# Density Functional Theory Study of Vibrational Spectra of Anthracene Neutral and Radical Cation 

Sang Yeon Lee ${ }^{\text {t§ }}$ and Bong Hyun Boo ${ }^{\text {¹* }}$<br>${ }^{\dagger}$ Department of Chemistry, Pusan National University, Pusan 609-735, Korea<br>${ }^{\mathbf{t}}$ Department of Chemistry, Chungnam National University, Tacjon 305-764, Korea<br>${ }^{8}$ Center for Molecular Science, Korea Advanced Institute of Science and. Technology. Taejon 305-701, Korea Received May 29, 1996


#### Abstract

Ab initio Hartree-Fock and Becke 3-Lee-Yang-Parr (B3LYP) density functional theory calculations using 6-31G* basis set were carried out to study the vibrational spectra of anthracene neutral ( $h_{10}$ and $d_{10}$ ) and radical cation ( $h_{10}$ ). We report results of the fundamental vibrational frequencies obtained on the basis of the calculations. The assignments of fundamentals show a one-to-one correspondence between the observed and calculated fundamentals.


## Introduction

Of growing interest is density functional theory (DFT) study since the methods predict relatively accurate molecular structures and vibrations with moderate computational effort. ${ }^{1-8}$ It is indicated in the published results that DFT method reproduces experimental vibrational frequencies with higher accuracy than do the Hartree-Fock (HF) and second order M $1 l$ er-Plesset perturbation theory (MP2) calculations. ${ }^{1 \sim 8}$ Even when a uniform scaling is performed for the computed vibrational frequencies, DFT methods show better agreement with experiment than do the HF and MP2 methods. ${ }^{1.2 .33}$ Some examples of the calculation on the vibrational frequencies of benzene are represented by Handy et al..$^{2}$ and Wheeless et al., ${ }^{4}$ in which MP2 calculations obtained by using $6-311 \mathrm{G}(\mathrm{d}$, p) basis set still underestimate the vibrational frequencies corresponding to the CCC out-of-plane bending, and overestimate that corresponding to a b2u mode which tends to dissociate into three acetylene molecules. On the other hand, these frequencies are accurately predicted by DFT methods using the Becke-Lee-Yang-Parr (BLYP) functional ${ }^{2,4}$

Rauhut and Pulay (RP) applied the BLYP and B3LYP functional methods to 20 small molecules such as benzene, ether and methanol, etc. whose vibrational frequencies are exactly assigned, and they then derived uniform scaling factors, 0.995 and 0.963 , having root mean square (rms) deviations of 26.2 and $18.5 \mathrm{~cm}^{-1}$, respectively ${ }^{5}$ When these workers also applied these scaling factors to another 11 molecules such as aniline, ethanol and oxetane, etc., the rms deviations turned out to be 26.9 and $19.7 \mathrm{~cm}^{-1}$ for the BLYP and B3LYP methods, respectively.

The electronic-vibrational structures of anthracene (AN) and its derivatives are the subject of ongoing spectroscopic ${ }^{9,10}$ and theoretical ${ }^{11,12}$ investigations. It is well known in the literature that AN molecule fluoresces efficiently upon excitation of $S_{0}$ to $S_{1}$. Also the ultrafast dynamics of the van der Waals complexes of anthracene derivatives has been studied by time-resolved fluorescence spectroscopy. ${ }^{10}$ Published results indicate that the vibrational frequencies appended to the $0_{0}^{0}$ band origin in AN consist of fundamentals, overtones, second overtones, and combination bands. ${ }^{9.10 a \cdot c}$ Thus the determination of fundamental vibrational frequencies and modes in
the ground-state is necessary to elucidate the vibronic spectra.

