as well as effective fast transient trap density. The results demonstrate that the pulsed I-V measurements are an essential technique for evaluating high- κ gate dielectric devices.

Index Terms—High- κ , pulse I-V, threshold voltage instability, fast transient charge trapping, mobility, bias temperature instability, hafnium, MOSFET

I. INTRODUCTION

Since silicon dioxide gate dielectric thickness continues to be scaled, gate leakage current may exceed the required specification limits for future technology nodes. To circumvent this issue, higher permittivity (κ) gate dielectrics, which allow thicker dielectric layers at the same equivalent oxide thickness (EOT), have been introduced [1]. This new class of gate dielectrics requires

new measurement methods and analysis. In particular, hafnium (Hf)-based oxides, which are being widely accepted as a leading candidate for high- κ gate dielectrics [1], exhibits high density of as-grown structural defects [2-16], some representing electron traps that complicate the evaluation of “intrinsic” performance in high- κ gate stacks.

The electron charge trapping process potentially causes threshold voltage instability, mobility degradation, and poor reliability and performance, which are some of the primary issues affecting the use of Hf-based materials in production [17-65]. While significant progress has been made in fabricating high- κ gate stacks with less charge trapping [66-69], pulse-based characterization techniques [30, 70-73] remain critically important in understanding charge trapping processes for low power applications, memory cells, etc. This review addresses pulsed current-voltage (I-V) characterization and analysis approaches.

II. FAST TRANSIENT CHARGE TRAPPING

Hf-based high- κ gate dielectrics exhibit a phenomenon known as the fast transient charging effect (FTCE) [37]. This effect is caused by substrate electrons injected to and captured by shallow traps (just below the Hf-based dielectric conduction band, Fig. 1) [74], where the trapping may occur within hundreds of nanoseconds to microseconds [50, 73]. The charges trapped during this process are well-localized within the dielectric, with their spatial distribution through the dielectric thickness being determined by the Fermi energy of the channel electrons and the trap energy (presumably, a neutral or single-charged negative oxygen vacancies) [4, 7, 9-11, 16]. The

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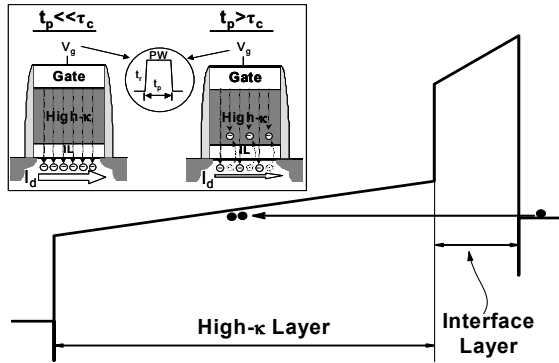


Fig. 1. Band diagram schematic for the electron trapping process during pulsed I-V measurements. Inset: Two scenarios are shown: when the pulse time t_p (sum of t_r and pulse width) is much shorter and longer than the trapping time τ_c . The latter defines the onset of electron trapping into the high- κ bulk.

fast trapping, which is reversible with no residual damage [30, 75, 76], can contribute significantly to threshold voltage (V_t) instability and degradation of device performance in Hf-based gate stack nMOS transistors.

pMOS Hf-based transistors, on the other hand, typically exhibit negligible fast transient charging (FTC) [33, 77] in inversion due to a lack of defect states with the appropriate energies. An example is shown in Fig. 2 for an inversion bias of $V_g = -2$ V with a standard mid-gap work function metal electrode illustrating that the FTC sites are far away from any resonant injection with the gate or substrate injection points. If significant FTC is detected, the test set-up or measurement may have issues [73] (should not be confused with NBTI measurements, which may have a fast component).

Because conventional DC characterization can be

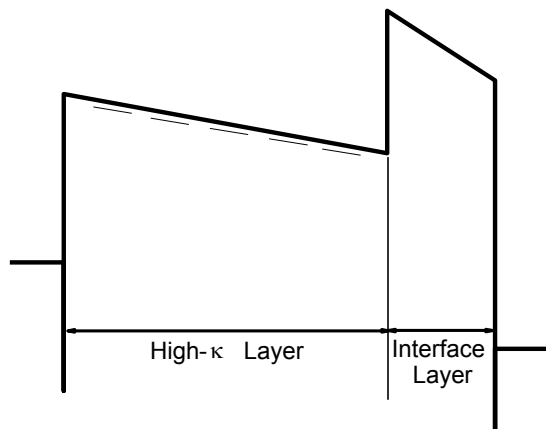


Fig. 2. Band diagram schematic for a pMOS gate stack illustrating injected electrons from the gate are not in resonance with the FTC trap sites, thereby demonstrating that negligible FTC occurs in pMOS devices.

adversely affected by FTC while the DC measurement is being conducted, thus lowering the drive current (Fig. 3(a)), the pulsed I-V methodology was introduced to quantify the transient charge trapping as well as to measure “intrinsic” (i.e., trapping-free) device characteristics [30, 70-72].

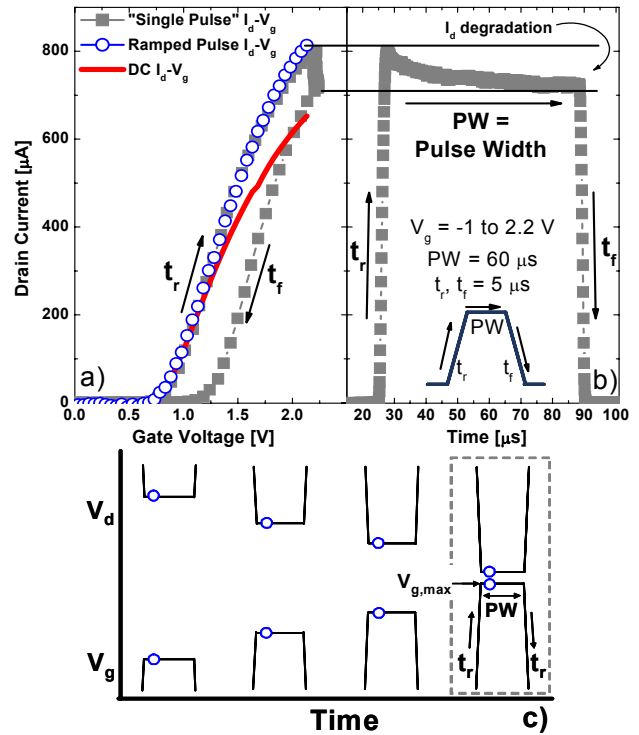


Fig. 3. Various forms of pulsed I_d data in the (a) pulse voltage, or (b) time domains based on the “ramped” pulse or “single” pulse (dashed box) illustration in (c). Results clearly show that I_d degrades at the top of the pulse (i.e., pulse width portion) due to FTC while the ramped pulse I_d and the t_r portion coincide to demonstrate trapping-free characteristics. DC I_d - V_g is included for comparison.

III. MEASUREMENT TECHNIQUE AND INSTRUMENTATION CONFIGURATIONS

1. Measurement Technique

Pulsed drain current-gate voltage (I_d - V_g) measurements are usually performed by applying a trapezoidal (e.g., rise time, pulse width, and fall time) or triangular (e.g., rise and fall time only) V_g pulse on the gate of a transistor configured as an inverter and then measuring I_d . While a voltage pulse is applied to the gate of the FET, the V_g and drain voltage (V_d) are simultaneously recorded on a digital storage oscilloscope. The drain

voltage response is converted to I_d and plotted against the pulse bias, V_g , or pulse time. From the V_d response, I_d (in the linear regime) can be calculated as [78]

$$I_d = \frac{V_{dd}}{V_d} \left(\frac{V_{dd} - V_d}{R_L} \right) \quad (1)$$

where V_{dd} is the drain voltage, V_d is the drain response voltage, and R_L is the load resistance.

