# Three Characteristic Beltrami System in Even Dimensions (I): $p$-Harmonic Equation 

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Abstract. This paper deals with space Beltrami system with three characteristic matrices in even dimensions, which can be regarded as a generalization of space Beltrami system with one and two characteristic matrices. It is transformed into a nonhomogeneous $p$-harmonic equation $d^{*} A\left(x, d f^{I}\right)=d^{*} B(x, D f)$ by using the technique of outer differential forms and exterior algebra of matrices. In the process, we only use the uniformly elliptic condition with respect to the characteristic matrices. The Lipschitz type condition, the monotonicity condition and the homogeneous condition of the operator $A$ and the controlled growth condition of the operator $B$ are derived.

## 1. Introduction and statement of result

For $\Omega$ a bounded domain in $\mathrm{R}^{n}(n \geq 2)$, we consider a mapping $f=$ $\left(f^{1}, f^{2}, \cdots, f^{n}\right) \in W_{l o c}^{1, n}\left(\Omega, \mathrm{R}^{n}\right)$. The differential $D f(x)=\left(\frac{\partial f^{i}}{\partial x_{j}}\right)_{1 \leq i, j \leq n}$ and its determinant $J_{f}(x)=\operatorname{det} D f(x)$ are, therefore, defined almost everywhere in $\Omega$. Throughout this paper, we assume that $J_{f}(x)$ is non-negative, that is, $f$ is sensepreserving. For a matrix $A$, we define the operator norm of it by $|A|=\sup _{|h|=1}|A h|$.

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Definition. A mapping $f=\left(f^{1}, f^{2}, \cdots, f^{n}\right) \in W_{l o c}^{1, n}\left(\Omega, \mathrm{R}^{n}\right)$ is said to be $K$ quasiregular, $1 \leq K<\infty$, if

$$
\begin{equation*}
|D f(x)|^{n} \leq K J_{f}(x) \tag{1.1}
\end{equation*}
$$

for almost every $x \in \Omega$. If, in addition, $f$ is homeomorphic, then it is called $K$ quasiconformal.

The theory of higher dimensional quasiregular mappings began with Yu. G. Rešhetnyak's theorem, stating that they are continuous, discrete and open, if they are nonconstant, see [9]. Development of the analytic theory of quasiregular mappings depends upon advances in PDEs, harmonic analysis and (in dimension 2) complex function theory. The first equation of particular relevance to the theory of quasiregular mappings is the $n$-dimensional Beltrami system

$$
\begin{equation*}
D^{t} f(x) D f(x)=J_{f}^{2 / n}(x) G(x) \tag{1.2}
\end{equation*}
$$

where $D^{t} f(x)$ is the transpose of $D f(x), G(x) \in \mathrm{GL}(n)$ is a positively-defined, symmetric, and determinant 1 matrix, and satisfies the following uniformly elliptic condition

$$
\alpha|\xi|^{2} \leq\langle G(x) \xi, \xi\rangle \leq \beta|\xi|^{2}, \quad \forall \xi \in \mathrm{R}^{n}, 0<\alpha \leq \beta<\infty
$$

Yu.G.Rešhetnyak obtained in [9] the following result: every component function $u=f^{i}, i=1,2, \cdots, n$ of equation (1.2) is a weak solution of the following divergence type elliptic equation (also called $A$-harmonic equation)

$$
\begin{equation*}
\operatorname{div} A(x, \nabla u)=0 \tag{1.3}
\end{equation*}
$$

where $A(x, \xi)=\left\langle G^{-1}(x) \xi, \xi\right\rangle^{\frac{n-2}{2}} G^{-1}(x) \xi$. See also [6] and [7]. The weak solution of (1.3) means that

$$
\begin{equation*}
\int_{\Omega}\langle A(x, \nabla u), \nabla \varphi\rangle d x=0 \tag{1.4}
\end{equation*}
$$

for arbitrary $\varphi \in W_{0}^{1, n}(\Omega)$. By this result, we know that there is a close relationship between space quasiregular mappings and the weak solution of the $A$-harmonic equation (1.3).

Consider the following Beltrami system with two characteristic matrices

$$
\begin{equation*}
D^{t} f(x) H(x) D f(x)=J_{f}^{2 / n}(x) G(x) \tag{1.5}
\end{equation*}
$$

where $H(x), G(x) \in \mathrm{GL}(n)$, and satisfy the following uniformly elliptic condition

$$
\begin{equation*}
\alpha_{1}|\xi|^{2} \leq\langle G(x) \xi, \xi\rangle \leq \beta_{1}|\xi|^{2}, \quad \forall \xi \in \mathrm{R}^{n}, 0<\alpha_{1} \leq \beta_{1}<\infty \tag{1.6}
\end{equation*}
$$

$$
\begin{equation*}
\alpha_{2}|\eta|^{2} \leq\langle H(x) \eta, \eta\rangle \leq \beta_{2}|\eta|^{2}, \quad \forall \eta \in \mathrm{R}^{n}, 0<\alpha_{2} \leq \beta_{2}<\infty \tag{1.7}
\end{equation*}
$$