Numerous vibrational spectroscopic data have been reported on ground-state AN neutral (AN- $h_{10}{ }^{14}$ and AN- $d_{10}{ }^{15.16}$ ), and on the radical cation ( $\mathrm{AN}^{+}-h_{10}$ ) in a matrix-isolated system. ${ }^{17.18}$ Although the experimental works are excellent in the assignment of symmetry to the observed vibrational frequencies, there are, however, no assignment to $a_{u}$ modes owing to the forbidden infrared and Raman selection rule. And also there is some ambiguity in the assignments of fundamental modes of AN $d_{10}$ and $\mathrm{AN}^{+}-h_{10}$. Many workers have endeavoured to unambiguosly reproduce fundamentals of $\mathrm{AN}-h_{10}{ }^{11,2,2,17}$ and $\mathrm{AN}^{+}-h_{10}{ }^{17}$ and to assign their modes with $a b$ initio HF calculations. Still large deviations are found between experiment and the calculations. Recently, the infrared spectra of $\mathrm{AN}-h_{10}$ and $\mathrm{AN}{ }^{+}-h_{10}$ are more reliably reproduced by Langhoff with the B3LYP calculation using the $4-31 G^{*}$ basis set, in which multiple scaling factors were employed to reduce rms deviations. ${ }^{7}$ Unfortuately, comparison between experiment and the calculations could be made for only some frequencies having relatively high intensities in the infrared spectra. ${ }^{?}$

By using the HF and B3LYP methods, we performed the normal vibrational mode analysis of $\mathrm{AN}-h_{10}, \mathrm{AN}-d_{10}$, and $\mathrm{AN}^{+}-h_{10}$ with the use of the expanded basis set $6-31 \mathrm{G}^{*}$ to predict the spectral positions of the missing lines and to select fundamentals from the various vibrational frequencies reported experimentally. By using the uniform scaling factor reported by RP, we could predict the spectral positions of the molecules with the similar accuracy as predicted by Langhoff with the multiple scaling factors. ${ }^{7}$ It is found in this study that DFT methods are sufficiently powerful to predict fundamentals and their intensities.

## Calculations

The molecular vibrations were calculated at the HF and B3LYP levels of theory with $6-31 G^{*}$ basis set by using the Gaussian 94 program. ${ }^{19}$ Vibrational frequencies were computed with the HF and B3LYP methods. They were scaled by 0.8929 (HF/6-31G*) and 0.963 (B3LYP/6-31G*). The frequencies and normal modes for AN were determined by dia-
gonalizing the mass-weighted force constant matrix.

## Results and Discussion

Our calculations show that anthracene radical cation turned out to have $D_{2 n}$ symmetry like anthracene neutral in the ground-state. The AN neutral and radical cation involve 66 fundamentals having the various symmetries of $12 a_{g}+5$ $a_{u}+4 b_{1 g}+11 b_{1 u}+6 b_{26}+11 b_{2 \mu}+11 b_{3 g}+6 b_{34}$. We seek to determine fundamentals making a one-to-one correspondence between the observed and calculated frequencies. Fundamentals having the same symmetries are well separated and thus the assignment is straight forward. The scaled vibrational frequencies and modes of $\mathrm{AN}-h_{10}$ and $\mathrm{AN}-d_{10}$ calculated with the HF and B3LYP methods are represented in Table 1. Our scaled frequencies of $\mathrm{AN}-h_{10}$ obtained from the B3LYP calculation are in excellent agreement with experiment and are found to have rms deviations of 30.98 and $19.30 \mathrm{~cm}^{-1}$ for the HF and B3LYP, respectively. Therefore we conclude that among the methods the B3LYP calculations are quite reliable in predicting the fundamentals.

However, the assignments of fundamentals of $\mathrm{AN}-d_{10}$ are less straightforward due to the missing lines. But reliable prediction of many fundamentals and modes of $\mathrm{AN}-d_{10}$ could be made by comparing the frequencies and modes of the protiated system with those of the deuterated counterpart. Our scaled frequencies obtained from the B3LYP calculation are in excellent agreement with experiments and are found to have rms deviations of $12.97 \mathrm{~cm}^{-1}$. As expected, the ring torsional and deformation modes are not so sensitive to the deuterium substitution. However, the $\mathrm{C}-\mathrm{H}$ stretching and bending frequencies correctly reflect the deuterium isotope ef-
fect.
For the $\mathrm{AN}^{+}-h_{10}$ system, we represent in Table 2 the calculated results along with a few observed vibrational peaks. The experimental frequencies are well correlated with the calculated values. It is quite interesting that the observed frequencies are also found to have the high intensities in the B3LYP and HF calculations.

