

라그란제 보간을 사용한 비선형 클라인 고든 미분방정식의 수치해

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요 약

비선형 클라인 고든 방정식의 수치해를 구하기 위해 라그란제 보간을 사용하는데 비선형 항을 계산하기 위해 보간식의 차이가 거의 없는 변형된 식을 사용하여 해의 안정성과 해의 수렴성을 밝히고 오차를 분석하였다. 즉 $I(x)^3$ 대신에 $f(x_i)^3 I_i(x)$ 을 사용하였으며 오차는 $c\left(\frac{1}{N}\right)^{N-1} h^{N(N-1)} \left(\frac{N}{2}\right)^{N-1} / \left(\frac{N}{2}\right)!$ 이하임을 보였고 여기서 N은 다항식의 차수이다.

Numerical Solution for Nonlinear Klein-Gordon Equation by Using Lagrange Polynomial Interpolation with a Trick

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ABSTRACT

In this paper, by using Lagrange polynomial interpolation with a trick such that for $f(x)^3$ we shall use $f(x_i)^3 I_i(x)$ instead of $I(x)^3$ where $I(x) = \sum f(x_i) I_i(x)$. We show the convergence and stability and calculate errors. These errors are approximately less than $c\left(\frac{1}{N}\right)^{N-1} h^{N(N-1)} \left(\frac{N}{2}\right)^{N-1} / \left(\frac{N}{2}\right)!$ where N is a polynomial degree.

키워드 : 비선형 클라인고든 미분방정식(Non-linear Klein-Gordon Equation), 라그란제 보간(Lagrange Interpolation)

1. Introduction

The nonlinear Klein Gordon equation

$$\frac{\partial^2 u}{\partial t^2} - \Delta u + V_u(u) = f \tag{1}$$

where Δ is the Laplacian operator in R^d ($d=1,2,3$), $V_u(u)$ is the derivative of the "Newtonian potential function" V , and f is a source term independent of the solution u , in various areas of mathematical physics. Among the particular cases which are the practical relevance, we take $V_u(u) = |u|^\alpha u$ with $\alpha > 0$ (quantum mechanics), refer to[5].

The convergence of the Galerkin finite element method for second order hyperbolic equations has been studied by

many authors : cf. among others Dupont[3], who obtained error estimates for time-discrete and time continuous approximations of linear problems, and Dendy[2], who examined nonlinear problems as well as various modified Galerkin methods. To compute the nonlinear term, the product approximation is used by Yves Tourigny[6]. This approximation is a technique which consists of replacing the nonlinear term by its interpolant in the finite-dimensional subspaces. This provides an interesting alternative to numerical quadrature and greatly eases the implementation of the Galerkin method.

In this paper, by using Lagrange polynomial interpolation with a trick such that let $I(x)$ be an interpolation function with n-node x_i of an arbitrary function $f(x)$, if we need an interpolation for $f(x)^3$, then we shall use $f(x_i)^3 I_i(x)$ instead of $I(x)^3$ where $I(x) = \sum_i f(x_i) I_i(x)$,

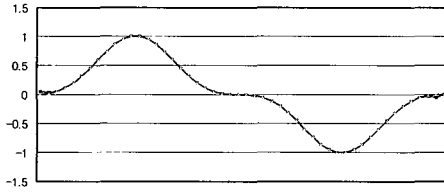
we get the numerical solutions of (1) when $\Delta u = \frac{\partial^2 u}{\partial x^2}$.

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We show an example about $f(x)^3, f(x_i)^3 I_i(x), I(x)^3$, where $f(x) = \sin x$ has 15 nodes in (Figure 1).



(Figure 1) $f(x)^3$ and $I(x)^3$ are same for every point but $f(x)^3 I_i(x)$ has some difference at near the boundary

To show the stability and convergence, we set as many distinct points

$$x_k \quad k \in J \text{ (a set of indices)}$$

in the domain Ω or in its boundary $\partial\Omega$, as the dimension of the space $Pol_N(\Omega)$. At the number of these points, located on $\partial\Omega$, the boundary conditions are imposed. The remaining points are used to enforce the differential equation.

