

Verification of Numerical Technique for Hydraulic Fracturing Stimulation - by Comparison with Analytical Solutions -

수압파쇄 설계를 위한 수치해석기법의 증명 - 해석식과의 비교를 중심으로 -

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ABSTRACT : Hydraulic fracturing technology has been widely applied in the industry for the recovery of the natural resources such as gas, oil and geothermal heat from hot dry rock. During hydraulic fracturing stimulation, multiple cracks are created resulting in mechanical interaction between cracks. Such an interaction influences obtaining hydraulic fracturing key parameters (crack opening, length, and borehole net pressure). The boundary collocation method (BCM) has been proved to be very effective in considering mechanical interaction. However, for better confidence, it needs to be verified by comparison with analytical solutions such as stress intensity factors. In this paper, three cases, single fracture in remote uniaxial tension, single fracture in remote shear stress field and two arbitrary segments in an infinite plane loaded at infinity are considered. As a result, the BCM is proved to be valid technique to consider mechanical interaction between cracks and can be used to estimate the hydraulic fracturing parameters such as opening of the fracture, and so on.

Keywords : Hydraulic fracturing, Boundary collocation method, Stress Intensity Factor

요지 : 수압파쇄기술은 가스나 석유, 지열 등 자원추출을 하기 위해 다양한 분야에서 전세계적으로 응용되고 있는 기술이다. 이러한 수압파쇄 작업 시 복수의 균열이 필수적으로 발생하여 균열간 기계적인 상호작용을 유발하는데 이러한 상호작용은 수압파쇄 시 얻어질 수 있는 결과(균열 폭, 균열 길이, 보어홀 내 압력)에 큰 영향을 끼치게 된다. 수치해석기법인 경계병치법은 이러한 균열 간의 역학적 상호작용을 고려하는데 유효한 수치해석적 기법으로 개발이 되고 있으나 응력확대계수를 계산하는 해석식과의 비교 등을 통한 검증이 필요하다. 이를 위해 무한평면에 일축 인장 응력과 전단응력이 작용하는 단일균열의 경우 및 임의의 두 균열이 존재하는 경우의 응력확대계수 및 균열폭 해석식과 본 수치해석기법을 통해 얻은 값을 비교하였다. 그 결과, 본 연구에서 제시한 경계병치법은 해석식과 상당히 근접한 결과를 나타내어, 균열간의 기계적인 상호작용을 고려하는데 유효함을 검증하였으며, 추후 수압파쇄 시 설계에 필요한 균열폭 등의 변수를 계산하는데 사용할 수 있음을 나타내었다.

주요어 : 수압파쇄, 경계병치법, 응력확대계수

1. Introduction

Hydraulic fracturing technology has been widely applied for the recovery of natural resources such as oil, gas, geothermal heat from hot dry rock (Fig. 1), waste disposal and remediation of contaminated water to create permeable reactive barrier (Fig. 2). During hydraulic fracturing stimulation, however, multiple fractures are expected to be created and mechanical interaction induced by fluid injection between fractures is essential. To consider mechanical interaction between fractures

(e.g., Naceur and Roegiers, 1990) is of great importance especially in the area of the recovery of natural resources because it affects fractures opening, length, and net pressure (difference between remote stress and internal pressure).

The boundary collocation method (BCM) is based on the primarily analytical derivations accounting for the mechanical interaction between fractures. This technique are, however, fairly complicated and, in many cases, not trivial. Sim et al. (2006) and Sim and Kim (2005) have improved this technique to visualize opening of the cracks induced by fluid

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injection and have verified it by comparing opening with that of finite element method. For better confidence in the version of the BCM developed in those papers, stress intensity factor (SIF) needs to be compared with other analytical solutions because SIF gives further information how crack will behave in the next. In this paper, therefore, SIF computed by the BCM are mainly compared with available analytical solutions.