Almost all the modes are delocalized over whole molecule and thus can not be assigned to several local bonds. This is a characteristic feature of cyclic compounds, particularly aromatic compounds. ${ }^{17}$ Therefore we represent in Table 1 only the approximate mode descriptions.
$\boldsymbol{a}_{g}$ symmetry. Under $D_{2 k}$ symmetry, $a_{g}$ modes are infrared inactive and Raman active. On the basis of our calculations and the reported infrared and Raman spectra of AN$h_{10}$, we made a reliable one-to-one correspondence between the fundamentals reported experimentally and our calculated frequencies with the B3LYP and HF methods. Twelve frequencies which we identify as $a_{1}$ fundamentals are 397,625 , $754,1007,1164,1264,1412,1480,1556,3027,3048$ and 3072 $\mathrm{cm}^{-1}$. It is shown that the HF values deviate from the experimental data, particularly in the $\mathrm{C}-\mathrm{H}$ stretching region.

For AN $-d_{10}$, among the twelve fundamentals, eleven fundamentals can be identified as $382,602,709,842,1156,1388$, $1402,1534,2258,2266$, and $2288 \mathrm{~cm}^{-1}$, the frequencies based on the reported Raman spectra by Bree and Kidd. ${ }^{15 b}$

For the radical cation, almost all of the frequencies are close to those of the protiated counterpart except those in the $\mathrm{C}-\mathrm{H}$ stretching region in which the frequencies are found to be higher than those of the neutral system.
It is shown in the B3LYP calculation that most of the fundamentals involve more or less ring deformation and in-

Table 1. Comparison of the observed and calculated vibrational spectra of anthracene ${ }^{*}$

| sym | No | Anthracene ( $h 10$ ) |  |  |  |  |  |  | Anthracene ( $d 10$ ) |  |  | Approx. ${ }^{\text {A }}$ mode |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | HF |  |  |  | B3LYP |  | exp | B3LYP |  | exp ${ }^{\prime}$ |  |
|  |  | freq | $I_{i k}$ | $I_{\text {sum }}$ | DP ${ }^{\text {s }}$ | freq | $I_{1 R}$ |  | freq | $I_{H}$ |  |  |
| $a_{g}$ | 1 | 376 | 0.00 | 15.99 | 0.24 | 384 | 0.00 | 397 | 370 | 0.00 | 382 | ring defm |
|  | 2 | 611 | 0.00 | 1.66 | 0.47 | 618 | 0.00 | 625 | 594 | 0.00 | 602 | ring defm |
|  | 3 | 735 | 0.00 | 42.94 | 0.10 | 738 | 0.00 | 754 | 693 | 0.00 | 709 | ring defm |
|  | 4 | 971 | 0.00 | 22.41 | 0.44 | 1002 | 0.00 | 1007 | 828 | 0.00 | - | ip C-H bend, ring defm |
|  | 5 | 1152 | 0.00 | 27.05 | 0.18 | 1155 | 0.00 | 1164 | 835 | 0.00 | 842 | ip C-H bend |
|  | 6 | 1225 | 0.00 | 203.89 | 0.19 | 1255 | 0.00 | 1264 | 1149 | 0.00 | 1156 | ring defm, ip $\mathrm{C}-\mathrm{H}$ bend |
|  | 7 | 1398 | 0.00 | 860.51 | 0.21 | 1391 | 0.00 | 1412 | 1378 | 0.00 | 1388 | ring defm, ip C-H bend |
|  | 8 | 1468 | 0.00 | 71.19 | 0.57 | 1478 | 0.00 | 1480 | 1399 | 0.00 | 1402 | ip C-H bend, ring defm |
|  | 9 | 1577 | 0.00 | 143.22 | 0.52 | 1550 | 0.00 | 1556 | 1526 | 0.00 | 1534 | ring defm, ip C-H bend |
|  | 10 | 2993 | 0.00 | 19.34 | 0.13 | 3056 | 0.00 | 3027 | 2257 | 0.00 | 2258 | sym $\mathrm{C}_{9.10}-\mathrm{H}$ str |
|  | 11 | 3002 | 0.00 | 305.06 | 0.61 | 3065 | 0.00 | 3048 | 2262 | 0.00 | 2266 | $\mathrm{C}-\mathrm{H}$ str |
|  | 12 | 3024 | 0.00 | 688.06 | 0.13 | 3089 | 0.00 | 3072 | 2290 | 0.00 | 2288 | tot sym C-H str |
| $a_{\text {x }}$ | 1 | 120 | 0.00 | 0.00 | 0.00 | 120 | 0.00 | $137{ }^{\text {d }}$ | 109 | 0.00 | 110 | ring tor |
|  | 2 | 486 | 0.00 | 0.00 | 0.00 | 489 | 0.00 | $552^{\text {d }}$ | 437 | 0.00 | - | ring tor |
|  | 3 | 734 | 0.00 | 0.00 | 0.75 | 731 | 0.00 | 743 | 621 | 0.00 | - | ring tor |
|  | 4 | 855 | 0.00 | 0.00 | 0.00 | 839 | 0.00 | 858 | 729 | 0.00 | - | oop C-H bend |
|  | 5 | 1001 | 0.00 | 0.00 | 0.75 | 953 | 0.00 | $958{ }^{\text {d }}$ | 796 | 0.00 | - | oop C-H bend |
| $b_{1 g}$ | 1 | 227 | 0.00 | 7.14 | 0.75 | 230 | 0.00 | 244 | 215 | 0.00 | 228 | ring tor |
|  | 2 | 474 | 0.00 | 3.80 | 0.75 | 471 | 0.00 | 479 | 405 | 0.00 | 415 | ring tor |

