# Structure, Modified Scaled Quantum Mechanical Force Field and A Priori Prediction of Vibrational Spectra and Their Assignment and Exponential Scaling of Frequencies of Triphenylene 

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#### Abstract

The stnucture, force field and vibrational spectra of triphenylene are studied by B3LYP/6-31G*(5d) level of theory. The results are compared to those of the related system. phenanthrene. The scale factors in nonredundant local coordinates obtained after fitting the DFT frequencies to the experimental numbers of phenanthrene- $\mathrm{d}_{6}$ and $-\mathrm{d}_{10}$ are transferred to predict the spectra and assignment of triphenylene for in-plane modes. The frequencies based on scaling methodology due to Lee $e t$ al. are also obtained. These frequencies are compared with the predicted numbers based on scale factors from phenanthrene. Probable assignment for out-of-plane modes is proposed based on simple scaling of Scott and Radom (scale factor 0.9614 ) as well as by scaling methodology by Lee et al.


Key Words : Density functional calculation. Triphenylene, Vibrational spectra

## Introduction

Triphenylene is an eighteen $\pi$ electron planar tetracyclic fused ring system with six membered rings. This system has been attracting much attention very recently. Ionization spectral bands from UV photoelectron measurements have been studied for triphenylene and its radical cation. ${ }^{1}$ The NMR properties of triphenylene have also been studied as molecular magnets. ${ }^{2.3}$ It is also important in host-guest chemistry as its use has been found in triphenylene scaffold. ${ }^{4}$ Unusual phosphorescence spectra of triphenylene are of particular interest both theoretically and experimentally: Polycyclic aromatic hydrocarbons (PAHs) are the largest known class of chemical carcinogens and mutagens. ${ }^{\text { }}$ In addition. triphenylene-based discotic (colunmar) liquid crystalline materials could be used as device materials. The high charge carrier mobilities photoconductivity and photophysical properties of tripheny lene derivatives are of special interest. ${ }^{6}$
Along with these studies the force field and vibratonal spectra of tripheny lene were also studied. The vibrational spectra of several PAHs and their cations and a few anions have been studied in great detail by several authors because of their importance as the origin of infrared emission bands in the emission bands in the interstellar radiation from many galactic sources. ${ }^{7}$ Presently the vibrational spectroscopy has become an indispensable tool to study various aspects of isolated molecules and clusters. The specific/complete vibrational analysis was done on various systems e.g. water clusters. ${ }^{8}$ sundry weak complexes of benzene and ethene. ${ }^{5}$ calix[4]hydroquinone nanotube bundles ${ }^{10}$ and many other systems. The nature of inter-ring couplings is also the focus of attention in some of the papers.
The first infrared and Raman spectra of triphenylene in the

[^0]crystalline state and solution was reported by Schettino. ${ }^{11}$ The low temperature ligh resolution phosphorescence spectrum and fundamentals deduced from the analysis of it was reported by Nishi and co-workers. ${ }^{12}$ The matrix-isolated IR spectrum in the range $500-3150 \mathrm{~cm}^{-1}$ was reported Hudgins and Sandford ${ }^{13}$ On the theoretical side. the scaled and calculated frequencies were reported by Ling and Lifshitz at the HF level using 3-21G basis set. Some reassignments of the fundamentals of triphenylene were made on the basis of the vibronic activity by QCFF/PI calculation in the phosphorescence spectnm of the disklike moleule. ${ }^{14}$ Later Keszthelyi et al. reported the structure and vibrational spectra based on HF. B-LYP and B3-LYP methods using $6-31 \mathrm{G}$ basis set and revisited the assignment. ${ }^{15}$ Langhoff computed the spectra at the B3LYP/4-31G level. ${ }^{16}$ He reported the scaled IR active vibrational frequencies (scale factor 0.958). The matrix-isolated IR spectrum in the range $500-3150 \mathrm{~cm}^{-1}$ was reported Hudgins and Sandford ${ }^{13}$ Inspite of these studies. there remain some uncertainties in all the empirical assiguments. It is to be noted that in the condensed phase interactions of this PAH with the surrounding medium and with triphnylene molecules leading to clusters can lower the symmetry and affect the force field. thus alter strongly the frequencies and intensities of the vibrational modes. With this view in consideration, the vibrational study on tripheny lene of $C_{2 v}$ symmetry is relevant. The approach isto simulate the slighly perturbed structure from the global minimum $D_{3 h}$ structure so as to reliably assign the experimetal frequencies. thus the calculated frequencies of both the structures will be marginally different.

To the best of the knowledge. no study exits on force field. complete vibrational spectra and assignment of this molecule in terms of non-redundant local coordinates. Here the objective is to predict the spectra from a related standard system. The scale factors of the related standard system. phenanthrene with $C_{2 v}$ symmetry were transferred to predict
and assign the spectra of triphenylene with $\mathrm{C}_{3 x}$ symmetry: The scaled frequencies were compared with the available experimental ones. One of the important aspects of vibrational spectroscopy is to predict accurately the imactive fundamentals $9 \mathrm{~A} 2^{*}$ and $5 \mathrm{~A}_{1}^{*}$. But they were observed due to symmetry lowering in condensed phase though they have very weak intensities. Thus experiments have proven to be difficult in this aspect. Theoretical prediction is of immense value and the predicted values of triphenylene with $\mathrm{C}_{\hat{2}}$ symmetry will accurately simulate the inactive fundamentals in $D_{\text {sh }}$ symmetry. The vibrational frequencies are usually overestimated. Thus there is a need to scale appropriately the frequencies. To this end a modified scaled quantum mechanical approach (SQM) was used in vibrational analysis for a number of hydrocarbons and heterocycles. ${ }^{17-19}$ Here this fitting methodology is also used. Previously a scaling method for frequencies was developed by Lee et $\alpha l$. and this method was successfully applied to s-trans-1,3-butadiene and all-trans-1,3,5.7-octatetraene ${ }^{20.21}$. Here the same scaling approach was used to predict the vibrational spectra. Previously we studied azulene. phenantluene with a view to evaluating inter-ring couplings and evaluate and analyze the transferability of non-redundant local force constants. In continuation of this work the results of triphenylene are presented here.