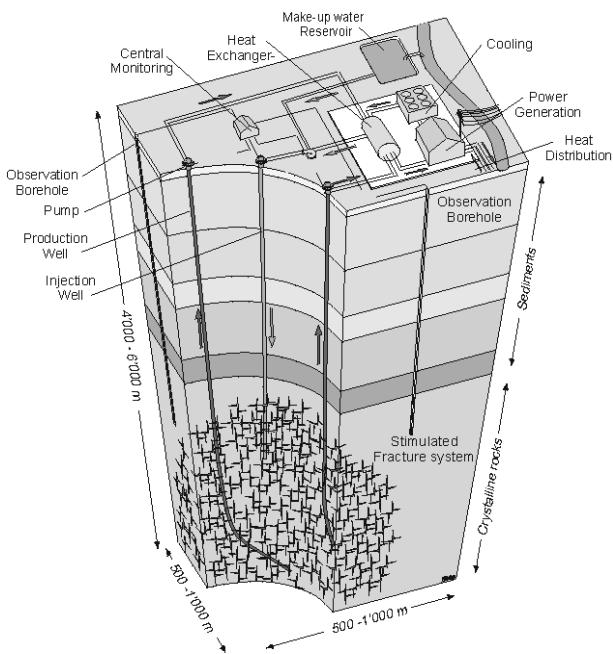


Fig. 1. Geothermal heat extraction (<http://ww2.mne.ksu.edu/~geocrack/>)

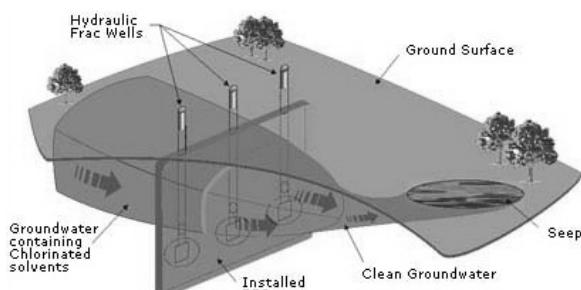


Fig. 2. Permeable reactive vertical barrier (Hocking, 2002)

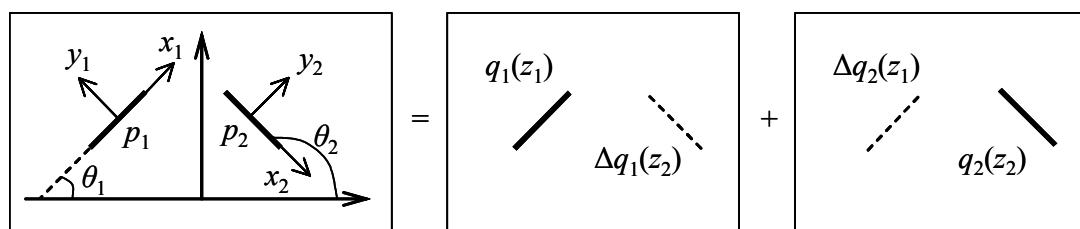


Fig. 3. Principle of numerical technique for interacting with hydraulic fractures

2. Boundary Collocation Method (BCM)

The principle of the BCM has been introduced by Sim et al. (2006) by introducing a pair of interacting fractures in an infinite plane that are located at z_1 and z_2 , respectively (Fig. 3). Known tractions, p_1 and p_2 are applied to the first and second fracture, respectively. Based on the superposition principle, this problem can be represented as a sum of two auxiliary problems for a single fracture. In the first auxiliary problem, unknown tractions acting on the first fracture q_1 induce stresses Δq_1 at the location of the second fracture. Likewise, in the second auxiliary problem, unknown tractions loading the second fracture q_2 generate stresses Δq_2 at the location of the first fracture. Therefore, the tractions in the original problem can be written as

$$q_1(z_1) + \Delta q_2(z_1) = p_1(z_1), \quad q_2(z_2) + \Delta q_1(z_2) = p_2(z_2) \quad (1)$$

In this method, the complex potential is assumed to be an infinite series truncated to a finite number of collocation points. These collocation points can be determined to satisfy prescribed boundary conditions. In this version of BCM, an approach using Chebyshev polynomials as an approximate function for the solution of multiple defects in an infinite homogeneous medium (e.g., Gladwell and England, 1977; McCartney and Gorley, 1987) are employed. As a result, stress intensity factors (SIF) and opening of the cracks, respectively, can be expressed as follows

$$K_I + iK_{II} = \sqrt{\pi c} \sum_{m=1}^M (\pm 1)^{m+1} (\alpha_m + i\beta_m) \quad (2)$$

$$\Delta v + i\Delta u = \frac{4c}{E'} \sqrt{1 - \xi^2} \sum_{m=1}^M \frac{1}{m} (\alpha_m + i\beta_m) U_{m-1}(\xi) \quad (3)$$

where K_I is mode I SIF, K_{II} is mode II SIF, c is the half size of the crack, α_m and β_m are real unknown coefficients to be determined, M is the number of collocation points, $U_m(\xi)$