We assume that for any $k \in J$, there exists a polynomial $\phi_k \in Pol_N(\Omega)$, necessarily unique, such that

$$\phi_k(x_m) = \begin{cases} 1 & \text{if } k = m, \\ 0 & \text{if } k \neq m. \end{cases}$$

The ϕ_k 's form a basis for the polynomials of degree N , since $v(x) = \sum_{k \in J} v(x_k) \phi_k(x)$ for all $v \in Pol_N(\Omega)$. Let J be divided into two disjoint subsets J_e and J_b , such that if $k \in J_b$, the x_k 's are on the part $\partial\Omega$ of the boundary. Moreover, let L_N be an approximation to the operator L in which derivatives are taken at the points x_k 's. The polynomial $u^N \in Pol_N(\Omega)$ is a solution that satisfies the equations

$$\begin{cases} L_N u^N(x_k) = f(x_k) & \text{for all } k \in J_e, \\ B u^N(x_k) = 0 & \text{for all } k \in J_b. \end{cases}$$

The unknowns in this method are the values of u^N at the points x_k 's, i.e., the coefficients of u^N with respect to the Lagrange polynomial. We consider a bilinear form $(u, v)_N$ on the space $C^0(\Omega)$ of the functions continuous up to the boundary of Ω by fixing a family of weights w_k and setting

$$(u, v) = \sum_{k \in J} u(x_k) \overline{v(x_k)} w_k.$$

The existence of the basis ensures that $(u, v)_N$ is an inner product on $Pol_N(\Omega)$. Consequently, we define a discrete norm on $Pol_N(\Omega)$ as

$$\|u\|_N = \{(u, u)_N\}^{\frac{1}{2}} \text{ for } u \in Pol_N(\Omega).$$

The basis of ϕ_k 's is orthogonal under the discrete inner product. We make the assumption that the nodes $\{x_k\}$ and the weights $\{w_k\}$ are such that

$$(u, v)_N = (u, v) \text{ for all } u, v \text{ such that } uv \in Pol_{2N-1}(\Omega).$$

In all the applications, this assumption is fulfilled since the x_k 's are the knots of quadrature formulas of Gaussian type.

Let X_N be the space of the polynomials of degree less than or equal to which satisfy the boundary conditions, i.e.,

$$X_N = \{v \in Pol_N(\Omega) \mid Bv(x_k) = 0 \text{ for all } k \in J_b\}.$$

Then this method is equivalently written as

$$\begin{cases} u^N \in X_N \\ (L_N u^N, \phi_k) = (f, \phi_k)_N \text{ for all } k \in J_e. \end{cases}$$

If Y_N is the space spanned by the ϕ_k 's with $k \in J_e$, i.e.,

$$Y_N = \{v \in Pol_N(\Omega) \mid v(x_k) = 0 \text{ for all } k \in J_b\},$$

then can be written as

$$\begin{cases} u^N \in X_N \\ (L_N u^N, v) = (f, v)_N \text{ for all } v \in Y_N. \end{cases}$$

2. Stability

Let Ω be an interval $[-1, 1]$. We would like to approximate the solution of the following problem

$$\frac{\partial^2 u}{\partial t^2} - \frac{\partial^2 u}{\partial x^2} + |u|^a u = f \text{ in } Q = \Omega \times [0, T]$$

$$u(\cdot, 0) = u_0 \text{ and } \left(\frac{\partial u}{\partial t}\right)(\cdot, 0) = u_1 \text{ in } \Omega$$

$$u(-1, t) = u(1, t) = 0 \text{ for } t \in [0, T]$$

where $u_0 \in H_0(\Omega), u_1 \in L(Q)$ and $f \in L(Q)$ are given functions, a small $T > 0$.

