

## ON QUASI EINSTEIN MANIFOLDS

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ABSTRACT. The object of the present paper is to study some properties of a quasi Einstein manifold. A non-trivial concrete example of a quasi Einstein manifold is also given.

### 1. Introduction

A Riemannian or a semi-Riemannian manifold  $(M^n, g)$ ,  $n = \dim M \geq 2$ , is said to be an Einstein manifold if the following condition

$$(1.1) \quad S = \frac{r}{n}g$$

holds on  $M$ , where  $S$  and  $r$  denote the Ricci tensor and the scalar curvature of  $(M^n, g)$  respectively. According to ([1], p. 432), (1.1) is called the Einstein metric condition. Einstein manifolds play an important role in Riemannian Geometry as well as in general theory of relativity. Also Einstein manifolds form a natural subclass of various classes of Riemannian or semi-Riemannian manifolds by a curvature condition imposed on their Ricci tensor ([1], pp. 432–433). For instance, every Einstein manifold belongs to the class of Riemannian manifolds  $(M^n, g)$  realizing the following relation :

$$(1.2) \quad S(X, Y) = ag(X, Y) + bA(X)A(Y),$$

where  $a, b \in \mathbf{R}$  and  $A$  is a non-zero 1-form such that

$$(1.3) \quad g(X, U) = A(X)$$

for all vector fields  $X$ .

A non-flat Riemannian manifold  $(M^n, g)$  ( $n > 2$ ) is defined to be a quasi Einstein manifold [7] if its Ricci tensor  $S$  of type  $(0, 2)$  is not identically zero and satisfies the condition (1.2). We shall call  $A$  the associated 1-form and  $U$  is called the generator of the manifold. In this paper we consider  $U$  as a unit vector field.

Quasi Einstein manifolds arose during the study of exact solutions of the Einstein field equations as well as during considerations of quasi-umbilical hypersurfaces of semi-Euclidean spaces. For instance, the Robertson-Walker

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spacetime are quasi Einstein manifolds [8]. Also quasi Einstein manifold can be taken as a model of the perfect fluid spacetime in general relativity [4]. So quasi Einstein manifolds have some importance in the general theory of relativity. Considering this aspect we are motivated to study such a manifold.

It is to be noted that M. C. Chaki and R. K. Maity [2] also introduced the notion of quasi Einstein manifolds which is different from that of R. Deszcz [7]. They took  $a$  and  $b$  as scalars and the generator  $U$  of the manifold as a unit vector field. The study of quasi Einstein manifolds was continued by M. C. Chaki [3], S. Guha [10], U. C. De and G. C. Ghosh [5], [6] and many others.

Section 2 of this paper contains a non-trivial concrete example of a quasi Einstein manifold. In Section 3 we enquire under what condition the associated 1-form  $A$  will be closed in a quasi Einstein manifold. In Section 4 we prove that if a quasi Einstein manifold satisfies cyclic Ricci tensor, then the integral curves of the vector field  $U$  are geodesic. In Section 5 it is shown that if  $U$  is a Killing vector field, then the quasi Einstein manifold satisfies cyclic Ricci tensor. Finally we study a quasi Einstein manifold admitting a conformal Killing vector field.

## 2. Example of a quasi Einstein manifold

In this section we construct a metric of quasi Einstein manifold.

Let  $E^5$  be a Euclidean space with Cartesian coordinates  $(x^1, x^2, y^1, y^2, z)$  or,  $(x^\alpha, y^\alpha, z)$  ( $\alpha = 1, 2$ ). Let us consider

$$(2.1) \quad A = \sqrt{\left(\frac{3}{2b}\right)}(dz - \sum_{\alpha=1}^2 y^\alpha dx^\alpha), \text{ (where } b \text{ is a constant).}$$

If we put

$$(2.2) \quad x^{\alpha^*} \equiv x^{2+\alpha} = y^\alpha, \quad x^\Delta = z, \quad \Delta = 5,$$

we have from (2.1) that

$$(2.3) \quad A_i = \left(-\sqrt{\left(\frac{3}{2b}\right)}y^\alpha, 0, \sqrt{\left(\frac{3}{2b}\right)}\right).$$