$= \sin((m+1)\arccos(\xi))/\sin(\arccos(\xi))$ is the m th order Chebyshev polynomial of the second kind, $\xi = x/c$ is the dimensionless coordinate along the crack, Δv is normal opening, Δu is shear opening, $E' = E/(1 - \nu^2)$ for plane strain, E is the Young's modulus, ν is the Poisson's ratio, and + and - in (2) indicate SIFs for the case of $\xi > 0$ and $\xi < 0$, respectively.

3. Verification by Comparison with Available Solutions

3.1 Single Fracture in Remote Uniaxial Tension

A single inclined fracture in remote uniaxial tension is considered as shown in Fig. 4. The remote stress field is

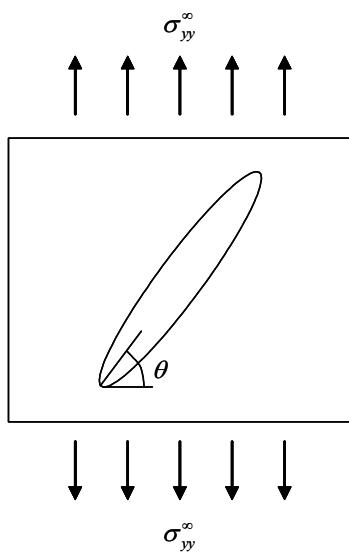


Fig. 4. Single inclined fracture in remote uniaxial tension

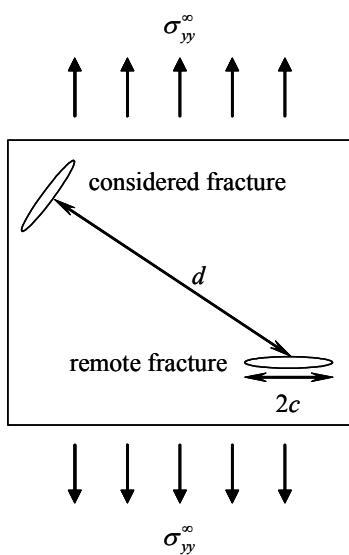


Fig. 5. Remotely spaced fracture ($d = 100c$)

given by $\sigma_{xx}^\infty = 0$, $\sigma_{yy}^\infty = 1$, and $\tau_{xy}^\infty = 0$ while the fracture is inclined at an angle of $\theta = 55^\circ$ with respect to the x-axis. Fracture half length is $c = 1$. Since the numerical code is developed for the case of more than two fractures ($N \geq 2$), an auxiliary fracture can be placed at a distance of $d = 100c$ such that their interaction becomes negligible as shown in Fig. 5. Therefore, both fractures can be considered as totally isolated from each other. Hereafter, the unit is not shown because it can be any type of unit system as long as it is used consistently in all parameters.

The analytical solution for SIF gives (Tada et al., 2000)

$$K_I = \sigma_{yy}^\infty \sin\left(\frac{\pi}{2} - \theta\right) \sin\left(\frac{\pi}{2} - \theta\right) \sqrt{\pi c} \quad (4)$$

$$K_{II} = \sigma_{yy}^\infty \cos\left(\frac{\pi}{2} - \theta\right) \sin\left(\frac{\pi}{2} - \theta\right) \sqrt{\pi c} \quad (5)$$

The SIFs computed by the BCM show good agreement with analytical solutions (Table 1) showing 0.02% error for both K_I and K_{II} .