The solution $u^N(x, t)$ of the Legendre Tau approx-

imation of this problem is for all $t > 0$ a polynomial of degree N in x , which is zero at $x = \pm 1$ and satisfies the equations,

$$\begin{aligned} & \int_{-1}^1 [u_{tt}^N(x, t) - u_{xx}^N(x, t) + |u^N(x, t)|^\alpha u^N(x, t)]v(x)dx \\ &= \int_{-1}^1 f(x, t)v(x)dx \quad t > 0, \text{ for all } t \in P_{N-2} \\ & \int_{-1}^1 [u^N(x, 0) - u_0(x)]v(x)dx = 0 \\ & \int_{-1}^1 [u_t^N(x, 0) - u_1(x)]v(x)dx = 0 \end{aligned} \quad (3)$$

Let we set $X_N = \{u \in P_N \mid u(-1) = u(1) = 0\}$, $Y_N = P_{N-2}$, and $(u, v) = \int_{-1}^1 u(x)v(x)dx$. For all $u \in X_N$ we have

$$\int_{-1}^1 u_{xx} P_{N-2} u dx = - \int_{-1}^1 u_{xx}^N u dx = \int_{-1}^1 (u_x)^2 dx$$

But, we know that the degree of $|u^N|^\alpha u^N$ is greater than $2N-1$. Here, we shall use the approximation of $|u^N|^\alpha u^N$ in (3). We substitute $I_N |u^N|^\alpha u^N$ instead of $|u^N|^\alpha u^N$ where $I_N : C(\Omega) \rightarrow X_N$ is the interpolation operator.

We shall find the approximate solution $u_N \in X_N$ such that

$$\begin{aligned} & \int_{-1}^1 [u_{tt}^N(x, t) - u_{xx}^N(x, t) + I_N |u^N(x, t)|^\alpha u^N(x, t)]v(x)dx \\ &= \int_{-1}^1 f(x, t)v(x)dx \quad t > 0, \text{ for all } t \in P_{N-2}, \\ & \int_{-1}^1 [u^N(x, 0) - u_0(x)]v(x)dx = 0 \\ & \int_{-1}^1 [u_t^N(x, 0) - u_1(x)]v(x)dx = 0. \end{aligned} \quad (4)$$

Theorem 1. For some $T > 0$,

$$\begin{aligned} & \|P_{N-2} u_t^N(t)\|_{L_\infty(-1,1)}^2 + \|u_x^N(t)\|^2 + (2\beta/P) \|u^N(t)\|_{L^p(-1,1)}^p \\ & \leq \{P_{N-2} u_t^N(0)\|_{L_\infty(-1,1)}^2 + \|u_x^N(0)\|^2 + (2\beta/P) \|u^N(0)\|_{L^p(-1,1)}^p \\ & + \int_0^T \|f(s)\|_{L_\infty(-1,1)}^2 ds\} e^T \end{aligned}$$

proof. Take $v \in Pol_{N-2} u_t^N$, from the left hand side first term in (4)

$$\begin{aligned} & \int_{-1}^1 u_t(x, t) P_{N-2} u_t^N(x, t) dx \\ &= \int_{-1}^1 P_{N-2} u_t(x, t) P_{N-2} u_t^N(x, t) (1-x^2) dx \\ &= (1/2) \frac{d}{dt} \|P_{N-2} u_t^N(t)\|_{L_\infty(-1,1)}^2, \end{aligned}$$

and the second term,

$$\begin{aligned} & - \int_{-1}^1 u_{xx}^N(x, t) P_{N-2} u_t^N(x, t) dx \\ &= \int_{-1}^1 u_{xx}^N(x, t) P_{N-2} \frac{d}{dt} u^N(x, t) dx \\ &= \int_{-1}^1 u_x^N(x, t) \frac{d}{dt} u_x^N(x, t) dx \\ &= (1/2) \frac{d}{dt} \|u_x^T(t)\|^2. \end{aligned}$$

Now, for $p = \alpha + 2$, refer to [4],

$$\begin{aligned} & (1/p) \frac{d}{dt} \|u^N(x, t)\|_{L^p(-1,1)}^p \\ &= \int_{-1}^1 |u^N(x, t)|^\alpha u^N(x, t) u_t^N(x, t) dx \\ &= \int_{-1}^1 |u^N(x, t)|^\alpha u^N(x, t) (1-x^2) P_{N-2} u_t^N(x, t) dx \end{aligned}$$

We can choose the β which satisfies $\|u^N - I_N u^N\|_w^2 = \|u^N - P_N u^N\|_w^2 + \|R_N u^N\|_w^2$.