## Methodology

The initial geometries of triphenylene and phenanthrene with $\mathrm{C}_{\hat{2}}$ symmetries were fully optimized using B3LYP/6$31 \mathrm{G}^{*}$ (5d) method available in Gaussian94. ${ }^{2-}$ The vibrational frequencies and cartesian force constants at the optimized geometries were obtained analytically and transformed to local coordinate space following the recommendation of Pulay et al. ${ }^{23}$ Previously we assigned the fundamentals of phenantlurene and phenanthrene-d $d_{10}$ based on HF/4-21G level calculations. ${ }^{18}$ Here the DFT frequencies are fitted to


Figure 1. Atom Numberings and Intemal Coordinates of Triphenylene.
the same set of experimental numbers of phenanthrene-d $d_{0}$ and- $\mathrm{d}_{10}$ using the modified SQM procedure. The scale factors were transferred from phenanthrene to triphenylene to predict the in-plane modes. The potential energy distribution (PED) of the in-plane normal modes in local coordinates for triphenylene was computed using force constants after scaling. For out-of-plane modes a simple scale factor of 0.9614 as recommended by Scott and Radom ${ }^{-5}$ was used to scale the frequencies. Thus the scaled in-plane force field in non-redundant local coordinates and the assignment were obtained. The scaling methodology of Lee et al is also employed for the present system ${ }^{20.21}$ and the results are compared with those of modified SQM.

## Results and Discussion

The units of various terms are as follows: The vibrational frequencies in this paper are expressed in $\mathrm{cm}^{-1}$. The stretching and their interaction force constants are in mdyne $\mathrm{A}^{-1}$. bending. wagging, torsions and their interactions are in mdyne $\mathrm{A} \mathrm{rad}^{-2}$. Interactions of stretching force constants with bendings. wags and torsions are in mdyne rad ${ }^{-1}$. The structure and internal coordinates of triphenylene are shown in Figure 1. The optimized geometrical parameters are presented in Table 1. The DFT frequencies for phenantlurene-dij and $-\mathrm{d}_{10}$ are fitted to the experimental ones and the fitting reproduces the experimental numbers well with average deviations of 3.1 and $2.5 \mathrm{~cm}^{-1}$ with $\mathrm{C}-\mathrm{H}$ stretches and without $\mathrm{C}-\mathrm{H}$ stretches respectively. The local and symmetry coordinates are given in Tables 2 and 3 , respectively:

The diagonal force constants at the B3LYP/6-31G* level of theory are listed in Table 4. The fundamental frequencies and their assignments of triphenylene along with their potential energy distribution (PED) given in Tables 5 and 6.
(i) Optimized Geometry: The geometrical parameters in Table 1 are considered up to the third decimal place. At this level of precision the geometry reflects the $\mathrm{D}_{3 \mathrm{~h}}$ symmetry: The geometrical parameters $\left(\mathrm{R}_{1.6}\right)$ of triphenylene (1.414.

Table 1. Optimized geometrical parameters of triphenylene at B3LYP/6-31G*(5d). The bond lengths are in $\AA$ and bond angles are in degrees ${ }^{a}$

|  | Theory | Expt. $^{b}$ |  | Theory | Expt. ${ }^{b}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{R}_{1}$ | 1.414 |  | $\alpha_{1}$ | 119.6 |  |
| $\mathrm{R}_{2}$ | 1.422 |  | $\alpha_{2}$ | 122.0 |  |
| $\mathrm{R}_{3}$ | 1.414 |  | $\alpha_{3}$ | 118.4 |  |
| $\mathrm{R}_{4}$ | 1.383 |  | $\alpha_{4}$ | 118.4 |  |
| $\mathrm{R}_{4}$ | 1.401 |  | $\alpha_{5}$ | 122.0 |  |
| $\mathrm{R}_{6}$ | 1.383 |  | $\alpha_{6}$ | 119.6 |  |
| $\mathrm{R}_{7}$ | 1.467 | 1.447 | $\alpha_{7}$ | 120.0 |  |
| $\mathrm{R}_{8}$ | 1.422 |  | $\alpha_{8}$ | 120.0 | 120.0 |
| $\mathrm{R}_{11}$ | 1.467 |  | $\alpha_{19}$ | 122.0 | 122.0 |
| $\mathrm{R}_{1}$ | 1.414 | 1.416 | $\alpha_{20}$ | 119.6 | 119.6 |
| $\mathrm{R}_{18}$ | 1.383 | 1.378 | $\alpha_{23}$ | 118.4 | 118.4 |
| $\mathrm{R}_{19}$ | 1.401 | 1.401 |  |  |  |

${ }^{4}$ For geometrical parameter, see Fig. 1. ${ }^{\text {R Refs. 11, 24. }}$

