# ABSOLUTELY STABLE EXPLICIT SCHEMES FOR REACTION SYSTEMS 

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#### Abstract

We introduce two numerical schemes for solving a system of ordinary differential equations which characterizes several kinds of linear reactions and diffusion from biochemistry, physiology, etc. The methods consist of sequential applications of the simple exact solver for a reversible reaction. We prove absolute stability and convergence of the proposed explicit methods. One is of first order and the other is of second order. Numerical results are included.


## 1. Introduction

Many phenomena of interest in physiology and biochemistry are characterized by reactions among several chemical species and diffusion in various mediums (see $[6-8,10]$ ). In a closed system, both reactions and diffusion are governed by a system of ordinary differential equations (ODEs)

$$
\begin{equation*}
\dot{\mathbf{y}}(t)=M \mathbf{y}(t) \tag{1.1}
\end{equation*}
$$

which guarantees conservation of the total amount of $\mathbf{y}(t)$ for any $t \geq 0$. Since we are concerned with the steady-state solution as well as the transient in simulations of very large systems of chemical reactions or molecular dynamics, we need to take the overall computational cost into consideration. Many physiologists and biochemists prefer explicit methods to implicit methods since implementation of the explicit methods is easier than the others. The popular methods for reaction systems are simple explicit schemes such as Euler's method, Runge-Kutta method, etc. However, it is well-known that conditional stability, the typical weak point of explicit methods, is very fatal for stiff problems. In the past few decades, many studies on numerical methods for stiff ODEs have been done in various aspects (see $[1,2,4,5,11]$ ).

The aim of this paper is to present two absolutely stable explicit schemes which are applicable to a linear reaction system (1.1). In 1978, Rush and

[^0]Larsen [9] introduced an iterative procedure for the Hodgkin-Huxley model for cell membrane behavior, which is composed of a circuit equation for currents and a coupled system of nonlinear ODEs for the ionic gates. An integration algorithm was suggested for a numerical solution to the ODEs for the ionic gates, which was based on the exact solution of a linearized ionic gate equation. Similarly, the methods in this paper are motivated by the simple exact solver for a reversible reaction. In spite of their explicitness, we have unconditional stability, that is, stability without any condition on the step size. Furthermore, we prove the convergence of the proposed methods; one is of first order and the other is of second order.

This paper is organized as follows. In Section 2, we introduce the reaction systems of our interest and propose two numerical methods for a linear reaction system. Section 3 provides theoretical results for convergence and stability of the proposed methods. In Section 4, we provide numerical experiments for typical reaction systems.

## 2. Reaction systems and numerical methods

We consider two typical types of reactions: reversible reactions and circular reactions. A reaction of the type

$$
A \xlongequal[k_{b}]{\stackrel{k_{f}}{\leftrightharpoons}} B
$$

is called the reversible reaction, where $k_{f}$ and $k_{b}$ are the rate constants for the forward and backward reactions. One interesting biochemical system to which the reversible first order equations apply is the carbonic acid system:

$$
\begin{equation*}
\mathrm{CO}_{2}+\mathrm{H}_{2} \mathrm{O} \underset{k_{-1}}{\stackrel{k_{1}}{\rightleftharpoons}} \mathrm{H}_{2} \mathrm{CO}_{3} \underset{\substack{\text { very } \\ \text { fast }}}{\stackrel{K_{a}}{\longrightarrow}} \mathrm{H}^{+}+\mathrm{HCO}_{3}^{-} \tag{2.1}
\end{equation*}
$$

Equation (2.1) reduces to

$$
\mathrm{CO}_{2} \underset{k_{b}}{\stackrel{k_{f}}{\leftrightarrows}} \mathrm{HCO}_{3}^{-}
$$

(see [8]). Then, the rate equations are written down as

$$
\begin{align*}
& \frac{d A}{d t}=-k_{b} A+k_{f} B \\
& \frac{d B}{d t}=k_{b} A-k_{f} B \tag{2.2}
\end{align*}
$$

where $A(t)$ and $B(t)$ are the concentrations of $\mathrm{CO}_{2}$ and $\mathrm{HCO}_{3}^{-}$as functions of time $t$. Another interesting example of such a type occurs frequently in metabolic studies. The other type is a circular reaction shown in Figure 1. We


Figure 1. A circular reaction system with 3 substances $A_{0}, A_{1}$ and $A_{2}$
may write the differential equations describing this process as

$$
\frac{d}{d t}\left[\begin{array}{l}
A_{0}(t)  \tag{2.3}\\
A_{1}(t) \\
A_{2}(t)
\end{array}\right]=\left[\begin{array}{ccc}
-b_{10}-b_{20} & f_{10} & f_{20} \\
b_{10} & -f_{10}-b_{21} & f_{21} \\
b_{20} & b_{21} & -f_{21}-f_{20}
\end{array}\right]\left[\begin{array}{l}
A_{0}(t) \\
A_{1}(t) \\
A_{2}(t)
\end{array}\right]
$$

From the fact that the total concentration $A(t)+B(t)$ remains constant for all $t \geq 0$ in a closed system, the exact solution of (2.2) is written in a form

$$
\begin{aligned}
{\left[\begin{array}{l}
A(t) \\
B(t)
\end{array}\right] } & =\left[\begin{array}{cc}
\frac{1}{k_{f}+k_{b}}\left(k_{f}+k_{b} e^{-\left(k_{f}+k_{b}\right) t}\right) & \frac{k_{f}}{k_{f}+k_{b}}\left(1-e^{-\left(k_{f}+k_{b}\right) t}\right) \\
\frac{k_{b}}{k_{f}+k_{b}}\left(1-e^{-\left(k_{f}+k_{b}\right) t}\right) & \frac{1}{k_{f}+k_{b}}\left(k_{b}+k_{f} e^{-\left(k_{f}+k_{b}\right) t}\right)
\end{array}\right]\left[\begin{array}{l}
A(0) \\
B(0)
\end{array}\right] \\
& =\exp \left(t\left[\begin{array}{cc}
-k_{b} & k_{f} \\
k_{b} & -k_{f}
\end{array}\right]\right)\left[\begin{array}{l}
A(0) \\
B(0)
\end{array}\right] .
\end{aligned}
$$

Similarly, we can find the exact solution for the circular reaction in Figure 1. In general, a reaction system is characterized by a coupled system of ODEs. To solve a relevant eigenvalue problem is the first step in solving such a coupled system exactly (see [12]). But as the number of substances increases, the exact solver suffers from typical difficulties in large scale eigenvalue problems. There are some numerical techniques in common use: the Euler method, which is the simplest one, but requires a small size of time step $\Delta t$; the Runge-Kutta method, which is more complicated, but allows much bigger time steps to be taken. Now for the circular reaction case, we propose new numerical methods motivated by the above process that is used to find the exact solution to a reversible reaction. For the sake of simplicity, we illustrate these algorithms for a simple circular reaction (2.3) although they are applicable to more complicated reaction systems.

## Algorithm 1: $\mathbf{C R}^{2}$

(1) For each $k \in \mathbb{N}$, let $A_{i, k}$ be the approximate solution to $A_{i}(t)$ at time $t_{k}=k \Delta t$.
(2) For $k=0,1,2, \ldots$ do:
(a) Set $A_{0}^{i c}=A_{0, k}, A_{1}^{i c}=A_{1, k}$ and $A_{2}^{i c}=A_{2, k}$.
(b) Find the exact solution $A_{0}^{\text {temp }}(t)$ and $A_{1}^{\text {temp }}(t)$ of the reversible reaction with the initial values $A_{0}^{i c}$ and $A_{1}^{i c}$ :

$$
A_{0} \underset{b_{10}}{\stackrel{f_{10}}{=}} A_{1}
$$

Set $A_{0}^{i c}=A_{0}^{\text {temp }}\left(t_{k+1}\right)$ and $A_{1}^{i c}=A_{1}^{\text {temp }}\left(t_{k+1}\right)$.
(c) Find the exact solution $A_{0}^{\text {temp }}(t)$ and $A_{2}^{\text {temp }}(t)$ of the reversible reaction with the initial values $A_{0}^{i c}$ and $A_{2}^{i c}$ :

$$
A_{0} \underset{\underset{b_{20}}{ }}{\stackrel{f_{20}}{=}} A_{2}
$$

Set $A_{0, k+1}=A_{0}^{\text {temp }}\left(t_{k+1}\right)$ and $A_{2}^{i c}=A_{2}^{\text {temp }}\left(t_{k+1}\right)$.
(d) Find the exact solution $A_{1}^{\text {temp }}(t)$ and $A_{2}^{\text {temp }}(t)$ of the reversible reaction with the initial values $A_{1}^{i c}$ and $A_{2}^{i c}$ :

$$
A_{1} \underset{b_{21}}{\stackrel{f_{21}}{\leftrightharpoons}} A_{2}
$$

Set $A_{1, k+1}=A_{1}^{\text {temp }}\left(t_{k+1}\right)$ and $A_{2, k+1}=A_{2}^{t e m p}\left(t_{k+1}\right)$.