In [5], the authors have derived the following result: let $f \in W_{l o c}^{1, n}\left(\Omega, \mathrm{R}^{n}\right)$ is a generalized solution of (1.5) with the conditions (1.6) and (1.7), then $d u=d f^{i_{1}} \wedge$ $\cdots \wedge d f^{i_{l}}\left(1 \leq i_{1}<\cdots<i_{l} \leq n, 1 \leq l \leq n\right)$ is a weak solution of the $p$-harmonic equation

$$
d^{*} A(x, d u)=d^{*} B(x, D f)
$$

where

$$
\begin{aligned}
A(x, d u) & =\left(\frac{\left\langle G_{\#}^{-1}(x) d u, d u\right\rangle}{H_{\#}^{I}(x)}\right)^{\frac{n-2 l}{2 l}} G_{\#}^{-1}(x) d u \\
& =\left(\frac{\left\langle G_{\#}^{-1}(x) d u, d u\right\rangle}{H_{\#}^{I}(x)}\right)^{\frac{p-2}{2}} G_{\#}^{-1}(x) d u, \\
B(x, D f) & =J_{f}(x) D_{\#}^{-1} f(x)\left[H_{\#}^{-1}(x)-H_{\#}^{-1}\left(x_{0}\right)\right] d x^{I},
\end{aligned}
$$

where $H_{\#}^{I}(x)$ is the element of $H_{\#}^{-1}(x)$ lies in the diagonal with respect to $I=$ $\left(i_{1}, \cdots, i_{l}\right)$. The weakly monotonicity result of every component function of (1.5) also been derived. For some related results, see [3], [4] and [12].

Remark 1. Let $n=2$. If we let $z=x_{1}+i x_{2}$ and $w(z)=f^{1}\left(x_{1}, x_{2}\right)+i f^{2}\left(x_{1}, x_{2}\right)$, then equations (1.2) and (1.5) are equivalent separately to the Beltrami equation with characteristic function $\mu(z)$

$$
\begin{equation*}
w_{\bar{z}}=\mu_{1}(z) w_{z}, \quad|\mu(z)| \leq k_{1}<1 \tag{1.8}
\end{equation*}
$$

and the Beltrami equation with two characteristic functions $\mu_{1}(z)$ and $\mu_{2}(z)$

$$
\begin{equation*}
w_{\bar{z}}=\mu_{1}(z) w_{z}+\mu_{2}(z) \overline{w_{z}}, \quad\left|\mu_{1}(z)\right|+\left|\mu_{2}(z)\right| \leq k_{2}<1 \tag{1.9}
\end{equation*}
$$

L.Bers ${ }^{[1]}$, I.N.Vekua ${ }^{[10]}$ and Wen Guochun ${ }^{[11]}$ have made lots of research on the complex equations (1.8) and (1.9), and have obtained some celebrated results on them.

In this paper, we continue to study a more general Beltrami system, i.e., the following Beltrami system with three characteristic matrices

$$
\begin{equation*}
D^{t} f(x) H(x) D f(x)+K(x) D^{t} f(x) D f(x)=J_{f}^{2 / n}(x) G(x) \tag{1.10}
\end{equation*}
$$

where $H(x)$ is a diagonal matrix, $G(x)$ and $K(x)$ are positively-defined, symmetric matrices, and satisfy the following uniformly elliptic conditions

$$
\begin{array}{ll}
\alpha_{1}|\xi|^{2} \leq\langle H(x) \xi, \xi\rangle \leq \beta_{1}|\xi|^{2}, & \forall \xi \in \mathrm{R}^{n}, 0<\alpha_{1} \leq \beta_{1}<\infty, \\
\alpha_{2}|\eta|^{2} \leq\langle G(x) \eta, \eta\rangle \leq \beta_{2}|\eta|^{2}, & \forall \eta \in \mathrm{R}^{n}, 0<\alpha_{2} \leq \beta_{2}<\infty, \tag{1.12}
\end{array}
$$

$$
\begin{equation*}
\alpha_{3}|\zeta|^{2} \leq\langle K(x) \zeta, \zeta\rangle \leq \beta_{3}|\zeta|^{2}, \quad \forall \zeta \in \mathrm{R}^{n}, 0 \leq \alpha_{3} \leq \beta_{3}<\infty . \tag{1.13}
\end{equation*}
$$

Remark 2. It is easy to see that if $f=\left(f^{1}, f^{2}, \cdots, f^{n}\right) \in W_{l o c}^{1, n}\left(\Omega, \mathrm{R}^{n}\right)$ is a generalized solution of (1.10), then $f$ is $\left(\frac{\beta_{2}}{\alpha_{1}+\alpha_{3}}\right)^{n / 2}$-quasiregular.