Now we consider a symmetric tensor field in  $E^5$  defined by

$$(2.4) \quad g_{ij} = \begin{pmatrix} \frac{1}{4}(1 + (y^1)^2) & \frac{y^1 y^2}{4} & 0 & 0 & -\frac{1}{4}y^1 \\ \frac{1}{4}y^1 y^2 & \frac{1}{4}(1 + (y^2)^2) & 0 & 0 & -\frac{1}{4}y^2 \\ 0 & 0 & \frac{1}{4} & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{4} & 0 \\ -\frac{1}{4}y^1 & -\frac{1}{4}y^2 & 0 & 0 & \frac{1}{4} \end{pmatrix}.$$

Then  $g_{ij}$  defines a positive definite Riemannian metric. The contravariant components of the tensor  $g^{ij}$  are given by

$$(2.5) \quad g^{ij} = \begin{pmatrix} 4 & 0 & 0 & 0 & 4y^1 \\ 0 & 4 & 0 & 0 & 4y^2 \\ 0 & 0 & 4 & 0 & 0 \\ 0 & 0 & 0 & 4 & 0 \\ 4y^1 & 4y^2 & 0 & 0 & 4(1 + y^1 + y^2) \end{pmatrix}.$$

We find the Christoffel symbols, by means of (2.4) and

$$[ij, r] = \frac{1}{2} \left[ \frac{\partial g_{jr}}{\partial x^i} + \frac{\partial g_{ir}}{\partial x^j} - \frac{\partial g_{ij}}{\partial x^r} \right].$$

We can verify that

$$(2.6) \quad [\alpha\beta^*, \gamma] = \frac{1}{8} (\delta_{\gamma\beta} y^\alpha + \delta_{\alpha\beta} y^\gamma),$$

$$(2.7) \quad [\alpha\beta, \gamma^*] = -\frac{1}{8} (\delta_{\gamma\alpha} y^\beta + \delta_{\beta\gamma} y^\alpha),$$

$$(2.8) \quad [\alpha\beta^*, \Delta] = -\frac{\delta_{\alpha\gamma}}{8}, \quad [\alpha^* \Delta, \gamma] = -\frac{\delta_{\alpha\gamma}}{8},$$

$$(2.9) \quad [\alpha\Delta, \gamma^*] = \frac{\delta_{\alpha\gamma}}{8},$$

the other components are zero. We first show that the symmetry of the indices  $\alpha, \beta^*$  of the Christoffel symbols. That is, we show that  $[\alpha\beta^*, \gamma] = [\beta^*\alpha, \gamma]$ . There are total 8 components which are  $[13, 1]$ ,  $[13, 2]$ ,  $[23, 1]$ ,  $[23, 2]$ ,  $[14, 1]$ ,  $[14, 2]$ ,  $[24, 1]$ , and  $[24, 2]$ .

We take any one of the above components, say,  $[23, 1]$ . Then

$$\begin{aligned} [23, 1] &= \frac{1}{2} \left[ \frac{\partial g_{31}}{\partial x^2} + \frac{\partial g_{21}}{\partial x^3} - \frac{\partial g_{23}}{\partial x^1} \right] \\ &= \frac{1}{2} \left( 0 + \frac{1}{4} y^2 - 0 \right), \text{ using (2.2) and (2.4)} \\ &= \frac{1}{8} y^2 \end{aligned}$$

and

$$\begin{aligned} [32, 1] &= \frac{1}{2} \left[ \frac{\partial g_{21}}{\partial x^3} + \frac{\partial g_{31}}{\partial x^2} - \frac{\partial g_{32}}{\partial x^1} \right] \\ &= \frac{1}{2} \left( \frac{1}{4} y^2 + 0 - 0 \right), \text{ using (2.2) and (2.4)} \\ &= \frac{1}{8} y^2, \end{aligned}$$

that is,

$$[23, 1] = [32, 1] = \frac{1}{8} y^2.$$

We also show that the symmetry of the Christoffel symbols  $[\alpha\delta, \gamma^*] = [\delta\alpha, \gamma^*]$ . There are total 4 components which are [15, 3], [15, 4], [25, 3], [25, 4].

