

## <sup>11</sup>B Nuclear Magnetic Resonance Study of Spin Structures in Terbium Tetraboride

B. J. Mean<sup>1</sup>, K. H. Kang<sup>1</sup>, J. H. Kim<sup>1</sup>, I. N. Hyun<sup>1</sup>, Moohee Lee<sup>1,\*</sup> and B. K. Cho<sup>2</sup>

Department of Physics, Konkuk University, Seoul 143-701, Korea <sup>2</sup>Center for Frontier Materials, Department of Materials Science and Engineering, GIST, Gwangju 500-712, Korea

\*\*Received October 1, 2006\*\*

**Abstract**:  $^{11}\text{B}$  nuclear magnetic resonance (NMR) measurements were performed on the single crystals of TbB<sub>4</sub> to investigate local electronic structure and 4f spin dynamics.  $^{11}\text{B}$  NMR spectrum, Knight shift, spin-lattice and spin-spin relaxation rates were measured down to 4 K at 8 T.  $^{11}\text{B}$  NMR shift and linewidth are huge and strongly temperature dependent due to the 4f moments. In addition, both are proportional to magnetic susceptibility, indicating that the hyperfine field at the boron site originates from the 4f spins of Tb. Below  $T_N$ , the single broad resonance peak of  $^{11}\text{B}$  NMR splits into several peaks reflecting the local magnetic fields due to antiferromagnetic spin arrangements. The longitudinal and the transverse relaxation rates,  $1/T_1$  and  $1/T_2$ , independent of temperature above  $T_N$ , decreases tremendously confirming huge suppression of spin fluctuation below  $T_N$ . Key words: TbB<sub>4</sub>,  $^{11}\text{B}$  NMR, antiferromagnetism, rare-earth tetraborides

Recently, the binary boride compounds have been extensively studied since these systems show a variety of fascinating and intriguing properties ranging from various magnetic phase transitions, heavy fermion behavior, mixed-valence phenomena to superconductivity. Those compounds include MgB<sub>2</sub>, hexaborides and tetraborides. Rare-earth tetraborides RB<sub>4</sub> have been investigated thoroughly over a few decades by using single crystals. Most RB<sub>4</sub> (R = Nd, Sm,

<sup>\*</sup> To whom correspondence should be addressed. E-mail: mhlee@konkuk.ac.kr

Gd, Tb, Dy, Ho, Er, Yb) exhibits antiferromagnetism  $^3$  in the range of  $T_N = 7$  - 57 K. The exceptional compounds are PrB4 with a ferromagnetic order  $^3$  at  $T_C = 25$  K, and CeB<sub>4</sub> and YbB<sub>4</sub> with no order. The magnetic interaction is known to be invoked by Rutherman-Kittel-Kasuya-Yoshida (RKKY) interaction.

TbB<sub>4</sub> is known as a metallic antiferromagnetic compound with  $T_N = 43$  K <sup>4-6</sup> and the effective magnetic moment  $\mu_{eff} = 9.8$   $\mu_B$ . The tetragonal to orthorhombic distortion occurs at 80 K <sup>4</sup> much above  $T_N$ . However, the distortion and relative volume change is reported to be small for TbB<sub>4</sub>. <sup>4</sup> According to neutron diffraction studies, magnetic structure is antiferromagnetic collinear with the moments perpendicular to the tetragonal c-aixs. <sup>5.6</sup> Below  $T_N$  another magnetic phase transition around 22 K shows up, as for DyB<sub>4</sub>, HoB<sub>4</sub> and TmB<sub>4</sub>. <sup>7</sup> Since nuclear magnetic resonance (NMR) measurements sensitively probe local magnetic field and electronic spin dynamics using nuclear spins, we have adopted <sup>11</sup>B NMR measurements to probe the spin configurations as well as local electronic structures and 4*f* spin dynamics in TbB<sub>4</sub>.

Pure single crystals of TbB<sub>4</sub> were synthesized by the solution growth method using Al flux, which is similar to the growth of rare-earth hexaborides. Magnetization and magnetoresistance measurements were carried out to characterize these crystals. Pulsed NMR measurements were carried out with our home-made pulsed spectrometers. The broad spectra were scanned by the point-by-point method changing the spectrometer frequency. The phase-alternating pulse sequences were employed to significantly reduce the electromechanical vibration (ringdown) after pulses. Relaxation rates,  $1/T_1$  and  $1/T_2$ , were measured by the solid-echo pulse sequence. <sup>11</sup>B NMR spectrum, shift,  $1/T_1$  and  $1/T_2$  were measured in the range of 4.0 - 295 K at 8.0 T.

Fig. 1 shows the <sup>11</sup>B NMR spectrum at 295 K (a) and 10 K (b). The spectrum at 295 K, above  $T_N = 43$  K, shows a single broad peak with the negative shift of K = -0.0040% and the large linewidth of  $\Delta v = 0.93$  MHz. The spectrum at 50 K shows a single peak with a much larger shift and a larger linewidth; K = -0.024% and  $\Delta v = 3.6$  MHz, as shown in Fig. 2. The negative shift of <sup>11</sup>B NMR indicates that the shift at the boron sites is dominated by corepolarization field. Fig. 2 confirms that the <sup>11</sup>B NMR shift as well as linewidth is strongly