In fact, we can obtain by (1.10) that for arbitrary $\xi \in \mathrm{R}^{n}$,

$$
\left\langle D^{t} f(x) H(x) D f(x) \xi, \xi\right\rangle+\left\langle K(x) D^{t} f(x) D f(x) \xi, \xi\right\rangle=J_{f}^{2 / n}(x)\langle G(x) \xi, \xi\rangle
$$

By (1.11) $\sim(1.13)$ we can derive that

$$
\begin{aligned}
\left(\alpha_{1}+\alpha_{3}\right)|D f(x) \xi|^{2} & \leq\langle H(x) D f(x) \xi, D f(x) \xi\rangle+\alpha_{3}\langle D f(x) \xi, D f(x) \xi\rangle \\
& \leq\left\langle D^{t} f(x) H(x) D f(x) \xi, \xi\right\rangle+\left\langle K(x) D^{t} f(x) D f(x) \xi, \xi\right\rangle \\
& =J_{f}^{2 / n}(x)\langle G(x) \xi, \xi\rangle \leq \beta_{2} J_{f}^{2 / n}(x)|\xi|^{2}
\end{aligned}
$$

Therefore

$$
|D f(x)|^{2} \leq \frac{\beta_{2}}{\alpha_{1}+\alpha_{3}} J_{f}^{2 / n}(x)
$$

Thus the desired result follows.
Remark 3. Combined with the result in Remark 2 and the regularity result for $K$ quasiregular mappings in [11], we can obtain the regularity result for the generalized solution of (1.10): There exists $p=p(n, K)>n$, such that any generalized solution of $(1.10)$ is in fact in the space $W_{\text {loc }}^{1, p}\left(\Omega, \mathrm{R}^{n}\right)$.

The system (1.10) has a strong background when $n=2$ and has been widely studied. But for the case when $n>2$, it is very difficult to study it directly since the system (1.10) is overdetermined and fully nonlinear. An effective research method is to transform it to an elliptic equation, and study the elliptic equaiton by using the method of differential geometry and the analytical method of Sobolev spaces. The aim of this paper is to transform system (1.10) to a nonhomogeneous $p$-equation, and derive the estimates of the operators $A$ and $B$. This make a bridge between system (1.10) and nonhomogeneous elliptic equations. We obtain
Theorem. Suppose that the space dimension $n$ is even and $f \in W_{l o c}^{1, n}\left(\Omega, R^{n}\right)$ is a generalized solution of (1.10), which satisfy conditions (1.11) ~ (1.13), then $d f^{I}=d f^{i_{1}} \wedge \cdots \wedge d f^{i_{l}}, 1 \leq i_{1}<i_{2}<\cdots<i_{l} \leq n, l=\frac{n}{2}$, is a weak solution of the following $p$-harmonic equation

$$
\begin{equation*}
d^{*} A\left(x, d f^{I}\right)=d^{*} B(x, D f), \tag{1.14}
\end{equation*}
$$

where

$$
\begin{aligned}
A\left(x, d f^{I}\right)= & G_{\#}^{-1}(x) d f^{I} \\
B(x, D f)= & G_{\#}^{-1}(x) D_{\#}^{t} f(x) H_{\#}(x)\left(H_{\#}^{-1}(x)\right. \\
& \left.-H_{\#}^{-1}\left(x_{0}\right)\right) d x^{I}-G_{\#}^{-1}(x) K_{\#}(x) D_{\#}^{t} f(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}
\end{aligned}
$$

The operator A satisfies the Lipschitz type condition

$$
\begin{equation*}
\left|A\left(x, h_{1}\right)-A\left(x, h_{2}\right)\right| \leq C_{1}\left|h_{1}-h_{2}\right|, \quad C_{1}=\frac{1}{\alpha_{3}^{l}} \tag{1.15}
\end{equation*}
$$

the monotonicity condition

$$
\begin{equation*}
\left\langle A\left(x, h_{1}\right)-A\left(x, h_{2}\right), h_{1}-h_{2}\right\rangle \geq C_{2}\left|h_{1}-h_{2}\right|^{2}, \quad C_{2}=\frac{1}{\beta_{3}^{l}} \tag{1.16}
\end{equation*}
$$

and the homogeneous condition

$$
\begin{equation*}
A(x, \lambda h)=\lambda A(x, h), \quad \forall \lambda \in R \tag{1.17}
\end{equation*}
$$

and the operator $B$ satisfies the controlled growth condition

$$
\begin{equation*}
|B(x, h)| \leq C_{3}|h|, \quad C_{3}=\left(\frac{2\left(\beta_{1}^{l}+\beta_{3}^{l}\right)}{\alpha_{2}^{l}}\right)^{\frac{1}{2}} \cdot \frac{1}{\alpha_{1}^{l}}+\frac{1}{\alpha_{3}^{l}} \tag{1.18}
\end{equation*}
$$

Remark 4. We can derive some useful properties of the generalized solutions of (1.10) by using the $p$-harmonic equation (1.14). These results will be published in other papers.