Normal and shear opening can be also calculated analytically based on (6) and (7) respectively. Poisson's ratio $\nu = 0.21$ and Young's modulus $E = 1$ are used as the elastic properties of the material. The analytical solutions for normal and shear opening (displacement discontinuities) are given by (Tada et al., 2000)

$$\Delta v = \frac{4(1-\nu^2)}{E} \sigma_y^\infty \sin\left(\frac{\pi}{2} - \theta\right) \sin\left(\frac{\pi}{2} - \theta\right) \sqrt{c^2 - x^2} \quad (6)$$

$$\Delta u = \frac{4(1-\nu^2)}{E} \sigma_y^\infty \cos\left(\frac{\pi}{2} - \theta\right) \sin\left(\frac{\pi}{2} - \theta\right) \sqrt{c^2 - x^2} \quad (7)$$

These normal and shear opening of the single inclined fracture (Fig. 4) are compared with the analytical solutions as shown in Fig. 6 and Fig. 7 and show good agreement.

3.2 Single Fracture in Remote Shear Stress Field

A single inclined fracture in a remote shear stress field is considered as shown in Fig. 8. The remote stress field is

Table 1. Comparison of computed SIFs with analytical solutions

	K_I	K_{II}
The BCM	0.5830	0.8330
Tada et al. (2000)	0.5831	0.8328

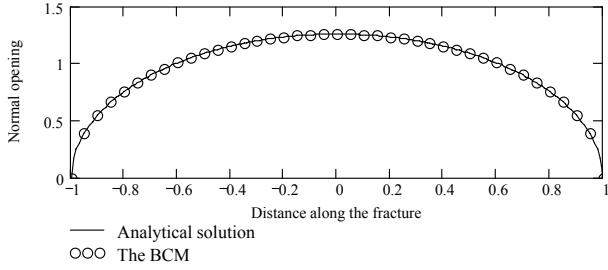


Fig. 6. Comparison of normal opening with analytical solution

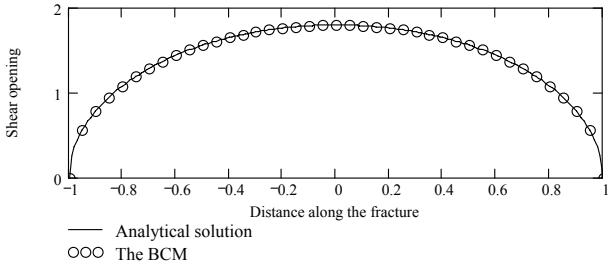


Fig. 7. Comparison of shear opening with analytical solution

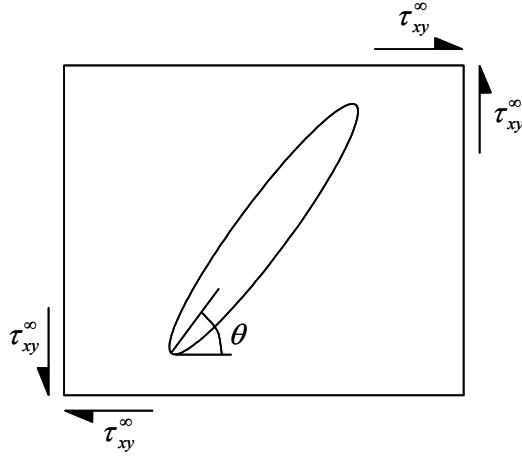


Fig. 8. Single inclined fracture in remote shear stress field

given by $\sigma_{xx}^\infty = 0$, $\sigma_{yy}^\infty = 0$, and $\tau_{xy}^\infty = 1$ while the fracture is inclined at the same angle of $\theta = 55^\circ$. Fracture half length is $c = 1$. The auxiliary fracture is also placed at a distance of $d = 100c$ such that their interaction become negligible.

An analytical solution for SIFs is (Tada et al., 2000)

$$K_I = -\tau_{xy}^\infty \sin(2\theta)\sqrt{\pi c} \quad (8)$$

$$K_{II} = \tau_{xy}^\infty \cos(2\theta)\sqrt{\pi c} \quad (9)$$

Comparing with the analytical solutions in Table 2, the computed results by the BCM show 0.02% error for K_I and 0.03% error for K_{II} , in good agreement with the analytical solutions.

