

Physicochemical Properties of ZrF₄-Based Fluoride Glasses Containing Rare-Earth Ions

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In the XRD study of 56ZrF₄ · 34BaF₂ · 4AlF₃ · (6-x)LnF₃ · xLnF₃ glasses (Ln=Ce, Nd, Gd, Tb), halo pattern characteristic of an amorphous sample appeared. When the halo peak angle (θ_p) was converted into a wavenumber with $Q_p=4\pi\sin\theta_p/\lambda$ (λ is the wavelength of the radiation used), it was found that the Q_p values varied almost linearly with the concentration of LnF₃. The emission spectra of Ce³⁺-containing fluoride glasses under 273 nm excitation had a peak maximum at ca. 300 nm (Ce³⁺ 5d-4f transition). The maximal intensity of the fluorescence was observed when the CeF₃ content was extremely low (ca. 1 mol%). DTA measurement revealed that these fluoride glasses had two crystallization temperatures. In 56ZrF₄ · 34BaF₂ · 4AlF₃ · (6-x)LnF₃ · xNdF₃ glasses, the activation energies of crystallization obtained from a Kissinger plot were 1.7 and 5.0 eV for the glass with x=2, and 1.9 and 5.6 eV for the glass with x=4.

Key words: Fluoride glass, Rare-earth ion, Fluorescence, Kissinger plot

I. Introduction

Fluorozirconate glasses (ZrF₄-based glasses) prepared by Poulain et al.¹⁾ have attracted much interest because of the high optical transparency in the visible to infrared (IR) region. Since ZrF₄ cannot be a network former by itself, the formation of glass samples becomes feasible when ZrF₄ is melted with other fluorides, e.g., BaF₂. Lucas²⁾ described that the introduction of AlF₃ into ZrF₄-based glasses was favorable for the glass formation. Also, small amounts of rare earth ions proved to be effective for stabilizing the glassy phase.²⁾

Especially, the ZrF₄-based fluoride glasses containing rare earth ions are attractive when developing rare earth-doped fiber amplifiers for use in optical communication system. The present study was carried out in order to obtain some physicochemical properties of ZrF₄ · BaF₂ · AlF₃ · LaF₃ · LnF₃ (Ln=Ce, Nd, Gd, Tb).

II. Experimental

The fluoride glasses, 56ZrF₄ · 34BaF₂ · 4AlF₃ · (6-x)LnF₃ · xLnF₃ glasses (Ln=Ce, Nd, Gd, Tb), were prepared by melting the desired mixture of the starting materials of analytical reagent grade under a dry nitrogen atmosphere. The melt was kept at 1123 K for 5 min and quenched by use of cold copper plate under in air. The formation of homogeneous glass samples was confirmed by X-ray diffraction (Rigaku Denki RAD2B) patterns. Fluorescence spectra were recorded on a Hitachi Fluorescence Spectrophotometer (F-4500). Glass transition tempera-

ture (T_g) and crystallization peak temperature (T_c) were determined by DTA (Shimadzu DTA-50).

III. Results and Discussion

1. Formation of glasses

Fig. 1. shows typical examples of X-ray diffraction patterns recorded for 56ZrF₄ · 34BaF₂ · 4AlF₃ · (6-x)LnF₃ · xCeF₃

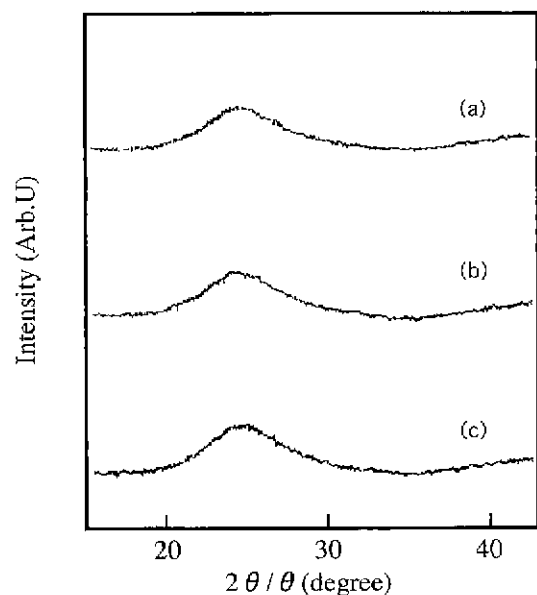


Fig. 1. X-Ray diffraction patterns of 56ZrF₄ · 34BaF₂ · 4AlF₃ · (6-x)LnF₃ · xCeF₃ glasses. (a) x=1.8, (b) x=2.4, (c) x=3.0.