## 2. Preliminaries

Let $e^{1}, e^{2}, \cdots, e^{n}$ denote the standard unit basis of $\mathrm{R}^{n}$. For $l=0,1, \cdots, n$, the linear space of $l$ differential forms, spanned by the exterior products $e^{I}=$ $e^{i_{1}} \wedge e^{i_{2}} \wedge \cdots \wedge e^{i_{l}}$ corresponding to all ordered $l$-tuples $I=\left(i_{1}, i_{2}, \cdots, i_{l}\right), 1 \leq i_{1}<$ $i_{2} \cdots<i_{l} \leq n$, is denoted by $\wedge^{l}=\wedge^{l}\left(\mathrm{R}^{n}\right)$. Thus, $\wedge^{0}\left(\mathrm{R}^{n}\right)=\mathrm{R}$ and $\wedge^{1}=\mathrm{R}^{n}$. We define the Hodge star operator

$$
*: \wedge^{l}\left(\mathrm{R}^{n}\right) \rightarrow \wedge^{n-l}\left(\mathrm{R}^{n}\right)
$$

by the rule

$$
* 1=e^{1} \wedge e^{2} \wedge \cdots \wedge e^{n}
$$

and

$$
\alpha \wedge * \beta=\beta \wedge * \alpha=\langle\alpha, \beta\rangle(* 1)
$$

for all $\alpha, \beta \in \wedge^{l}, l=1,2, \cdots n$. The norm of $\alpha \in \wedge^{l}$ is then

$$
|\alpha|^{2}=\langle\alpha, \alpha\rangle=*(\alpha \wedge * \alpha) \in \wedge^{0}=\mathrm{R} .
$$

The Hodge star is an isometric isomorphism on $\wedge^{l}$ with

$$
*: \wedge^{l} \rightarrow \wedge^{n-l}
$$

and

$$
* *=(-1)^{l(n-l)}: \wedge^{l} \rightarrow \wedge^{l} .
$$

A differential $l$-form $\omega$ on $\Omega$ is simply a Schwarz distribution on $\Omega$ with values in $\wedge^{l}=\wedge^{l}\left(\mathrm{R}^{n}\right)$. We write $\omega \in D^{\prime}\left(\Omega, \wedge^{l}\right)$. Therefore every $l$-from $\omega$ may be written uniquely as

$$
\alpha(x)=\sum_{1 \leq i_{1}<\cdots<i_{l} \leq n} \alpha_{i_{1}, \cdots, i_{l}}(x) d x_{i_{1}} \wedge \cdots \wedge d x_{i_{l}},
$$

where the coefficients $\alpha_{i_{1}, \cdots, i_{l}}(x)$ are functions or distributions.
Of fundamental concern to us will be the exterior derivative

$$
d: D^{\prime}\left(\Omega, \wedge^{l}\right) \rightarrow D^{\prime}\left(\Omega, \wedge^{l+1}\right)
$$

it is uniquely determined by the following three conditions:
(i) if $l=0$, then $d f$ is the differential of $f$;
(ii) for $\alpha \in D^{\prime}\left(\Omega, \wedge^{l}\right), \beta \in D^{\prime}\left(\Omega, \wedge^{k}\right)$, we have

$$
d(\alpha \wedge \beta)=d \alpha \wedge \beta+(-1)^{l} \alpha \wedge d \beta
$$

(iii) $d(d \alpha)=0$.

The formal adjoint ( also called the Hodge codifferential ) of $d$ is the operator

$$
d^{*}=(-1)^{n l+1} * d *: D^{\prime}\left(\Omega, \wedge^{l+1}\right) \rightarrow D^{\prime}\left(\Omega, \wedge^{l}\right)
$$

For any forms $\alpha \in L^{p}\left(\Omega, \wedge^{l}\right), \beta \in L^{q}\left(\Omega, \wedge^{l}\right), 1 \leq p, q \leq \infty, \frac{1}{p}+\frac{1}{q}=1$, define

$$
(\alpha, \beta)=\int_{\Omega}\langle\alpha(x), \beta(x)\rangle d x
$$

Definition. Let $\beta \in L^{q}\left(\Omega, \wedge^{l}\right)$. We call the codifferential of $\beta$ is zero in the weak sense, if for any $\alpha \in L_{1}^{p}\left(\Omega, \wedge^{l-1}\right), 1 \leq p, q<\infty, \frac{1}{p}+\frac{1}{q}=1$ with compact support in $\Omega$, we have

$$
(d \alpha, \beta)=0
$$

We can also define the differential of $\alpha$ is zero in the weak sense in the same manner.

Let $G$ be an $n \times n$ matrix. The $l^{t h}$ exterior power of $G$ is the linear operator

$$
G_{\#}: \wedge^{l} \rightarrow \wedge^{l}
$$

defined by

$$
G_{\#}\left(\alpha_{1} \wedge \alpha_{2} \wedge \cdots \wedge \alpha_{l}\right)=G \alpha_{1} \wedge G \alpha_{1} \wedge \cdots \wedge G \alpha_{l}
$$

for $\alpha_{1}, \alpha_{2}, \cdots, \alpha_{l} \in \wedge^{1}\left(\mathrm{R}^{n}\right)$ and then extend linearly to all of $\wedge^{l}$. Thus $G_{\#}$ is a $C_{n}^{l} \times C_{n}^{l}$ matrix.