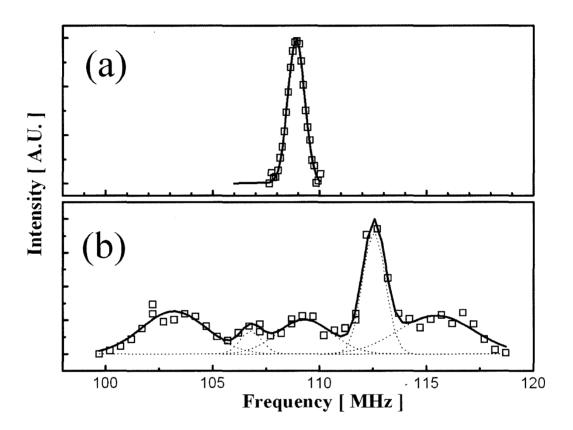


Fig. 1. <sup>11</sup>B NMR spectrum at 8 T; (a) At 295 K above  $T_N$  (b) At 10 K below  $T_N$ . The lines represent a Gaussian fit to respective resonance peaks.

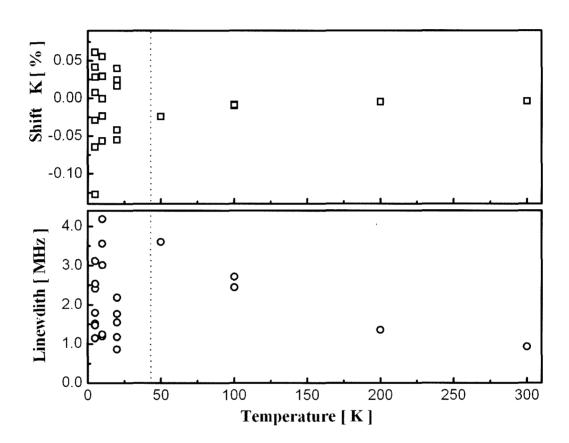


Fig. 2. Shift and linewidth of  $^{11}$ B NMR at 8 T. The dotted line stands for  $T_N$ .

temperature dependent increasing in proportion to the susceptibility down to  $T_N$ . This confirms that the shift and the linewidth of <sup>11</sup>B NMR are dominated by the hyperfine field from the Tb 4f spins.

Fig. 1(b) shows the <sup>11</sup>B NMR spectrum at 10 K below  $T_N = 43$  K. This contrasts with Fig. 1(a) above  $T_N$ . The *single* broad peak at 295 K in Fig. 1(a) splits into *six* even broader resonance peaks in the antiferromagnetic state. Since a magnetic order of electron spins generates local magnetic field in addition to the static external magnetic field for NMR measurements, the splitting of NMR peaks are often a clear signature of magnetic order. Furthermore, the local field, obtained from the respective peak frequency, at the respective nuclear position is utilized to find out spin configurations in the ordered state. In particular, the antiferromagnetic order is often evidenced by doubling up the number of resonance peaks due to the two alternating spin orientations, as shown below  $T_N$  in Fig. 2.

At this point, it should be pointed out that RB<sub>4</sub> has a tetragonal crystal structure with P4/mbm point symmetry at room temperature. Thus the R ions are at the 4g symmetry positions and the boron sublattice forms up chains of B<sub>6</sub> octahedra parallel to the c-axis, which are connected by pairs of boron atoms. Therefore, there are three different boron sites, which are magnetically nonequivalent. The difference is direction and distance to the Tb 4f moments. The temperature evolution of four resonance peaks is crucial to determine spin orientations. However, the detailed analysis using computer simulation is required to pin down correct spin configuration. Far above  $T_N$ , relaxation rates  $1/T_1$  and  $1/T_2$  are very large, respectively,  $\sim 10^4$  and  $\sim 10^5$  s<sup>-1</sup> and temperature-independent suggesting that the spin correlation between Tb 4f spins is weak and the spin fluctuation is independent. Below  $T_N$ , not only  $1/T_1$  but also  $1/T_2$  is significantly reduced to, respectively,  $\sim 10^{-1}$  and  $\sim 10^4$  s<sup>-1</sup> at 5 K, indicating that the spins are ordered and the spin dynamics are suppressed dramatically.

In summary, we have performed <sup>11</sup>B NMR measurements on the single crystals of TbB<sub>4</sub> to investigate local electronic structures and 4f spin dynamics. <sup>11</sup>B NMR shift and linewidth are huge and strongly temperature dependent due to the 4f moments. In addition, both are proportional to the magnetic susceptibility, indicating that the hyperfine field at the boron site

originates from the 4f spins of Tb. Below  $T_N$ , the single broad resonance peak of <sup>11</sup>B NMR splits into several peaks reflecting the local magnetic fields due to antiferromagnetic spin arrangements.  $1/T_1$  and  $1/T_2$ , independent of temperature above  $T_N$ , decrease dramatically below  $T_N$  confirming huge suppression of spin fluctuation.

## **REFERENCES**

- 1. J. Etourneau, J. Less-Common Met. 110, 267 (1985).
- 2. W. Schäfer and G. Will, Z. Kristallogr. 144, 217 (1976).
- 3. J. H. Buschow et al., J. Chem. Phys. 57, 3910 (1972).
- 4. Z. Heiba et al., J. Phys. Chem. Solids 47, 651 (1986).
- 5. G. Will et al., J. Less-Common Met. 82, 349 (1981).
- 6. F. Elf et al., Solid State Commun. 40, 579 (1981).
- 7. Z. Fisk et al., Solid State Commun. 39, 1189 (1981).
- 8. J. S. Rhyee et al., Phys. Rev. B 65, 205112 (2002).