Table 2. Comparisons of computed SIFs with analytical solution

	K_I	K_{II}
The BCM	-1.6656	-0.6062
Tada et al. (2000)	-1.6660	-0.6060

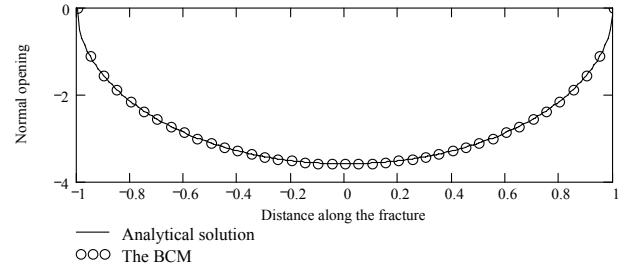


Fig. 9. Comparison of normal opening with analytical solution

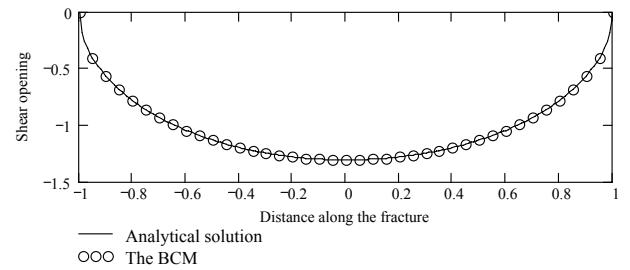


Fig. 10. Comparison of shear opening with analytical solution

The analytical solutions for normal and shear openings are given by (Tada et al., 2000)

$$\Delta v = -\frac{4(1-\nu^2)}{E} \tau_{xy}^\infty \sin(2\theta)\sqrt{c^2 - x^2} \quad (10)$$

$$\Delta u = \frac{4(1-\nu^2)}{E} \tau_{xy}^\infty \cos(2\theta)\sqrt{c^2 - x^2} \quad (11)$$

These normal and shear openings of the single inclined fracture (Fig. 8) are compared with analytical results as shown in Fig. 9 and Fig. 10 and show good agreement.

3.3 Two Arbitrary Segments in an Infinite Plane Loaded at Infinity

The asymptotic solution by Savruk and Datsyshin (1973) considers a fracture consisting of two equal, arbitrarily oriented segments in an infinite plane loaded at infinity. Stresses at infinity are $\sigma_{xx}^\infty = 0.5$, $\sigma_{yy}^\infty = 1$, and $\tau_{xy}^\infty = 0$. The fracture configuration is summarized in Table 3 and plotted in Fig. 11. Two segments are located fairly far ($d > c$) from each other to make sure that the asymptotic solution is accurate.

The form of the solution presented by Savruk and Datsyshin (1973) depends upon two indices, n and k . One of them

Table 3. Configuration of two arbitrary segments that are located fairly far ($d > c$) from each other

Segment	Center of segment (x, y)	Inclination angle θ	Half length c
1	(0, 0)	-74°	1
2	(1, 5)	30°	1

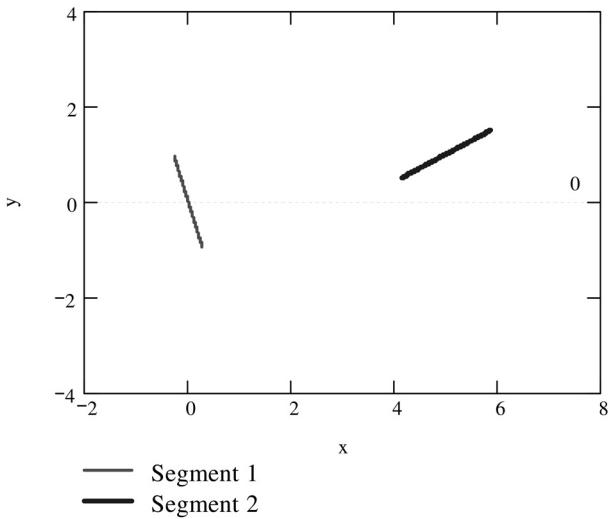


Fig. 11. Two arbitrary segments in an infinite plane loaded at infinity

should be equal to one and the other, two. Therefore, the segment number is identified by a pair of digits (n, k). For example, (1, 2) represents the first segment while (2, 1) represents the second one. Savruk and Datsyshin's (1973) solutions for K_I on the left (K_{I_l}) and right (K_{I_r}) crack tips can be expressed as