Remark 2.1. $\mathrm{CR}^{2}$ stands for Consecutive Reversible Reactions.
The key idea of Algorithm 1 is that we approximately regard a circular reaction as a consecutive reaction which consists of three separated reversible reactions. Although we solve three reversible reactions in the following order

$$
A_{0} \leftrightharpoons A_{1}, A_{0} \leftrightharpoons A_{2}, A_{1} \leftrightharpoons A_{2}
$$

the algorithm does not depend upon the ordering we adopt to split a circular reaction into a chain of reversible reactions. In a form of matrix exponential, Algorithm 1 is represented by the following linear system:

$$
\begin{aligned}
{\left[\begin{array}{l}
A_{0, k+1} \\
A_{1, k+1} \\
A_{2, k+1}
\end{array}\right] } & =\exp \left(\Delta t \bar{M}_{21}\right) \exp \left(\Delta t \bar{M}_{20}\right) \exp \left(\Delta t \bar{M}_{10}\right)\left[\begin{array}{l}
A_{0, k} \\
A_{1, k} \\
A_{1, k}
\end{array}\right] \\
& \stackrel{\text { def }}{=} L_{2}\left[\begin{array}{l}
A_{0, k} \\
A_{1, k} \\
A_{2, k}
\end{array}\right],
\end{aligned}
$$

where $\bar{M}_{i j}$ is a 3 by 3 matrix relevant to the reversible reaction between $A_{j}$ and $A_{i}$ such that

$$
\left(\bar{M}_{i j}\right)_{l m}= \begin{cases}-b_{i j} & \text { if }(l, m)=(j, j),  \tag{2.4}\\ f_{i j} & \text { if }(l, m)=(j, i), \\ b_{i j} & \text { if }(l, m)=(i, j), \\ -f_{i j} & \text { if }(l, m)=(i, i), \\ 0 & \text { otherwise }\end{cases}
$$

We interpret the Algorithm 1 for $n$ substances inductively, i.e., at $t=t_{k}$, we assume that the reaction system with $j$ substances $A_{0}, A_{1}, \ldots, A_{j-1}(1 \leq j \leq$ $n-1)$ is solved, then we solve the parallel reaction that consists of $j$ reversible reactions:

$$
\begin{array}{ccl}
A_{j} & \stackrel{f_{j 0}}{\stackrel{f_{j 0}}{\rightleftharpoons}} & A_{0} \\
A_{j} & \stackrel{f_{j 1}}{\stackrel{f_{j 1}}{\rightleftharpoons}} & A_{1} \\
& \vdots & \\
A_{j} & \underset{b_{j(j-1)}}{b_{j(j-1)}} & A_{j-1}
\end{array}
$$

## Algorithm 2: $\mathbf{S C R}^{2}$

(1) For each $k \in \mathbb{N}$, let $A_{i, k}^{s}$ be the approximate solution to $A_{i}(t)$ at time $t_{k}=k \Delta t$.
(2) For $k=0,1,2, \ldots$ do:
(a) Set $A_{0}^{i c}=A_{0, k}, A_{1}^{i c}=A_{1, k}$ and $A_{2}^{i c}=A_{2, k}$.
(b) As in Algorithm 1, solve three reversible reactions according to chosen ordering: $A_{0} \leftrightharpoons A_{1}, A_{0} \leftrightharpoons A_{2}, A_{1} \leftrightharpoons A_{2}$. Define $A_{i, k+1}^{s_{1}}$ by the obtained approximation $A_{i, k+1}$ for $i=1,2,3$.
(c) Solve three reversible reactions in reverse order: $A_{2} \leftrightharpoons A_{1}, A_{2} \leftrightharpoons$ $A_{0}, A_{1} \leftrightharpoons A_{0}$. Define $A_{i, k+1}^{s_{2}}$ by the obtained approximation $A_{i, k+1}$ for $i=1,2,3$.
(d) Set

$$
A_{i, k+1}^{s}=\frac{A_{i, k+1}^{s_{1}}+A_{i, k+1}^{s_{2}}}{2} \text { for } i=1,2,3 .
$$

Remark 2.2. For reaction systems with more than three substances, there are two different ways to implement the reverse ordering in the step (2)-(c) of the Algorithm 2; one is to reverse the order of reversible reactions in the step (2)-(b) and the other is to put the substances used in the step (2)-(b) in reverse order. In other words, we suppose that in a reaction system with four substances, we solve the six reversible reactions in the following order

$$
A_{0} \leftrightharpoons A_{1}, A_{0} \leftrightharpoons A_{2}, A_{1} \leftrightharpoons A_{2}, A_{0} \leftrightharpoons A_{3}, A_{1} \leftrightharpoons A_{3}, A_{2} \leftrightharpoons A_{3}
$$

Then, at our convenience we adopt one of two reverse ordering

$$
\begin{equation*}
A_{2} \leftrightharpoons A_{3}, A_{1} \leftrightharpoons A_{3}, A_{0} \leftrightharpoons A_{3}, A_{1} \leftrightharpoons A_{2}, A_{0} \leftrightharpoons A_{2}, A_{0} \leftrightharpoons A_{1} \tag{2.5}
\end{equation*}
$$

and
(2.6) $\quad A_{3} \leftrightharpoons A_{2}, A_{3} \leftrightharpoons A_{1}, A_{2} \leftrightharpoons A_{1}, A_{3} \leftrightharpoons A_{0}, A_{2} \leftrightharpoons A_{0}, A_{1} \leftrightharpoons A_{0}$.

Remark 2.3. Algorithm 2 is a symmetrized version of Algorithm 1 and can be applied to $n$ substances inductively as in the Algorithm 1.

## 3. Stability and convergence results

Now, we consider a linear reaction system with $(n+1)$ substances:

$$
\frac{d}{d t}\left[\begin{array}{c}
A_{0}(t)  \tag{3.1}\\
A_{1}(t) \\
\vdots \\
A_{n}(t)
\end{array}\right]=M_{n}\left[\begin{array}{c}
A_{0}(t) \\
A_{1}(t) \\
\vdots \\
A_{n}(t)
\end{array}\right],
$$

where the matrix $M_{n}$ with entries $\left\{m_{i j}\right\}_{i, j=1}^{n+1}$ satisfies

$$
m_{i i}<0 \quad \forall i, \quad m_{i j} \geq 0 \quad \forall i \neq j, \quad \sum_{i=1}^{n+1} m_{i j}=0 \quad \forall j .
$$

Remark 3.1. The assumption $\sum_{i} m_{i j}=0$ enforces the concerned reaction system to be conserved. In fact, the requirement of the conserved system can be dropped since the non-conserved system can be converted into the conserved one by adding an artificial substance. In Section 4, we state in more details how to extend the proposed algorithm to the non-conserved systems.

Focusing on the individual reversible reactions between two substances yields a detailed form of $M_{n}$ as follows:

$$
M_{n}=\left[\begin{array}{ccccc}
-\sum_{i=1}^{n} b_{i 0} & f_{10} & \cdots & f_{(n-1) 0} & f_{n 0} \\
b_{10} & -f_{10}-\sum_{i=2}^{n} b_{i 1} & \cdots & f_{(n-1) 1} & f_{n 1} \\
\vdots & \vdots & \ddots & \vdots & \vdots \\
b_{(n-1) 0} & b_{(n-1) 1} & \cdots & -\sum_{j=0}^{n-2} f_{(n-1) j}-b_{n(n-1)} & f_{n(n-1)} \\
b_{n 0} & b_{n 1} & \cdots & b_{n(n-1)} & -\sum_{j=0}^{n-1} f_{n j}
\end{array}\right] .
$$

Note that the linear reaction system (3.1) includes a chain of reversible reactions and circular reactions. We can easily confirm that the real parts of all eigenvalues of $M_{n}$ are nonpositive. Hence the solution for the problem (3.1) is bounded for any $t>0$.

Definition. A numerical method for the problem like (3.1) is said to be stable if it does not blow up as $t \rightarrow \infty$. Furthermore, a numerical method is called absolutely stable if it is stable without any condition on the step size $\Delta t$.

### 3.1. Stability and convergence of $\mathbf{C R}^{2}$ for linear reaction systems

When applying the $\mathrm{CR}^{2}$ algorithm to (3.1), we get the approximate solution $\left(A_{i, k+1}\right)_{i=0}^{n}$ such that

$$
\left[\begin{array}{c}
A_{0, k+1}  \tag{3.2}\\
A_{1, k+1} \\
\vdots \\
A_{n, k+1}
\end{array}\right]=L_{n}\left[\begin{array}{c}
A_{0, k} \\
A_{1, k} \\
\vdots \\
A_{n, k}
\end{array}\right],
$$

where

$$
\begin{equation*}
L_{n}=\left(e^{\Delta t \bar{M}_{n(n-1)}} \cdots e^{\Delta t \bar{M}_{n 0}}\right) \cdots\left(e^{\Delta t \bar{M}_{21}} e^{\Delta t \bar{M}_{20}}\right) e^{\Delta t \bar{M}_{10}} \tag{3.3}
\end{equation*}
$$

and $\bar{M}_{i j}$ is an $(n+1) \times(n+1)$ matrix with entries in (2.4). The expression of the form

$$
\begin{equation*}
\prod_{\substack{i, j=0 \\ i>j}}^{n} \exp \left(\Delta t \bar{M}_{i j}\right) \tag{3.4}
\end{equation*}
$$

formally implies $L_{n}$ in (3.3).
Let $\|\mathbf{x}\|_{1}=\sum_{i=1}^{n}\left|x_{i}\right|$ be a norm on $\mathbb{R}^{n}$.
Theorem 3.2. The algorithm $C R^{2}$ applied to linear reaction systems is absolutely stable.
Proof. The $\mathrm{CR}^{2}$ algorithm gives an approximate solution to (3.1) such that $\mathbf{y}^{k+1}=L_{n} \mathbf{y}^{k}$ with $\mathbf{y}^{k}=\left(A_{0, k} \cdots A_{n, k}\right)^{T}$. Note that each matrix operator $\exp \left(\Delta t \bar{M}_{i j}\right)$ in (3.3) conserves the total concentration of reaction substances since the action of $\exp \left(\Delta t \bar{M}_{i j}\right)$ is characterized by a reversible reaction between two substances in a closed system. Hence, it follows that

$$
\begin{equation*}
\left\|\mathbf{y}^{k+1}\right\|_{1}=\left\|\mathbf{y}^{k}\right\|_{1} \tag{3.5}
\end{equation*}
$$

that is, the approximation does not blow up independently of $\Delta t$.
Remark 3.3. We shall compare the stability of $\mathrm{CR}^{2}$ with that of typical numerical ODE solvers in more details. In order to analyze numerical methods for ODEs, a simple scalar ODE