Table 2. Definition of local coordinates for triphenylene ${ }^{2}$

| Coordinate | Description |
| :---: | :---: |
| In-plane |  |
| $\mathrm{l}_{1-21}=\mathrm{R}_{1-21}$ | CC stretch |
| $1_{22-39}=\mathrm{I}_{1-1}$ | CH stretch |
| $1_{3-45}=\beta_{1-12}$ | CH rock |
| $1_{46}\left(\delta_{1}\right)=\alpha_{1}-\alpha_{4}+\alpha_{3}-\alpha_{4}+\alpha_{4}-\alpha_{6}$ | Ring bending |
| $1_{4}\left(\delta_{2}\right)=2 \alpha_{1}-\alpha_{2}-\alpha_{3}+2 \alpha_{4}-\alpha_{4}-\alpha_{6}$ | Ring bending |
| $1_{48}\left(\delta_{3}\right)=\alpha_{2}-\alpha_{3}+\alpha_{3}-\alpha_{6}$ | Ring bending |
|  | Ring bending ${ }^{2}$ |
| Out-of-plane |  |
| $1_{1-12}=\gamma_{1-8}$ | CH wag |
| $113\left(\tau_{1}\right)=t_{1}-t_{2}+t_{3}-t_{4}+t_{3}-t_{6}$ | Ring Torsion |
| $1_{14}\left(\tau_{2}\right)=2 t_{1}-t_{2}-t_{3}+2 t_{4}-t_{s}-t_{6}$ | Ring Torsion |
| $l_{15}\left(\tau_{3}\right)=t_{2}-t_{3}+t_{s}-t_{s}$ | Ring Torsion |
| $1_{10-24}\left(\tau_{+-\varphi)}\right)$ | Ring Torsion ${ }^{3}$ |
|  | Torsion |
| $\mathrm{l}_{26}\left(\tau_{11}\right)=\mathrm{t}_{1.2 .4 .14}-\mathrm{t}_{\text {S. } 2.4 .6}$ | Torsion |
|  | Torsion |

${ }^{\text {"See Fig. I: the }} \mathrm{CH}$ rock is detined as. $\beta_{\mathrm{h}}=2^{-1}\left(\phi_{1}-\mathrm{f}_{\mathrm{i}}\right) . \gamma_{i}$ numberings are same as $\mathrm{r}_{1},{ }^{1} \delta_{1.6,4} \delta_{: .9}$ and $\delta_{00.12}$ are detined similarly with $\alpha_{7,12}, \alpha_{13.18}$ and $\alpha_{19,24}$ respectively: $\tau_{4.6}, \tau_{7.9}, \tau_{10.12}$ are defined similarly as $\tau_{1.3}$ with $t_{7.15}$. $\mathrm{t}_{12}$-is and $\mathbf{t}_{19}$-主 respectively.
1.422. $1.414,1.383,1.401,1.383$ A) match well with those of phenantlurene $(1.415,1.427,1.415,1.383,1.408,1.381 \mathrm{~A})$. The major differences arise in the case of the sandwitched ring. The values for $\mathrm{R}_{7}$. $\mathrm{R}_{\mathrm{s}}$ and $\mathrm{R}_{11}$ are 1.467. 1.422 and 1.467 A for triphenylene, whereas they are 1.435. 1.340 and 1.458 A for phenanthrene. The calculated geometrical parameters agree well with the experimental values. ${ }^{11-24}$
(ii) Force Field: Triphenylene could be regarded as consisting of a phenanthrene moiety and a benzene ring. Hence we expect similarities between the phenanthrene ring moiety and triphenylene. Based on the diagonal force constants of these two molecules given in Table 4 and some selected interaction force constants, the following inferences could be made regarding the force fields of tripheny lene and phenantlurene.
(1) $\mathrm{C}-\mathrm{H}$ diagonal force constants of triphenylene are larger compared to phenanthrene. Interaction force constants involving $\mathrm{C}-\mathrm{H}$ ( $\sigma$ bonds) are small and only the first neighbours contribute significantly for all the two sy'stems.
(2) In the case of $\mathrm{C}-\mathrm{C}$. the diagonal force conatants of triphenylene are larger than those of phenanthrene although some are similar. The interaction force constants between $\mathrm{C}-\mathrm{C}$ bonds are much larger compared to phenanthrene and similar to phenanthrene they extend over the entire ring.
(3) Diagonal $f_{\beta \beta \beta}$ s are in general higher for triphenylene than for phenanthrene. $\mathrm{C}-\mathrm{H}-\beta_{i}$ interactions are small and negligible in most cases. $\mathrm{C}-\mathrm{C}-\beta_{\mathrm{i}}$ interactions are substantial only for the connecting bonds.
(4) The force constants for ring bendings ( $\delta$ ) are higher for triphenylene compared to phenanthrene.
These conclusions clearly indicate that the force constants of triphenylene are larger than those of phenantlurene although some of them are similar.

Table 3. Detinition of symunetry coordinates for in-plane modes of triphenylene in terms of local coordinates ${ }^{a}$