$$K_{I_l}(n, k) = \frac{1}{2} \sigma_{yy}^{\infty} \sqrt{\pi c_n} \left[(1 + \varepsilon) + (1 - \varepsilon) \cos(2\theta_n) + \left(\frac{\lambda^2}{16} \right) I_1(n, k) + (-1)^{k+1} \left(\frac{\lambda^3}{32} \right) I_2(n, k) + \frac{\lambda^4}{256} I_3(n, k) \right] \quad (12)$$

$$K_{I_r}(n, k) = \frac{1}{2} \sigma_{yy}^{\infty} \sqrt{\pi c_n} \left[(1 + \varepsilon) + (1 - \varepsilon) \cos(2\theta_n) + \left(\frac{\lambda^2}{16} \right) I_1(n, k) + (-1)^k \left(\frac{\lambda^3}{32} \right) I_2(n, k) + \frac{\lambda^4}{256} I_3(n, k) \right] \quad (13)$$

where

$$I_1(n, k) = 2(1 + \varepsilon) (\cos(2\theta_k - \beta) + \cos(2\beta - \theta_n) - \cos(2(2\beta - \theta_k - \theta_n))) + (1 - \varepsilon) \left(2 \cos(2\beta) + \cos(2(\beta - \theta_k - \theta_n)) + \cos(2(\beta - \theta_n + \theta_k)) \right) \quad (14)$$

$$I_2(n, k) = (1 + \varepsilon) \left(3 \cos(2\theta_k + \theta_n - 3\beta) + 2 \cos(3(\beta - \theta_n)) \right) - 3 \cos(5\beta - 2\theta_k - 3\theta_n) + (1 - \varepsilon) \left(3 \cos(\theta_n - 3\beta) + \cos(3\beta - 3\theta_n - 2\theta_k) \right) + \cos(3\beta - 3\theta_n + 2\theta_k) - 3 \cos(5\beta - 3\theta_n) \quad (15)$$

$$I_3(n, k) = 2(1 + \varepsilon) \begin{cases} 3 \cos(4(\theta_k - \beta)) + 20 \cos(2(2\beta - \theta_n - \theta_k)) \\ - 6 \cos(2(3\beta - 2\theta_k - \theta_n)) + 7 \cos(4(\beta - \theta_n)) \\ - 14 \cos(2(3\beta - 2\theta_n - \theta_k)) + 5 + 2 \cos(2(\theta_n - \theta_k)) \\ - 4 \cos(2(\theta_n - \theta_k)) - 2 \cos(2(\theta_k - \beta)) \end{cases} \\ + (1 - \varepsilon) \begin{cases} 9 \cos(2(\theta_k - 2\beta)) + 3 \cos(2(2\beta - 2\theta_k - \theta_n)) \\ + 35 \cos(2(2\beta - \theta_n)) - 16 \cos(2(3\beta - \theta_n - \theta_k)) \\ + 7 \cos(2(2\beta - 2\theta_n - \theta_k)) + 6 \cos(2(2\beta - 2\theta_n + \theta_k)) \\ - 24 \cos(2(3\beta - 2\theta_n)) + 10 \cos(2\theta_n) + 3 \cos(2\theta_k) \\ - 6 \cos(2\beta) + \cos(2(\theta_k - 2\theta_n)) - 2 \cos(2(\theta_k - \beta - \theta_n)) \\ - 2 \cos(2(\beta - \theta_n - \theta_k)) - 2 \cos(2(\beta - 2\theta_n)) \end{cases} \quad (16)$$

$$\varepsilon = \frac{\sigma_{xx}^{\infty}}{\sigma_{yy}^{\infty}} \quad (17)$$

$$\lambda = \frac{2c}{d} \quad (18)$$

$$\beta = \arctan \left(\frac{y_2 - y_1}{x_2 - x_1} \right) \quad (19)$$

Their solutions for K_{II} on the left (K_{II_l}) and right (K_{II_r}) crack tips can be expressed as

$$K_{II_l}(n, k) = \frac{1}{2} \sigma_{yy}^{\infty} \sqrt{\pi c_n} \left[(1 - \varepsilon) \sin(2\theta_n) - \left(\frac{\lambda^2}{16} \right) II_1(n, k) + (-1)^k \left(\frac{\lambda^3}{32} \right) II_2(n, k) - \frac{\lambda^4}{256} II_3(n, k) \right] \quad (20)$$