Figs. 3(a) and 3(b) illustrate I_d - V_g (pulse voltage domain) and I_d -time (pulse time domain) characteristics, respectively, wherein the change of I_d during a “single” pulse measurement can be determined. The dashed box in Fig. 3(c) illustrates the “single” pulse wherein the V_g pulse with its corresponding V_d response is captured and converted into I_d - V_g or I_d -time. From this, the hysteresis in the I_d - V_g curve (Fig. 3(a)) or the decrease of I_d in the I_d -time plot (Fig. 3(b)) illustrates the threshold voltage change due to charge trapping.

Another approach to the pulsed I-V technique is the “ramped pulse” I-V (Fig. 3(a)), which generates a complete I_d - V_g curve within the total “single” pulse V_g range by applying multiple pulses with incrementally rising amplitude V_g up to $V_g = V_{g,max}$ and extracting I_d at each V_g value (Fig. 3(c)) [33, 72]. The advantage of this approach is that each pulse can have a very short characteristic time (around 37 ns = rise time + pulse width with ~0.1% duty cycle in this work), allowing “intrinsic” I_d values to be obtained. It also provides better resolution by averaging the response pulse V_d values during a portion of the pulse width time, thereby reducing the effect of fluctuations. Comparing this approach to the “single” pulse method (see Fig. 3(a)) revealed that a 5 μ s rise time in the single pulse measurement is short enough to obtain the trapping-free I_d characteristic in this particular gate stack due to the 1 nm SiO₂ interfacial layer since both techniques coincide on the “trapping-free” portion of the I_d - V_g .

In summary, the “ramped pulse” approach executes a truly trap-free I_d - V_g curve, which is vital, for instance, for extracting intrinsic mobility [79]. However, because it is ultra-fast, this methodology does not allow the amount of trapping that can occur in the gate stack to be quantified. Therefore, the “single” pulse measurement should be used to study charge trapping. To obtain a

correct charge trapping quantity, the rise time portion in the “single” pulse measurement must coincide with the ultra-short, “ramped” pulse one. In the (I_d -time) domain, the “single” pulse approach also delivers information on the charge trapping kinetics. This is the method used to obtain ΔV_t , explained in section V-1.

2. Instrumentation Configurations

Four basic test set-up configurations have been employed to execute many of the pulse-based methodologies reported in the literature [30, 70-72]. In two of these configurations, the MOSFET is operated in an inverter circuit with a load resistor arranged in two ways: 1) a standard resistor placed in series between V_{dd} and the drain of the MOSFET as shown in Fig. 4(a) [30], or 2) the 50 Ω environment of a radio frequency (RF) set-up (Fig. 4(b)) [72]. The third configuration, shown in Fig. 4(c), incorporates an operational amplifier (op-amp) [71] that takes advantage of the virtual short circuit property of op-amp operation. This allows the drain current to be measured across the sense resistor, R. The fourth measurement incorporates a current amplifier in series between common ground and the source of the MOSFET (Fig. 4(d)) [70].

A. “Conventional” set-up

This configuration is a straightforward approach to setting up a pulsed I_d - V_g measurement because it can be built using “standard” lab equipment and cables (Fig. 4(a)) [30]. This configuration can facilitate measurements typically within tens of microseconds when conducting the I_d - V_g form of the measurement [30]. The time limitation is caused by not having a true radio frequency (RF)-capable configuration. Explicit and quantifiable limitations can be found in [71]. However, simply using a small dimension device (e.g., W/L = 10/1 μ m or smaller) and relatively short cables will enable measurements in the tens of microseconds regime.

B. Ultra-Short Pulsed I-V set-up

This configuration requires RF-capable instrumentation and test structures to achieve optimum measurement performance with nanosecond regime timing. The

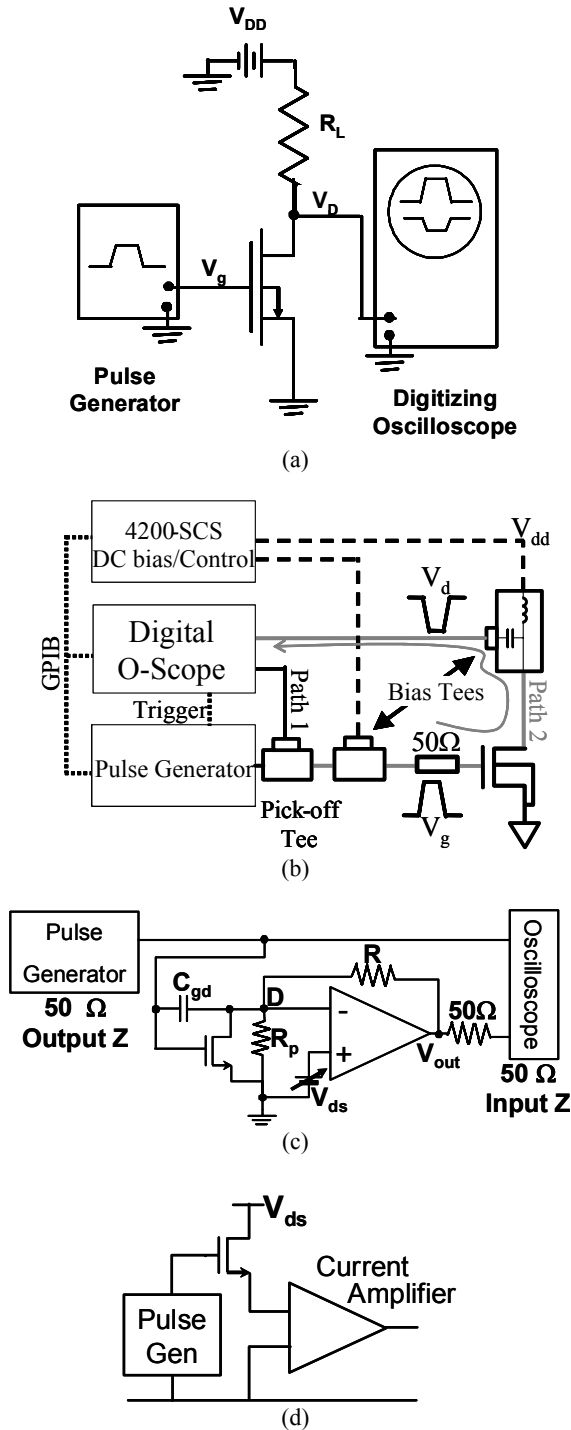


Fig. 4. (a) Schematic illustration of the measurement setup used for pulsed I_D - V_G experiments in the μs range; (b) Experimental set-up for the ultra-short pulsed I - V measurement illustrating signal propagation paths (1 and 2) that should be equal in length and the internal circuitry of a conventional bias tee consisting of an inductor and a capacitor; (c) custom op-amp circuit that takes advantage of the “virtual short” property of the op-amp to mitigate the charging and discharging of parasitic capacitance; (d) current amplifier configuration with the current amplifier placed between source and ground with V_{dd} applied to the drain.

approach still operates with the device under test in an inverter circuit, but uses the $50\ \Omega$ environment as the load resistor. The difference in the experimental set-up is shown in Fig. 4(b) [72]. K. Jenkins, et al., first introduced this type of measurement configuration [80]. Bias tees present a critical component in this RF set-up due to the $50\ \Omega$ environment. The schematically shown bias tee in Fig. 4(b) isolates the AC and DC signals intended for the oscilloscope and transistor drain, respectively, from each other. Note that the characteristic charging time of the internal capacitor in the bias tee may limit the upper range of the pulse width while the minimum current resolution is limited by the voltage resolution of the scope at the smallest range divided by $50\ \Omega$. The set-up discussed here also permits a pulsed I_D - V_d capability due to the two bias tees used in the measurement set-up (Fig. 4(b)) without having to change the test set-up configuration [33, 73]. The “pick-off” tee used in the gate circuit in Fig. 4(b) allows the gate pulse voltage to be monitored non-intrusively. More configuration details can be found in [73]. A potential limiter of this technique is the need for RF timescale equipment and test structures. In addition, this ultra-fast measurement typically measures a completely trap-free result, which does not allow levels of fast transient trapping to be quantified. However, it should serve as a benchmark for all other pulse measurements to ensure an accurate FTCE analysis (see section V. 1.).