For the properties of $G_{\#}$ see [7] and [8].
Let $f: \Omega \rightarrow \mathrm{R}^{n}, f=\left(f^{1}, f^{2}, \cdots, f^{n}\right) \in W_{l o c}^{1, p}\left(\Omega, \mathrm{R}^{n}\right), p \geq 1$. Then $f$ induces a homomorphism

$$
f^{*}: C^{\infty}\left(\mathrm{R}^{n}, \wedge^{l-1}\right) \rightarrow L_{l o c}^{1, p}\left(\Omega, \wedge^{l-1}\right)
$$

called the pull back. More precisely, let $\alpha \in C^{\infty}\left(\mathrm{R}^{n}, \wedge^{l-1}\right), \alpha=\sum_{I} \alpha^{I} d x^{I}$. Then

$$
\left(f^{*} \alpha\right)(x)=\sum_{I} \alpha^{I}(f(x)) d f^{i_{1}} \wedge d f^{i_{2}} \wedge \cdots \wedge d f^{i_{l}}
$$

For the properties of the pull back operator see $[6] \sim[8]$.
Other unstated symbol in this paper can be found in [6]~[8].

## 3. $p$-harmonic equation for the Beltrami system (1.10)

From the system (1.10) we can obtain

$$
D^{-1} f(x)=J_{f}^{-2 / n}(x) G^{-1}(x) D^{t} f(x) H(x)+J_{f}^{-2 / n}(x) G^{-1}(x) K(x) D^{t} f(x)
$$

take the $l^{\text {th }}$ exterior power in the above equality, we derive that
$D_{\#}^{-1} f(x)=J_{f}^{-2 l / n}(x) G_{\#}^{-1}(x) D_{\#}^{t} f(x) H_{\#}(x)+J_{f}^{-2 l / n}(x) G_{\#}^{-1}(x) K_{\#}(x) D_{\#}^{t} f(x): \wedge^{l} \rightarrow \wedge^{l}$.
Hence

$$
\begin{align*}
J_{f}(x) D_{\#}^{-1} f(x)= & J_{f}^{(n-2 l) / n}(x) G_{\#}^{-1}(x) D_{\#}^{t} f(x) H_{\#}(x)  \tag{3.1}\\
& +J_{f}^{(n-2 l) / n}(x) G_{\#}^{-1}(x) K_{\#}(x) D_{\#}^{t} f(x): \wedge^{l} \rightarrow \wedge^{l}
\end{align*}
$$

Since $f(x)$ is differentiable almost everywhere, it is no loss of generality to assume $J_{f}(x) \neq 0, x$ is a differentiable point of $f(x)$, and a Lebesgue point of $J_{f}(x)$. Then by [2] we know that there exists a neighborhood $U$ of $x$, such that $f$ is homeomorphic on $U$. Let $V=f(U)$ and $d x^{I}=d x^{i_{1}} \wedge \cdots \wedge d x^{i_{l}}$. Fix $x_{0} \in \Omega$. By (3.1), we have

$$
\begin{align*}
& d^{*}\left[J_{f}(x) D_{\#}^{-1} f(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}\right]  \tag{3.2}\\
= & d^{*}\left[J_{f}^{(n-2 l) / n}(x) G_{\#}^{-1}(x) D_{\#}^{t} f(x) H_{\#}(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}\right] \\
& +d^{*}\left[J_{f}^{(n-2 l) / n}(x) G_{\#}^{-1}(x) K_{\#}(x) D_{\#}^{t} f(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}\right] .
\end{align*}
$$

Our nearest goal is to prove that the left hand side in the above equality is zero in the weak sense. To this end, take an arbitrary

$$
\psi(\zeta)=\sum_{J} \psi^{J}(\zeta) d \zeta^{J}=\sum_{J} \psi^{J}(\zeta) d \zeta^{j_{1}} \wedge \cdots \wedge d \zeta^{j_{l-1}} \in C_{0}^{\infty}\left(V, \wedge^{l-1}\right)
$$

Let

$$
\varphi(x)=\left(f^{*} \psi\right)(x)=\sum_{j} \psi^{J}(f(x)) d f^{j_{1}} \wedge \cdots \wedge d f^{j_{l-1}}
$$

Then by

$$
d f^{j} \wedge d f^{j_{1}} \wedge \cdots \wedge d f^{j_{l-1}}=D_{\#}^{t} f(x) d x^{j} \wedge d x^{j_{1}} \wedge \cdots \wedge d x^{j_{l-1}}
$$

(For this formula see [7] or [8]) we have

$$
\begin{aligned}
d \varphi(x) & =\sum_{J} d \psi^{J}(f(x)) \wedge d f^{j_{1}} \wedge \cdots \wedge d f^{j_{l-1}} \\
& =\left.\sum_{J} \sum_{j=1}^{n} \frac{\partial \psi^{J}(\zeta)}{\partial \zeta^{j}}\right|_{\zeta=f(x)} d f^{j} \wedge d f^{j_{1}} \wedge d f^{j_{l-1}} \\
& =D_{\#}^{t} f(x)\left(\left.\sum_{J} \sum_{j=1}^{n} \frac{\partial \psi^{J}(\zeta)}{\partial \zeta_{j}}\right|_{\zeta=f(x)} \mathrm{d} x^{j}\right) \wedge d x^{j_{1}} \wedge \cdots \wedge d x^{j_{l-1}}
\end{aligned}
$$