$$K_{II_r}(n, k) = \frac{1}{2} \sigma_{yy}^{\infty} \sqrt{\pi c_n} \left[(1 - \varepsilon) \sin(2\theta_n) - \left(\frac{\lambda^2}{16} \right) II_1(n, k) + (-1)^{k+1} \left(\frac{\lambda^3}{32} \right) II_2(n, k) - \frac{\lambda^4}{256} II_3(n, k) \right] \quad (21)$$

where

$$II_1(n, k) = 2(1 + \varepsilon) (\sin(2(\beta - \theta_n)) - \sin(2(2\beta - \theta_k - \theta_n))) + (1 - \varepsilon) \left(\sin 2(\beta - \theta_k - \theta_n) + \sin(2(\beta - \theta_n + \theta_k)) \right) - 2 \sin(2(2\beta - \theta_n)) \quad (22)$$

$$II_2(n, k) = (1 + \varepsilon) \left(\sin(3\beta - 2\theta_k - \theta_n) + 2 \sin(3(\beta - \theta_n)) \right) - 3 \sin(5\beta - 2\theta_k - 3\theta_n) + (1 - \varepsilon) \left(\sin(3\beta - \theta_n) + \sin(3\beta - 3\theta_n - 2\theta_k) \right) + \sin(3\beta - 3\theta_n + 2\theta_k) - 3 \sin(5\beta - 3\theta_n) \quad (23)$$

$$II_3(n, k) = (1 + \varepsilon) \left(26 \sin(2(2\beta - \theta_n - \theta_k)) - 12 \sin(2(3\beta - 2\theta_k - \theta_n)) \right) + 14 \sin(4(\beta - \theta_n)) - 28 \sin(2(3\beta - 2\theta_n - \theta_k)) + 2 \sin(2(\theta_k - \theta_n)) - 4 \sin(2(\beta - \theta_n)) + \sin(2(2\beta - 2\theta_k - \theta_n)) + 3 \sin(2(2\beta - 2\theta_k - \theta_n)) + 23 \sin(2(2\beta - \theta_n)) - 16 \sin(2(3\beta - \theta_k - \theta_n)) + 7 \sin(2(2\beta - 2\theta_n - \theta_k)) + 6 \sin(2(2\beta - 2\theta_n + \theta_k)) - 24 \sin(2(3\beta - 2\theta_n)) - 6 \sin(2\theta_n) + \sin(2\theta_k) - 2 \sin(2\beta) + \sin(2(\theta_k - 2\theta_n)) - 2 \sin(2(\theta_k - \beta - \theta_n)) - 2 \sin(2(\beta - \theta_n - \theta_k)) - 2 \sin(2(\beta - 2\theta_n)) \quad (24)$$

As shown in Table 4 and 5, the computed results are close to those of the asymptotic solutions by Savruk and Datsyshin (1973). These results are also compared with those calculated by Galybin's (Astakhov et al., 2000) code, which numerically

Table 4. Comparisons of K_I with asymptotic solutions by Savruk and Datsyshin (1973)

Segment	Asymptotic solution		The BCM	
	Left	Right	Left	Right
1	$K_{I_l} = 0.966$	$K_{I_r} = 0.983$	$K_{I_l} = 0.967$	$K_{I_r} = 0.983$
2	$K_{I_l} = 1.557$	$K_{I_r} = 1.559$	$K_{I_l} = 1.555$	$K_{I_r} = 1.557$

Table 5. Comparisons of K_{II} with asymptotic solutions by Savruk and Datsyshin (1973)

Segment	Asymptotic solution		The BCM	
	Left	Right	Left	Right
1	$K_{II_l} = -0.221$	$K_{II_r} = -0.230$	$K_{II_l} = -0.222$	$K_{II_r} = -0.229$
2	$K_{II_l} = 0.419$	$K_{II_r} = 0.405$	$K_{II_l} = 0.419$	$K_{II_r} = 0.405$