$$
\begin{equation*}
y^{\prime}(t)=\lambda y(t) \tag{3.6}
\end{equation*}
$$

is commonly-used since it is very helpful in predicting the stability behavior of numerical schemes. In such a typical stability analysis with (3.6), two stability properties, $A$-stability and L-stability are used especially for the stiff problems [3]. The stability function $R(z)$ is a major tool to measure the stability behavior of numerical solutions. But, since our algorithm is first aimed for a numerical solution to a reaction system not a scalar ODE, it is not easy to adopt directly the analysis techniques based on the stability function. Note that the $A$-stability and $L$-stability properties were originally introduced as criteria which show how well numerical solutions mimic important properties of the exact solution in form of exponential function. In this context, the algorithm $\mathrm{CR}^{2}$ might be regarded as an $A$-stable method in the sense that the numerical solution by $\mathrm{CR}^{2}$ is bounded independently of both the step size $\Delta t$ and the spectrum of $M_{n}$. In addition, the L-stability is essential to a numerical method for stiff ODEs, which makes stiff components of a numerical solution damped out rapidly. A numerical method is said to be L-stable if $\lim _{z \rightarrow \infty}|R(z)|=0$. Focusing on the main role of L-stability, it is sufficient to verify how well the proposed
$\mathrm{CR}^{2}$ maintains the rapidly decaying property of the exact solution to the stiff problems. In a reaction system (3.1), the stiffness ratio of the problem is in connection with difference in magnitude of reaction parameters. Assuming that the only one specific pair $A_{j_{s}} \leftrightharpoons A_{i_{s}}$ of reversible reaction reacts very fast relative to the other parts of reactions in (3.1), that is, the problem is stiff, the exact solution includes such a rapidly decaying component. Based on the structure of the $\mathrm{CR}^{2}$ in (3.3), the operator $\exp \left(\Delta t \bar{M}_{i_{s} j_{s}}\right)$ corresponding to the reversible reaction between $A_{j_{s}}$ and $A_{i_{s}}$ makes the approximation by $\mathrm{CR}^{2}$ preserve well the stiff component of the exact solution.

Let $\|A\|_{1}=\max _{1 \leq j \leq n} \sum_{i=1}^{n}\left|a_{i j}\right|$ be a matrix norm on $\mathbb{R}^{n \times n}$.
Theorem 3.4. The algorithm $C R^{2}$ applied to (3.1) has the order of convergence 1.

Proof. We denote the exact solution of (3.1) by $\mathbf{y}(t)=\left(A_{0}(t) \cdots A_{n}(t)\right)^{T}$. The $\mathrm{CR}^{2}$ gives an approximation to (3.1) such that $\mathbf{y}^{k+1}=L_{n} \mathbf{y}^{k}$ with $\mathbf{y}^{k}=$ $\left(A_{0, k} \cdots A_{n, k}\right)^{T}$. We will show that

$$
\left\|\mathbf{y}\left(t_{k+1}\right)-\mathbf{y}^{k+1}\right\|_{1} \leq C \Delta t
$$

where $C$ is a constant independent of $\Delta t$. Using Taylor's theorem yields

$$
\begin{equation*}
\mathbf{y}\left(t_{k+1}\right)=\left(I+\Delta t M_{n}+\frac{(\Delta t)^{2}}{2} M_{n}^{2}\right) \mathbf{y}\left(t_{k}\right)+O\left((\Delta t)^{3}\right) . \tag{3.7}
\end{equation*}
$$

Recalling the definition of the matrix exponential in terms of power series, $L_{n}$ is rewritten as

$$
\begin{aligned}
L_{n} & =\prod_{\substack{i, j=0 \\
i>j}}^{n}\left(I+\Delta t \bar{M}_{i j}+O\left((\Delta t)^{2}\right)\right) \\
& =I+\Delta t M_{n}+(\Delta t)^{2} L_{n, 2}+O\left((\Delta t)^{3}\right)
\end{aligned}
$$

where $\Pi$ is formally used the same as in (3.4) and the second equality is given by the fact that

$$
M_{n}=\sum_{\substack{i, j=0 \\ i>j}}^{n} \bar{M}_{i j} .
$$

Hence, it follows that

$$
\begin{equation*}
\mathbf{y}^{k+1}=\left(I+\Delta t M_{n}+(\Delta t)^{2} L_{n, 2}\right) \mathbf{y}^{k}+O\left((\Delta t)^{3}\right) \tag{3.8}
\end{equation*}
$$

Let $\mathbf{e}_{k}=\mathbf{y}\left(t_{k}\right)-\mathbf{y}^{k}$ denote the numerical error. Subtracting (3.8) from (3.7) gives

$$
\begin{aligned}
\mathbf{e}_{k+1} & =\left(I+\Delta t M_{n}+(\Delta t)^{2} L_{n, 2}\right) \mathbf{e}_{k}+\left(\frac{M_{n}^{2}}{2}-L_{n, 2}\right)(\Delta t)^{2} \mathbf{y}\left(t_{k}\right)+O\left((\Delta t)^{3}\right) \\
& =L_{n} \mathbf{e}_{k}+O\left((\Delta t)^{2}\right)
\end{aligned}
$$

Since $\left\|L_{n}\right\|_{1}=1$ by (3.5), it follows that

$$
\left\|\mathbf{e}_{k+1}\right\|_{1} \leq\left\|\mathbf{e}_{k}\right\|_{1}+C(\Delta t)^{2}
$$

with a constant $C$ independent of $\Delta t$ and $k$. We can check by induction on $k$ that

$$
\left\|\mathbf{e}_{k}\right\|_{1} \leq k C(\Delta t)^{2} \quad \forall k=0,1, \ldots
$$

If we restrict $t$ to the finite interval $\left[0, t_{F}\right]$, then we have that

$$
\begin{aligned}
\left\|\mathbf{e}_{k}\right\|_{1} & \leq\left\lfloor t_{F} / \Delta t\right\rfloor C(\Delta t)^{2} \\
& \leq C t_{F} \Delta t, \quad \forall k=0,1, \ldots,\left\lfloor t_{F} / \Delta t\right\rfloor
\end{aligned}
$$

Since $C$ is independent of $\Delta t$, it follows that

$$
\lim _{\substack{\Delta t \rightarrow 0 \\ 0 \leq k \leq\left\lfloor t_{F} / \Delta t\right\rfloor}}\left\|\mathbf{e}_{k}\right\|_{1}=0 .
$$

In other words, the algorithm $\mathrm{CR}^{2}$ is convergent and has the order of convergence 1.

### 3.2. Convergence of $\mathrm{SCR}^{2}$ for linear reaction systems

Applying the $\mathrm{SCR}^{2}$ to (3.1) gives

$$
\left[\begin{array}{c}
A_{0, k+1}^{s_{1}} \\
A_{1, k+1}^{s_{1}} \\
\vdots \\
A_{n, k+1}^{s_{1}}
\end{array}\right]=L_{n}^{s_{1}}\left[\begin{array}{c}
A_{0, k}^{s} \\
A_{1, k}^{s} \\
\vdots \\
A_{n, k}^{s}
\end{array}\right], \quad\left[\begin{array}{c}
A_{0, k+1}^{s_{2}} \\
A_{1, k+1}^{s_{2}^{2}} \\
\vdots \\
A_{n, k+1}^{s_{2}}
\end{array}\right]=L_{n}^{s_{2}}\left[\begin{array}{c}
A_{0, k}^{s} \\
A_{1, k}^{s} \\
\vdots \\
A_{n, k}^{s}
\end{array}\right] .
$$

Then, we have

$$
\left[\begin{array}{c}
A_{0, k+1}^{s}  \tag{3.9}\\
A_{1, k+1}^{s} \\
\vdots \\
A_{n, k+1}^{s}
\end{array}\right]=L_{n}^{s}\left[\begin{array}{c}
A_{0, k}^{s} \\
A_{1, k}^{s} \\
\vdots \\
A_{n, k}^{s}
\end{array}\right]=\frac{1}{2}\left(L_{n}^{s_{1}}+L_{n}^{s_{2}}\right)\left[\begin{array}{c}
A_{0, k}^{s} \\
A_{1, k}^{s} \\
\vdots \\
A_{n, k}^{s}
\end{array}\right]
$$

As mentioned in Remark 2.2, for reaction systems with more than three substances, there are two possible ways to choose the reverse order used in implementing $\mathrm{SCR}^{2}$. Note that how to reverse the order of reactions makes a difference in the relationship between $L_{n}^{s_{1}}$ and $L_{n}^{s_{2}}$. Hence, we shall analyze the algorithm $\mathrm{SCR}^{2}$ in two different ways.

First, let us handle the algorithm with the reverse order (2.5).
Lemma 3.5. If the algorithm $S C R^{2}$ is in company with the reverse order (2.5), the matrix $L_{n}^{s}$ in (3.9) has the Taylor expansion

$$
\begin{aligned}
L_{n}^{s} & =L_{n, 0}^{s}+L_{n, 1}^{s} \Delta t+L_{n, 2}^{s}(\Delta t)^{2}+O\left((\Delta t)^{3}\right) \\
& =I+M_{n} \Delta t+\frac{1}{2} M_{n}^{2}(\Delta t)^{2}+O\left((\Delta t)^{3}\right),
\end{aligned}
$$

where $M_{n}$ is the matrix in (3.1).