Therefore

$$
\begin{aligned}
& \int_{U}\left\langle J_{f}(x) D_{\#}^{-1} f(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}, d \varphi(x)\right\rangle d x \\
= & \int_{U}\left\langle H_{\#}^{-1}\left(x_{0}\right) d x^{I},\left(D_{\#}^{t} f(x)\right)^{-1} d \varphi(x)\right\rangle J_{f}(x) d x \\
= & \int_{U}\left\langle H_{\#}^{-1}\left(x_{0}\right) d x^{I},\left.\sum_{J} \sum_{j=1}^{n} \frac{\partial \psi^{J}(\zeta)}{\partial \zeta^{j}}\right|_{\zeta=f(x)} d x^{j} \wedge d x^{j_{1}} \wedge \cdots \wedge d x^{j_{l-1}}\right\rangle J_{f}(x) d x \\
= & \int_{V}\left\langle H_{\#}^{-1}\left(x_{0}\right) d \zeta^{I}, \sum_{J} \sum_{j=1}^{n} \frac{\partial \psi^{J}(\zeta)}{\partial \zeta^{j}} d \zeta^{j} \wedge d \zeta^{j_{1}} \wedge \cdots \wedge d \zeta^{j_{l-1}}\right\rangle d \zeta \\
= & \int_{V}\left\langle H_{\#}^{-1}\left(x_{0}\right) d \zeta^{J}, d \psi(\zeta)\right\rangle d \zeta=\int_{V}\left\langle d^{*}\left[H_{\#}^{-1}\left(x_{0}\right) d \zeta^{J}\right], \psi(\zeta)\right\rangle d \zeta=0
\end{aligned}
$$

The last equality holds since the differential form $H_{\#}^{-1}\left(x_{0}\right) d \zeta^{J}$ has constant coefficients. Thus we obtain from (3.2) that

$$
\begin{aligned}
& d^{*}\left[J_{f}^{(n-2 l) / n}(x) G_{\#}^{-1}(x) D_{\#}^{t} f(x) H_{\#}(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}\right] \\
= & -d^{*}\left[J_{f}^{(n-2 l) / n}(x) G_{\#}^{-1}(x) K_{\#}(x) D_{\#}^{t} f(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}\right]
\end{aligned}
$$

in the weak sense.
Let $d f^{I}=d f^{i_{1}} \wedge \cdots \wedge d f^{i_{l}}$. By $d f^{I}=D_{\#}^{t} f(x) d x^{I}$, we have

$$
\begin{aligned}
& d^{*}\left\{J_{f}^{(n-2 l) / n}(x) G_{\#}^{-1}(x)\left[d f^{I}+D_{\#}^{t} f(x)\left(H_{\#}(x) H_{\#}^{-1}\left(x_{0}\right)-\operatorname{Id}_{\#}\right) d x^{I}\right]\right\} \\
= & -d^{*}\left[J_{f}^{(n-2 l) / n}(x) G_{\#}^{-1}(x) K_{\#}(x) D_{\#}^{t} f(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}\right] .
\end{aligned}
$$

That is

$$
\begin{align*}
& d^{*}\left\{J_{f}^{(n-2 l) / n}(x) G_{\#}^{-1}(x) d f^{I}\right\}  \tag{3.3}\\
= & d^{*}\left\{J_{f}^{(n-2 l) / n}(x) G_{\#}^{-1}(x) D_{\#}^{t} f(x) H_{\#}(x)\left(H_{\#}^{-1}(x)-H_{\#}^{-1}\left(x_{0}\right)\right) d x^{I}\right\} \\
& -d^{*}\left[J_{f}^{(n-2 l) / n}(x) G_{\#}^{-1}(x) K_{\#}(x) D_{\#}^{t} f(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}\right] .
\end{align*}
$$

If the dimension $n$ is even, then take $l=\frac{n}{2}$, (3.3) becomes

$$
\begin{aligned}
d^{*}\left\{G_{\#}^{-1}(x) d f^{I}\right\}= & d^{*}\left\{G_{\#}^{-1}(x) D_{\#}^{t} f(x) H_{\#}(x)\left(H_{\#}^{-1}(x)-H_{\#}^{-1}\left(x_{0}\right)\right) d x^{I}\right\} \\
& -d^{*}\left[G_{\#}^{-1}(x) K_{\#}(x) D_{\#}^{t} f(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}\right]
\end{aligned}
$$

This is nothing but (1.14).