Table 6. Stress intensity factors calculated by Galybin's code

Segment	K _I		K _{II}	
	Left	Right	Left	Right
1	$K_{I_l} = 0.967$	$K_{I_r} = 0.983$	$K_{II_l} = -0.222$	$K_{II_r} = -0.229$
2	$K_{I_l} = 1.555$	$K_{I_r} = 1.557$	$K_{II_l} = 0.419$	$K_{II_r} = 0.405$

Table 7. Configuration of two arbitrary segments that are located closely

Segment	Center of segment (x, y)	Inclination angle θ	Half length c
1	(1, 0.9)	40°	1.2
2	(0.9, 2.1)	-30°	0.7

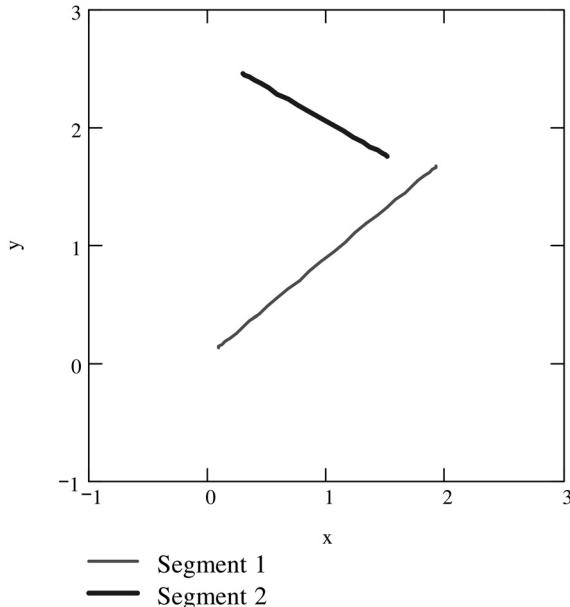


Fig. 12. Two arbitrary segments in an infinite plane loaded by internal pressure

calculates SIFs for the same fracture configuration. SIFs by Galybin's (Astakhov et al., 2000) code are very similar to those obtained by the BCM and the asymptotic solutions

Table 8. Comparisons of K_I with the results by Galybin's (Astakhov et al., 2000) code

Segment	Galybin's code		The BCM	
	Left	Right	Left	Right
1	$K_{I_l} = 3.863$	$K_{I_r} = 4.128$	$K_{I_l} = 3.862$	$K_{I_r} = 4.123$
2	$K_{I_l} = 1.426$	$K_{I_r} = 2.996$	$K_{I_l} = 1.427$	$K_{I_r} = 2.997$

Table 9. Comparisons of K_{II} with the results by Galybin's (Astakhov et al., 2000) code

Segment	Galybin's code		The BCM	
	Left	Right	Left	Right
1	$K_{II_l} = -6.075$	$K_{II_r} = -6.154$	$K_{II_l} = -6.075$	$K_{II_r} = -6.155$
2	$K_{II_l} = 2.422$	$K_{II_r} = 3.203$	$K_{II_l} = 2.422$	$K_{II_r} = 3.203$

(Savruk and Datsyshin, 1973) as shown in Table 6.

Although BCM results for SIFs match the asymptotic solution and Galybin's (Astakhov et al., 2000) code results closely enough, the difference between SIFs on different segment ends is not too large because the segments are located fairly far from each other. To have an independent comparison, the BCM is tested against Galybin's (Astakhov et al., 2000) code for a configuration (Table 7 and Fig. 12) of non-equal segments loaded by different internal pressure and with more significantly varying SIFs on both ends. The internal pressure for two segments are $p_1 = 2$ and $p_2 = 1$ respectively. As a result, the agreement between the two sets of results for SIFs obtained by independent codes is nearly ideal as shown in Table 8 and Table 9.

4. Conclusions

The boundary collocation method is verified by comparison of computed SIFs and openings with those of analytical solutions. Three cases, single fracture in remote uniaxial tension, single fracture in remote shear stress field and two arbitrary segments in an infinite plane loaded at infinity are considered. As a result, the SIFs and opening of the cracks are good agreement with analytical solutions. Therefore, the BCM can be extensively applied to estimate hydraulic fracturing parameters such as opening of the fracture where the mechanical interaction between fractures is expected.

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