| $\mathrm{A}_{1}$ | $\mathrm{S}_{1}=2^{-1:}\left(\mathrm{R}_{1}+\mathrm{R}_{12}\right)$ | $\mathrm{B}_{2}$ | $\mathrm{S}_{1}=2^{-12}\left(\mathrm{R}_{1}-\mathrm{R}_{12}\right)$ |
| :---: | :---: | :---: | :---: |
|  | $\mathrm{S}_{2}=2^{-1: 1}\left(\mathrm{R}_{2}+\mathrm{R}_{10}\right)$ |  | $\mathrm{S}_{2}=2^{-12}\left(\mathrm{R}_{2}-\mathrm{R}_{10}\right)$ |
|  | $\mathrm{S}_{3}=2^{-12}\left(\mathrm{R}_{3}+\mathrm{R}_{16}\right)$ |  | $\mathrm{S}_{3}=2^{-1 / 2}\left(\mathrm{R}_{3}-\mathrm{R}_{16}\right)$ |
|  | $\mathrm{S}_{4}=2^{-1 \cdot 2}\left(\mathrm{R}_{4}+\mathrm{R}_{15}\right)$ |  | $\mathrm{S}_{4}=2^{-1 / 2}\left(\mathrm{R}_{4}-\mathrm{R}_{15}\right)$ |
|  | $\mathrm{S}_{5}=2^{-1 / 2}\left(\mathrm{R}_{5}+\mathrm{R}_{14}\right)$ |  | $\mathrm{S}_{5}=2^{-1 / 2}\left(\mathrm{R}_{5}-\mathrm{R}_{14}\right)$ |
|  | $\mathrm{S}_{6}=2^{-1: 2}\left(\mathrm{R}_{6}+\mathrm{R}_{13}\right)$ |  | $\mathrm{S}_{6}=2^{-1 / 2}\left(\mathrm{R}_{6}-\mathrm{R}_{13}\right)$ |
|  | $\mathrm{S}_{7}=2^{-1 / 2}\left(\mathrm{R}_{7}+\mathrm{R}_{9}\right)$ |  | $\mathrm{S}_{7}=2^{-1 / 2}\left(\mathrm{R}_{7}-\mathrm{R}_{9}\right)$ |
|  | $\mathrm{S}_{8}=\mathrm{R}_{3}$ |  | $S_{8}=2^{-1 / 2}\left(R_{17}-R_{11}\right)$ |
|  | $\mathrm{S}_{9}=\mathrm{R}_{11}$ |  | $\mathrm{S}_{9}=2^{-1 / 2}\left(\mathrm{R}_{18}-\mathrm{R}_{20}\right)$ |
|  | $\mathrm{S}_{10}=2^{-1 \cdot 2}\left(\mathrm{R}_{17}+\mathrm{R}_{11}\right)$ |  | $\mathrm{S}_{10}=2^{-1 / 2}\left(\mathrm{r}_{1}-\mathrm{r}_{2}\right)$ |
|  | $\mathrm{S}_{11}=2^{-1: 2}\left(\mathrm{R}_{13}+\mathrm{R}_{20}\right)$ |  | $S_{11}=2^{-1 / 2}\left(r_{3}-r_{4}\right)$ |
|  | $\mathrm{S}_{12}=\mathrm{R}_{19}$ |  | $\mathrm{S}_{11}=2^{-1 / 2}\left(\mathrm{r}_{5}-\mathrm{r}_{6}\right)$ |
|  | $S_{13}=2^{-1: 2}\left(r_{1}+r_{2}\right)$ |  | $\mathrm{S}_{13}=2^{-1 / 2}\left(\mathrm{r}_{5}-\mathrm{r}_{3}\right)$ |
|  | $S_{14}=2^{-1 / 2}\left(r_{3}+r_{4}\right)$ |  | $\mathrm{S}_{14}=2^{-1 / 2}\left(\mathrm{r}_{9}-\mathrm{r}_{10}\right)$ |
|  | $S_{15}=2^{-1:}\left(r_{5}+r_{6}\right)$ |  | $S_{15}=2^{-12}\left(\mathrm{I}_{11}-r_{12}\right)$ |
|  | $S_{16}=2^{-1: 1}\left(r^{-}+r_{3}\right)$ |  | $S_{16}=2^{-12}\left(\beta_{1}+\beta_{11}\right)$ |
|  | $S_{17}=2^{-1: 2}\left(r_{19}+r_{10}\right)$ |  | $S_{17}=2^{-12}\left(\beta_{2}+\beta_{11}\right)$ |
|  | $S_{18}=2^{-1: 7_{( }\left(r_{11}+r_{12}\right)}$ |  | $\mathrm{S}_{18}=2^{-12}\left(\beta_{3}+\beta_{10}\right)$ |
|  | $\mathrm{S}_{19}=2^{-1: 1}\left(\beta_{1}-\beta_{12}\right)$ |  | $\mathrm{S}_{19}=2^{-12}\left(\beta_{4}+\beta_{9}\right)$ |
|  | $S_{20}=2^{-1: 1}\left(\beta_{2}-\beta_{11}\right)$ |  | $S_{20}=2^{-12}\left(\beta_{s}-\beta_{8}\right)$ |
|  | $S_{21}=2^{-1: 1}\left(\beta_{3}-\beta_{(1)}\right)$ |  | $S_{21}=2^{-12}\left(\beta_{5}-\beta_{-}\right)$ |
|  | $S_{22}=2^{-1:}\left(\beta_{4}-\beta_{4}\right)$ |  | $\mathrm{S}_{22}=2^{-12}\left(\delta_{1}-\delta_{-}\right)$ |
|  | $S_{23}=2^{-1: 1}\left(\beta_{5}-\beta_{8}\right)$ |  | $\mathrm{S}_{23}=2^{-12}\left(\delta_{2}+\delta_{8}\right)$ |
|  | $S_{24}=2^{-1: 1}\left(\beta_{4}-\beta_{-}\right)$ |  | $\mathrm{S}_{24}=2^{-12}\left(\delta_{3}-\delta_{0}\right)$ |
|  | $S_{25}=2^{-1: 1}\left(\delta_{1}-\delta_{5}\right)$ |  | $\mathrm{S}_{25}=\delta_{4}$ |
|  | $S_{26}=2^{-1:}\left(\delta_{2}+\delta_{8}\right)$ |  | $\mathrm{S}_{26}=\delta_{6}$ |
|  | $\mathrm{S}_{2}=2^{-1: 2}\left(\delta_{3}-\delta_{4}\right)$ |  | $\mathrm{S}_{2}=\delta_{10}$ |
|  | $S_{28}=\delta_{5}$ |  | $\mathrm{S}_{28}=\delta_{12}$ |
|  | $\mathrm{S}_{29}=\delta_{11}$ |  |  |