C. Operational Amplifier/Current Amplifier Set-up

This configuration has been achieved in two ways: a) a custom design using a standard op-amp circuit (Fig. 4(c)) [71], and b) a current amplifier (Fig. 4(d)) [70].

The custom-designed op-amp circuitry takes advantage of the short circuit property of the op-amp in which the voltage of the two inputs is expected to be equal. The fixed V_{dd} applied to the drain is provided by the voltage source; therefore, the parasitic capacitance, C_o , does not charge or discharge [71]. The measured output voltage captured from the oscilloscope is linked to I_d through the following expression [71]:

$$V_{out} = (I_d - I_{gd}) \cdot R + V_{ds} \quad (2)$$

where I_{gd} is the gate-to-drain current through the parasitic

capacitor (C_{gd}), V_{ds} is the drain voltage, and R is the sense resistance. Because the C_{gd} is weak in short channel devices, I_{gd} is negligible compared to the drain current, thereby allowing I_d to be extracted. The limiting factor for the time resolution of this configuration is determined by the bandwidth of the op-amp. Further set-up details can be found in [71].

The current amplifier, placed in series between ground and the source terminal while V_{dd} is applied to the drain, provides a relatively simple way to make the pulsed I_d - V_g measurement when a V_g pulse is applied to the gate of the MOSFET [70]. Similar to the op-amp circuit, the time resolution for this configuration is determined by the bandwidth of the op-amp.

IV. MEASUREMENT REQUIREMENTS

Several key components must be implemented to ensure robust measurement execution, otherwise significant errors can occur that will impact the extracted signals that are captured and stored on the digital oscilloscope. First, the trapezoidal or triangular pulse should start and end at a discharge condition. For the FTCE, reversibility of this effect was found to be at a discharge voltage of -1V [30, 75, 76]. In addition, a very low duty cycle (e.g., $\leq 0.1\%$) should be implemented. This allows the device under test to essentially be in an “off”/discharge condition most of the time and “on” for extremely short times. This way the “single” pulse can actually be a standard pulse train, allowing averaging if better resolution is desired and/or required to avoid fluctuations.

When using the measurement in the I_d - V_g form (Fig. 3(a)), the input pulse (i.e., V_g pulse) must be synchronized with the output response (i.e., response at the V_d node, dashed box, of Fig. 3(c)). To achieve this, the signal “paths” must take the same propagation time to traverse the path and simultaneously arrive at the oscilloscope. This is achieved by having the same cable length path for the V_g input and V_d response and by having an oscilloscope that can effectively remove any remaining propagation delay (e.g., paths 1 and 2 in Fig. 4(b)) [73, 81]. If this is not done, a measurement error could occur at the minimum frequency/time range of the particular pulse set-up. The error can manifest itself as a false hysteresis as shown in Fig. 5. This 2 nm HfO_2 gate stack,

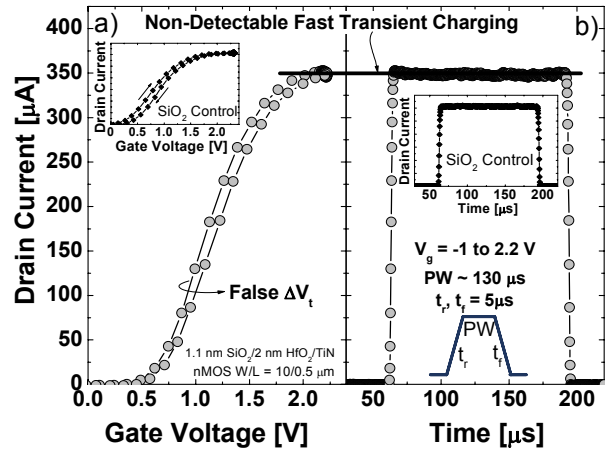


Fig. 5. Single pulsed I-V data illustrating a hysteresis artifact (ΔV_t) in the I_d - V_g domain (a) while no detectable charging (i.e., I_d degradation) is observed in the time domain, (b). The measurement set-up limitation is verified by the artificial hysteresis on conventional SiO_2 , which does not exhibit the FTCE (insets).

which has not suffered from significant FTC [66, 69, 82], should not exhibit any hysteresis. However, Fig. 5a clearly illustrates a small hysteresis, although the I_d -time data (Fig. 5(b)) shows no decay in the pulse width portion. The inset results exhibit a V_t shift artifact for a fast trapping-free SiO_2 sample (insets, Fig. 5) as well, confirming the hysteresis is not real. Therefore, estimating ΔV_t as a shift between I_d - V_g traces corresponding to the rise and fall of the pulse bias may not be accurate [70, 83, 84] when a non-optimized measurement set-up is used.

V. APPLICATIONS

The pulsed I-V methodology can be used for more than just measuring the “intrinsic” drive current performance in the absence of FTC. Proper analysis techniques can provide information on charge trapping [32, 70-74, 81, 83, 85] and the effect on device performance, separation of fast and slow trapping processes [20, 32, 50, 54, 55, 74, 86], and bias temperature instability [50, 87-90].

1. Quantifying Trapped Charge and Its Effects on Device Performance

Using pulsed I_d -time [30, 33, 73, 83] or a correction technique for pulsed I_d - V_g can circumvent the issues discussed in section IV [81]. In the case of I_d -time, very

fast rise and fall times can be used since propagation delay is not an issue. The V_d response is plotted against time instead of the input V_g pulse. The degradation in I_d with time is reflected in the fast transient charge trapping. The threshold voltage shift, ΔV_t , can be determined from the change in I_d through the expression [83, 91]

$$\Delta V_t = \frac{\Delta I_d}{I_d} (V_g - V_t) \quad (3)$$

where ΔI_d is the I_d degradation between the start and end of the V_g pulse width, I_d is the maximum I_d before transient charging, V_g is the pulse amplitude applied to the gate, and V_t is the threshold voltage. The approximation used in Eq. (3) is applicable when mobility does not decrease significantly due to a high vertical field in the channel caused by high gate bias (i.e., when V_g is so high that the slope of the $I_d(V_g)$ curve decreases significantly from its maximum value corresponding to maximum transconductance, $g_{m,max}$) [81].