Table 1. Continued

${ }^{4}$ Vibrational frequencies in $\mathrm{cm}^{-1}$. ${ }^{b}$ Depolarization ratio. ${ }^{\text {'Otherwise noted elsewhere, the frequencies from the result of Bakke } e t}$ al. (ref 14 d ). ${ }^{d}$ ref $14 \mathrm{a}-\mathrm{c}$. ${ }^{\text {ref }} 17$. 'ref 15 , the numbers noted by bold face refer to our new assignment based on the deuterium isotope effect in the vibrational shift, ${ }^{g}$ ref 16 , ${ }^{h}$ Mode descriptions are from the B3LYP/6-31G ${ }^{*}$, ip: in-plane; oop: out-of-plane, the modes assigned are sorted in the order of their contributions to the vibrational motions.

Table 2. Comparison of observed and calculated vibrational spectra of anthracene radical cation ${ }^{a}$

| sym | No. | HF/6-31G* |  |  |  | B3LYP/6-31G** |  | exp ${ }^{\text {c }}$ | $\exp ^{\text {d }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | freq | $I_{L R}$ | $I_{\text {Ran }}$ | DP ${ }^{6}$ | freq | $I_{I R}$ |  |  |
| $a_{k}$ | 1 | 371 | 0.00 | 269.14 | 0.14 | 381 | 0.00 |  |  |
|  | 2 | 590 | 0.00 | 375.24 | 0.50 | 604 | 0.00 |  |  |
|  | 3 | 726 | 0.00 | 770.74 | 0.17 | 738 | 0.00 |  |  |
|  | 4 | 1006 | 0.00 | 2745.65 | 0.50 | 1020 | 0.00 |  |  |
|  | 5 | 1153 | 0.00 | 351.16 | 0.16 | 1167 | 0.00 |  |  |
|  | 6 | 1209 | 0.01 | 5687.58 | 0.29 | 1249 | 0.00 |  |  |
|  | 7 | 1309 | 0.00 | 741.86 | 0.72 | 1374 | 0.00 |  |  |
|  | 8 | 1490 | 0.00 | 1869.33 | 0.58 | 1493 | 0.90 |  |  |
|  | 9 | 1569 | 0.00 | 269.45 | 0.38 | 1553 | 0.00 |  |  |
|  | 10 | 3021 | 0.00 | 88.69 | 0.30 | 3083 | 0.00 |  |  |
|  | 11 | 3028 | 0.00 | 153.91 | 0.75 | 3093 | 0.00 |  |  |
|  | 12 | 3049 | 0.00 | 591.86 | 0.40 | 3114 | 0.00 |  |  |
| $a_{u}$ | 1 | 117 | 0.00 | 0.00 | 0.75 | 117 | 0.00 |  |  |
|  | 2 | 471 | 0.00 | 0.00 | 0.75 | 467 | 0.00 |  |  |
|  | 3 | 708 | 0.00 | 0.00 | 0.75 | 711 | 0.00 |  |  |
|  | 4 | 892 | 0.00 | 0.00 | 0.75 | 871 | 0.00 |  |  |
|  | 5 | 1028 | 0.00 | 0.00 | 0.75 | 986 | 0.00 |  |  |
| $b_{4}$ | 1 | 214 | 0.00 | 5.17 | 0.75 | 214 | 0.00 |  |  |
|  | 2 | 448 | 0.00 | 1.80 | 0.75 | 447 | 0.00 |  |  |
|  | 3 | 769 | 0.00 | 0.01 | 0.75 | 748 | 0.00 |  |  |