Proof. Since $\mathrm{SCR}^{2}$ is a symmetrized version of $\mathrm{CR}^{2}$, each $L_{n}^{s_{i}}$ has the same property as $L_{n}$ in $\mathrm{CR}^{2}$. Thus, it is obvious that

$$
L_{n, 0}^{s}=I, \quad L_{n, 1}^{s}=M_{n} .
$$

For a linear reaction system with $(n+1)$ substances, recall

$$
L_{n}^{s_{1}}=\left(e^{\Delta t \bar{M}_{n(n-1)}} \cdots e^{\Delta t \bar{M}_{n 0}}\right) \cdots\left(e^{\Delta t \bar{M}_{21}} e^{\Delta t \bar{M}_{20}}\right) e^{\Delta t \bar{M}_{10}}
$$

According to the reverse order in (2.5), we have

$$
L_{n}^{s_{2}}=e^{\Delta t \bar{M}_{10}}\left(e^{\Delta t \bar{M}_{20}} e^{\Delta t \bar{M}_{21}}\right) \cdots\left(e^{\Delta t \bar{M}_{n 0}} \cdots e^{\Delta t \bar{M}_{n(n-1)}}\right) .
$$

Based on the fact that

$$
M_{n}=\sum_{\substack{i, j=0 \\ i>j}}^{n} \bar{M}_{i j}
$$

the power series expansion of each $L_{n}^{s_{i}}$ yields

$$
L_{n, 2}^{s}=\frac{1}{2} M_{n}^{2} .
$$

Next, we consider the $\mathrm{SCR}^{2}$ equipped with the reverse order (2.6), where the reaction substances are put in reverse order. In order to look into the procedure of the $\mathrm{SCR}^{2}$ algorithm applied to a reaction system (3.1), we first consider a circular reaction with 3 substances shown in Figure 1. Two parallel solution steps in $\mathrm{SCR}^{2}$ applied to the reaction system is represented as

$$
\left[\begin{array}{l}
A_{0, k+1}^{s_{1}} \\
A_{1, k+1}^{s_{1}} \\
A_{2, k+1}^{s_{1}}
\end{array}\right]=L_{2}^{s_{1}}\left[\begin{array}{l}
A_{0, k}^{s} \\
A_{1, k}^{s} \\
A_{2, k}^{s}
\end{array}\right], \quad\left[\begin{array}{c}
A_{2, k+1}^{s_{2}} \\
A_{1, k+1}^{s_{2}} \\
A_{0, k+1}^{s_{2}}
\end{array}\right]=\bar{L}_{2}^{s_{2}}\left[\begin{array}{c}
A_{2, k}^{s} \\
A_{1, k}^{s} \\
A_{0, k}^{s}
\end{array}\right] .
$$

Then, we have

$$
\left[\begin{array}{c}
A_{0, k+1}^{s} \\
A_{1, k+1}^{s} \\
A_{2, k+1}^{s}
\end{array}\right]=\frac{1}{2}\left(L_{2}^{s_{1}}+L_{2}^{s_{2}}\right)\left[\begin{array}{c}
A_{0, k}^{s} \\
A_{1, k}^{s} \\
A_{2, k}^{s}
\end{array}\right],
$$

where

$$
L_{2}^{s_{2}}=P_{2} \bar{L}_{2}^{s_{2}} P_{2}, \quad P_{2}=\left[\begin{array}{lll}
0 & 0 & 1 \\
0 & 1 & 0 \\
1 & 0 & 0
\end{array}\right] .
$$

Considering Taylor expansions of $L_{2}^{s_{1}}$ and $\bar{L}_{2}^{s_{2}}$ yields

$$
L_{2,2}^{s_{1}}-\frac{1}{2} M_{2}^{2}=\frac{1}{2}\left[\begin{array}{cc}
0 & -b_{20} f_{10}+b_{21} f_{10}-b_{21} f_{20}  \tag{array}\\
b_{10} b_{20}-b_{10} b_{21}+b_{20} f_{21} & 0 \\
-b_{10} b_{20}+b_{10} b_{21}-b_{20} f_{21} & b_{20} f_{10}-b_{21} f_{10}+b_{21} f_{20}
\end{array}\right.
$$

and
$\bar{L}_{2,2}^{s_{2}}-\frac{1}{2} P_{2} M_{2}^{2} P_{2}=\frac{1}{2}\left[\begin{array}{ccc}0 & -b_{20} f_{10}+b_{21} f_{10}-b_{21} f_{20} & \text { (III) } \\ b_{10} f_{20}-f_{10} f_{21}+f_{20} f_{21} & \text { (IV) } \\ -b_{10} f_{20}+f_{10} f_{21}-f_{20} f_{21} & b_{20} f_{10}-b_{21} f_{10}+b_{21} f_{20} & 0\end{array}\right]$,
where

$$
\begin{array}{ll}
\text { (I) }=b_{10} f_{20}-f_{10} f_{21}+f_{20} f_{21} & \text { (II) }=-b_{10} f_{20}+f_{10} f_{21}-f_{20} f_{21} \\
\text { (III) })=b_{10} b_{20}-b_{10} b_{21}+b_{20} f_{21} & \text { (IV) }=-b_{10} b_{20}+b_{10} b_{21}-b_{20} f_{21} .
\end{array}
$$

Then, we have

$$
\begin{equation*}
L_{2,2}^{s_{1}}+L_{2,2}^{s_{2}}-M_{2}^{2}=0 \tag{3.10}
\end{equation*}
$$

In general, the approximation $\left(A_{i, k+1}^{s}\right)_{i=0}^{n}$ obtained by $\mathrm{SCR}^{2}$ applied to (3.1) is characterized as

$$
\left[\begin{array}{c}
A_{0, k+1}^{s}  \tag{3.11}\\
A_{1, k+1}^{s} \\
\vdots \\
A_{n, k+1}^{s}
\end{array}\right]=L_{n}^{s}\left[\begin{array}{c}
A_{0, k}^{s} \\
A_{1, k}^{s} \\
\vdots \\
A_{n, k}^{s}
\end{array}\right]=\frac{1}{2}\left(L_{n}^{s_{1}}+L_{n}^{s_{2}}\right)\left[\begin{array}{c}
A_{0, k}^{s} \\
A_{1, k}^{s} \\
\vdots \\
A_{n, k}^{s}
\end{array}\right]
$$

where

$$
L_{n}^{s_{2}}=P_{n} \bar{L}_{n}^{s_{2}} P_{n}
$$

and $P_{n}$ is the $(n+1) \times(n+1)$ permutation matrix such that

$$
\left(P_{n}\right)_{i j}= \begin{cases}1 & \text { if } i+j=n+2 \\ 0 & \text { otherwise }\end{cases}
$$

Lemma 3.6. Consider a parallel reaction with $(n+1)$ substances, that is composed of $n$ reversible reactions:

$$
\begin{array}{ccc}
A_{n} & \stackrel{b_{n 0}}{\stackrel{f_{n 0}}{\rightleftharpoons}} A_{0} \\
A_{n} & \stackrel{f_{n 1}}{\rightleftharpoons} & A_{1} \\
& \vdots  \tag{3.12c}\\
& \stackrel{b_{n 1}}{\rightleftharpoons} & \\
A_{n} & \stackrel{f_{n(n-1)}}{\rightleftharpoons} & A_{n(n-1)}
\end{array} A_{n-1} .
$$

Applying the $C R^{2}$ algorithm to this problem in the following ordering:

$$
A_{0} \leftrightharpoons A_{n}, A_{1} \leftrightharpoons A_{n}, \ldots, A_{n-1} \leftrightharpoons A_{n}
$$

gives

$$
\left[\begin{array}{c}
A_{0, k+1} \\
A_{1, k+1} \\
\vdots \\
A_{n, k+1}
\end{array}\right]=\Lambda_{n}\left[\begin{array}{c}
A_{0, k} \\
A_{1, k} \\
\vdots \\
A_{n, k}
\end{array}\right],
$$

where $\left(A_{i, k+1}\right)_{i=0}^{n}$ is an approximation to $\left(A_{i}(t)\right)_{i=0}^{n}$ at $t=t_{k+1}$ and $\Lambda_{n}$ is an $(n+1) \times(n+1)$ matrix. Then, we have
$\Lambda_{n}=\left[\begin{array}{cccc} & & & f_{n 0} \Delta t \\ & \operatorname{diag}\left(1-b_{n j} \Delta t\right)_{j=0}^{n-1} & & \vdots \\ b_{n 0} \Delta t & \ldots & b_{n(n-1)} \Delta t & 1-\left(\sum_{j=0}^{n-1} f_{n j}\right) \Delta t\end{array}\right]+O\left((\Delta t)^{2}\right)$.
Proof. It is easy to show that the statement holds for $n=2$ by using Taylor's theorem. Suppose that the statement is true for a parallel reaction which consists of $n$ substances. Consider a parallel reaction with $(n+1)$ substances. By the inductive hypothesis, the approximate solution $\left(A_{i, k+1}\right)_{i=0}^{n}$ to the solution $\left(A_{i}(t)\right)_{i=0}^{n}$ at $t=t_{k+1}$ obtained by applying $\mathrm{CR}^{2}$ to (3.12) is written as follows:

$$
\left[\begin{array}{c}
A_{0, k+1}  \tag{3.13a}\\
A_{n}^{t e m p}
\end{array}\right]=\left[\begin{array}{cc}
1-b_{n 0} \Delta t+O\left((\Delta t)^{2}\right) & f_{n 0} \Delta t+O\left((\Delta t)^{2}\right) \\
b_{n 0} \Delta t+O\left((\Delta t)^{2}\right) & 1-f_{n 0} \Delta t+O\left((\Delta t)^{2}\right)
\end{array}\right]\left[\begin{array}{c}
A_{0, k} \\
A_{n, k}
\end{array}\right]
$$

$$
\left[\begin{array}{c}
A_{1, k+1} \\
\vdots \\
A_{n-1, k+1} \\
A_{n, k+1}
\end{array}\right]=\Lambda_{n-1}\left[\begin{array}{c}
A_{1, k} \\
\vdots \\
A_{n-1, k} \\
A_{n}^{\text {temp }}
\end{array}\right]
$$

where
(3.14)

$$
\left.\Lambda_{n-1}=\left[\begin{array}{cccc} 
& & & f_{n 1} \Delta t \\
& \operatorname{diag}\left(1-b_{n j} \Delta t\right)_{j=1}^{n-1} & & \vdots \\
b_{n 1} \Delta t & \ldots & & b_{n(n-1)} \Delta t
\end{array}\right]+\left(\sum_{n(n-1)}^{n-1} f_{n j}\right) \Delta t\right]\left[O\left((\Delta t)^{2}\right)\right.
$$

By (3.13), we have

$$
\left.\Lambda_{n}=\left[\begin{array}{cccc}
1 & 0 & \cdots & 0  \tag{3.15}\\
0 & & & \\
\vdots & & \Lambda_{n-1} \\
0 & & & 0
\end{array}\right]\left[\begin{array}{cccc}
1-b_{n 0} \Delta t+O\left((\Delta t)^{2}\right) & \cdots & 0 & f_{n 0} \Delta t+O\left((\Delta t)^{2}\right) \\
0 & 1 & \cdots & 0
\end{array}\right) 0 \begin{array}{c}
\vdots \\
\vdots \\
0
\end{array} \begin{array}{c}
0 \\
0
\end{array}\right]
$$

By substituting (3.14) into (3.15), we have the expansion form of $\Lambda_{n}$, that is, the statement is true for a parallel reaction with $(n+1)$ substances.