## 4. The estimates for the operators $A$ and $B$

In this section, we derive the estimates (1.15) (1.18) for the operators $A$ and $B$. We first derive the estimates for the operator

$$
A(x, h)=G_{\#}^{-1}(x) h
$$

(i) the Lipschitz type condition

$$
\left|A\left(x, h_{1}\right)-A\left(x, h_{2}\right)\right| \leq C_{1}\left|h_{1}-h_{2}\right|, \quad C_{1}=\frac{1}{\alpha_{3}^{l}}
$$

By (1.13), we know that

$$
\begin{equation*}
\frac{1}{\beta_{3}^{l}}|\zeta|^{2} \leq\left\langle G_{\#}^{-1}(x) \zeta, \zeta\right\rangle \leq \frac{1}{\alpha_{3}^{l}}|\zeta|^{2} \tag{4.1}
\end{equation*}
$$

Therefore

$$
\left|G_{\#}^{-1}(x)\right| \leq \frac{1}{\alpha_{3}^{l}}
$$

Hence

$$
\left|A\left(x, h_{1}\right)-A\left(x, h_{2}\right)\right| \leq\left|G_{\#}^{-1}(x)\right|\left|h_{1}-h_{2}\right| \leq \frac{1}{\alpha_{3}^{l}}\left|h_{1}-h_{2}\right| .
$$

(ii) the monotonicity inequality

$$
\left\langle A\left(x, h_{1}\right)-A\left(x, h_{2}\right), h_{1}-h_{2}\right\rangle \geq C_{2}\left|h_{1}-h_{2}\right|^{2}, \quad C_{2}=\frac{1}{\beta_{3}^{l}} .
$$

By (4.1) we can derive that

$$
\begin{aligned}
& \left\langle G_{\#}^{-1}(x) h_{1}-G_{\#}^{-1}(x) h_{2}, h_{1}-h_{2}\right\rangle \\
= & \left\langle G_{\#}^{-1}(x)\left(h_{1}-h_{2}\right), h_{1}-h_{2}\right\rangle \geq \frac{1}{\beta_{3}^{l}}\left|h_{1}-h_{2}\right|^{2}
\end{aligned}
$$

(iii) the homogeneous condition

$$
A(x, \lambda h)=\lambda A(x, h), \quad \forall \lambda \in \mathrm{R} .
$$

This can be easily derived by the definition of $A(x, h)$.

We next derive the controlled growth condition (1.18) of the operator $B$. By (3.2), we know that

$$
B(x, D f)=A\left(x, d f^{I}\right)+J_{f}(x) D_{\#}^{-1} f(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I},
$$

and from the Lipschitz type condition (1.15), we know that

$$
\begin{equation*}
\left|A\left(x, d f^{I}\right)\right| \leq C_{1}\left|d f^{I}\right|, \quad C_{1}=\frac{1}{\alpha_{3}^{l}} \tag{4.2}
\end{equation*}
$$

Hence we need only to estimate $\left|J_{f}(x) D_{\#}^{-1} f(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}\right|$.
By (1.10) we have

$$
H(x)+\left(D^{t} f(x)\right)^{-1} K(x) D^{t} f(x)=J_{f}^{2 / n}(x)\left(D^{t} f(x)\right)^{-1} G(x)(D f(x))^{-1}
$$

Take the $l^{t h}\left(l=\frac{n}{2}\right)$ exterior power, we have

$$
H_{\#}(x)+\left(D_{\#}^{t} f(x)\right)^{-1} K_{\#}(x) D_{\#}^{t} f(x)=J_{f}(x)\left(D_{\#}^{t} f(x)\right)^{-1} G_{\#}(x)\left(D_{\#} f(x)\right)^{-1} .
$$

Thus for any $\eta \in \mathrm{R}^{C_{2 l}^{l}}, \eta \neq 0$, we have

$$
\begin{align*}
& \left\langle H_{\#}(x) \eta, \eta\right\rangle+\left\langle K_{\#}(x) D_{\#}^{t} f(x) \eta,\left(D_{\#} f(x)\right)^{-1} \eta\right\rangle  \tag{4.3}\\
= & J_{f}(x)\left\langle G_{\#}(x) D_{\#}^{-1} f(x) \eta, D_{\#}^{-1} f(x) \eta\right\rangle .
\end{align*}
$$

By the uniformly elliptic conditions (1.11) (1.13), we know that

$$
\begin{align*}
& \alpha_{1}^{l}|\tau|^{2} \leq\left\langle H_{\#}(x) \tau, \tau\right\rangle \leq \beta_{1}^{l}|\tau|^{2}  \tag{4.4}\\
& \alpha_{2}^{l}|\tau|^{2} \leq\left\langle G_{\#}(x) \tau, \tau\right\rangle \leq \beta_{2}^{l}|\tau|^{2}  \tag{4.5}\\
& \alpha_{3}^{l}|\tau|^{2} \leq\left\langle K_{\#}(x) \tau, \tau\right\rangle \leq \beta_{3}^{l}|\tau|^{2} . \tag{4.6}
\end{align*}
$$

Therefore

$$
\begin{equation*}
\left\langle K_{\#}(x) D_{\#}^{t} f(x) \eta,\left(D_{\#} f(x)\right)^{-1} \eta\right\rangle \leq \beta_{3}^{l}\left\langle D_{\#}^{t} f(x) \eta,\left(D_{\#} f(x)\right)^{-1} \eta\right\rangle=\beta_{3}^{l}|\eta|^{2} \tag{4.7}
\end{equation*}
$$

and

$$
\begin{equation*}
\alpha_{2}^{l}\left|D_{\#}^{-1} f(x) \eta\right|^{2} \leq\left\langle G_{\#}(x) D_{\#}^{-1} f(x) \eta, D_{\#}^{-1} f(x) \eta\right\rangle \leq \beta_{2}^{l}\left|D_{\#}^{-1} f(x) \eta\right|^{2} \tag{4.8}
\end{equation*}
$$

