

An Analytical Expression for Current Gain of an IGBT

Jin-Woo Moon[†] and Sang-Koo Chung^{*}

Abstract – A simple analytical expression for a current gain of IGBT is derived in terms of the device parameters as well as a gate length dependent parameter, which allows for the determination of the current components of the device as a function of its gate length. The analytical results are compared with those from simulation results. A good agreement is found.

Keywords: IGBT, Current Gain, Analytical Expression

1. Introduction

IGBTs (Insulated Gate Bipolar Transistors) are widely used in applications of high voltage and current, enjoying the high current handling capabilities of BJTs (Bipolar Junction Transistors) and the fast switching characteristic of the MOSFET [1]. The current gain defined as the product of the injection efficiency and the base transport factor is a fundamental quantity for the bipolar action of an IGBT. The well-known expression for the base transport factor of BJTs has been commonly used for the bipolar action of the device [2]. However, due to the conductivity modulation in the drift region of the device, the factor for IGBTs is expected to be different from that for BJTs.

In this work, based on the ambipolar diffusion equation, a simple analytical expression for the current gain of an IGBT is derived in terms of the device parameters including a gate length dependent one. Two-dimensional simulations are carried out using ATLAS. The simulation results of the gate length dependent electron current of an IGBT are compared with those obtained from the derived expression for the current gain of the device. Limitations of the value variation for the current gain are discussed. Physical principles involved in the present analysis can be applied to all types of IGBT structure.

2. Analytical Expression

At high-level injection the hole and electron current densities are given, respectively, by $J_p(x) = c'J(x) - qD_a dp/dx$ and $J_n(x) = cJ(x) + qD_a dp/dx$, where $J(x)$ is the total current density, $p(x)$ is the carrier density with x for the vertical distance measured from the anode ($x = 0$) towards the gate edge in the middle of the accumulation re-

gion as shown in Fig.1, q is the electronic charge, and D_a is the ambipolar diffusion coefficient while $c' = 1/(1+b)$ and $c = 1 - c'$ are used with b for the ratio between the electron and hole mobilities. The carrier distribution in the drift region along the vertical path can be readily obtained from the ambipolar diffusion equation as $p(x) = p_0[\sinh(w-x)/L + k \sinh(x/L)]/\sinh[w/L]$, where L is the ambipolar diffusion length, p_0 is the carrier density at the anode junction, w is the thickness of the drift region, and $k = p(w)/p_0$ with $p(w)$ for the gate dependent carrier density near the gate edge.

The hole current density at $x = 0$ and $x = w$ may be written as

$$J_p(0) = c'J(0) + \bar{p}_0 J_{pd} \quad (1)$$

$$J_p(w) = c'J(w) + \bar{p}_0 J_{pd} [1/\cosh r - k]/[1 - k/\cosh r] \quad (2)$$

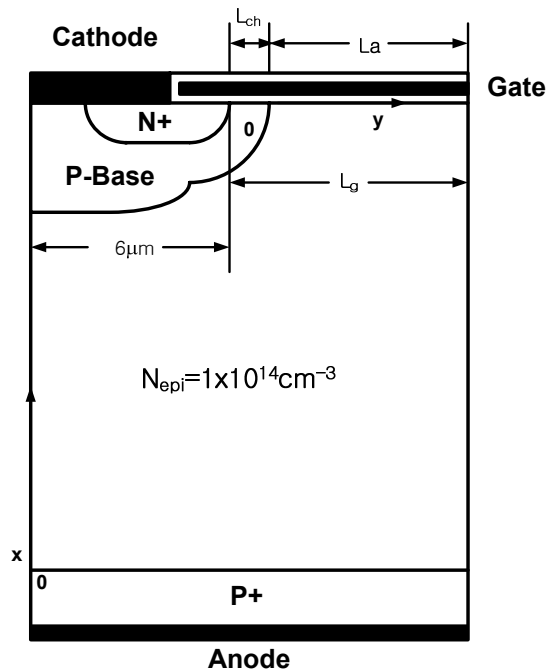


Fig. 1. The Cross section of a half unit cell of an IGBT with parameters

[†] Corresponding Author : Process Technology Development Korea, Analog Technology Development, Fairchild Semiconductor. (jinwoo.moon@fairchildsemi.com)

^{*} Division of Computer, Electronic and Communication Engineering, Yanbian University of Science & Technology (skchung@ybust.edu.cn)

with

$$J_{pd} = [qD_a N_{epi} / (L \tanh r)](1 - k / \cosh r) \quad (3)$$

where $r = w/L$ and $\bar{p}_0 = p_0 / N_{epi}$ are used. Assuming that $J(w) \approx J(0)$ satisfies in Eq. (2), combination of Eqs. (1) and (2) gives an expression for the base transport factor of the IGBT, resulting in

$$\alpha_T = J_p(w)/J_p(0) = c' / \gamma + [1 - c' / \gamma] (1 / \cosh r - k) / [1 - k / \cosh r] \quad (4)$$

where $\gamma = J_p(0)/J(0)$ is the anode injection efficiency of the IGBT treated previously [3].