Lemma 3.7. Under the same assumptions as in Lemma 3.6, we assume that

$$
\Lambda_{n}=\Lambda_{n, 0}+\Lambda_{n, 1} \Delta t+\Lambda_{n, 2}(\Delta t)^{2}+O\left((\Delta t)^{3}\right)
$$

Then,

$$
\Lambda_{n, 2}=\left[\begin{array}{cccccc}
\frac{b_{n 0}}{2}\left(b_{n 0}+f_{n 0}\right) & 0 & 0 & \cdots & 0 & \vdots \\
b_{n 0} f_{n 1} & \frac{b_{n 1}}{2}\left(b_{n 1}+f_{n 1}\right) & 0 & \ddots & \vdots & \vdots \\
b_{n 0} f_{n 2} & b_{n 1} f_{n 2} & \ddots & \ddots & 0 & \text { (I) } \\
\vdots & \vdots & & \ddots & \\
b_{n 0} f_{n(n-1)} & b_{n 1} f_{n(n-1)} & \cdots & \cdots & \frac{b_{n(n-1)}^{2}}{b_{n}}\left(b_{n(n-1)}+f_{n(n-1)}\right) & \text { (III) } \\
(\mathrm{IV}) & (\mathrm{V}) & \cdots & \cdots & -\frac{b_{n(n-1)}^{2}}{2}\left(b_{n(n-1)}+f_{n(n-1)}\right) & \text { (VI)}
\end{array}\right]
$$

where
$(\mathrm{I})=-\frac{f_{n 0}}{2}\left(b_{n 0}+f_{n 0}\right)$
$(\mathrm{II})=-\frac{f_{n 1}}{2}\left(2 f_{n 0}+b_{n 1}+f_{n 1}\right)$
$(\mathrm{III})=-\frac{f_{n(n-1)}}{2}\left(2 \sum_{i=0}^{n-2} f_{n i}+b_{n(n-1)}+f_{n(n-1)}\right)$
$(\mathrm{IV})=-\frac{b_{n 0}}{2}\left(b_{n 0}+f_{n 0}+2 \sum_{i=1}^{n-1} f_{n i}\right)$
$(\mathrm{V})=-\frac{b_{n 1}}{2}\left(b_{n 1}+f_{n 1}+2 \sum_{i=2}^{n-2} f_{n i}\right)$
$(\mathrm{VI})=\frac{1}{2} \sum_{i=0}^{n-1} f_{n i}\left(b_{n i}+f_{n i}\right)+\sum_{i<j} f_{n i} f_{n j}$.

Proof. By the Taylor expansion of $\Lambda_{2}$, the statement holds for $n=2$. Suppose that the statement is true for $\Lambda_{n-1,2}$. By Lemma 3.6, we have

$$
\Lambda_{n-1}=I+\left[\begin{array}{cccc}
-b_{n 1} & & & f_{n 1} \\
& \ddots & & \vdots \\
& & -b_{n(n-1)} & f_{n(n-1)} \\
b_{n 1} & \cdots & b_{n(n-1)} & -\sum_{i=1}^{n-1} f_{n i}
\end{array}\right] \Delta t+\Lambda_{n-1,2}(\Delta t)^{2}+O\left((\Delta t)^{3}\right) .
$$

It follows

$$
\begin{aligned}
& \Lambda_{n, 2}=\left[\begin{array}{cc}
0 & 0 \\
0 & \Lambda_{n-1,2}
\end{array}\right] \\
& +\left[\begin{array}{ccccc}
0 & 0 & \cdots & 0 & 0 \\
0 & -b_{n 1} & & & f_{n 1} \\
\vdots & & \ddots & & \vdots \\
0 & & & -b_{n(n-1)} & f_{n(n-1)} \\
0 & b_{n 1} & \cdots & b_{n(n-1)} & -\sum_{i=1}^{n-1} f_{n i}
\end{array}\right]\left[\begin{array}{ccccc}
-b_{n 0} & & & & f_{n 0} \\
& 0 & & & \\
& & \ddots & \\
& & & 0 & \\
b_{n 0} & & & & -f_{n 0}
\end{array}\right] \\
& +\frac{1}{2}\left[\begin{array}{ccccc}
b_{n 0}\left(b_{n 0}+f_{n 0}\right) & & & & -f_{n 0}\left(b_{n 0}+f_{n 0}\right) \\
& 0 & & & \\
& & \ddots & & \\
-b_{n 0}\left(b_{n 0}+f_{n 0}\right) & & & 0 & f_{n 0}\left(b_{n 0}+f_{n 0}\right)
\end{array}\right] \\
& =\left[\begin{array}{cc}
0 & 0 \\
0 & \Lambda_{n-1,2}
\end{array}\right]+\left[\begin{array}{ccccc}
\frac{b_{n 0}}{2}\left(b_{n 0}+f_{n 0}\right) & 0 & \cdots & 0 & -\frac{f_{n 0}}{2}\left(b_{n 0}+f_{n 0}\right) \\
b_{n 0} f_{n 1} & 0 & \cdots & 0 & -f_{n 0} f_{n 1} \\
\vdots & \vdots & \ddots & \vdots & \vdots \\
b_{n 0} f_{n(n-1)} & 0 & \cdots & 0 & -f_{n 0} f_{n(n-1)} \\
(\mathrm{I}) & 0 & \cdots & 0 & (\mathrm{II})
\end{array}\right],
\end{aligned}
$$

where
$(\mathrm{I})=-\frac{b_{n 0}}{2}\left(b_{n 0}+f_{n 0}+2 \sum_{i=1}^{n-1} f_{n i}\right)$,
$(\mathrm{II})=\frac{f_{n 0}}{2}\left(b_{n 0}+f_{n 0}+2 \sum_{i=1}^{n-1} f_{n i}\right)$.

Hence, by the inductive hypothesis, we conclude that the statement is true for $\Lambda_{n, 2}$.

Lemma 3.8. If the algorithm $S C R^{2}$ is in company with the reverse order (2.6), the matrix $L_{n}^{s}$ in (3.11) has the Taylor expansion

$$
\begin{aligned}
L_{n}^{s} & =L_{n, 0}^{s}+L_{n, 1}^{s} \Delta t+L_{n, 2}^{s}(\Delta t)^{2}+O\left((\Delta t)^{3}\right) \\
& =I+M_{n} \Delta t+\frac{1}{2} M_{n}^{2}(\Delta t)^{2}+O\left((\Delta t)^{3}\right),
\end{aligned}
$$

where $M_{n}$ is the matrix in (3.1).
Proof. As mentioned in the proof of Lemma 3.5, it is obvious that

$$
L_{n, 0}^{s}=I, \quad L_{n, 1}^{s}=M_{n}
$$

Let $D_{n}^{s_{i}}=L_{n, 2}^{s_{i}}-\frac{1}{2} M_{n}^{2}$ for $i=1,2$. To complete the proof, it suffices to show that

$$
D_{n}^{s_{1}}+D_{n}^{s_{2}}=0
$$

Because $D_{n}^{s_{2}}$ can be associated with $D_{n}^{s_{1}}$ by the permutation matrix $P_{n}$, we only need to characterize entries of $D_{n}^{s_{1}}$.