${ }^{\circ}$ Abbreviations. SS: Symmetry Species.
Table 4. Scaled(Sc.) force constants of triphenylene and phenanthrene at the B3LYP/6-31G* level

| Temm | Triphenylene Sc . | Phenanthrene Sc. | Term | Triphenylene Sc . | Phenanthrene Sc . |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{R}_{1} \mathrm{R}_{1}$ | 5.484 | 5.475 | $\mathrm{r}_{11} \mathrm{r}_{11}$ | 5.254 | 5.205 |
| $\mathrm{R}_{2} \mathrm{R}_{2}$ | 5.993 | 5.583 | $\beta_{1} \beta_{1}$ | 1.099 | 1.045 |
| $\mathrm{R}_{3} \mathrm{R}_{3}$ | 6.632 | 6.545 | $\beta_{2} \beta_{2}$ | 1.038 | 1.058 |
| $\mathrm{R}_{4} \mathrm{R}_{4}$ | 7.028 | 6.968 | $\beta_{3} \beta_{3}$ | 1.063 | 1.024 |
| $\mathrm{R}_{5} \mathrm{R}_{5}$ | 6.480 | 6.233 | $\beta_{4} \beta_{4}$ | 1.098 | 1.077 |
| $\mathrm{R}_{6} \mathrm{R}_{6}$ | 7.281 | 7.334 | $\beta_{5} \beta_{5}$ | 1.094 |  |
| $\mathrm{R}_{7} \mathrm{R}_{7}$ | 6.270 | 6.126 | $\beta_{0} \beta_{6}$ | 1.038 |  |
| $\mathrm{R}_{8} \mathrm{R}_{8}$ | 5.714 | 7.418 | $\delta_{1} \delta_{1}$ | 1.644 | 1.527 |
| $\mathrm{R}_{11} \mathrm{R}_{11}$ | 6.280 |  | $\delta_{1} \delta_{1}$ | 1.807 | 1.622 |
| $\mathrm{R}_{1}=\mathrm{R}_{1} \%$ | 6.622 |  | $\delta_{3} \delta_{3}$ | 1.592 | 1.453 |
| $\mathrm{R}_{18} \mathrm{R}_{18}$ | 7.026 |  | $\delta_{4} \delta_{4}$ | 1.736 | 1.616 |
| $\mathrm{R}_{19} \mathrm{R}_{19}$ | 6.477 |  | $\delta_{5} \delta_{5}$ | 2.125 | 1.825 |
| $\mathrm{r}_{1} \mathrm{r}_{\mathrm{l}}$ | 5.234 |  | $\delta_{6} \delta_{6}$ | 2.117 | 1.512 |
| $\mathrm{r}_{3} \mathrm{r}_{3}$ | 5.246 |  | $\delta_{10} \delta_{10}$ | 1.640 |  |
| $\mathrm{r}_{5} \mathrm{r}_{5}$ | 5.246 | 5.005 | $\delta_{11} \delta_{11}$ | 1.481 |  |
| $\mathrm{r}_{7} \mathrm{r}_{7}$ | 5.127 | 5.128 | $\delta_{12} \delta_{12}$ | 1.902 |  |
| $\mathrm{rag}_{9}$ | 5.234 | 5.227 |  |  |  |

Table 5. Vibrational Spectra, IR intensities ( $\mathrm{km} / \mathrm{mol}$ ) and Assignunent for In-plane modes of Triphenylene ${ }^{a}$