|  | 4 | 988 | 0.00 | 2.69 | 0.75 | 951 | 0.00 |  |  |
| $b_{14}$ | 1 | 236 | 14.41 | 0.00 | 0.74 | 225 | 0.42 |  |  |
|  | 2 | 631 | 0.50 | 0.00 | 0.74 | 643 | 0.01 |  |  |
|  | 3 | 876 | 4.82 | 0.00 | 0.73 | 881 | 0.53 |  |  |
|  | 4 | 1112 | 4.64 | 0.00 | 0.39 | 1123 | 0.12 |  |  |
|  | 5 | 1232 | 3.29 | 0.00 | 0.73 | 1266 | 21.42 | 1291(0.07) | 1290(0.06) |
|  | 6 | 1269 | 27.46 | 0.00 | 0.74 | 1277 | 5.77 |  | 1315(0.06) |
|  | 7 | 1442 | 84.41 | 0.00 | 0.74 | 1442 | 19.41 |  | 1430(0.01) |
|  | 8 | 1582 | 18.52 | 0.00 | 0.35 | 1570 | 58.05 |  | 1586(0.14) |
|  | 9 | 3019 | 0.11 | 0.00 | 0.75 | 3081 | 0.32 |  |  |
|  | 10 | 3026 | 31.98 | 0.00 | 0.75 | 3090 | 0.48 |  |  |
|  | 11 | 3039 | 5.32 | 0.00 | 0.44 | 3105 | 1.73 |  |  |
| $b_{24}$ | 1 | 266 | 0.00 | 7.94 | 0.75 | 264 | 0.00 |  |  |
|  | 2 | 540 | 0.00 | 1.15 | 0.75 | 541 | 0.00 |  |  |
|  | 3 | 716 | 0.00 | 3.77 | 0.75 | 732 | 0.00 |  |  |
|  | 4 | 863 | 0.00 | 1.33 | 0.75 | 841 | 0.00 |  |  |
|  | 5 | 926 | 0.00 | 50.31 | 0.75 | 919 | 0.00 |  |  |
|  | 6 | 1029 | 0.00 | 19.17 | 0.75 | 988 | 0.00 |  |  |
| $b_{2 u}$ | 1 | 580 | 2.29 | 0.00 | 0.60 | 587 | 1.87 |  |  |
|  | 2 | 793 | 32.86 | 0.00 | 0.22 | 805 | 3.36 |  |  |
|  | 3 | 1008 | 31.12 | 0.00 | 0.16 | 1020 | 6.01 | $1034(0.20)$ |  |
|  | 4 | 1086 | 30.21 | 0.00 | 0.18 | 1155 | 0.75 |  | 1183(0.01)? |
|  | 5 | 1136 | 2030.51 | 0.00 | 0.13 | 1179 | 161.32 | 1188(0.98) | $1189(0.70)$ |
|  | 6 | 1209 | 1138.72 | 0.04 | 0.29 | 1341 | 310.77 | 1341(1.00) | 1341(1.00) |
|  | 7 | 1322 | 449.39 | 0.00 | 0.26 | 1395 | 143.87 | 1418(0.97) | 1418(0.86) |
|  | 8 | 1447 | 31.07 | 0.00 | 0.53 | 1447 | 55.86 | 1457(0.05) | 1457(0.07) |
|  | 9 | 1534 | 41.51 | 0.00 | 0.16 | 1528 | 83.14 | 1540(0.04) | 1540(0.15) |
|  | 10 | 3027 | 8.80 | 0.00 | 0.19 | 3092 | 1.42 |  |  |
|  | 11 | 3048 | 18.09 | 0.00 | 0.45 | 3114 | 3.16 |  |  |
| $b_{38}$ | 1 | 401 | 0.00 | 5284.88 | 0.75 | 371 | 0.00 |  |  |
|  | 2 | 537 | 0.00 | 7358.39 | 0.75 | 499 | 0.00 |  |  |
|  | 3 | 885 | 0.00 | 17.43 | 0.75 | 891 | 0.00 |  |  |
|  | 4 | 1082 | 0.00 | 207.36 | 0.75 | 1086 | 0.00 |  |  |