Fig. 6 illustrates pulsed mobility extracted (to be discussed in section V-1, part A) from the rise (i.e., trapping-free) and fall (i.e., after trapping) time portions showing similar effective mobility values. Even after 100 μs of charging time, the falling I_d - V_g curve appears to be parallel to the rising curve suggesting no significant impact due to remote coulomb scattering from the trapped electrons (Fig. 6, inset), which confirms the

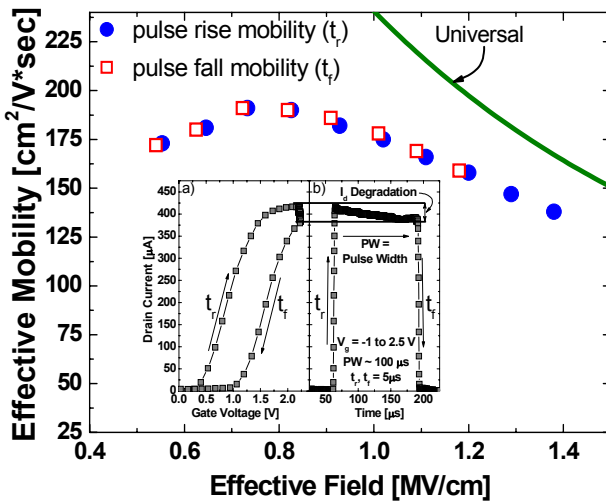


Fig. 6. Similar mobility values extracted from the rise (trap-free) and fall (after trapping) portions of the single pulse measurement (inset, a) demonstrating that remote coulomb scattering is negligible in high- κ gate stacks due to bulk trapping beyond the 1 nm SiO_2 -like interfacial layer.

applicability of Eq. (3) to this particular gate stack. The negligible effect of the trapped electrons on intrinsic mobility can be understood when one considers that the trapped charges are located in the bulk of the high- κ [76, 79], relatively far from the channel.

Therefore, similar pulse data in Fig. 3 can be used as an example of measurements where Eq. (3) is applicable: ΔV_t obtained from the data in Fig. 3(b) using Eq. (3) is equal to the hysteresis value extracted from Fig. 3(a). A precise expression for ΔV_t , which takes into account variation of the mobility value with changing threshold voltage, is given in [83]. With this extraction procedure for ΔV_t , the onset of FTCE can be evaluated [33]. Fig. 7 demonstrates the two extremes, with ($t_p > \tau_c$) and without ($t_p \ll \tau_c$) FTC, for a sample with significant trapping, as illustrated in Fig. 1. By applying pulses with a t_p magnitude between these extreme values, the onset of trapping, τ_c , can be determined.

For any pulsed I-V measurement, the pulse rise time plus its width represents the charging time, t_p . To ensure the absence of FTCE, the pulse charging time must be shorter than the onset of the trapping time, τ_c . This is schematically illustrated in Fig. 1 (inset), which includes two extreme cases of $I_{d,sat}$ current measured with ultra-short ($t_p \ll \tau_c$) and long pulse times. Fig. 8 shows an example of the effect of the pulse rising time (charging time) on I_d [73]. Dependence on the pulse rising times reflects on the charging at higher V_g values. Significant charging during the pulse rising time may result in FTCE being underestimated. Therefore, a measurement capability

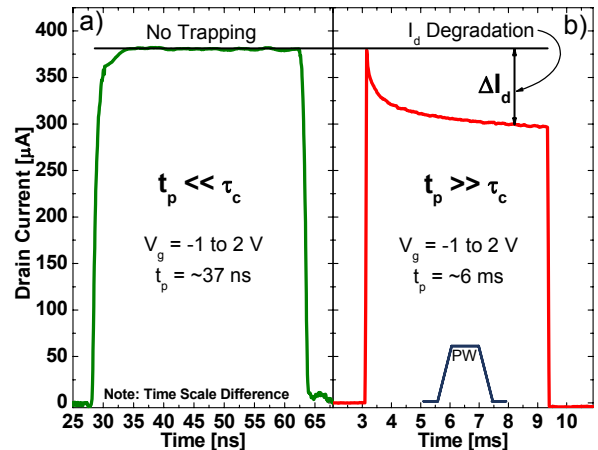


Fig. 7. Pulse I_d dependence on the charging time t_p when (a) $t_p \ll \tau_c$ (no electron trapping), and (b) $t_p \gg \tau_c$ (electron trapping is proportional to ΔI_d).

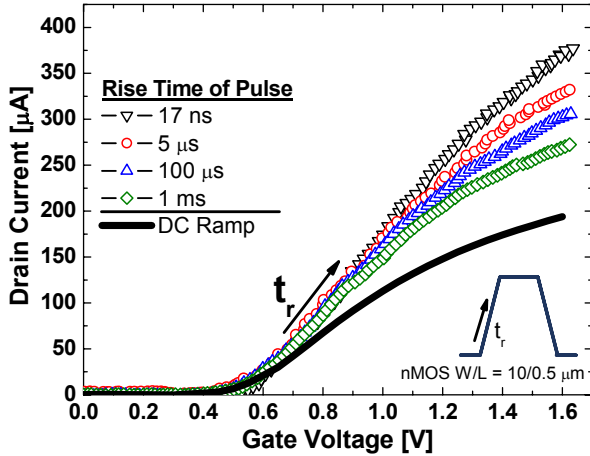


Fig. 8. Single pulsed I-V measurements on high- κ demonstrating the impact of the pulse rise time on charging. A fast enough rise time (i.e., short t_r) ensures a “trap-free” drain current.

that ensures a trapping-free characteristic throughout the entire pulse measurement is required to guarantee no FTCE. In addition, highly scaled interfacial layers in high- κ gate stacks may exhibit very fast trapping, which would require pulse measurements below the microsecond regime (Fig. 8). This can be achieved by using the ultra-short pulsed I-V set-up in Fig. 4(b). Furthermore, this is a potentially powerful measurement technique for characterizing high- κ gate stacks on silicon-on-insulators when self-heating effects should also be eliminated to achieve intrinsic device performance [80].

Another approach, which uses the pulsed I_d - V_g hysteresis, can be employed when the limitations of being able to use eq. (3) have been exceeded (Fig. 9) [81]. This approach effectively negates the propagation delay by intentionally introducing a phase delay time between

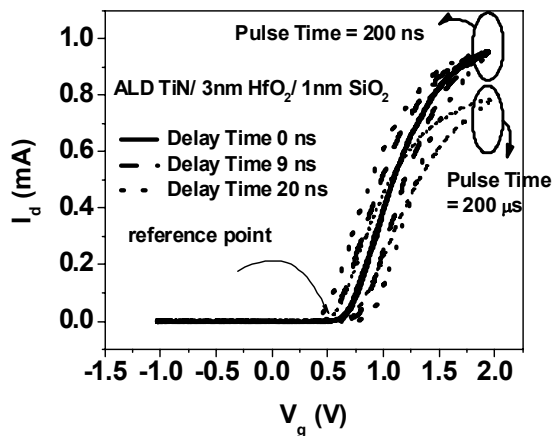


Fig. 9. The effect of phase delay time between the two input channels of the oscilloscope on the measured hysteresis of the I_d - V_g curves in a single pulse measurement.

the two channels of the oscilloscope. To obtain the accurate phase delay time to compensate for the propagation delay, I_d - V_g curves generated by “short” and “long” triangle pulses with equal rise and fall times, should be compared (Fig. 9). Three different phase delay times ranging from zero to 20 ns were used in both short (thick curves) and long (thin curves) pulses. It can be seen that the I_d - V_g curves generated by the short pulses are strongly affected by the phase delay times, which result in different I_d - V_g hysteresis values. On the other hand, the complete overlap of the long pulse curves obtained with different delay times indicates that the phase delay times in the ns range have no effect on the pulse in the μ s range. Therefore, by varying the phase delay time in the short pulse measurement to match the I_d - V_g curves of the short and long pulses within the voltage range for V_g less than V_t (where no FTCE takes place), the correct phase delay time that compensates the signal delay due to different propagation paths for the given pulse measurement set-up can be determined.

A. Impact of Dielectric Properties on Fast Transient Charge Trapping

Typical high- κ gate stack structures are comprised of two layers: an interfacial layer and a high- κ layer (Fig. 1). During gate stack fabrication, the interfacial layer can be engineered or it can grow spontaneously [33, 78, 92, 93]. The FTCE can be influenced by these two layers as well as by the amount of hafnium in the high- κ layer.