We first claim that

$$
\begin{equation*}
D_{n}^{s_{1}}=\sum_{i<j<k} D_{(i, j, k)}, \quad 0 \leq i, j, k \leq n, \tag{3.16}
\end{equation*}
$$

where

$$
D_{(i, j, k)}
$$

$$
=\begin{array}{ccc}
(i+1) \text { th } & (j+1) \text { th } & (k+1) \text { th } \\
\downarrow & \downarrow & \downarrow \\
0 & -b_{k i} f_{j i}+b_{k j} f_{j i}-b_{k j} f_{k i} & \text { (I) } \\
\frac{1}{2}\left[\begin{array}{ccc} 
\\
\frac{1}{2} \\
b_{j i} b_{k i}-b_{j i} b_{k j}+b_{k i} f_{k j} & (\text { II }) \\
-b_{j i} b_{k i}+b_{j i} b_{k j}-b_{k i} f_{k j} & b_{k i} f_{j i}-b_{k j} f_{j i}+b_{k j} f_{k i} & 0
\end{array}\right] & \leftarrow(i+1) \text { th } \\
\leftarrow(j+1) \text { th } \\
\leftarrow(k+1) \text { th }
\end{array}
$$

with

$$
\text { (I) }=b_{j i} f_{k i}-f_{j i} f_{k j}+f_{k i} f_{k j}, \quad \text { (II) }=-b_{j i} f_{k i}+f_{j i} f_{k j}-f_{k i} f_{k j}
$$

and other entries are zeros. The claim is verified by induction on $n$. By the previous argument, we already know that

$$
D_{2}^{s_{1}}=D_{(0,1,2)}
$$

We assume that the claim is true for a general reaction system with $n$ substances $A_{0}, A_{1}, \ldots, A_{n-1}$, i.e.,

$$
D_{n-1}^{s_{1}}=\sum_{i<j<k} D_{(i, j, k)}, \quad 0 \leq i, j, k \leq n-1
$$

Let us see how $L_{n}^{s_{1}}$ and $M_{n}$ change when adding $A_{n}$ into the general reaction system with $n$ substances. Since

$$
\left.\begin{array}{rl}
L_{n}^{s_{1}} & =\Lambda_{n}\left[\begin{array}{cc}
L_{n-1}^{s_{1}} & 0 \\
0 & 1
\end{array}\right] \\
& =\left(I+\Lambda_{n, 1} \Delta t+\Lambda_{n, 2}(\Delta t)^{2}+O\left((\Delta t)^{3}\right)\left[\begin{array}{c}
I+M_{n-1} \Delta t+L_{n-1,2}^{s_{1}}(\Delta t)^{2}+O\left((\Delta t)^{3}\right) \\
0
\end{array} 0\right.\right. \\
& 0
\end{array}\right],
$$

it follows that

$$
L_{n, 2}^{s_{1}}=\left[\begin{array}{cc}
L_{n-1,2}^{s_{1}} & 0 \\
0 & 0
\end{array}\right]+\Lambda_{n, 1}\left[\begin{array}{cc}
M_{n-1} & 0 \\
0 & 0
\end{array}\right]+\Lambda_{n, 2}
$$

Also, noting the fact that

$$
M_{n}=\left[\begin{array}{cccc} 
& & & f_{n 0} \\
& M_{n-1}-\operatorname{diag}\left(b_{n j}\right)_{j=0}^{n-1} & & \vdots \\
b_{n 0} & \ldots & b_{n(n-1)} & -\sum_{j=0}^{n-1} f_{n j}
\end{array}\right],
$$

Lemma 3.6 gives

$$
\begin{aligned}
M_{n}^{2} & =\left(\left[\begin{array}{cc}
M_{n-1} & 0 \\
0 & 0
\end{array}\right]+\Lambda_{n, 1}\right)^{2} \\
& =\left[\begin{array}{cc}
M_{n-1}^{2} & 0 \\
0 & 0
\end{array}\right]+\left[\begin{array}{cc}
M_{n-1} & 0 \\
0 & 0
\end{array}\right] \Lambda_{n, 1}+\Lambda_{n, 1}\left[\begin{array}{cc}
M_{n-1} & 0 \\
0 & 0
\end{array}\right]+\Lambda_{n, 1}^{2} .
\end{aligned}
$$

Then, we have that

$$
\begin{aligned}
D_{n}^{s_{1}} & =L_{n, 2}^{s_{1}}-\frac{1}{2} M_{n}^{2} \\
& =\left[\begin{array}{cc}
L_{n-1,2}^{s_{1}}-\frac{1}{2} M_{n-1}^{2} & 0 \\
0 & 0
\end{array}\right]+\frac{1}{2} \Lambda_{n, 1}\left[\begin{array}{cc}
M_{n-1} & 0 \\
0 & 0
\end{array}\right]-\frac{1}{2}\left[\begin{array}{cc}
M_{n-1} & 0 \\
0 & 0
\end{array}\right] \Lambda_{n, 1}+\Lambda_{n, 2}-\frac{1}{2} \Lambda_{n, 1}^{2} .
\end{aligned}
$$

Looking at $M_{n-1}, \Lambda_{n, 1}$ and $\Lambda_{n, 2}$ carefully, all $(i, j)$ entries contain subindex $n,(i-1)$ or $(j-1)$. It implies that we can write
$\frac{1}{2} \Lambda_{n, 1}\left[\begin{array}{cc}M_{n-1} & 0 \\ 0 & 0\end{array}\right]-\frac{1}{2}\left[\begin{array}{cc}M_{n-1} & 0 \\ 0 & 0\end{array}\right] \Lambda_{n, 1}+\Lambda_{n, 2}-\frac{1}{2} \Lambda_{n, 1}^{2}=\sum_{i<j} B_{i j n} \quad 0 \leq i, j<n$,
where $B_{i j n}$ is a matrix whose entries consist of products of $b_{n i}, f_{n i}, b_{n j}, f_{n j}, b_{j i}$, and $f_{j i}$. Now let us compute $B_{i j n}$ in details. For a matrix $M$, we define $(M)_{i j n}$ by a matrix whose entries consist of products of $b_{n i}, f_{n i}, b_{n j}, f_{n j}, b_{j i}$ and $f_{j i}$ among entries of $M$. We detail $B_{i j n}$ for all $i, j$ with $0 \leq i, j<n$ as follows:

$$
\begin{aligned}
& E_{1}=\left(\Lambda_{n, 1}\left[\begin{array}{cc}
M_{n-1} & 0 \\
0 & 0
\end{array}\right]\right)_{i j n} \\
&(i+1) \text { th }(j+1) \text { th } \\
& \downarrow \downarrow \\
&(n+1) \text { th } \\
&=\left[\begin{array}{ccc}
b_{j i} b_{n i} & -b_{n i} f_{j i} & \downarrow \\
-b_{j i} b_{n j} & b_{n j} f_{j i} & 0 \\
b_{j i}\left(b_{n j}-b_{n i}\right) & f_{j i}\left(b_{n i}-b_{n j}\right) & 0
\end{array}\right] \leftarrow(i+1) \text { th } \\
& \leftarrow(j+1) \text { th } \\
& \leftarrow(n+1) \text { th }
\end{aligned}
$$

$$
\begin{aligned}
& E_{2}=\left(\left[\begin{array}{cc}
M_{n-1} & 0 \\
0 & 0
\end{array}\right] \Lambda_{n, 1}\right)_{i j n} \\
& \left.=\begin{array}{ccc}
(i+1) \text { th } & (j+1) \text { th } & (n+1) \text { th } \\
\downarrow & \downarrow & \downarrow \\
b_{j i} b_{n i} & -f_{j i} b_{n j} & -b_{j i} f_{n i}+f_{j i} f_{n j} \\
-b_{j i} b_{n i} & f_{j i} b_{n j} & b_{j i} f_{n i}-f_{j i} f_{n j} \\
0 & 0 & 0
\end{array}\right] \leftarrow(i+1) \text { th } \\
& E_{3}=\left(\Lambda_{n, 2}\right)_{i j n} \\
& =\left[\begin{array}{ccc}
(i+1) \text { th } & (j+1) \text { th } & (n+1) \text { th } \\
\downarrow & \downarrow & \downarrow \\
{\left[\begin{array}{cc}
2(n-1) \\
b_{n i}\left(b_{n i}+f_{n i}\right) & 0 \\
b_{n i} f_{n j} & \frac{1}{2(n-1)} b_{n j}\left(b_{n j}+f_{n j}\right)
\end{array}\right.} & \text { (II) } \\
-\frac{1}{2(n-1)} b_{n i}\left(b_{n i}+f_{n i}\right)-b_{n i} f_{n j} & -\frac{1}{2(n-1)} b_{n j}\left(b_{n j}+f_{n j}\right) & (\text { III })
\end{array}\right] \leftarrow(i+1) \text { th } \begin{array}{c}
\leftarrow(j+1) \text { th } \\
\leftarrow(n+1) \text { th }
\end{array} \\
& \text { (I) }=-\frac{1}{2(n-1)} f_{n i}\left(b_{n i}+f_{n i}\right) \\
& \text { (II) }=-\frac{1}{2(n-1)} f_{n j}\left(b_{n j}+f_{n j}\right)-f_{n i} f_{n j} \\
& (\mathrm{III})=\frac{1}{2(n-1)} f_{n i}\left(b_{n i}+f_{n i}\right)+\frac{1}{2(n-1)} f_{n j}\left(b_{n j}+f_{n j}\right)+f_{n i} f_{n j} .
\end{aligned}
$$

Note that $(n-1)$ appears in the denominators. In the $(i+1, i+1)$ entry, $\frac{1}{2} b_{n i}\left(b_{n i}+f_{n i}\right)$ is independent of $j$. Therefore, we distribute it equally to all $\left(\Lambda_{n, 2}\right)_{i j n}$ of which $(i+1)$ th column is nonzero. In other entries, $(n-1)$ appears with the same reason.

$$
\begin{aligned}
E_{4} & =\left(\left(\Lambda_{n, 1}\right)^{2}\right)_{i j n} \\
& =\left[\begin{array}{ccc}
(i+1) \text { th } & (j+1) \text { th } & (n+1) \text { th } \\
\downarrow & \downarrow & \downarrow \\
\frac{1}{(n-1)} b_{n i}\left(b_{n i}+f_{n i}\right) & b_{n j} f_{n i} & \text { (I) } \\
b_{n i} f_{n j} & \frac{1}{(n-1)} b_{n j}\left(b_{n j}+f_{n j}\right) & \text { (II) } \\
-\frac{1}{(n-1)} b_{n i}\left(b_{n i}+f_{n i}\right)-b_{n i} f_{n j} & (\mathrm{III}) & \leftarrow(i+1) \text { th } \\
& (\mathrm{I})=-\frac{1}{(n-1)} f_{n i}\left(b_{n i}+f_{n i}\right)-f_{n i} f_{n j} & \leftarrow(n+1) \text { th } \\
& (\mathrm{II})=-\frac{1}{(n-1)} f_{n j}\left(b_{n j}+f_{n j}\right)-f_{n i} f_{n j} \\
& (\mathrm{III})=-\frac{1}{(n-1)} b_{n j}\left(b_{n j}+f_{n j}\right)-b_{n j} f_{n i} \\
& (\mathrm{IV})=\frac{1}{(n-1)} f_{n i}\left(b_{n i}+f_{n i}\right)+\frac{1}{(n-1)} f_{n j}\left(b_{n j}+f_{n j}\right)+2 f_{n i} f_{n j},
\end{array}\right.
\end{aligned}
$$

where for each $E_{l}, E_{l}\left(m_{1}, m_{2}\right)=0 \quad \forall m_{1}, m_{2} \notin\{i+1, j+1, n+1\}$. Then, we have

$$
\begin{aligned}
& B_{i j n}=\frac{1}{2} E_{1}-\frac{1}{2} E_{2}+E_{3}-\frac{1}{2} E_{4}
\end{aligned}
$$

$$
\begin{aligned}
& =D_{(i, j, n)},
\end{aligned}
$$

where

$$
(\mathrm{I})=b_{j i} f_{n i}-f_{j i} f_{n j}+f_{n i} f_{n j}, \quad(\mathrm{II})=-b_{j i} f_{n i}+f_{j i} f_{n j}-f_{n i} f_{n j}
$$