Thus, by (4.3) (4.8) we can derive that

$$
\alpha_{2}^{l} J_{f}(x)\left|D_{\#}^{-1} f(x) \eta\right|^{2} \leq\left(\beta_{1}^{l}+\beta_{3}^{l}\right)|\eta|^{2}
$$

That is

$$
\begin{equation*}
J_{f}(x)\left|D_{\#}^{-1} f(x)\right|^{2} \leq \frac{\beta_{1}^{l}+\beta_{3}^{l}}{\alpha_{2}^{l}} \tag{4.9}
\end{equation*}
$$

By using the system (1.10) again, we have

$$
J_{f}(x) H_{\#}^{-1}(x)=D_{\#} f(x) G_{\#}^{-1}(x) D_{\#}^{t} f(x)+D_{\#} f(x) G_{\#}^{-1}(x) K_{\#}(x) D_{\#}^{t} f(x) H_{\#}^{-1}(x)
$$

Thus

$$
\begin{aligned}
& J_{f}(x)\left\langle H_{\#}^{-1}(x) d x^{I}, d x^{I}\right\rangle \\
= & \left\langle D_{\#} f(x) G_{\#}^{-1}(x) D_{\#}^{t} f(x) d x^{I}, d x^{I}\right\rangle+\left\langle D_{\#} f(x) G_{\#}^{-1}(x) K_{\#}(x) D_{\#}^{t} f(x) H_{\#}^{-1}(x) d x^{I}, d x^{I}\right\rangle .
\end{aligned}
$$

Since the matrix $H(x)$ is diagonal, so is $H^{-1}(x)$ and $H_{\#}^{-1}(x)$. Therefore

$$
\begin{align*}
& J_{f}(x)\left(H_{\#}^{-1}(x)\right)^{I}  \tag{4.10}\\
= & \left\langle G_{\#}^{-1}(x) D_{\#}^{t} f(x) d x^{I}, D_{\#}^{t} f(x) d x^{I}\right\rangle \\
& +\left(H_{\#}^{-1}(x)\right)^{I}\left\langle G_{\#}^{-1}(x) K_{\#}(x) D_{\#}^{t} f(x) d x^{I}, D_{\#}^{t} f(x) d x^{I}\right\rangle \\
= & \left\langle G_{\#}^{-1}(x) d f^{I}, d f^{I}\right\rangle+\left(H_{\#}^{-1}(x)\right)^{I}\left\langle G_{\#}^{-1}(x) K_{\#}(x) d f^{I}, d f^{I}\right\rangle,
\end{align*}
$$

where $\left(H_{\#}^{-1}(x)\right)^{I}$ is the element lies in the diagonal of $H_{\#}^{-1}(x)$ with respect to $I=\left(i_{1}, \cdots, i_{l}\right)$. By the condition (1.11) we know that

$$
0<\frac{1}{\beta_{1}^{l}} \leq\left|\left(H_{\#}^{-1}(x)\right)^{I}\right| \leq \frac{1}{\alpha_{1}^{l}}
$$

Hence from (4.10) we have

$$
\begin{equation*}
J_{f}(x)=\frac{\left\langle G_{\#}^{-1}(x) d f^{I}, d f^{I}\right\rangle}{\left(H_{\#}^{-1}(x)\right)^{I}}+\left\langle G_{\#}^{-1}(x) K_{\#}(x) d f^{I}, d f^{I}\right\rangle \tag{4.11}
\end{equation*}
$$

By (4.11) and (4.4) (4.6), we can easily derive that

$$
\begin{equation*}
J_{f}(x) \leq \frac{\beta_{1}^{l}+\beta_{3}^{l}}{\alpha_{2}^{l}}\left|d f^{I}\right|^{2} \tag{4.12}
\end{equation*}
$$

Combining (4.9) with (4.12), we have

$$
J_{f}^{2}(x)\left|D_{\#}^{-1} f(x)\right|^{2} \leq \frac{2\left(\beta_{1}^{l}+\beta_{3}^{l}\right)}{\alpha_{2}^{l}}\left|d f^{I}\right|^{2}
$$

It follows that

$$
J_{f}(x)\left|D_{\#}^{-1} f(x)\right| \leq\left(\frac{2\left(\beta_{1}^{l}+\beta_{3}^{l}\right)}{\alpha_{2}^{l}}\right)^{\frac{1}{2}}\left|d f^{I}\right|
$$

Hence

$$
\begin{equation*}
\left|J_{f}(x) D_{\#}^{-1} f(x) H_{\#}^{-1}\left(x_{0}\right) d x^{I}\right| \leq\left(\frac{2\left(\beta_{1}^{l}+\beta_{3}^{l}\right)}{\alpha_{2}^{l}}\right)^{\frac{1}{2}} \cdot \frac{1}{\alpha_{1}^{l}}\left|d f^{I}\right| \tag{4.13}
\end{equation*}
$$

The desired controlled growth condition of $B$ is obtained by combining (4.2) with (4.13).

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