Table 2. Continued

| sym | No. | HF/6-31G* |  |  |  | B3LYP/6-31G* |  | exp ${ }^{\text {c }}$ | $\exp ^{\text {a }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | freq | $I_{\text {IR }}$ | $I_{\text {Raw }}$ | DP ${ }^{\text {b }}$ | freq | $I_{t}$ |  |  |
| $b_{3 s}$ | 5 | 1213 | 0.00 | 9619.26 | 0.75 | 1176 | 0.00 |  |  |
|  | 6 | 1266 | 0.00 | 4159.98 | 0.75 | 1264 | 0.00 |  |  |
|  | 7 | 1391 | 0.00 | 19.15 | 0.75 | 1387 | 0.00 |  |  |
|  | 8 | 1531 | 0.00 | 1033.58 | 0.75 | 1509 | 0.00 |  |  |
|  | 9 | 2401 | 0.00 | >9999.99 | 0.75 | 1530 | 0.00 |  |  |
|  | 10 | 3025 | 0.00 | 5777.89 | 0.75 | 3090 | 0.00 |  |  |
|  | 11 | 3039 | 0.00 | 3915.57 | 0.75 | 3105 | 0.00 |  |  |
|  | 1 | 83 | 1.59 | 0.00 | 0.75 | 85 | 1.62 |  |  |
|  | 2 | 372 | 2.71 | 0.00 | 0.75 | 372 | 0.12 |  |  |
|  | 3 | 423 | 31.54 | 0.00 | 0.75 | 432 | 22.79 | 432(0.07) |  |
|  | 4 | 759 | 102.49 | 0.00 | 0.75 | 740 | 77.64 |  | $748(0.26)$ |
|  | 5 | 907 | 38.37 | 0.00 | 0.75 | 895 | 31.05 | 912(0.15) | 912(0.09) |
|  | 6 | 992 | 5.64 | 0.00 | 0.75 | 959 | 8.19 |  |  |

"Vibrational frequencies in $\mathrm{cm}^{-1}$. ${ }^{\text {D }}$ Depolarization ratio, 'ref 17 . ${ }^{\text {t }}$ ref 18. "The numbers in the parentheses are relative intensities.
plane $\mathrm{C}-\mathrm{H}$ bending.
$\boldsymbol{a}_{u}$ symmetry. The modes are infrared and Raman inactive vibrations. Five fundamentals at 137, 552, 743, 858 and $958 \mathrm{~cm}^{-1}$ are observed in the Raman spectrum. ${ }^{14}$ For the perdeuterated anthracene, only one fundamental is found at $110 \mathrm{~cm}^{-1}$. Most of the fundamentals include more or less ring torsions and out-of-plane $\mathrm{C}-\mathrm{H}$ bending. The calculated fundamentals in the deuterated and radical cation systems can not be compared with the experimental values due to the missing lines.
$\boldsymbol{b}_{1 g}$ symmetry. Four fundamentals attributable to $b_{l q}$ mode are $244,479,747$, and $956 \mathrm{~cm}^{-1}$. For the deuterated molecule, only three fundamentals are identified as 228,415 , and $613 \mathrm{~cm}^{-1}$. In the ring torsional modes, the isotope effect is quite small, reflecting the skeletal vibration of the molecule. The frequencies of the radical cation are very close to those of the neutral molecule. Almost all of the vibrations include out-of-plane bending modes, and some of them involve the ring torsion.
$\boldsymbol{b}_{1 \mu}$ symmetry. Eleven modes which we identify as $b_{14}$ are $234,653,906,1147,1272,1317,1448,1620,3007,3053$, and $3084 \mathrm{~cm}{ }^{1}$. Except fundamental $9 b_{1 u}\left(3007 \mathrm{~cm}^{-1}\right)$, the B3LYP calculation reproduces quite accurately the experimental data. It is shown that the HF values greatly underestimate the experimental frequencies particularly in the $\mathrm{C}-\mathrm{H}$ stretching region. In the deuterated analog, only ten fundamentals are identified as $b_{14}$. These are $220,592,822,879$, 1258, 1406, 1584, 2248, 2264, and $2283 \mathrm{~cm}^{-1}$. These fundamentals are well correlated with our B3L.YP calculation, correctly reflecting the isotope effect. The B3LYP method is also found to correctly reproduce the vibrational frequencies of the radical cation.