The inductive hypothesis yields

$$
\begin{aligned}
D_{n}^{s_{1}} & =\sum_{i<j<k<n} D_{(i, j, k)}+\sum_{i<j<n} D_{(i, j, n)} \\
& =\sum_{i<j<k} D_{(i, j, k)} .
\end{aligned}
$$

Similarly, we obtain the relationship

$$
\bar{L}_{n, 2}^{s_{2}}-\frac{1}{2} P_{n} M_{n}^{2} P_{n}=\sum_{i<j<k} D_{(n-k, n-j, n-i)} .
$$

Note that

$$
D_{(i, j, k)}+P_{n} D_{(n-k, n-j, n-i)} P_{n}=0 .
$$

Consequently, we find that

$$
\begin{aligned}
D_{n}^{s_{2}} & =L_{n, 2}^{s_{2}}-\frac{1}{2} M_{n}^{2} \\
& =P_{n}\left(\bar{L}_{n, 2}^{s_{2}}-\frac{1}{2} P_{n} M_{n}^{2} P_{n}\right) P_{n} \\
& =P_{n}\left(\sum_{i<j<k} D_{(n-k, n-j, n-i)}\right) P_{n} \\
& =-\sum_{i<j<k} D_{(i, j, k)},
\end{aligned}
$$

that is,

$$
D_{n}^{s_{1}}+D_{n}^{s_{2}}=0
$$

The proof is complete.
By using the same arguments as in Theorem 3.4 based on Lemma 3.5 and Lemma 3.8, we can easily get the following theorem.

Theorem 3.9. The algorithm $S C R^{2}$ applied to (3.1) has the order of convergence 2, which is in company with the reverse order either (2.5) or (2.6).

## 4. Numerical experiments

In this section, we present numerical results which verify the theoretical results in the previous sections. In addition, we provide an example of the non-conserved system which can be converted to the conserved one.

Figure 2-(a) depicts a circular reaction with three substances $A, B$ and $C$ with initial values $A(0)=1, B(0)=2$ and $C(0)=3$. The system of ODEs

(a)

$f_{3}$

(b)

Figure 2. (a) A circular reaction (b) A general reaction system


Figure 3. Stability of the $\mathrm{CR}^{2}$ algorithm

TABLE 1. Convergence behaviors of $\mathrm{CR}^{2}, \mathrm{SCR}^{2}$ and BDF2 by means of $\left\|\mathbf{y}\left(t_{F}\right)-\mathbf{y}^{N}\right\|_{1}$ and $\sum_{i=1}^{N}\left\|\mathbf{y}\left(t_{i}\right)-\mathbf{y}^{i}\right\|_{1} \Delta t$ with $\Delta t$ varying from $10^{-1}$ to $10^{-6}$ where $N=t_{F} / \Delta t$ with $t_{F}=3$

| $\Delta t$ | $\left\\|\mathbf{y}\left(t_{F}\right)-\mathbf{y}^{N}\right\\|_{1}$ |  |  | $\sum_{i=1}^{N}\left\\|\mathbf{y}\left(t_{i}\right)-\mathbf{y}^{i}\right\\|_{1} \Delta t$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | CR $^{2}$ | SCR $^{2}$ | BDF2 | CR $^{2}$ | SCR $^{2}$ | BDF2 |
| $10^{-1}$ | $3.4182 \mathrm{e}-01$ | $1.6979 \mathrm{e}-01$ | $3.0723 \mathrm{e}-12$ | 1.0295 | $5.2196 \mathrm{e}-01$ | $6.4610 \mathrm{e}-02$ |
| $10^{-2}$ | $3.2857 \mathrm{e}-02$ | $1.4643 \mathrm{e}-02$ | $1.5698 \mathrm{e}-13$ | $9.9615 \mathrm{e}-02$ | $4.4658 \mathrm{e}-02$ | $3.5786 \mathrm{e}-03$ |
| $10^{-3}$ | $2.1366 \mathrm{e}-03$ | $3.0403 \mathrm{e}-04$ | $2.3378 \mathrm{e}-12$ | $6.5804 \mathrm{e}-03$ | $9.3141 \mathrm{e}-04$ | $4.9409 \mathrm{e}-04$ |
| $10^{-4}$ | $1.8653 \mathrm{e}-04$ | $3.0979 \mathrm{e}-06$ | $2.4701 \mathrm{e}-11$ | $5.7888 \mathrm{e}-04$ | $9.5901 \mathrm{e}-06$ | $8.9323 \mathrm{e}-06$ |
| $10^{-5}$ | $1.8376 \mathrm{e}-05$ | $3.1126 \mathrm{e}-08$ | $4.2676 \mathrm{e}-11$ | $5.7083 \mathrm{e}-05$ | $9.6640 \mathrm{e}-08$ | $9.6306 \mathrm{e}-08$ |
| $10^{-6}$ | $1.8348 \mathrm{e}-06$ | $3.7633 \mathrm{e}-10$ | $7.4494 \mathrm{e}-10$ | $5.7002 \mathrm{e}-06$ | $1.0620 \mathrm{e}-09$ | $2.5923 \mathrm{e}-09$ |

becomes

$$
\frac{d}{d t}\left[\begin{array}{l}
A(t)  \tag{4.1}\\
B(t) \\
C(t)
\end{array}\right]=\left[\begin{array}{ccc}
-1001 & 10 & 1 \\
1000 & -15 & 10 \\
1 & 5 & -11
\end{array}\right]\left[\begin{array}{l}
A(t) \\
B(t) \\
C(t)
\end{array}\right]=M\left[\begin{array}{l}
A(t) \\
B(t) \\
C(t)
\end{array}\right] .
$$

Many explicit numerical methods such as Euler's method and the fourth order Runge-Kutta method (RK4 method) in common use have their pros and cons. The structure of explicit methods is so simple that they are popular and may
diminish computing cost at each time step. But, due to its conditional stability, the restriction on the choice of the size of time step is troublesome. For a stiff problem such as (4.1), the drawback of explicit methods is more severe. Note that Euler's method for (4.1) is stable only if

$$
\Delta t \leq \frac{S_{1}}{\lambda_{\min }} \approx 1.9782 \times 10^{-3} \text { with } S_{1}=-2
$$

and the RK4 method for (4.1) is stable only if

$$
\Delta t \leq \frac{S_{2}}{\lambda_{\min }} \approx 2.7531 \times 10^{-3} \text { with } S_{2} \approx-2.7853
$$

where $\lambda_{\text {min }}$ is the minimum nonzero eigenvalue of $M$ and

$$
S_{i}=\min _{z \in \partial D_{i} \subset \mathbb{C}} \operatorname{Re}(z), \quad i=1,2
$$

with the stability domain $D_{i}$ of Euler's method and the RK4 method, respectively. Figure 3 shows that the $\mathrm{CR}^{2}$ algorithm gives a stable approximation even if a large $\Delta t$ is used. For a stiff problem, Gear's method is widely used in order to avoid the problems of instability due to the stiffness. We investigate the difference between two proposed methods and one of Gear's method, a twostep backward differentiation formula (BDF2) which is an $A$-stable method of order 2. BDF2 takes a long time to solve a linear system. Moreover, we should pay special attention during the process of solving the relevant linear system with BDF2 because it may be ill-conditioned for a large $\Delta t$. Hence, two explicit algorithms $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$ are superior to BDF2 in view of efficiency and simplicity of implementation. To observe the convergence speed of three methods, two different errors $\left\|\mathbf{y}\left(t_{F}\right)-\mathbf{y}^{N}\right\|_{1}$ and $\sum_{i=1}^{N}\left\|\mathbf{y}\left(t_{i}\right)-\mathbf{y}^{i}\right\|_{1} \Delta t$ are shown in Table 1, where $\left\|\mathbf{y}\left(t_{F}\right)-\mathbf{y}^{N}\right\|_{1}$ is computed at the terminal point $t_{F}=3$ and $\sum_{i=1}^{N}\left\|\mathbf{y}\left(t_{i}\right)-\mathbf{y}^{i}\right\|_{1} \Delta t$ is calculated over all discrete time steps $t=t_{i}$ with $N=t_{F} / \Delta t$. Note that the performance of $\mathrm{SCR}^{2}$ is comparable to that of BDF2 over all discrete time steps while BDF2 shows more accurate results at the terminal point $t_{F}$.