Effects of the interfacial layer thickness, physical thickness of the high- κ dielectric, and its Hf content on the fast electron trapping process were evaluated by using a combination of the ultra-fast pulse technique described above (covering the range up to 100 ns) and the conventional “single” pulse (covering the range between 100 ns to 6 ms) [30, 76, 94]. The effect of interfacial layer thickness is demonstrated in Figs. 10-12. Since the trap sites in the high- κ layer are filled by tunneling from the inversion layer through the interfacial layer (Fig. 1), a thinner interface leads to a stronger tunneling current and, hence, greater electron trapping. Accordingly, critical time (the pulse width at which the onset of trapping occurs), τ_c , decreases and the V_t shift increases for thinner interfacial oxides (Fig. 10). The impact of electron trapping on DC mobility is shown in

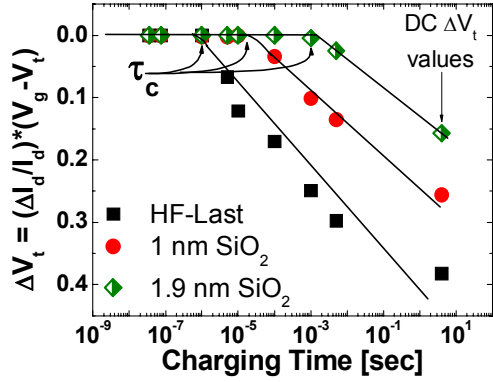


Fig. 10. Normalized ΔV_t for varying IL thicknesses. Thicker IL increases the tunneling distance (see Fig. 1) to available traps, thereby increasing τ_c .

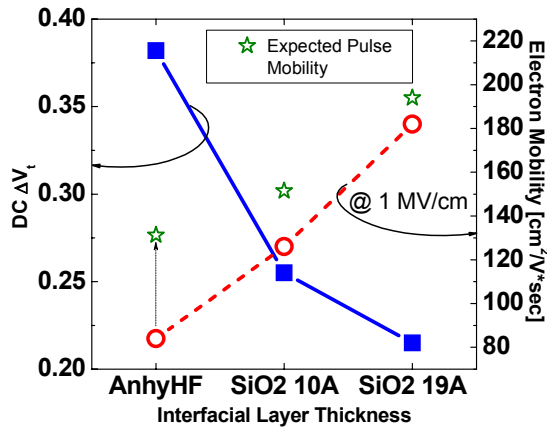


Fig. 11. Correlation of V_t shift (trapped charge) to the extracted DC mobility illustrating degraded mobility as the IL becomes thinner.

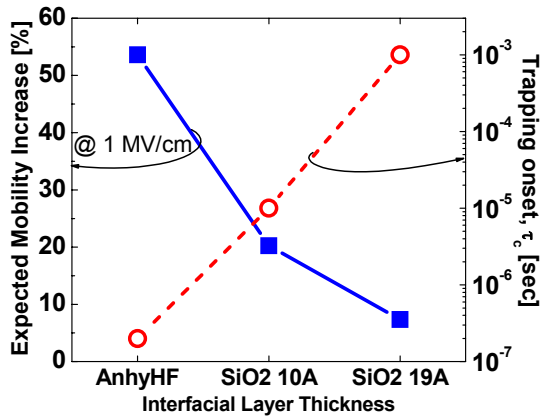


Fig. 12. Expected mobility increase (thinner IL \rightarrow larger pulse/DC $I_{d,lin}$ ratio) as a function of interfacial layer thickness compared with the onset of the critical time, τ_c .

Fig. 11 where the DC mobility increases as the ΔV_t values, calculated based on the ΔI_d [83, 91], diminish in proportion to the thickness of interfacial layer.

A mobility increase (which may be expected if the fast

charge trapping is eliminated), called “expected mobility,” was estimated from the ratio of $I_{d,lin}$ values measured in the pulse and DC regimes (Figs. 11 and 12).

For samples of different $HfSi_xO_y$ (20% SiO_2) physical thicknesses of 3, 5, and 7 nm, all with a TiN electrode, the interfacial layer was fixed at ~ 1 nm SiO_2 . As the physical thickness of the dielectric increases, the total number of traps accessible at the given applied bias also increases, or the filled trap location shifts towards the device channel, resulting in a greater V_t shift [76] (Fig. 13) that correlates with the degradation of the electron mobility (Fig. 14).

Hafnium content has an apparent influence on both the electron trapping and τ_c (Figs. 15 and 16). The ΔV_t increases with the Hf content indicating that electron traps may be associated with the Hf atoms.

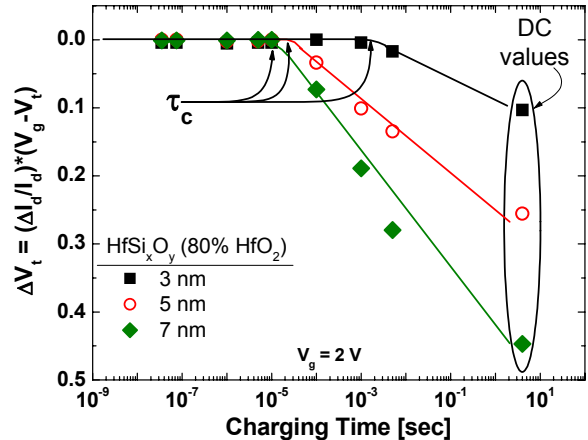


Fig. 13. ΔV_t for various high- κ physical thicknesses. Physically thicker high- κ is subject to faster and more severe charge trapping.

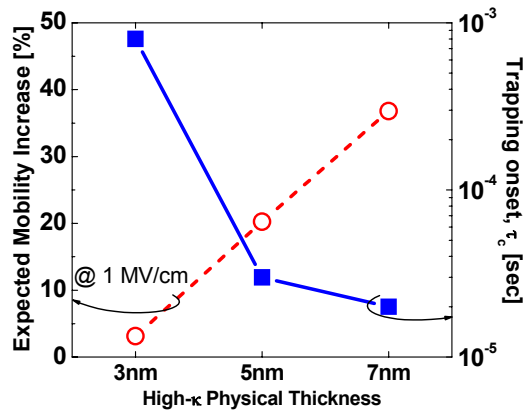


Fig. 14. Expected mobility increase (thinner high-k \rightarrow larger pulse/DC $I_{d,lin}$ ratio) as a function of high- κ layer thickness compared with the onset of the critical time, τ_c .

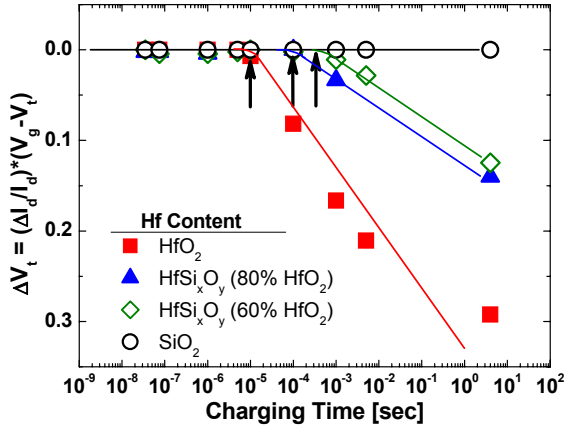


Fig. 15. Normalized ΔV_t for various hafnium contents. More Hf results in faster and greater charge trapping.

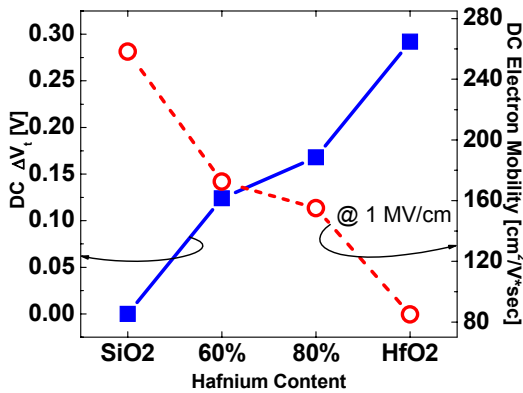


Fig. 16. An increase in DC ΔV_t and the corresponding decrease in DC mobility as the Hf content increases.