Eq. (4) reveals that the well-known expression of $\alpha_T \approx [1 / \cosh r]$ for the base transport factor of BJTs is not valid for IGBTs and reduces to it only for the case of $c' \approx 0$ and $k \approx 0$, indicating no drift field for holes in the base region and a vanishing hole concentration in the proximity of the collector junction assumed generally for a P-N-P in active-region operation.

The current gain of an IGBT is found accordingly to be

$$\alpha_{pnp} = \gamma \alpha_T = c' + [\gamma - c'] (1 / \cosh r - k) / (1 - k / \cosh r) \quad (5)$$

Eq. (5) reduces to $\alpha_{pnp} = c' [1 - 1 / \cosh r] + [\gamma / \cosh r]$

for $k = 0$, showing that the current gain of the IGBT in this case is larger than that of the BJT given by $\alpha_{pnp} = \gamma / \cosh r$ and then decreases with the increase of k , approaching ultimately to $\alpha_{pnp} = c'$ for $k = 1 / \cosh r$. The anode injection efficiency, γ in Eq. (5) depends also on k as will be discussed shortly. In the base of the P-N-P structure of an IGBT, both the diffusion and the drift of the carriers are so effective as to result in α_{pnp} as a function of the holes injected from the anode and the electrons injected through the accumulation layer under the gate as well as the carrier mobility.

3. Simulation Results and Discussions

A cross section of the half unit cell with a thickness of $w = 95 \mu\text{m}$ for an IGBT with a blocking capability of 600V used in the simulation is shown in Fig. 1, where the drift region is assumed to be uniformly doped with an impurity concentration of $N_{epi} = 10^{14} \text{cm}^{-3}$ with $\tau_{n0} = \tau_{p0} = 1 \mu\text{s}$ for the electron and hole lifetime, respectively, and a double diffused p body region has a diffusion depth of $5 \mu\text{m}$ and peak concentrations of $2 \times 10^{19} \text{cm}^{-3}$ and $1 \times 10^{17} \text{cm}^{-3}$ at the cathode contact and the inversion layer, respectively. The p+ anode has a thickness of $W_A = 30 \mu\text{m}$ and a uniform distribution of the impurity with a concentration of $N_A = 5 \times 10^{18} \text{cm}^{-3}$. The device has a total width of $(6 + L_g) \mu\text{m}$ with a fixed channel length of $L_{ch} = 2 \mu\text{m}$ and an accumulation layer length of L_a varied from $4 \mu\text{m}$ to $48 \mu\text{m}$ in the simulation, resulting in $L_g = L_a + L_{ch}$ for

the gate length, which allows investigation of the current components of the device as a function of its gate length.

Fig. 2 shows the carrier distribution in the drift region at $V_g = 15[V]$ as a function of the distance measured from the gate at $y = L_g$ along the vertical direction towards the anode with the gate length as a parameter. The carrier density remains almost unvaried in the drift region of the anode side with $p_0 \approx 6.6 \times 10^{16} [\text{cm}^{-3}]$, but increases significantly in the proximity of the gate with increase of the gate length. The values are found to be $p(w) \approx 5.3 \times 10^{15}$, 8.7×10^{15} , $1.4 \times 10^{16} [\text{cm}^{-3}]$ with the corresponding values of $k = 0.08, 0.13, 0.21$ for $L_g = 6, 14, 26 [\mu\text{m}]$, respectively. The simulation results are fitted by $p(w) = 2.7 \times 10^{15} + 0.43 L_g (\mu\text{m}) \times 10^{15} [\text{cm}^{-3}]$.

The n-channel MOSFET supplies electrons through the gate of an IGBT that replaces electrons recombined with the injected holes in the drift region as well as that reach the anode junction. The values for $p(w)$ thus represent the electron injection capability of the accumulation layer.