| SS | $\mathrm{SS}\left(\mathrm{D}_{\text {sh }}\right)$ | Expt ${ }^{\text {b }}$ | sc | unse | IR int. | $\omega^{\text {bc }}$ | PED |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{A}_{1}$ | $\mathrm{A}_{1}$ | 3117 | \$107 | 3243 | 359 | 3117 | $37 \mathrm{~S}_{18}+21 \mathrm{~S}_{14}+14 \mathrm{~S}_{15}$ |
|  | E | 3098 | 3105 | 3241 | 4620 | 3115 | $40 \mathrm{~S}_{13}+22 \mathrm{~S}_{14}+20 \mathrm{~S}_{13}$ |
|  | E | 3 [91 | 3094 | 3222 | 110 | 3008 | $52 \mathrm{~S}_{13}+40 \mathrm{~S}_{15}$ |
|  | $\mathrm{A}_{1}$ | 3085 | 3086 | 3206 | 002 | 3083 | $69 \mathrm{~S}_{17}+21 \mathrm{~S}_{18}$ |
|  | E | 3075 | 3076 | 3203 | 4685 | 3080 | $50 \mathrm{~S}_{14}+31 \mathrm{~S}_{15}$ |
|  | E | 305 | 3052 | 3189 | 244 | 3067 | $86 \mathrm{~S}_{16}$ |
|  | E | 1612 | 1613 | 1666 | 025 | $15 \%$ | $21 \mathrm{~S}_{6}+16 \mathrm{~S}_{3}+14 \mathrm{~S}_{4}$ |
|  | E | 1579 | 1567 | 1634 | 007 | 1507 | $21 \mathrm{~S}_{12}$ |
|  | $\mathrm{A}_{1}{ }^{\text {a }}$ | 1550 | 1548 | 1606 | 000 | 1541 | 258s |
|  | E | 1502 | 1509 | 1550 | 1118 | 1490 | $14 S_{33}+13 S_{20}$ |
|  | $\mathrm{A}_{1}$ | 1434 | 1435 | 1496 | 001 | 1440 | $24 S_{21}+16 S_{23}+15 S_{22}$ |
|  | E | 1434 | 1430 | 1482 | 2033 | 1427 | $14 \mathrm{~S}_{22}+12 \mathrm{~S}_{2}$ |
|  | E | 1341 | 1340 | 1383 | 020 | 1335 | $22 \mathrm{~S}_{99}+21 \mathrm{~S}_{19}+16 \mathrm{~S}_{15}$ |
|  | $\mathrm{A}_{1}$ | 1300 | 1310 | 1377 | 000 | 1330 | $25 \mathrm{~S}+\mathrm{S}^{2} 2 \mathrm{~S}_{2}$ |
|  | E | $1300{ }^{-}$ | 1284 | 1337 | 154 | 1292 | $33 \mathrm{~S}_{111}+18 \mathrm{~S}_{11}+17 \mathrm{~S}_{8}$ |
|  | E | 1256 | 1258 | 1285 | 382 | 124 | $26 \mathrm{~S}_{3}+18 \mathrm{~S}_{19}+13 \mathrm{~S}_{22}$ |
|  | A ${ }^{\text {a }}$ | 1230 | 1211 | 1263 | 000 | 1223 | $24 \mathrm{~S}_{1}+16 \mathrm{~S}_{36}+12 \mathrm{~S}_{3}$ |
|  | A | 1180 | 1191 | 1221 | 000 | 1184 | $21 S_{21}+16 S_{21}$ |
|  | E | 1173 | 1171 | 1202 | 027 | 1166 | $49 \mathrm{~S}_{3+}+17 \mathrm{~S}_{23}$ |
|  | E | 1106 | 1106 | 1142 | 032 | 1109 | $21 S_{1}+21 S_{15}+14 S_{4}$ |
|  | $\mathrm{A}_{5}$ | 1662 | 1063 | 1100 | 000 | 1089 | 518. |
|  | E | 1052 | 1051 | 1086 | 369 | 1057 | $40 \mathrm{~S}_{12}+17 \mathrm{~S}_{23}$ |
|  | E | 1603 | 998 | 1023 | 158 | 996 | $34 \mathrm{~S}_{35}+12 \mathrm{~S}_{3}$ |
|  | E | 776 | 784 | 789 | 000 | 773 | $40 \mathrm{~S}_{26}+23 \mathrm{~S}_{25}$ |
|  | $\mathrm{A}_{1}{ }^{\text {a }}$ | \%00 | 700 | 713 | 000 | 700 | $31 S_{27}+18 S_{24}+13 S_{2}+13 S_{26}$ |
|  | E | (12) | 633 | 635 | 559 | 615 | $49 \mathrm{~S}_{29}+31 \mathrm{~S}_{27}$ |
|  | $\mathrm{A}_{1}$ | 418 | 419 | 424 | 000 | 419 | $24 S^{-}+17 S_{2}=$ |
|  | E | 406. | 405 | 410 | 080 | 406 | $59 \mathrm{~S}_{28}+15 \mathrm{~S}_{26}$ |
|  | E | 264 | 264 | 265 | 028 | 263 | $41 \mathrm{~S}_{26}$ |
| $\mathrm{B}_{1}$ | E | 3098 | 3102 | 3240 | 5013 | 3114 | $46 \mathrm{~S}_{11}+43 \mathrm{~S}_{12}$ |
|  | A | 3091 | 3095 | 3224 | 047 | 3099 | $67 \mathrm{~S}_{14}+24 \mathrm{~S}_{15}$ |
|  | E | 3085 | 3084 | 3220 | 023 | 3095 | $44 \mathrm{~S}_{111}+38 \mathrm{~S}_{12}$ |
|  | E | 3075 | 3075 | 3203 | 4659 | 3080 | $71 \mathrm{~S}_{15}+22 \mathrm{~S}_{1+}$ |
|  | E | 1075 | 3072 | 3189 | 263 | 3067 | $45 \mathrm{~S}_{10}+36 \mathrm{~S}_{11}$ |
|  | A | 3055 | 3052 | 3188 | 000 | 3066 | $85 \mathrm{~S}_{13}$ |
|  | $\mathrm{A}_{2}$ | 1612 | 1623 | 1670 | 000 | 1601 | $29 \mathrm{~S}_{9}+27 \mathrm{~S}_{8}$ |
|  | E | 1612 | 1616 | 1665 | 026 | 1596 | $23 \mathrm{~S}_{6}+16 \mathrm{~S}_{3}$ |
|  | E | 1579 | 1567 | 1634 | 007 | 1568 | $23 \mathrm{~S}_{4}+20 \mathrm{~S}_{5}+16 \mathrm{~S}_{2}$ |
|  | E | 1501 | 1501 | 1550 | 1136 | 1490 | $17 \mathrm{~S}_{[0}+13 \mathrm{~S}_{-}+13 \mathrm{~S}_{21}$ |
|  | $\mathrm{A}_{2}$ | 1471 | 1472 | 1505 | 000 | 1449 | $23 \mathrm{~S}_{18}+16 \mathrm{~S}_{35}+13 \mathrm{~S}_{3}$ |
|  | E | 1434 | 1446 | 1482 | 2019 | 1427 | $24 \mathrm{~S}_{21}+17 \mathrm{~S}_{16}$ |
|  | E | 1341 | 1341 | 1383 | 017 | 1335 | $28 S^{-}+1 \cdot 1 S_{20}$ |
|  | E | 1300 | 1294 | 1336 | 147 | 1291 | $15 \mathrm{~S}_{16}+14 \mathrm{~S}_{3}+13 \mathrm{~S}_{19}+12 \mathrm{~S}_{25}$ |
|  | A | 1256 | 1279 | 1322 | 000 | 1278 | $21 S_{20}+14 S^{-}$ |
|  | E | 1245 | 1239 | 1285 | 388 | 1243 | $23 \mathrm{~S}_{1}+18 \mathrm{~S}_{6}+14 \mathrm{~S}_{2}$ |
|  | E | 1173 | 1175 | 1203 | 027 | 1166 | $31 \mathrm{~S}_{18}+27 \mathrm{~S}_{19}+14 \mathrm{~S}_{16}$ |
|  | A ${ }^{\text {a }}$ | 1136 | 1150 | 1179 | 000 | 1144 | $14 \mathrm{~S}_{17}+13 \mathrm{~S}_{4}+12 \mathrm{~S}_{9}+12 \mathrm{~S}_{21}$ |
|  | E | 1110 | 1114 | 1142 | 034 | 1109 | $16 S_{9}+14 S_{8}+14 S_{21}+13 S_{2}+12 S_{4}$ |
|  | E | 1052 | 1054 | 1087 | 362 | 1057 | 418 |
|  | E | 1005 | 1015 | 1023 | 159 | 997 | $60 \mathrm{~S}_{22}+33 \mathrm{~S}_{27}$ |
|  | A | (9) 3 | 993 | 1015 | 000 | 989 | $26 \mathrm{~S}_{27}+17 \mathrm{~S}_{1}+13 \mathrm{~S}_{22}$ |
|  | E | 76 | 779 | 790 | 000 | 774 | $29 S_{28}+14 S_{27}+13 S$ |
|  | E | 62) | 633 | 636 | 561 | 626 | $48 \mathrm{~S}_{3+}+27 \mathrm{~S}_{13}$ |
|  | Az | 607 | 614 | 617 | 000 | 607 | $40 \mathrm{~S}_{33}+29 \mathrm{~S}_{23}$ |
|  | A ${ }^{\text {c }}$ | 530 | 560 | 566 | 000 | 558 | $56 \mathrm{~S}_{25}$ |
|  | E | 410 | 409 | 411 | 078 | 406 | $64 \mathrm{~S}_{26}$ |
|  | E | 264 | 262 | 264 | 028 | 262 | $33 \mathrm{~S}_{28}+17 \mathrm{~S}$ |