B. Factors Controlling V_t shift

In the pulsed I-V measurement, the effect of the trapped charge on ΔV_t is primarily determined by the V_g pulse conditions. The V_g pulse amplitude affects both the magnitude of the injection current and the point of injection (the charge trapping location) in the dielectric. To understand the major factors controlling ΔV_t —charge supply or trapping location with respect to the gate electrode—the pulse width was fixed at 100 μs and the pulse V_g (i.e., pulse height) was varied to study the bias dependence [73].

It has been demonstrated that the fast trapping process, with a characteristic time of about 100 μs , is independent of temperature, implying that the pre-existing traps are filled by the resonant capture of the tunneling electron injected from the substrate [20, 74]. The charges trapped by this process are well-localized within the dielectric thickness, with their location being determined by the

Fermi energy of the channel electrons and the trap energy, whose value is close to the HfO_2 conduction band edge, Fig. 1 (presumably, a neutral or single-charged negative oxygen vacancy) [4, 7, 9-11, 16]. Therefore, we assume for simplicity that the traps are characterized by a single energy value that is close to the bottom of the conduction band [4, 20, 74]. For longer stress times, over 1 second, the trapped charge distribution widens due to slow temperature-activated electron capture processes [50, 74].

The obtained dependences of ΔV_t vs. injected charge, Fig. 17, show no correlation, indicating that fast transient charging is not a charge supply-driven phenomenon. To evaluate the impact of the position of the trapped charge on ΔV_t , band diagrams were constructed, as shown in the examples in Fig. 18, based on the physical thicknesses and κ values of each layer in the different bi-layer stacks under investigation. The band diagrams were generated for various V_g pulses using the software developed by Boise State University [95]. Then, at the given V_g , the location of the trapped charge is defined by the resonance between the Fermi level in the substrate and the energy level of the pre-existing traps in the Hf-based dielectric [50, 74]. Knowing the distance of the trapped charge from the gate electrode allows the trap charge density to be calculated by

$$N_{trap} = \frac{\epsilon_o \cdot \kappa \cdot \Delta V_t}{q \cdot x} \tag{4}$$

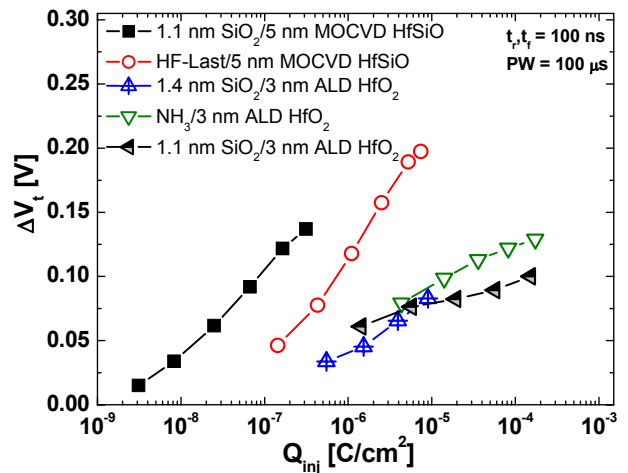


Fig. 17. ΔV_t vs. the injected charge, Q_{inj} illustrating no correlation between the ΔV_t and injected charge.

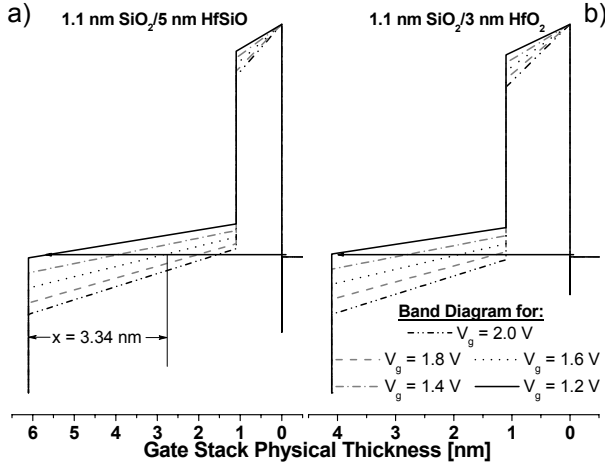


Fig. 18. The gate stack energy band diagrams for (a) 1.1 nm $\text{SiO}_2/5$ nm HfSiO and (b) 1.1 nm $\text{SiO}_2/3$ nm HfO_2 for various pulsed V_g values. An example of the resonant electron trapping location with respect to the Fermi level for $V_g = 1.6$ V is shown.

where ΔV_t is the extracted threshold voltage shift, x is the experimentally extracted distance from the gate electrode (example, Fig. 18(a)), q is the fundamental electron charge, ϵ_0 is the permittivity of free space, and κ is the dielectric constant of the high- κ layer. (For more details on the assumptions and limitations of this trap density extraction procedure, see [73].)

By changing V_g , one can obtain x and thus the trap density profile across the thickness of the high- κ film using Eq. (4). While the position x at each given pulse voltage may not be defined exactly, the measurements still allow the relative distribution of x values to be extracted in a self-consistent way. Fig. 19 shows the

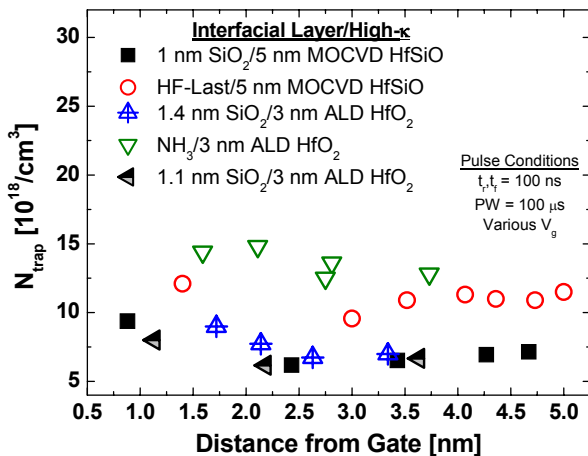


Fig. 19. Defect density profile of different high- κ gate stacks calculated using the resonant trap location distance referenced from the gate electrode (Fig. 18) for the corresponding V_g pulse value.

defect profiles for several gate stacks. Within the referenced assumptions, uniform trap density profiles were obtained in all gate stacks analyzed in this study. As follows, the observed increase of ΔV_t with the pulse amplitude is due to the trapped charge being closer to the substrate rather than the amount of the trapped charge being greater.

In summary, this approach probes each sample through its high- κ thickness by changing the measurement pulse voltage, thereby changing the band bending (i.e., resonant trapping point), which allows trap density information to be directly extracted at each point in the given sample [73].