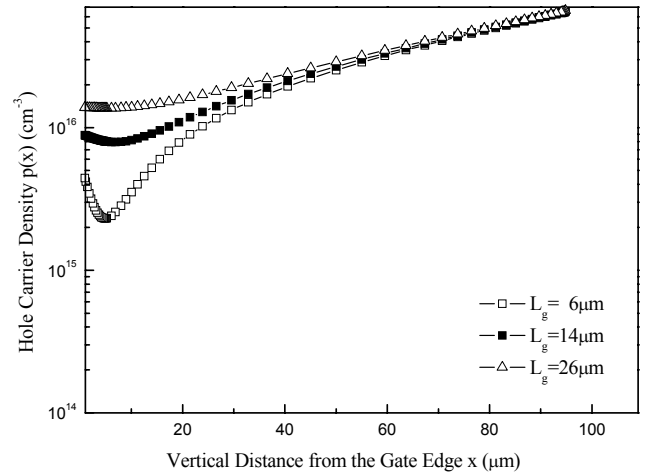


Fig. 2. Carrier distribution in drift region with gate length of IGBT as a parameter at $V_g = 15V$.

The injected carrier density, \bar{p}_0 can be calculated [3] by equating the expressions for the electron current density on either side of the anode junction given by $J_n(0) = cJ(0) - \bar{p}_0 J_{pd}$ from Eq. (1) and $J_n(0) = J_{n0} \bar{p}_0^2$. The result leads an expression for $\bar{p}_0 = \sqrt{(J_{pd} / 2J_{n0})^2 + cJ(0) / J_{n0}} - J_{pd} / 2J_{n0}$, where $J_{n0} = qgD_A N_{epi}^2 / [N_A L_A \tanh(W_A / L_A)]$ with D_A and L_A for the diffusion constant and length for electrons in the anode, respectively, and with g for the band gap narrowing [4] in the heavily doped anode. The present anode structure with $D_A = 3.54 [\text{cm}^2 / \text{s}]$, $L_A = 1.88 [\mu\text{m}]$, and

$g = 15.5$ gives $J_{n0} = 0.93 \times 10^{-4} [A/cm^2]$. In the present device structure $\mu_p = 493 [cm^2/V \cdot s]$ is used for the hole mobility, giving $c' = 0.266$, $D_a = 18.74 [cm^2/s]$ with $\tau_a = 2 \mu s$ leads to $L = 61 \mu m$, resulting in $1/\cosh r \cong 0.4$ and $J_{pd} \cong 53.7(1-0.4k) [mA/cm^2]$ obtained from Eq. (3). Thus we obtain $J_{pd} \cong 53.7, 50.4, 48.7 [mA/cm^2]$ and $\bar{p}_0 = 645, 658, 664$, which yield in turn $\gamma = 0.61, 0.6, 0.59$ for $k = 0, 0.13, 0.21$, respectively, showing that the effects of the gate length on the anode injection efficiency are limitedly small as expected. These values are found to be in excellent agreements with the numerical ones obtained from Fig. 2 within 0.5% in error. The results show that the electron injection increases with an increase of the gate length, calling an increase of the hole injection from the anode junction, which results in an increase in the electron current density accordingly as $J_n(0) = J_{n0} \bar{p}_0^2$ and decrease in $J_p(0)$ by the same amount. This brings about a decrease in the anode injection efficiency and the current gain as well.

The current gain for the present IGBT structure calculated from Eq.(5) is found to be $\alpha_{pnp} \approx 0.4, 0.354, 0.33$ for $k \cong 0, 0.13, 0.21$, respectively, indicating that α_{pnp} decreases with increase of L_g , reducing to $\alpha_{pnp} \approx 0.33$ at $k \approx 0.21$ for $L_g \geq 26 [\mu m]$. An approximate expression for $\alpha_{pnp} \cong 0.266 + 0.334(0.4 - k)/(1 - 0.4k)$ may be found for the present IGBT when, for simplicity, $\gamma \cong 0.6$ is taken in Eq.(5), which would yield $\alpha_{pnp} \cong 0.24$ for the BJT case of $c' \cong 0$ and $k \cong 0$.

Assuming that $I_{mos} = (1 - \alpha_{pnp}) I_{anode}$ holds in the bipolar transistor-MOSFET model of IGBT[2], we finally arrive at $I_{mos} \cong 3.9 + 0.65 L_g (\mu m) [\mu A]$ for $k = 0.13$ and $I_{mos} \cong 4 + 0.67 L_g (\mu m) [\mu A]$ for $k = 0.21$, whereas $I_{mos} \cong 4.6 + 0.76 L_g (\mu m)$ for the BJT case.

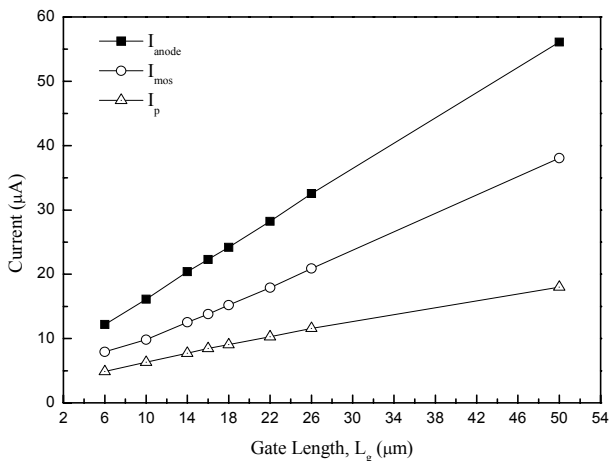


Fig. 3. Variation of electron, hole, and total current vs. gate length of IGBT at $V_g = 15V$.