$$\begin{aligned} & (1/p) \frac{d}{dt} \|u^N(x, t)\|_{L^p(-1,1)}^p \\ & \leq \beta \int_{-1}^1 I_N \{|u^N(x, t)|^\alpha u^N(x, t)\} P_{N-2} u_t^N(x, t) dx \end{aligned}$$

Therefore, from the equation

$$\begin{aligned} & \int_{-1}^1 [u_{tt}^N(x, t) - u_{xx}^N(x, t) + I_N |u^N(x, t)|^\alpha u^N(x, t)] \\ & \quad P_{N-2} u_t^N(x, t) dx \\ &= \int_{-1}^1 f(x, t) P_{N-2} u_t^N(x, t) dx \end{aligned}$$

we obtain

$$\begin{aligned} & (1/2) \frac{d}{dt} \|P_{N-2} u_t^N(t)\|_{L_\infty(-1,1)}^2 + (1/2) \frac{d}{dt} \|u_x^N(t)\|^2 \\ & + (1/p) \beta \frac{d}{dt} \|u^N(t)\|_{L^p(-1,1)}^p \\ & \leq \int_{-1}^1 [u_{tt}^N(x, t) - u_{xx}^N(x, t) + I_N |u^N(x, t)|^\alpha u^N(x, t)] P_{N-2} u_t^N(x, t) dx \\ & \leq (1/2) \|f(t)\|_{L_\infty(-1,1)}^2 + (1/2) \|P_{N-2} u_t^N(x)\|_{L_\infty(-1,1)}^2 \\ & \|P_{N-2} u_t^N(t)\|_{L_\infty(-1,1)}^2 + \|u_x^N(t)\|^2 + (2/p\beta) \|u^N(t)\|_{L^p(-1,1)}^p \\ & \leq \|P_{N-2} u_t^N(0)\|_{L_\infty(-1,1)}^2 + \|u_x^N(0)\|^2 + (2/p\beta) \|u^N(0)\|_{L^p(-1,1)}^p \\ & + \int_0^t \|f(s)\|_{L_\infty(-1,1)}^2 ds + \int_0^t \|P_{N-2} u_t^N(s)\|^2 ds \end{aligned}$$

Applying Gronwall's inequality we complete the proof.

This theorem shows the stability of the approximate solution of u^N for

$$\begin{aligned}
 0 &= \int_{-1}^1 (u^N(x, 0) - u_0(x)) u_{0_x}^N dx \\
 &= - \int_{-1}^1 (u_x^N(x, 0) - u_{0_x}(x)) u_{0_x}^N dx \\
 \int_{-1}^1 u_x^N(x, 0) u_{0_x}^N dx &= \int_{-1}^1 u_{0_x}(x) u_{0_x}(x) u_{0_x}^N dx \\
 &\leq c \int_{-1}^1 u_{0_x}(x) u_{0_x}(x) dx \leq c \|u_0\|_{H_0^1(\Omega)}^2
 \end{aligned}$$

3. Convergence

Let R_N be a projection operator from a dense subspace W of D_B upon X_N , where D_B is a set which satisfies the boundary condition of (2). For each $u \in W$, we further require $R_N u$ to satisfy the exact boundary conditions, i.e.,

$$R_N : W \rightarrow X_N \cap D_B.$$

We define the norm $\|g\|_{E^*} = \sup_{u \in E, u \neq 0} \frac{(g, u)}{\|u\|_E}$

for all $g \in E^*$ that is dual of E .

Let $e(x, t) = u^N(x, t) - R_N u$. We obtain the following theorem.

Theorem 2. Assume that $|u|^a u \in H^1(-1, 1)$.

$$\begin{aligned}
 &\|P_{N-2} e_t(t)\|_{L_w(-1,1)} + \|e_x(t)\|^2 \\
 &\leq (\|P_{N-2} e_t(0)\|_{L_w(-1,1)}^2 + \|e_x(0)\|^2 + M^2 T) e^T \\
 &\leq (\|P_{N-2} e_t(0)\|_{L_w(-1,1)}^2 + c \|e_0\|_{H_0^1(\Omega)}^2 + M^2 T) e^T
 \end{aligned}$$

proof. From (3), we have

$$\begin{aligned}
 \int_{-1}^1 [u_{tt}^N(x, t) - u_{xx}^N(x, t) + I_N |u^N(x, t)|^a u^N(x, t)] v(x) dx &= 0 \\
 t > 0, \text{ for all } v \in P_{N-2}
 \end{aligned}$$