[^1](iii) Vibrational Spectra: Triphenylene in $\mathrm{C}_{3 v}$ symmetry has 57 in-plane and 27 out-of-plane normal modes of vibration spanning irreducible representations $29 \mathrm{~A}_{1}$ (IR.R) + $14 \mathrm{~A}_{2}(\mathrm{R})+13 \mathrm{~B}_{1}$ (IR.R) $+28 \mathrm{~B}_{2}$ (IR.R). The spectral activities are given in parentheses e.g. $29 \mathrm{~A}_{1}$ modes are both IR and Raman active and so on. The vibrational modes of triphenylene in $D_{s h}$ symmetry span the irreducible representations $10 \mathrm{~A}_{1}(\mathrm{R})+9 \mathrm{~A}_{2}()+19 \mathrm{E}(\mathrm{IR} . \mathrm{R})$ for in-plane and $5 \mathrm{~A}_{1}^{\prime \prime}()+4 \mathrm{~A}_{2}^{*}(\mathrm{IR})+9 \mathrm{E}^{*}(\mathrm{R})$ for out-of-plane motions where the parentheses enclosing no content are both IR and Raman inactive. The B3LYP/6-31G* frequencies obtained for phenanthrene with $\mathrm{C}_{24}$ symmetry from Gaussian94 program are fitted to the experimental ones of phenanthrene$\mathrm{d}_{(1)}$ and $-\mathrm{d}_{\mathrm{l}_{1}}$ for the in-plane fundamentals. The agreement between the fitted and the experimental frequencies is very good for both the isotopomers with an average deviation of $3.1 / 2.5 \mathrm{~cm}^{-1}$ with/without C-H stretches.

Most of the assignments are very straightforward as the frequencies agree with the predicted numbers from very well-studied system. phenanthrene and scaled frequencies by the scaling method of Lee et af ${ }^{2021}$ Some assiguments merit some comments. The bands at 1300 and $1341 \mathrm{~cm}^{-1}$ (belonging to $\mathrm{E}^{*}$ and $\mathrm{A}_{1}{ }^{\prime}$ respectively) in both the assignments of Schettino ${ }^{11}$ and Nishi et al ${ }^{12}$ are interchanged.

These assignments are in good agreement with the predicted numbers ( 1340 and $1310 \mathrm{~cm}^{-1}$ ). Noting that the band at 1300 $\mathrm{cm}^{-1}$ has a weak intensity. it is evident that the present assigument is justified.

The bands at 1246 and $1256 \mathrm{~cm}^{-1}$ considered as a single fundamental in the work of Hudgins et al. ${ }^{13}$ are regarded here as two fundamentals. This is also supported by the theoretical predictions. The band at $947 \mathrm{~cm}^{-1}$ was reported by Hudgins et al. However, in view of the absence of suitable match with the predicted frequencies. the band is uncertain as a fundamental. It is to be noted that the alternative assigument is at $1025 \mathrm{~cm}^{-1}$ as reported by Schettino. ${ }^{11}$ The bands at 1560 and $1422 \mathrm{~cm}^{-1}$ in ref. 11 appear to be doubtful as they are not in agreement with the predicted numbers. This fact was further supported by Keszthelyi et al. ${ }^{15}$ However, the band at $1180 \mathrm{~cm}^{-1}$ suspected by Keszthelyi et al. as a fundamental is considered as a fundamental in the present work.