Almost all of the vibrations include the ring deformation, in-plane $\mathrm{C}-\mathrm{H}$ bending and stretching.
$\boldsymbol{b}_{2 g}$ symmetry. Five fundamentals attributable to $b_{2 g}$ mode are $287,580,773,896$, and $952 \mathrm{~cm}^{-1}$. The previously assigned fundamentals corresponding to the $4 b_{2 g}$ and $5 b_{2 g}$, 896 and $916 \mathrm{~cm}^{-1}$ are not correlated with the calculated frequencies. But the corresponding frequencies in the deu-
terated molecules are well predicted. As seen in Table 1, we leave the $4 b_{26}$ fundamental unassigned and assume that the peak at $896 \mathrm{~cm}^{-1}$ is attributed to the $5 b_{24}$.

For AN- $d_{10,}$ only four fundamentals are identified as $b_{28}$. These are $261,644,659$, and $761 \mathrm{~cm}^{-1}$. However, it is found in the deuteated system that the $4 b_{28}$ value ( $659 \mathrm{~cm}^{-1}$ ) is found to be very close to the B3LYP value.
Almost all of the vibrations include the out-of-plane bending and ring torsional modes.
$\boldsymbol{b}_{2 u}$ symmetry. Eleven modes which we identify as $b_{2 n}$ are 601, 809, 998, 1124, 1167, 1346, 1397, 1460, 1534, 3021, and $3048 \mathrm{~cm}^{-1}$. Two fundamentals, 1346 ( $6 b_{24}$ ) and $1460 \mathrm{~cm}^{-1}$ $\left(8 b_{2 u}\right)$, are from the result of Szczepanski. ${ }^{17}$ The Bakke's value, $1494 \mathrm{~cm}^{-1}$, is found to deviate very much from our calculated value, $1445 \mathrm{~cm}^{-1}$. For AN $-d_{10}$, nine fundamentals could be identified as $\delta_{24}$. These are $575,703,824,941,1245$, $1298,1335,1493$, and $2267 \mathrm{~cm}^{-1}$. It is assumed that the $3 b_{2 n}$ peak is immersed in the $4 b_{2 \mu}$ peak ( $824 \mathrm{~cm}^{-1}$ ).

Here it is also found that the C-H frequencies of $\mathrm{AN}^{+}$$h_{10}$ are higher than those of AN- $h_{10}$.

Most of the vibrations include in-plane C-H bending, ring deformation, and $\mathrm{C}-\mathrm{H}$ stretching modes.
$\boldsymbol{b}_{3 g}$ symmetry. Ten fundamentals attributable to the $b_{3 k}$ mode are 397, 522, 903, 1102, 1187, 1273, 1574, 1632, 3017 and $3054 \mathrm{~cm}^{-1}$. Like the $4 b_{2 k}$ fundamental, the previously assigned fundamental $7 b_{38}, 1433 \mathrm{~cm}^{-1}$ deviate very much from the calculated value, and thus we leave it unassigned.

For AN- $d_{10}$, only eight fundamentals are identified as $b_{38}$. These are $367,502,816,886,943,1233,1615$, and $2276 \mathrm{~cm}^{-1}$ whose values are very close to our B3LYP frequencies. The fundamentals at 367 and $886 \mathrm{~cm}^{-1}$ are newly assigned.
In the skeletal ring deformation, the frequencies such as fundamentals $8 \mathrm{~b}_{3 k}$ and $9 b_{3 g}$ of the cation are much lower than those corresponding to the neutral system, reflecting the decrease of the bond orders in the rings.
Almost all the vibrations include the out-of-plane bending and ring torsional modes.
$\boldsymbol{b}_{34}$ symmetry. Six modes which we identify as $b_{3 u}$ are $106,380,474,737,883$, and $956 \mathrm{~cm}^{-1}$. For AN- $d_{10}$, fundamen-
tals we identify as $b_{34}$ are $102,349,397,560,722$, and 784 $\mathrm{cm}^{-1}$ which correctly match the calculated fundamentals. Most of the vibrations include in-plane C-H bending, ring deformation, and $\mathrm{C}-\mathrm{H}$ stretching modes.

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