The simulation results for the electron, hole, and total current denoted as I_{mos} , I_p and I_{anode} , respectively, are represented as data points in Fig.3 as a function of the gate length at $V_g = 15[V]$. With $J(0) = 100 [A/cm^2]$ fixed, the total current for the present device is given by $I_{anode} = I_p + I_{mos} = 6 + L_g (\mu m) [\mu A]$, whereas the electron current is found to be fitted by $I_{mos} = 4 + 0.66 L_g (\mu m) [\mu A]$ at least for $L_g \geq 6 [\mu m]$, which shows a fair agreement with the calculated results of $I_{mos} \cong 3.9 + 0.65 L_g (\mu m) [\mu A]$ for $k = 0.13$ and $I_{mos} \cong 4 + 0.67 L_g (\mu m) [\mu A]$ for $k = 0.21$, given above, supporting the validity of Eq.(5). The similar simulation results are also obtained at $V_g = 10[V]$ and $V_g = 20[V]$. The V_g effects on I_{mos} , I_p and I_{anode} are found to be negligible.

4. Conclusion

In conclusion, it has been demonstrated that a simple analytical expression is derived for the current gain of the IGBT, which gives a good agreement with the simulation results for the gate length dependent MOSFET current of the IGBT. The analytical expression for the current gain of the IGBT provided in the present work, though approximate with one-dimensional treatment, will be useful in IGBT design calculations.

References

- [1] D. S. Kuo and C. Hu, "An analytical model for the power Bipolar-MOS Transistor", *Solid-State Electronics*, Vol. 29, No. 12, pp 1229-1237, 1986.
- [2] B. J. Baliga, *Modern Power Devices*, pp. 363, John Wiley & Sons, 1987.
- [3] S. K. Chung, "Injection currents analysis of p+/n-buffer junction", *IEEE Trans. Electron Devices*, vol. 45, pp. 1850-1854, Aug. 1998.
- [4] J. W. Slotboom and H. C. de Graaff, "Measurements of bandgap narrowing in Si bipolar transistors", *Solid-State Electronics*, Vol. 19, pp. 857-862, 1976.



Jin-Woo Moon received his B.S. and M.E. degree in electronics engineering from Ajou University in 2000 and 2002, respectively. Since 2002, he has been a candidate for a Ph. D. degree in Ajou University. Since December 2004, he has joined Fairchild Semiconductor, Bucheon, Korea. His research interests

include design, fabrication and process development of power semiconductor devices.



Sang-Koo Chung received his B.S. degree in electronics engineering in 1962 from Seoul National University, Seoul, Korea, and his M.S. and D.Sc. degrees in electrical engineering in 1972 and 1974, respectively, from Washington University, St. Louis, MO., where he was also a post-doctoral fellow from 1974 to 1975. He then worked on garnet films for bubble memory application. He also attended the Technical University, Berlin, Germany, from 1964 to 1966.

In late 1975, he joined the Institute for Applied Physics, University Regensburg, Germany, where he worked in the area of magnetic thin films and magneto-optics. Since 1978 he has been a professor of Electronics Engineering at Ajou University, Suwon, Korea, where he served as Dean of the Engineering School from 1981 to 1983. During the school year 1983-1984, he was a visiting scholar at U.C.L.A., Los Angeles, CA, where he was engaged in research on HEMT devices. From 1990 to 1994, he has served as Dean of the Graduate School, Ajou University. In 1996 he spent a sabbatical year as a guest researcher at ABB Semiconductor AG., Switzerland, where he was engaged in power silicon device analysis and modeling. His current research interests include device modeling, simulation, and fabrication technology of power semiconductor devices. In February 2003, he retired, becoming a professor emeritus at Ajou University. In September 2003, he joined Yanbian University of Science and Technology (YUST) as a professor.

Dr. Chung is a senior member of the IEEE Electron Device Society, the IEEE Microwave Theory and Techniques Society, the New York Academy of Sciences, the Korean Institute of Telematics and Electronics, and the Korean Institute of Electrical Engineers. He is listed in *Who's who in the World* (1998), *Who's who in finance and industry* (1999), and *Who's who in Science and Engineering* (2000). He received the Kyung Gi Do Cultural Award in Natural Science in 2001.