We take any one of the above components, say, [15, 3]. Then

$$\begin{aligned} [15, 3] &= \frac{1}{2} \left[ \frac{\partial g_{53}}{\partial x^1} + \frac{\partial g_{13}}{\partial x^5} - \frac{\partial g_{15}}{\partial x^3} \right] \\ &= \frac{1}{2} \left( 0 + 0 - \left(-\frac{1}{4}\right) \right), \text{ using (2.2) and (2.4)} \\ &= \frac{1}{8} \end{aligned}$$

and

$$\begin{aligned} [51, 3] &= \frac{1}{2} \left[ \frac{\partial g_{13}}{\partial x^5} + \frac{\partial g_{53}}{\partial x^1} - \frac{\partial g_{51}}{\partial x^3} \right] \\ &= \frac{1}{2} \left( 0 + 0 - \left(-\frac{1}{4}\right) \right), \text{ using (2.2) and (2.4)} \\ &= \frac{1}{8}, \end{aligned}$$

that is,

$$[15, 3] = [51, 3] = \frac{1}{8}.$$

The equations from (2.5) to (2.9) with the help of  $\left\{ \begin{matrix} h \\ i & j \end{matrix} \right\} = g^{hr} [i, j, r]$  implies that

$$(2.10) \quad \left\{ \begin{matrix} \mu \\ \alpha & \beta^* \end{matrix} \right\} = \frac{1}{2} \delta_{\mu\beta} y^\alpha, \quad \left\{ \begin{matrix} \mu^* \\ \alpha & \beta \end{matrix} \right\} = -\frac{1}{2} (\delta_{\alpha\mu} y^\beta + \delta_{\mu\beta} y^\alpha),$$

$$(2.11) \quad \left\{ \begin{matrix} \mu^* \\ \alpha & \Delta \end{matrix} \right\} = \frac{1}{2} \delta_{\alpha\mu}, \quad \left\{ \begin{matrix} \Delta \\ \alpha & \beta^* \end{matrix} \right\} = \frac{1}{2} (y^\alpha y^\beta - \delta_{\alpha\beta}),$$

$$(2.12) \quad \left\{ \begin{matrix} \Delta \\ \Delta & \beta^* \end{matrix} \right\} = -\frac{1}{2} y^\beta, \quad \left\{ \begin{matrix} \mu \\ \alpha^* & \Delta \end{matrix} \right\} = -\frac{1}{2} \delta_{\alpha\mu},$$

the other components are zero. After straightforward calculations, we obtain the independent components of the curvature tensor  $R_{ijkl}$  as follows

$$(2.13) \quad R_{\delta\gamma\beta\alpha} = \frac{1}{16} (\delta_{\alpha\beta} y^\gamma y^\delta + \delta_{\beta\gamma} y^\alpha y^\delta - \delta_{\alpha\gamma} y^\beta y^\delta - \delta_{\beta\delta} y^\alpha y^\gamma),$$

$$(2.14) \quad R_{\delta^* \gamma^* \beta \alpha} = \frac{1}{16} (\delta_{\alpha\gamma} \delta_{\beta\delta} - \delta_{\alpha\delta} \delta_{\beta\gamma}), \quad R_{\delta\gamma\Delta\alpha} = (\delta_{\alpha\gamma} y^\delta - \delta_{\alpha\delta} y^\gamma),$$

$$(2.15) \quad R_{\delta\gamma^* \beta^* \alpha} = \frac{1}{16} (\delta_{\beta\gamma} y^\alpha y^\delta - 2\delta_{\alpha\beta} \delta_{\gamma\delta} - \delta_{\alpha\gamma} \delta_{\beta\delta}),$$

$$(2.16) \quad R_{\delta\Delta\Delta\alpha} = \frac{1}{16} \delta_{\alpha\delta}, \quad R_{\Delta\gamma^* \Delta\alpha^*} = -\frac{1}{16} \delta_{\alpha\gamma}, \quad R_{\delta^* \Delta \beta^* \alpha} = \frac{1}{16} \delta_{\beta\delta} y^\alpha,$$

the other independent components are zero. From (2.5), (2.13), (2.14), (2.15) and (2.16), it follows that the Ricci tensor has the following non-zero components.