Next, we consider a somewhat complicated general reaction system in Figure 2-(b) with initial values

$$
A(0)=1, \quad B(0)=0, \quad C(0)=0, \quad D(0)=0, \quad E(0)=0 \quad \text { and } \quad F(0)=0,
$$

where

$$
\begin{aligned}
& f_{1}=0.5, \quad f_{2}=0.01, \quad f_{3}=5.0, \quad f_{4}=0.1, \quad f_{5}=0.1, \quad f_{6}=1.0 \\
& b_{1}=0.05, \quad b_{2}=0.001, \quad b_{3}=0.5, \quad b_{4}=0.01, \quad b_{5}=0.01, \quad b_{6}=1.0
\end{aligned}
$$

In order to illustrate convergence properties of two algorithms $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$, two kinds of errors measured in the $l_{1}$-norm and estimates of the order of convergence are summarized in Tables 2 and 3. Table 2 gives errors $\| \mathbf{y}\left(t_{F}\right)$ $\mathbf{y}^{N} \|_{1}$ for varing $\Delta t$ and estimates of the convergence order corresponding to $\left\|\mathbf{y}\left(t_{F}\right)-\mathbf{y}^{N}\right\|_{1}$ where $N=t_{F} / \Delta t$ with the terminal point $t_{F}=10$. The numbers presented in Table 2 imply that when the time step $\Delta t$ is halved, the errors

TABLE 2. Convergence behaviors of the algorithms $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$ by means of $\left\|\mathbf{y}\left(t_{F}\right)-\mathbf{y}^{N}\right\|_{1}$ with $\Delta t$ varying from $1 / 2$ to $1 / 256$ where $N=t_{F} / \Delta t$ with $t_{F}=10$

|  | $\mathrm{CR}^{2}$ |  | $\mathrm{SCR}^{2}$ |  |
| :---: | :---: | :---: | :---: | :---: |
| $\Delta t$ | $\left\\|\mathbf{y}\left(t_{F}\right)-\mathbf{y}^{N}\right\\|_{1}$ | order | $\left\\|\mathbf{y}\left(t_{F}\right)-\mathbf{y}^{N}\right\\|_{1}$ | order |
| $1 / 2$ | $5.7923 \mathrm{e}-02$ |  | $1.0089 \mathrm{e}-02$ |  |
| $1 / 4$ | $2.7704 \mathrm{e}-02$ | 1.0640 | $3.4552 \mathrm{e}-03$ | 1.5460 |
| $1 / 8$ | $1.3454 \mathrm{e}-02$ | 1.0421 | $1.2101 \mathrm{e}-03$ | 1.5137 |
| $1 / 16$ | $6.6209 \mathrm{e}-03$ | 1.0229 | $3.6181 \mathrm{e}-04$ | 1.7418 |
| $1 / 32$ | $3.2836 \mathrm{e}-03$ | 1.0118 | $9.9202 \mathrm{e}-05$ | 1.8668 |
| $1 / 64$ | $1.6350 \mathrm{e}-03$ | 1.0059 | $2.5992 \mathrm{e}-05$ | 1.9323 |
| $1 / 128$ | $8.1584 \mathrm{e}-04$ | 1.0030 | $6.6537 \mathrm{e}-06$ | 1.9659 |
| $1 / 256$ | $4.0750 \mathrm{e}-04$ | 1.0015 | $1.6833 \mathrm{e}-06$ | 1.9829 |

TABLE 3. Convergence behaviors of the algorithms $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$ by means of $\sum_{i=1}^{N}\left\|\mathbf{y}\left(t_{i}\right)-\mathbf{y}^{i}\right\|_{1} \Delta t$ with $\Delta t$ varying from $1 / 2$ to $1 / 256$ where $N=t_{F} / \Delta t$ with $t_{F}=10$

|  | $\mathrm{CR}^{2}$ |  | $\mathrm{SCR}^{2}$ |  |
| :---: | :---: | :---: | :---: | :---: |
| $\Delta t$ | $\sum_{i=1}^{N}\left\\|\mathbf{y}\left(t_{i}\right)-\mathbf{y}^{i}\right\\|_{1} \Delta t$ | order | $\sum_{i=1}^{N}\left\\|\mathbf{y}\left(t_{i}\right)-\mathbf{y}^{i}\right\\|_{1} \Delta t$ | order |
| $1 / 2$ | 1.4005 |  | $2.6305 \mathrm{e}-01$ |  |
| $1 / 4$ | $6.6803 \mathrm{e}-01$ | 1.0680 | $7.0304 \mathrm{e}-02$ | 1.9036 |
| $1 / 8$ | $3.2096 \mathrm{e}-01$ | 1.0575 | $1.8225 \mathrm{e}-02$ | 1.9477 |
| $1 / 16$ | $1.5670 \mathrm{e}-01$ | 1.0344 | $4.7130 \mathrm{e}-03$ | 1.9512 |
| $1 / 32$ | $7.7354 \mathrm{e}-02$ | 1.0185 | $1.2067 \mathrm{e}-03$ | 1.9656 |
| $1 / 64$ | $3.8422 \mathrm{e}-02$ | 1.0095 | $3.0594 \mathrm{e}-04$ | 1.9797 |
| $1 / 128$ | $1.9147 \mathrm{e}-02$ | 1.0048 | $7.7070 \mathrm{e}-05$ | 1.9890 |
| $1 / 256$ | $9.5573 \mathrm{e}-03$ | 1.0024 | $1.9344 \mathrm{e}-05$ | 1.9943 |

$\left\|\mathbf{y}\left(t_{F}\right)-\mathbf{y}^{N}\right\|_{1}$ of $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$ decay linearly and quadratically, respectively. On the other hand, Table 3 displays the errors $\sum_{i=1}^{N}\left\|\mathbf{y}\left(t_{i}\right)-\mathbf{y}^{(i)}\right\|_{1} \Delta t$ and estimates of the convergence order. According to Table 3, we confirm that the algorithms $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$ converge of order 1 and 2, respectively, over the whole interval $(0,10]$.

Finally, we slightly relax the requirement of the conserved system. Let us consider a non-conserved system

$$
\frac{d}{d t}\left[\begin{array}{l}
A(t)  \tag{4.2}\\
B(t)
\end{array}\right]=\left[\begin{array}{cc}
-10 & 0.5 \\
1 & -1
\end{array}\right]\left[\begin{array}{l}
A(t) \\
B(t)
\end{array}\right]=M_{v}\left[\begin{array}{l}
A(t) \\
B(t)
\end{array}\right]
$$

Table 4. Convergence behaviors of the algorithms $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$ by means of $\sum_{i=1}^{N}\left\|\mathbf{y}_{v}\left(t_{i}\right)-\mathbf{y}_{v}^{i}\right\|_{1} \Delta t$ with $\Delta t$ varying from $10^{-1}$ to $10^{-5}$ where $\mathbf{y}_{v}(t)=[A(t), B(t)]^{T}$ and $N=t_{F} / \Delta t$ with $t_{F}=5$

| $\Delta t$ | $\mathrm{CR}^{2}$ |  | $\mathrm{SCR}^{2}$ |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $\sum_{i=1}^{N}\left\\|\mathbf{y}_{v}\left(t_{i}\right)-\mathbf{y}_{v}^{i}\right\\|_{1} \Delta t$ | order | $\sum_{i=1}^{N}\left\\|\mathbf{y}_{v}\left(t_{i}\right)-\mathbf{y}_{v}^{i}\right\\|_{1} \Delta t$ | order |
| $10^{-1}$ | $2.7803 \mathrm{e}-01$ |  | $1.3788-01$ |  |
| $10^{-2}$ | $2.6722 \mathrm{e}-02$ | 1.0172 | $1.7209 \mathrm{e}-03$ | 1.9038 |
| $10^{-3}$ | $2.6594 \mathrm{e}-03$ | 1.0021 | $1.7650 \mathrm{e}-05$ | 1.9890 |
| $10^{-4}$ | $2.6580 \mathrm{e}-04$ | 1.0002 | $1.7695 \mathrm{e}-07$ | 1.9989 |
| $10^{-5}$ | $2.6579 \mathrm{e}-05$ | 1.0000 | $1.7660 \mathrm{e}-09$ | 2.0009 |

where $A(0)=1$ and $B(0)=10$. We introduce an artificial substance $C$ which plays a role as a sink substance in the following conserved system

$$
\frac{d}{d t}\left[\begin{array}{l}
A(t)  \tag{4.3}\\
B(t) \\
C(t)
\end{array}\right]=\left[\begin{array}{ccc}
-10 & 0.5 & 0 \\
1 & -1 & 0 \\
9 & 0.5 & 0
\end{array}\right]\left[\begin{array}{l}
A(t) \\
B(t) \\
C(t)
\end{array}\right]=M_{c}\left[\begin{array}{c}
A(t) \\
B(t) \\
C(t)
\end{array}\right],
$$

where $C(0)=0$. It is easily expected that the modification of the non-conserved system into the conserved system can be done provided $M_{v}$ is columnwisediagonally dominant. Table 4 shows the errors for the approximations obtained by applying the algorithms $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$ to (4.3).

Remark 4.1. Let us consider a non-conserved system where the relevant matrix $M_{v}$ violates the columnwise-diagonal dominance. Similarly to (4.3), the nonconserved system can be converted into the conserved one by adding an artificial substance which makes each column of the resultant matrix $M_{g}$ add up to 0 . But, $M_{g}$ differs from $M_{c}$ in (4.3) in the fact that $M_{g}$ has negative off-diagonal entries. Unfortunately, our mathematical verification of the proposed methods does not cover such a case. But in practice, we numerically observed that the algorithms $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$ applied to the above case give the same performance as in a linear reaction system (3.1) in view of both stability and convergence.

## 5. Conclusions

We proposed two numerical schemes $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$ applicable to linear reaction systems. Since the $\mathrm{CR}^{2}$ algorithm is motivated by the exact solver for a reversible reaction with two substances, it is very easy to implement. On the other hand, the $\mathrm{SCR}^{2}$ algorithm, a symmetrized version of the $\mathrm{CR}^{2}$ algorithm, is enhanced in view of the convergence speed. Unlike most of explicit methods, $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$ are absolutely stable in spite of their explicitness. We
analyzed stability and convergence properties of $\mathrm{CR}^{2}$ and $\mathrm{SCR}^{2}$ and confirmed the theoretical results by the numerical experiments.

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