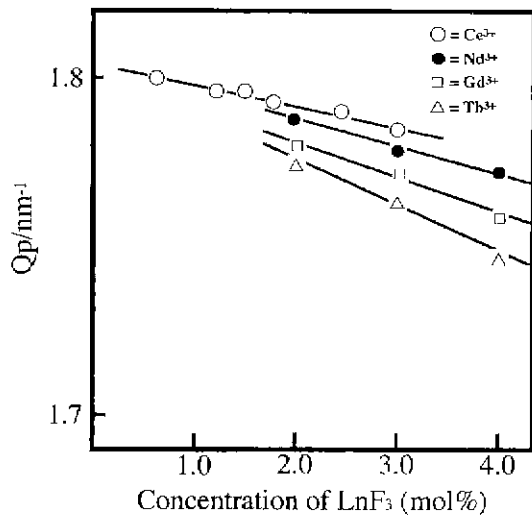


Fig. 2. Plot of the wavenumber, Q_p/nm^{-1} , vs. LnF_3 concentration in $56\text{ZrF}_4 \cdot 34\text{BaF}_2 \cdot 4\text{AlF}_3 \cdot (6-x)\text{LaF}_3 \cdot x\text{LnF}_3$ glasses.

glasses. A broad diffuse diffraction peak which is attributed to an amorphous sample is observed at around 20-30 degrees. The peak maximum of these halo patterns shifts to a lower angle with the increase of "x". Halo patterns were also obtained in the ZrF_4 -based fluoride glasses containing other rare earth ions. After the halo peak angle (Q_p) was converted into a wavenumber with $Q_p = 4\pi \sin\theta/\lambda$ (λ is the wavelength of the radiation used), Q_p values are plotted against the concentration of LnF_3 , as shown in Fig. 2. It was found that the Q_p values varied almost linearly with the concentration of Ln, while it diminished with the decrease of the ionic radius of Ln^{3+} . ZrF_4 - BaF_2 glasses with low BaF_2 contents are known to have a layer structure composed of zigzag chains of ZrF_6 units, whereas they have a pseudo-chain structure composed of ZrF_7 units when the $\text{BaF}_2/\text{ZrF}_4$ ratio is close to unity.³¹ Ba^{2+} ions occupy so-called "interstitial" sites, as a network modifier, surrounded by ZrF_6 and/or ZrF_7 units,³¹ just like several alkali and alkaline earth metal ions in several oxide glasses. Thus, it can be considered that the relationship between Q_p values and the concentration of Ln^{3+} is attributed to a gradual increase in the nearest-neighbour distance of the plane including Ln^{3+} ions. The Q_p values decrease linearly with increase of concentration.

2. Fluorescence spectra

A typical emission spectrum for $56\text{ZrF}_4 \cdot 34\text{BaF}_2 \cdot 4\text{AlF}_3 \cdot 5.4\text{LaF}_3 \cdot 0.6\text{CeF}_3$ glass under 273 nm excitation is shown in Fig. 3. This photoluminescence can be assigned to the electron transition in the $4f^1 \rightarrow 4f^05d^1$.⁴¹ Generally, the wavelength of the emission peak for Ce^{3+} in many compounds is located at ca. 350-360 nm.⁴¹ Hence, it is assumed that the Ce^{3+} ions in these fluoride glasses occupy distorted sites; the peak position of this emission band has shifted to a lower wavelength region.

It is well known that the emission intensity in various

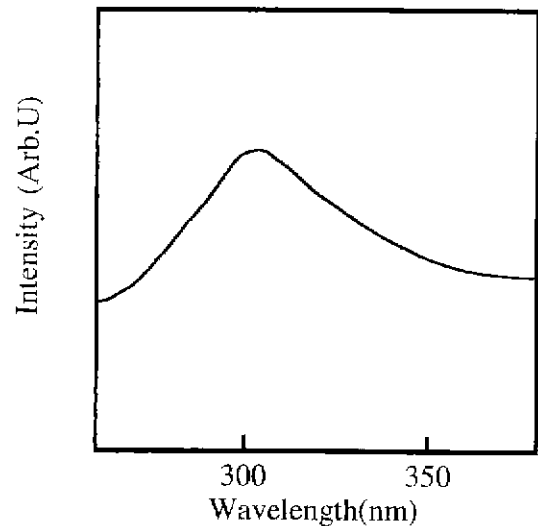


Fig. 3. Emission spectra of $56\text{ZrF}_4 \cdot 34\text{BaF}_2 \cdot 4\text{AlF}_3 \cdot 4\text{LaF}_3 \cdot 2\text{CeF}_3$ glass excited with the UV light with the wavelength of 273 nm.