Take $v = e_t(x, t)$

$$\begin{aligned}
 0 &= \int_{-1}^1 [u_{tt}^N - u_{xx}^N + I_N |u^N|^a u^N] - (u_{tt} - u_{xx} + |u|^a u) e_t dx \\
 &= \int_{-1}^1 (u_{tt}^N - R_N u_{tt}^N + R_N u_{tt}^N - u_{tt}) e_t dx \\
 &\quad - \int_{-1}^1 (u_{xx}^N - R_N u_{xx}^N + R_N u_{xx}^N - u_{xx}) e_t dx \\
 &\quad + \int_{-1}^1 (I_N |u^N|^a u^N - |u|^a u) e_t dx
 \end{aligned}$$

We get

$$\begin{aligned}
 &\frac{1}{2} \frac{d}{dt} \|P_{N-2} e_t(t)\|_{L_w(-1,1)}^2 + \frac{1}{2} \frac{d}{dt} \|e_x(t)\|^2 \\
 &= \int_{-1}^1 (u_{tt} - R_N u_{tt}^N) e_t + (R_N u_{tt}^N - u_{xx}) e_t \\
 &\quad + (|u|^a u - I_N |u^N|^a u^N) e_t dx
 \end{aligned}$$

We refer to [1] : For each $v \in H_0^1(-1, 1)$

$$\begin{aligned}
 &(P_{N-2}(u_{tt} - R_N u_{tt}^N), v) \\
 &= (u_{tt} - R_N u_{tt}^N, v) - (u_{tt} - R_N u_{tt}^N, v - P_{N-2} v) \\
 &= ((u_{tt} - R_N u_{tt}^N)_x, (\phi - R_N \phi)_x) - (u_{tt} - R_N u_{tt}^N, v - P_{N-2} v)
 \end{aligned}$$

where ϕ is the only function in $H_0^1(-1, 1)$ satisfying $-\phi_{xx} = v$,

then we obtain,

$$\|P_{N-2}(u_{tt} - R_N u_{tt}^N)\|_{E^*} \leq CN^{1-m} \|u_{tt}\|_{H^{m-2}(-1,1)}.$$

For each $v \in H_0^1(-1, 1)$

$$\begin{aligned}
 &(P_{N-2}(u - R_N u)_{xx}, v) \\
 &= -((u - R_N u)_x, v_x) - ((u - R_N u)_{xx}, v - P_{N-2} v) \\
 &= -((u - R_N u)_x, v_x) - (u_{xx} - P_{N-2} u_{xx}, v - P_{N-2} v)
 \end{aligned}$$

here we have used the fact that both $P_{N-2} u_{xx}$ and $(R_N u)_{xx}$ are orthogonal to $v - P_{N-2} v$. Using the same approximation results as before, we deduce

$$\|P_{N-2}(u - R_N u)_{xx}\|_{E^*} \leq CN^{1-m} \|u\|_{H^{m-2}(-1,1)}.$$

In Legendre approximations, for all $u \in H^m(-1, 1)$

$$\|u - I_N u\|_{H^l(-1,1)} \leq CN^{2l + \frac{1}{2} - m} \|u\|_{H^m(-1,1)}$$

for $0 \leq l \leq m$ with $m > \frac{1}{2}$.

Assume that $|u|^a u \in H^1(-1, 1)$, let $l = 0$. We get

$$\begin{aligned}
 &\| |u|^a u - I_N |u^N|^a u^N \|_{L^2(-1,1)} = \| |u|^a u - I_N |u|^a u \|_{L^2(-1,1)} \\
 &\leq CN^{\frac{1}{2} - m} \| |u|^a u \|_{H^m(-1,1)}
 \end{aligned}$$

We may assume that $m > 2$.