$$(2.17) \quad R_{\delta\alpha} = -\frac{1}{2}(\delta_{\alpha\delta} - 2y^\alpha y^\delta), \quad R_{\delta^*\alpha} = 0, \quad R_{\Delta\Delta} = 1,$$

$$(2.18) \quad R_{\delta^*\alpha^*} = -\frac{\delta_{\alpha\delta}}{2}, \quad R_{\gamma^*\Delta} = 0, \quad R_{\gamma\Delta} = -y^\gamma.$$

Now from the equations (2.3), (2.4), (2.17) and (2.18) we have

$$(2.19) \quad R_{ij} = -2g_{ij} + bA_i A_j.$$

Hence  $E^5$  with the metric (2.4) is a quasi Einstein manifold. Summing up we can state the following:

**Theorem 2.1.** *The  $E^5$  with the metric (2.4) is a quasi Einstein manifold where the associated 1-form  $A$  is given by (2.1).*

### 3. Nature of the associated 1-form of a quasi Einstein manifold

A Riemannian manifold is said to satisfy Codazzi type of Ricci tensor [9] if its Ricci tensor  $S$  satisfies the following condition

$$(3.1) \quad (\nabla_X S)(Y, Z) = (\nabla_Y S)(X, Z).$$

Now from (1.2) we obtain

$$(\nabla_X S)(Y, Z) = b[(\nabla_X A)(Y)A(Z) + (\nabla_X A)(Z)A(Y)].$$

Hence from (3.1) we have

$$b[(\nabla_X A)(Y)A(Z) + (\nabla_X A)(Z)A(Y)] = b[(\nabla_Y A)(X)A(Z) + (\nabla_Y A)(Z)A(X)]$$

or,

$$(3.2) \quad (\nabla_X A)(Y)A(Z) - (\nabla_Y A)(X)A(Z) + (\nabla_X A)(Z)A(Y) - (\nabla_Y A)(Z)A(X).$$

Since equation (3.2) holds for all  $Z$ , putting  $Z = U$  and using  $(\nabla_X A)(U) = 0$ , since  $U$  is a unit vector, we have

$$(3.3) \quad (\nabla_X A)(Y) - (\nabla_Y A)(X) = 0,$$

or,

$$(3.4) \quad dA(X, Y) = 0.$$

Thus we can state the following theorem:

**Theorem 3.1.** *If a quasi Einstein manifold satisfies the Codazzi type of Ricci tensor, then the associated 1-form  $A$  is closed.*

#### 4. Quasi Einstein manifold satisfying cyclic Ricci tensor

Let us suppose that the manifold satisfies cyclic Ricci tensor. Then

$$(4.1) \quad (\nabla_X S)(Y, Z) + (\nabla_Y S)(Z, X) + (\nabla_Z S)(X, Y) = 0.$$

It is known [11] that Cartan hypersurfaces are manifolds, with non-parallel Ricci tensor, satisfying the form (4.1). From (1.2) we obtain

$$(\nabla_X S)(Y, Z) = b[(\nabla_X A)(Y)A(Z) + (\nabla_X A)(Z)A(Y)].$$

Hence from (4.1) we have

$$(4.2) \quad (\nabla_X A)(Y)A(Z) + (\nabla_X A)(Z)A(Y) + (\nabla_Y A)(X)A(Z) \\ + (\nabla_Y A)(Z)A(X) + (\nabla_Z A)(Y)A(X) + (\nabla_Z A)(X)A(Y) = 0.$$

Contracting  $Y$  and  $Z$  in (4.2) we have

$$(4.3) \quad 2(\nabla_U A)(X) = 0, \quad \text{since } (\nabla_X A)(U) = 0, \\ \text{or, } (\nabla_U A)(X) = 0 \\ \text{or, } g(X, \nabla_U U) = 0 \quad \text{for all } X.$$

Equation (4.3) implies  $\nabla_U U = 0$ . This leads to the following theorem:

**Theorem 4.1.** *If a quasi Einstein manifold satisfies cyclic Ricci tensor, then the integral curves of the vector field  $U$  are geodesic.*