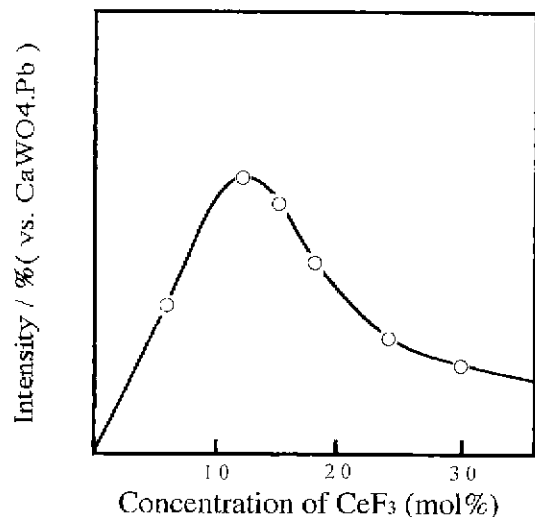


Fig. 4. Dependence of the intensity of the emission band for Ce^{3+} on the concentration of CeF_3 .

phosphors is affected by the concentration of Ce^{3+} . The effect of Ce^{3+} ion concentration on the emission band is presented in Fig. 4. The photoluminescence outputs (PL) of the samples were compared with that of $\text{CaWO}_4 \cdot \text{Pb}$ (NBS 1026). The intensity of PL showed a specific pattern; initially, there was a sharp increase in the intensity along with the increase in the Ce^{3+} ions and then it reached a maximum at ca. 1 mol% of CeF_3 . This effect is ascribed to the well-known "concentration quenching effect".

Thus, because the peak position of the Ce^{3+} emission band for $56\text{ZrF}_4 \cdot 34\text{BaF}_2 \cdot 4\text{AlF}_3 \cdot (6-x)\text{LaF}_3 \cdot x\text{CeF}_3$ was located at a lower wavelength and the quenching effect of the emission intensity was observed at lower Ce^{3+} concentration, the rare earth ions in $56\text{ZrF}_4 \cdot 34\text{BaF}_2 \cdot 4\text{AlF}_3 \cdot (6-x)\text{LaF}_3 \cdot x\text{LnF}_3$ are considered to be located at "interstitial

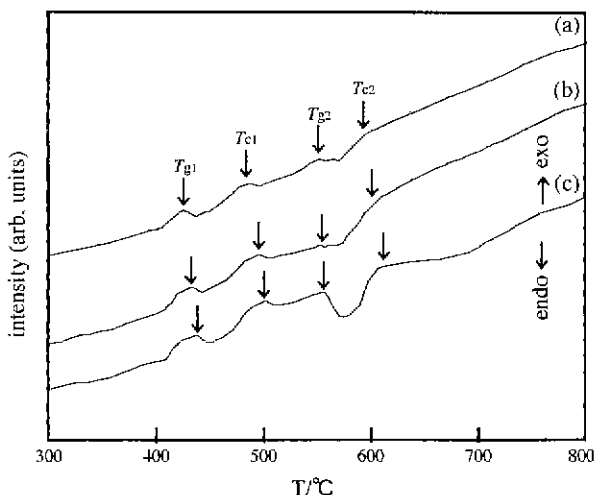


Fig. 5. DTA curves of 56ZrF₄ · 34BaF₂ · 4AlF₃ · 4LaF₃ · 2CeF₃ glass recorded at heating rates of (a) 20, (b) 25 and (c) 30 K/min.

sites" which a surrounded by ZrF₆ octa field, as will be described later.

3. Crystallization energy

The short-range structure of ZrF₄-based fluoride glasses doped with rare earth ions was discussed with the values of activation energy (E_a) of crystallization. A Kissinger plot in the differential thermal analysis (DTA) is generally used for estimating the E_a of crystallization. The Kissinger equation⁵⁾ is expressed by

$$\ln(T_c^2/\alpha)E_a/RT_c + \text{const} \quad (1)$$

in which T_c, α and R are the crystallization peak temperature (in K), heating rate of the sample, and the gas constant, respectively. DTA measurements were carried out for the Ce³⁺- and Nd³⁺-doped fluoride glasses. Fig. 5 shows a DTA curve of 56ZrF₄ · 34BaF₂ · 4AlF₃ · (6-x)LaF₃ · xNdF₃ glass recorded at heating rates of 20, 25 and 30 K. This fluoride glass had two crystallization temperatures.

And the peak temperatures of the crystallization increased with an elevating of the heating rates. Two crystallization temperatures were observed for 56ZrF₄ · 34BaF₂ · 4AlF₃ · (6-x)LaF₃ · xNdF₃ glasses. The slope of the straight line in the Kissinger plot of 56ZrF₄ · 34BaF₂ · 4AlF₃ · (6-x)LaF₃ · xCeF₃ (Fig. 6) yielded the E_a values of 1.4 and 4.0 eV for the glass with x=2, and 1.6 and 4.8 eV for the glass with x=3. A Kissinger plot for 56ZrF₄ · 34BaF₂ · 4AlF₃ · (6-x)LaF₃ · xNdF₃ is shown in Fig. 7 for comparison, which yielded E_a values of 1.7 and 5.0 eV for the glass with x=2, and 1.9 and 5.6 eV for the glass with x=4.