$$\begin{aligned}
 \text{Let } M &= CN^{1-m} \|u_{tt}\|_{H^{m-2}(-1,1)} + CN^{1-m} \|u\|_{H^m(-1,1)} \\
 &+ CN^{-\frac{1}{2}} \| |u|^a u \|_{H^{m-2}(-1,1)}
 \end{aligned}$$

clearly $M \rightarrow 0$ as $N \rightarrow \infty$. From (5)

$$\begin{aligned}
 &\frac{1}{2} \frac{d}{dt} \|P_{N-2} e_t(t)\|_{L_w(-1,1)}^2 + \frac{1}{2} \frac{d}{dt} \|e_x(t)\|^2 \\
 &\leq \frac{1}{2} M^2 + \frac{1}{2} \|P_{N-2} e_t(t)\|_{L_w(-1,1)}^2 \\
 &\|P_{N-2} e_t(t)\|_{L_w(-1,1)}^2 + \|e_x(t)\|^2 \\
 &\leq \|P_{N-2} e_t(0)\|_{L_w(-1,1)}^2 + \|e_x(0)\|^2 \\
 &+ \int_0^t M^2 ds + \int \|P_{N-2} e_t(x, s)\|_{L_w(-1,1)}^2 ds
 \end{aligned}$$

We know that $\|e_x(0)\|^2 \leq c \|e_0\|_{H^1(\Omega)}^2$ and applying Gronwall's inequality we conclude the proof.

4. Numerical results

Set $u^N(x, t) = \sum_{i=0}^N a_i(t) l_i(x)$ where $l_i(x)$ is a N -degree

Lagrange polynomial with $N+1$ nodes as $-1 = x_0 < x_1 < x_2 < \dots < x_N = 1, \alpha = 2$. We substitute $u^N(x, t)$ into (2), we get

$$\begin{aligned} \frac{d^2 a_i(t)}{dt^2} &- (l_0''(x_i) a_0(t) + \dots + l_N''(x_i) a_N(t)) \\ &- |a_i(t)|^2 a_i(t) = f(x_i, t) \\ &i = 0, 1, 2, \dots, N. \end{aligned}$$

in here we use the trick at $|a_i(t)|^2 a_i(t)$.

Applying the boundry condition and the difference equation with

$$\begin{aligned} \frac{d^2 a_i(t)}{dt^2} &= \frac{a_i(t_{j+1}) - 2a_i(t_j) + a_i(t_{j-1}))}{h^2} \\ a_i(t_0) &= 0 \\ a_i(t_1) &= hu_1(x_i) \end{aligned}$$

where h is a mesh size and $t_j = jh$. For one example, let

$$\begin{aligned} f(x, t) &= -2 \sin(\pi x) + |(t-t^2) \sin(\pi x)|^2 (t-t^2) \sin(\pi x) \\ &+ \pi^2 (t-t^2) \sin(\pi x) \end{aligned}$$

$$\frac{\partial u}{\partial t}(x, 0) = \sin(\pi x),$$

then we obtain the numerical solution as follows.

Practically, this example has exact solution such that $(t-t^2) \sin(\pi x)$. We can calculate errors. These errors

are approximately less than $C \left(\frac{1}{N}\right)^{N-1} hN(N-1) \left(\frac{N}{2}\right)^{N-1} /$

$\left(\frac{N}{2}\right)!$ in that $C \left(\frac{1}{N}\right)^{N-1}$ is estimated from M which is

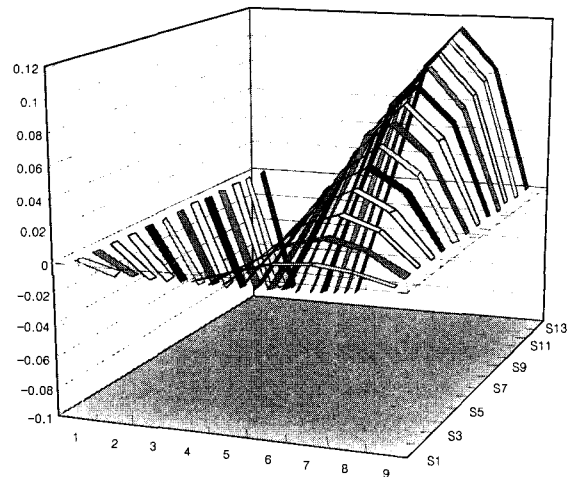
in theorem 2., and $N(N-1) \left(\frac{N}{2}\right)^{N-1} / \left(\frac{N}{2}\right)!$ is calculated by a second order differentiation of N -degree Lagrange polynomial in <Table 1>. Briefly, the errors are less than

$\left(\frac{1}{2}\right)^{N-1} hN(N-1) \left(\frac{N}{2}\right)^{N-1} / \left(\frac{N}{2}\right)!$ and are independent of α .