#### 5. The generator $U$ as a Killing vector field

In this section let us consider the generator  $U$  of the manifold is a Killing vector field. Then we have

$$(5.1) \quad (\mathcal{L}_U g)(X, Y) = 0,$$

where  $\mathcal{L}$  denotes the Lie derivative. From which we get

$$(5.2) \quad g(\nabla_X U, Y) + g(X, \nabla_Y U) = 0.$$

Again since  $g(\nabla_X U, Y) = (\nabla_X A)(Y)$ , we get from (5.2)

$$(5.3) \quad (\nabla_X A)(Y) + (\nabla_Y A)(X) = 0 \quad \text{for all } X, Y.$$

Similarly, we have

$$(5.4) \quad (\nabla_X A)(Z) + (\nabla_Z A)(X) = 0 \quad \text{for all } X, Z$$

and

$$(5.5) \quad (\nabla_Z A)(Y) + (\nabla_Y A)(Z) = 0 \quad \text{for all } Z, Y.$$

Now from (1.2) we get

$$(5.6) \quad (\nabla_X S)(Y, Z) + (\nabla_Y S)(Z, X) + (\nabla_Z S)(X, Y) \\ = \{(\nabla_X A)(Y) + (\nabla_Y A)(X)\}A(Z) + \{(\nabla_X A)(Z) + (\nabla_Z A)(X)\}A(Y) \\ + \{(\nabla_Z A)(Y) + (\nabla_Y A)(Z)\}A(X).$$

Putting (5.3), (5.4) and (5.5) in (5.6) we have

$$(\nabla_X S)(Y, Z) + (\nabla_Y S)(Z, X) + (\nabla_Z S)(X, Y) = 0.$$

Thus we can state the following:

**Theorem 5.1.** *If the generator of the quasi Einstein manifold is a Killing vector field then the manifold satisfies cyclic Ricci tensor.*

**6. Quasi Einstein manifold admitting a conformal Killing vector field**

We need the following theorem which is due to Obata [12], [13], [14].

**Theorem A.** *If a complete Riemannian manifold  $M$  of dimension  $n \geq 2$  admits a non-constant function  $\rho$  such that*

$$\nabla \nabla \rho = -c^2 \rho g,$$

where  $c$  is a positive constant, then  $M$  is isometric to a sphere of radius  $\frac{1}{c}$  in  $(n + 1)$ -dimensional Euclidean space.

In this section we suppose that a quasi Einstein manifold admits a conformal Killing vector field [15]. Then we get

$$(6.1) \quad (\mathcal{L}_X g)(Y, Z) = 2\rho g(Y, Z),$$

where  $\rho$  is a non-zero scalar. If  $\rho = \text{constant}$ , then the vector field is called a homothetic vector field. If  $X$  is a conformal Killing vector, then we have [15]

$$(6.2) \quad (\mathcal{L}_X S)(Y, Z) = -(n - 2)(\nabla d\rho)(Y, Z) + \nabla \rho g(Y, Z).$$

Also from (1.2) we have

$$(6.3) \quad (\mathcal{L}_X S)(Y, Z) = a(\mathcal{L}_X g)(Y, Z) + b\mathcal{L}_X \{A(Y)A(Z)\}.$$

Now we suppose that

$$(6.4) \quad \mathcal{L}_X \{A(Y)A(Z)\} = \frac{2\rho}{n - 1}g(Y, Z) \quad \text{and} \quad a > r.$$

Using (6.1), (6.2) and (6.4) in (6.3) we obtain

$$(6.5) \quad -(n - 2)(\nabla d\rho)(Y, Z) + \nabla \rho g(Y, Z) = 2a\rho g(Y, Z) + \frac{2b\rho}{n - 1}g(Y, Z).$$

Contracting (6.5) we have

$$(6.6) \quad \nabla \rho = \frac{(r - a)n}{n - 1}\rho.$$

Putting the value of  $\nabla \rho$  in (6.5) we get

$$(6.7) \quad (\nabla d\rho)(Y, Z) = -\frac{(a - r)}{n - 1}\rho g(Y, Z).$$

Now using the theorem of Obata from (6.7) we see that a complete quasi Einstein manifold of dimension  $n \geq 2$  is isometric to a sphere.

But a sphere is an Einstein manifold. Hence we can state the following:

**Theorem 6.1.** *A complete quasi Einstein manifold of dimension  $n \geq 2$  can not admit a nonhomothetic conformal vector field.*

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