If a rare earth ion is cross-linked to a F⁻ ion coordinating to a Zr⁴⁺ ion in layer structure of composed of zigzag chains of ZrF₆ units, the average binding energy is estimated as follows; 1/8 (ΔH_{Nd-F}+ΔH_{Zr-F})=4.7 and 1/8 (ΔH_{Ce-F}+ΔH_{Zr-F})=4.7 eV.

Since 4~5 eV of the activation energy of crystallization

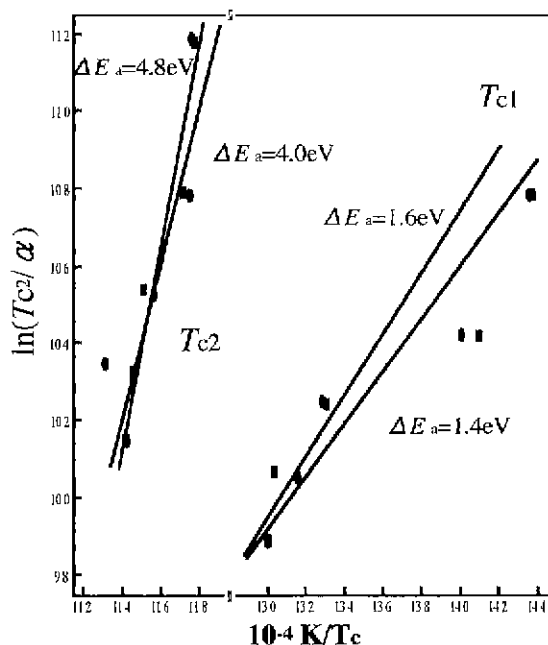


Fig. 6. Kissinger plot for (●)56ZrF₄ · 34BaF₂ · 4AlF₃ · 4LaF₃ · 2CeF₃ and (■)56ZrF₄ · 34BaF₂ · 4AlF₃ · 3LaF₃ · 3CeF₃ glasses.

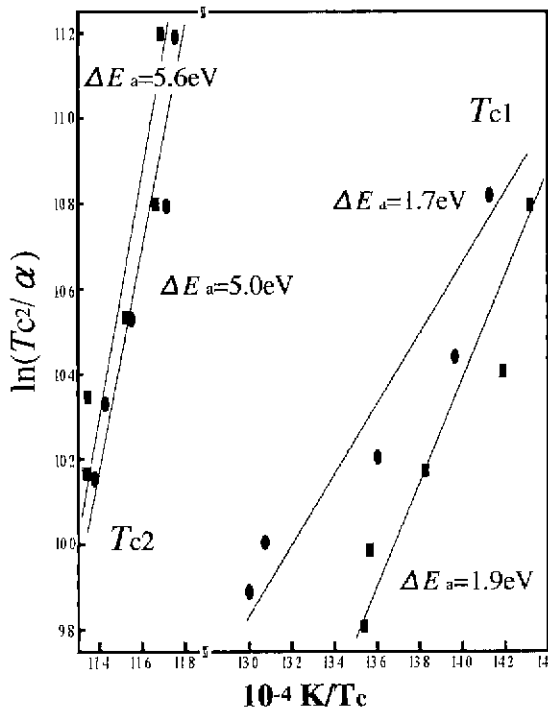


Fig. 7. Kissinger plot for (●)56ZrF₄ · 34BaF₂ · 4AlF₃ · 4LaF₃ · 2NdF₃ and (■)56ZrF₄ · 34BaF₂ · 4AlF₃ · 2LaF₃ · 4NdF₃ glasses.

was ranged in this average binding energy, it was considered that cleavage Ln-F and Zr-F bonds triggered the crystallization.

IV. Summary

We have studied some physicochemical properties of

rare earth doped ZrF_4 -based fluoride glasses. From $Q_p = 4\pi \sin \theta_p / \lambda$ (θ_p : a 1/2 of halo peak angle in the XRD $56ZrF_4 \cdot 34BaF_2 \cdot 4AlF_3 \cdot (6-x)LaF_3 \cdot xLnF_3$ glasses (Ln=Ce, Nd, Gd, Tb), λ : wavelength of the radiation used), it was found that the θ_p values varied almost linearly with the concentration of LnF_3 . In the photoluminescence of Ce^{3+} containing fluoride glasses, the wavelength of emission peak was smaller than that of normal phosphors and the concentration quenching effect was observed at extreme low concentration of CeF_3 . 2~5 eV of the activation energy of crystallization obtained from a Kissinger plot was ranged in the average binding energy which was due to the bridged rare earth ion.

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