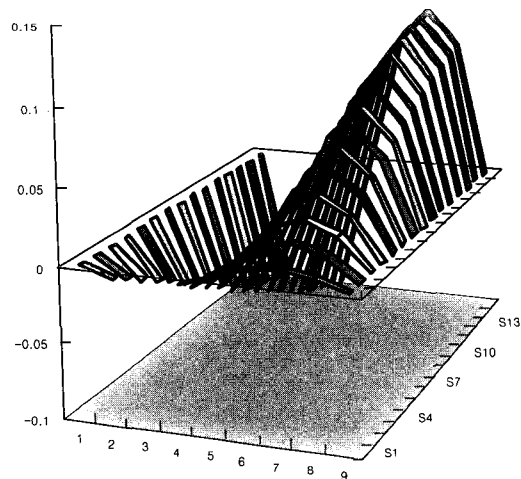
<Table 1> The numerical estimation of $u(1/2, 0.01)$, error and error bound. For time value t we show results by

10th iteration, where $K = \left(\frac{1}{2}\right)^{N-1} hN(N-1) \left(\frac{N}{2}\right)^{N-1} / \left(\frac{N}{2}\right)!$ $h=0.001$

Node	Numerical	Exact	Error	K
N=8	9.913096E-3	9.9E-3	1.3096E-5	1.8229E-5
N=12	9.900053E-3	9.9E-3	5.3E-8	8.9518E-8



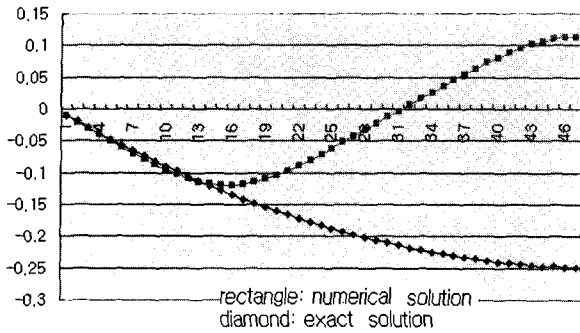
(a) Exact Solution $t \in [0, 0.15]$



(b) Numerical Solution $t \in [0, 0.15]$

(Figure 2) Exact solution and numerical solution when t is from 0 to 0.15 mesh $h=0.001$

Until the time variable t is small, the numerical solution is stable in (Figure 2), but if t is greater than 0.15, this numerical solution is unstable. In (Figure 3), we show that the numerical solution is not stable when t is greater than 0.15. Nevertheless, this numerical solution is not bad for some small t .



(Figure 3) The difference between numerical solution and exact solution is big when t is greater than 0.15

References

[1] Claudio Canuto M. Yousuff Hussaini Alfio Quarteroni Thomas A. Zang, Spectral Methods in Fluid Dynamics. Springer_Verlag, 1988.
 [2] Dendy, J. E., An analysis of some Galerkin schemes for the solution of nonlinear time-dependent problems. SIAM J. Numer. Anal.12, pp.541-565, 1975.
 [3] Dupont, T. L., estimates for Galerkin methods for sec-

ond-order hyperbolic equations. SIAM J. Numer. Anal. 10, pp.392-410, 1973.

[4] Masanori Hosoya and Yoshio Yamada, On Some nonlinear wave equations I : local existence and regularity of solutions. J. Fac. Sci. univ. Tokyo Sect. IA, Math. 38, pp.225-238, 1991.
 [5] Perring, J. K. and Skyrme, T. R. H, A model unified field equation. Nucl. Phys. 31, pp.550-555, 1962.
 [6] Yves Truginy, Product approximation for nonlinear Klein_Gordon equations., IMA journal of Numerical Analysis. 9, pp.449-462, 1990.



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