C. The Effect of FTC on Channel Mobility

In the linear regime, the FTCE can significantly impact mobility [31, 63, 64, 66, 79, 83, 91, 92, 96-98]. In a conventional DC sweep of the I_D - V_g measurement, the threshold voltage shifts as the measurement progresses, thereby decreasing the drive current at each bias sweep point resulting in decreased mobility. Several FTC correction techniques discussed in the literature can correct for this effect [31, 63, 79, 83]. One technique that requires only two measurements and minimal data analysis was presented in [79]. In this approach, pulsed I_D - V_g is combined with North Carolina State University's CVC [99] and mob2d [100] modeling parameter extraction software. To extract mobility, mob2d performs a non-linear, least squares fit to experimental pulsed I_D data in conjunction with the EOT and substrate doping supplied from CVC. EOT and substrate doping are negligibly impacted by FTC [101]. Therefore, the result is a trap-free mobility extraction. Fig. 20 illustrates the impact of trapped charge in the DC mobility data compared to the trap-free, pulsed-based extraction for the I_D - V_g data in the inset where the pulse-based mobility is higher than its DC-based counterpart [32, 80]. Fig. 21 compares the pulse and DC mobility extractions using CVC and mob2d for high- κ gate stacks with different interfacial layer (IL) thicknesses. The sample with a 1.9 nm SiO_2 interfacial layer exhibits essentially the same DC and pulsed mobility since fast transient charging is negligible in thick interfaces (quite similar DC and pulsed I_D - V_g curves – not shown). As the interface becomes thinner, the difference between the pulsed and

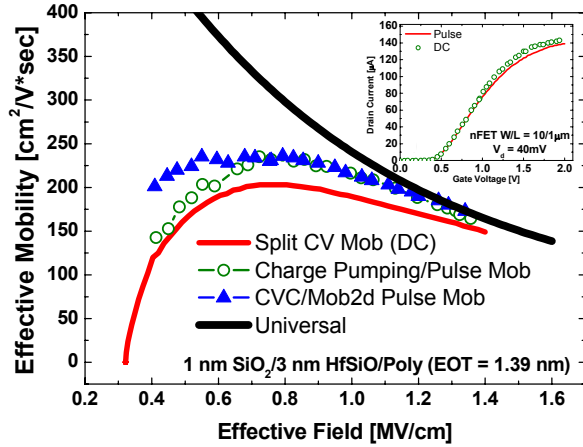


Fig. 20. Comparison of DC mobility to the novel pulse mobility technique for the I_d - V_g data in the inset. The mobility values extracted using the proposed mob2d pulse methodology and the direct measurement of N_{inv} technique [32] show excellent agreement.

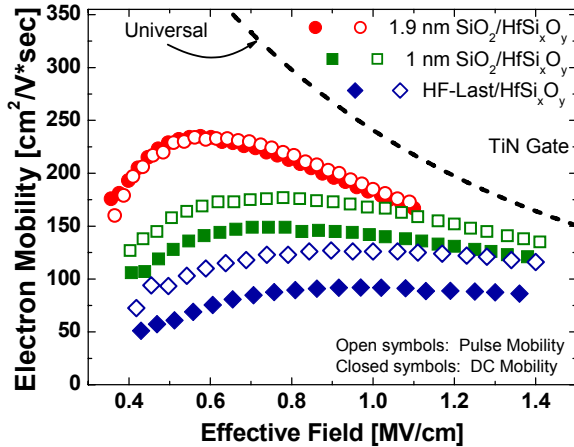


Fig. 21. Application of the pulse mobility extraction method [80] to gate stacks with interfacial layers (IL) of varying thicknesses and fixed $HfSiO_x$ dielectric layer.

DC mobility for a given gate stack becomes greater. The overall degradation in the pulsed (intrinsic) mobility as the interface layer becomes thinner can be attributed to poorer interfacial layer quality [102] and greater effects of soft optical phonons [64] and crystallinity [21, 103].

D. Fast Transient Trapped Charge Relaxation

Reports have shown that, in addition to FTC, high- κ gate dielectrics also demonstrate fast relaxation of the trapped charge, although on a slightly longer time scale [88, 104]. Fig. 22 illustrates this fast relaxation phenomenon. After a two-level V_g pulse (Fig. 22 inset: the high level is in strong inversion, and the low level is just

above V_t) the drain current recovers due to relaxation (charge detrapping) at this low V_g level in $\sim 150 \mu s$. Fig. 23 demonstrates that the trapped charge can return to the DC measurement results in as little as 250 μs .

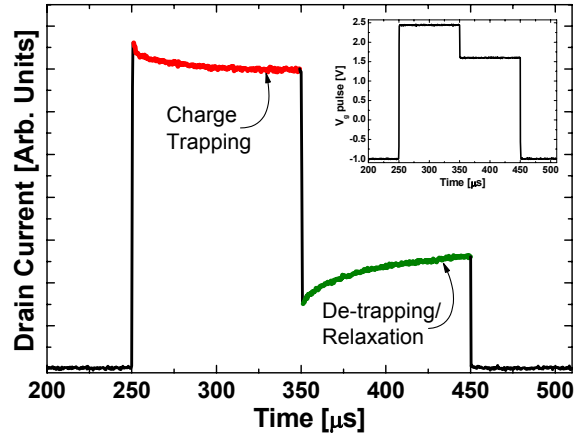


Fig. 22. Illustration of fast transient charge trapping (I_d decay) and detrapping (I_d recovery) for the V_g pulse in the inset.

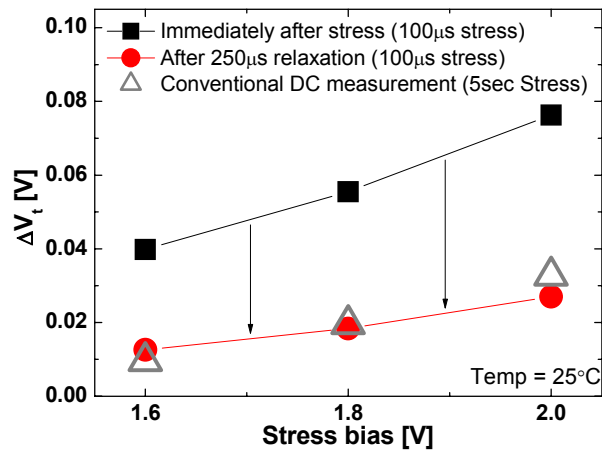


Fig. 23. V_t shifts immediately after the 100 μs stress and after 250 μs of post-stress relaxation compared to conventional “stress and sense” approach.

2. Implications of Fast Transient Trapping and Detrapping During Stress Measurements

Note that since fast trapping and recovery can occur in high- κ gate stack transistors, it significantly impacts positive bias temperature instability with an adverse effect on the lifetime projection [87, 88]. Fast recovery can occur between the stress and “sense” (i.e., monitoring parameter measurement) of conventional constant voltage stress with interspersed DC I_d - V_g measurements (V_t extraction) due to the interval between the removal of the stress and the subsequent “sense” I_d - V_g

measurement. This can lead to significant underestimation of the V_t shift [22, 25, 32, 84, 89, 90, 105, 106]. Therefore, the time delay between the stress and measurement must be shortened.

A. Effect of Relaxation and Its Impact on the Lifetime Projection

Several techniques have been introduced in the literature to mitigate the effects of fast relaxation [88, 90, 107]. One of particular interest employs the pulsed I-V methodology [88, 90]. Fig. 24 illustrates the measurement schematic for the pulsed approach vs. conventional DC stress and sense. As Fig. 24(b) shows, the pulsed I-V is measured immediately in the downward pulse from the stress condition, thereby removing any relaxation [88]. To study the impact of fast relaxation on the extraction of V_t , the I_d of the rising trace on the inverted triangular pulse can be extracted after modulating the fall time of the downward portion, which can effectively quantify the relaxation time [88]. The results are shown in Fig. 25 where the quantity of the trapped charge, reflected in the V_t shift, is similar to the results of the conventional DC approach as the relaxation time increases [88]. Moreover, the slopes (i.e., power law exponents) change significantly with changes in the relaxation time. Therefore, it is worthwhile to separate the effects of the fast relaxation on the time evolution of degradation with stress.