For out-of-plane modes. the strategy of scaling due to Scott and Radom ${ }^{25}$ employed before to azulene ${ }^{17}$ is used here. The frequencies are predicted based on simple scaling recommended by Scott and Radom25 and exponential scaling methodology due to Lee et al. ${ }^{2(2)}$ and compared with the results of Hudgins et al. ${ }^{13}$

Table 6. Vibrational Spectra, IR intensities $(\mathrm{km} / \mathrm{mol})$ and Assignment of Tripheny lene tor the out-of-plane modes ${ }^{a}$

| $\mathrm{SS}\left(\mathrm{C}_{2 \mathrm{v}}\right)$ | $\mathrm{SS}\left(\mathrm{D}_{31}\right)$ | Expt [13] | sc | unsc | IR int. | $\omega^{* *}$ | PED ${ }^{\text {² }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{A}_{2}$ | E | 947 | 943 | 981 | 0.00 | 957 | $22 l_{1}+221_{2}$ |
|  | $\mathrm{Al}_{1}$ |  | 934 | 971 | 0.00 | 947 | $12 l_{1}+121_{2}+121$ |
|  | $E$ |  | 890 | 926 | 0.00 | 904 | $1410+141_{10}+1315$ |
|  | $A_{1}$ |  | 844 | 878 | 0.00 | 859 | $161_{16}+121_{22}$ |
|  | E. |  | 833 | 866 | 0.00 | 847 | $21 l_{21}+21 l_{21}$ |
|  | E |  | 755 | 785 | 0.00 | 769 | $221_{22}$ |
|  | A. ${ }^{\text {I }}$ |  | 754 | 785 | 0.00 | 769 | $221_{16}+111_{13}$ |
|  | E |  | 694 | 722 | 0.00 | 709 | 37122 |
|  | A. |  | 557 | 579 | 0.00 | 571 | $161_{13}+161_{19}+161_{22}$ |
|  | E* |  | 536 | 558 | 0.00 | 550 | $11 l_{22}+111_{24}$ |
|  | E: |  | 421 | 438 | 0.00 | 433 | $2311_{4}+231_{21}+161_{24}$ |
|  | E* |  | 264 | 275 | 0.00 | 273 | $171_{14}+201_{1}+16 l_{15}$ |
|  | A. |  | 106 | 110 | 0.00 | 110 | $211_{15}+151_{15}+151_{21}$ |
|  | E: |  | 52 | 54 | 0.00 | 54 | $331_{18}+15 l_{24}$ |
| $\mathrm{B}_{1}$ | E | 947 | 943 | 981 | 0.00 | 957 | $1810+181_{10}+151-$ |
|  | $\mathrm{Al}^{\prime}$ |  | 911 | 948 | 1.14 | 925 | $1013+104_{4}$ |
|  | $\mathrm{R}^{\prime \prime}$ |  | 891 | 927 | 0.00 | 905 | $161_{3}+161_{4}+121_{1}+121_{2}$ |
|  | E |  | 832 | 865 | 0.00 | 846 | $201_{3}+201_{6}+141_{11}$ |
|  | $E^{\prime}$ |  | 755 | 785 | 0.01 | 770 | $141_{13}+141_{19}$ |
|  | $\mathrm{As}^{\prime}$ | 741 | 730 | 759 | 142.40 | 744 | 2010 |
|  | $\mathrm{E}^{\prime}$ |  | 694 | 722 | 0.01 | 709 | $2811_{13}+281_{19}$ |
|  | E' |  | 536 | 558 | 0.00 | 550 | $15 l_{25}$ |
|  | P |  | 423 | 440 | 0.00 | 435 | $201_{15}+201_{21}+201_{23}$ |
|  | $\mathrm{A}_{4}$ |  | 417 | 434 | 2.93 | 429 | $181_{14}+181_{20}+171_{23}$ |
|  | $\mathrm{E}^{\prime \prime}$ |  | 265 | 276 | 0.00 | 274 | $31 l_{25}+171_{20}+17 l_{14}$ |
|  | $\mathrm{Al}^{\prime \prime}$ |  | 120 | 125 | 1.73 | 125 | $251_{25}+251_{26}+25 l_{2}$ |
|  | E |  | 51 | 53 | 0.00 | 53 | 321: |

"Only contributions to the potential energy distributions in symmetry coordinates equal to or larger than $10^{\circ} \%$ are included. In-plane and out-of-plane
 ${ }^{6}$ PEDs of out-of-plane modes are based on unscaled local force field.

## Conclusions

The structure. force field and vibrational spectra. and their assignment of triphenylene of $\mathrm{C}_{2 y}$ symmetry are studied by B3LYP/6-31G* level of density functional theory: The results are compared to the related system. phenantlurene. The force field of tripheny lene is higher from its constituent ring system. although some of the force constants are similar. The earlier assignment was confirmed or reassigned for in-plane modes. For in-plane modes the frequencies scaled by two independent methods match very well. and for out-of-plane modes the agreement is reasonable. The out-ofplane frequencies predicted compare reasonably with the latest assignment obtained from IR spectra.

Acknowledgements. This study was supported by BK21. I thank the anonymous referees for their valuable criticisms/ comments.

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[^1]:    "Only contributions to the potential energy distnbutions (PED) in symmetry coordinates equal to or larger than $10{ }^{\circ}$. are included. In-plane and out-ofplane modes are treated separately. When one experimental band has large deviation from the predicted number. the experimental frequency is marked
     C-H stretches, the scale factor $\left(\alpha=2.548 \times 10^{-5}\right)$ from the pair ( $\omega^{\omega \mathrm{Bic}}=758 ; \omega^{\circ}=744$ ) was used. Abbreviations: sc: scaled: unsc: unscaled; SS: Symmetry Species. ${ }^{3}$ From refs. 11.12.13.