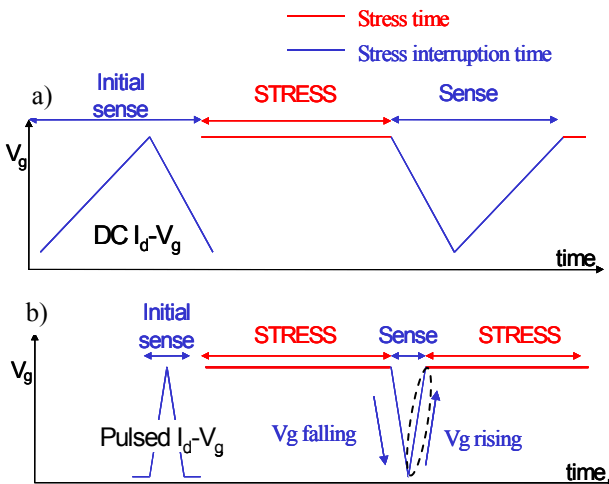


Fig. 24. The schematic comparison of stress-sense schemes using the stress-interrupted DC I_d - V_g method and single pulse method.

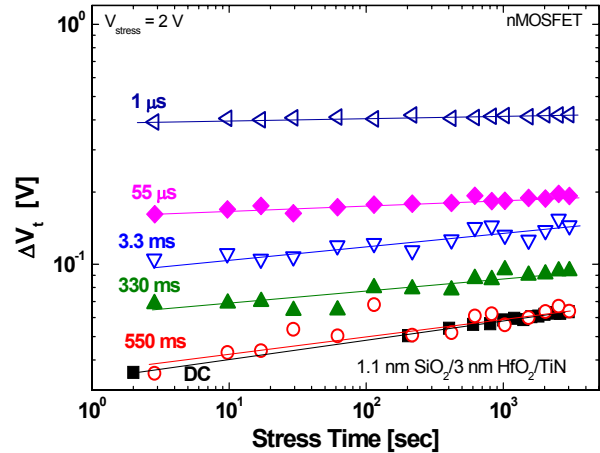


Fig. 25. ΔV_t dependence on stress time when extracting V_t from the rising portion of V_g (see Fig. 24(b)) for different pulse times.

B. Separation of Fast and Slow Trapping

In addition to fast transient mechanisms, slow trapping and/or trap generation can impact V_t instability in high- k transistors when constant voltage stress is measured [22, 25, 30, 34, 84, 88, 89, 105]. The oxygen vacancy defects responsible for the FTCE have also been found to exhibit a temperature-activated slow migration of trapped electrons to unoccupied traps [50, 74]. Fig. 26 illustrates these two components over a long time scale [50]. A temperature-independent feature is detected in fast measurement times (pulsed I-V regime), indicating a direct tunneling mechanism to the neutral or single-charged negative oxygen vacancy fast transient traps. Then, at longer stress times, a temperature effect is seen when a larger ΔV_t is extracted at the higher temperature. The fast relaxation detraps in a similar manner [50].

To remove the effect of the FTC so that only the slow migration and/or trap generation can be studied, the previously mentioned pulsed I-V combined with constant voltage stress and subsequent data analysis was used [88]. A physical model-based subtraction of the initial V_t (contains only the fast relaxation portion) from each subsequent V_t point in the time dependence effectively removes the fast transient detrapping component [88]. Fig. 27 shows parallel curves after the initial data point subtraction provides a “universal” slope, which suggests that the dependence of the power law exponent on pulse time is due to fast transient charge detrapping [87]. In addition, Fig. 28 further confirms this universality because different measurement techniques have a very similar

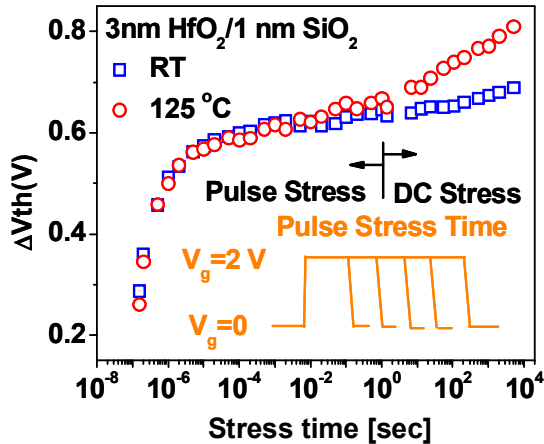


Fig. 26. Fast (temperature-independent) and slow (temperature-dependent) charge trapping processes in typical HfO₂ gate stacks where the fast process pseudo-saturates within 100 μs and transitions to the slow process around 10⁰ sec of stressing.

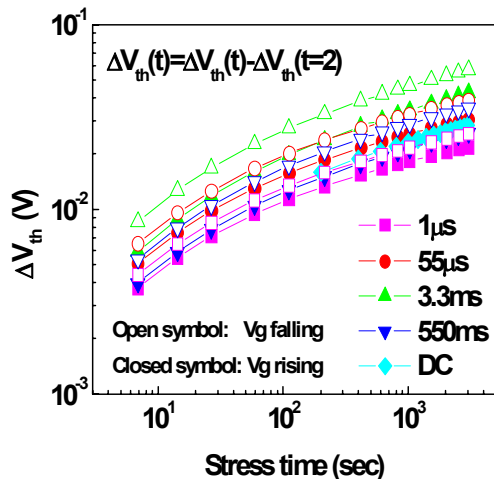


Fig. 27. The power law dependence after removal of the FTC portion of the V_t shift (i.e., subtraction of the first ΔV_t point from all subsequent points) demonstrating that all curves fall on an intrinsic curve regardless of the sense measurement time.

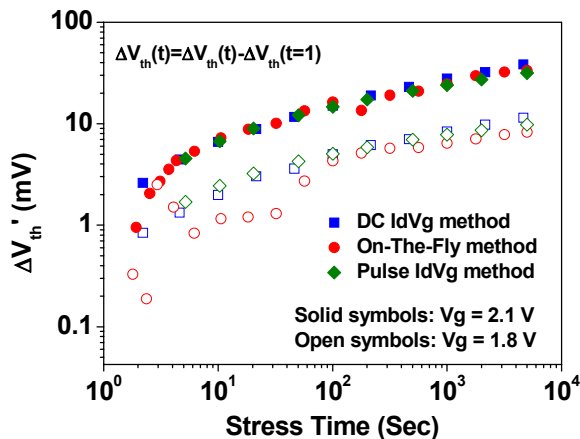


Fig. 28. The correlation of ΔV_t measured using different characterization methods.

slope (power law exponent) after removing the fast transient effects when measured on the same high-κ gate dielectric stack [87].

VI. CONCLUSIONS

Several pulsed I-V measurement configurations have been critically evaluated. The ultra-short pulsed I-V method is found to be the benchmark for determining truly trap-free characteristics of high-κ gate dielectric devices. The pulsed I-V methodology is shown to be capable of obtaining “intrinsic” (i.e., trap-free) performance of hafnium-based gate dielectrics. To evaluate reliability, the pulsed-based approach is required to extract device lifetime, which can be impacted by the fast transient phenomena. However, to ascertain the trap-free results requires a rigorous evaluation of the measurement configuration along with its corresponding analysis. Failing to calibrate the measurement set-up properly (or employing a set-up that is too slow) results in an underestimation of the V_t shift, effective mobility, and trap density.

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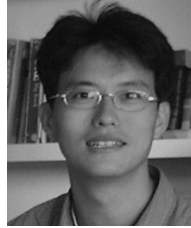
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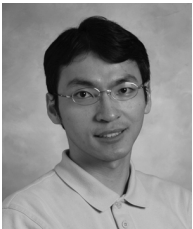
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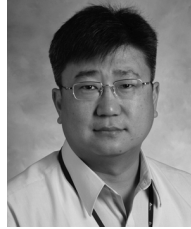
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