

A SURVEY ON SYMPLECTIC GEOMETRY

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1. Introduction

A *symplectic manifold* is a pair (M, ω) consisting of a smooth manifold M and a non-degenerate closed 2-form ω on M . Locally, $\omega = \sum_{i,j=1}^n \omega_{ij} dx^i \wedge dx^j$ and $d\omega = 0$, where $n = \dim M$. The condition $d\omega = 0$ implies that locally $\omega = d\alpha$ with $\alpha = \sum_{k=1}^n \alpha_k dx^k$. There are three main sources of symplectic manifolds.

(A) Phase space for classical mechanical systems

$M = T^*(N)$ is the cotangent bundle of a smooth manifold N , i.e., the configuration space. $T(N)$ is the position-velocity space and $T^*(N)$ is the impulse-coordinate space. If q^1, \dots, q^n are local coordinate in N , $\partial_1, \dots, \partial_n$ the corresponding tangent vector fields, p_1, \dots, p_n the corresponding coordinates in the fibres of $T^*(N)$ so that $\langle p, \partial_i \rangle = p_i$, then the 1-form $\alpha = \sum_{i=1}^n p_i dq^i$ is a canonically defined form which is invariant under $\text{Diff}(N)$ and $\omega = d\alpha$ is closed (even exact) and non-degenerate.

(B) Complex projective algebraic varieties

These are subvarieties of complex projective varieties $P_n(\mathbb{C})$ defined by algebraic equations with a canonical Kähler form obtained by restricting the canonical Kähler

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form ω on $P_n(\mathbb{C})$. If $\omega = \operatorname{Re}(\omega) + i\operatorname{Im}(\omega)$, the real part $\operatorname{Re}(\omega)$ of ω defines a Riemannian metric and the imaginary part $\operatorname{Im}(\omega)$ defines a symplectic structure. These can be thought of as a *superization* of an even and an odd part. The even parts are Riemannian manifolds and the odd parts are symplectic manifolds.

(C) Coadjoint orbits of a Lie group

The study of coadjoint orbits of a Lie group is quite closely related to that of irreducible unitary representations of G . The relation between these was given by A. Kirillov [K1] in the case when G is a *nilpotent* Lie group. For a wide class of Lie groups including but not restricted to nilpotent groups, we get a similar relation between the coadjoint orbits and the unitary dual \hat{G} (cf. [K2]). For a general theory of a Lie group, we refer to [He] and [V].

Let G be a simply connected Lie group with Lie algebra \mathfrak{g} . Then we have the coadjoint mapping

$$\operatorname{Ad}^*(x) : \mathfrak{g}^* \rightarrow \mathfrak{g}^*, \quad x \in G.$$

Here $\operatorname{Ad}^*(x)$ is the *contragredient* of the adjoint mapping $\operatorname{Ad}(x) = \mathfrak{g} \rightarrow \mathfrak{g}$. Therefore

$$(\operatorname{Ad}^*(x)l)(X) = l(\operatorname{Ad}(x^{-1})X), \quad l \in \mathfrak{g}^*, X \in \mathfrak{g}. \quad (1.1)$$

For the present time being, we fix a \mathbb{R} -linear form $l \in \mathfrak{g}^*$ on \mathfrak{g} once and for all. We let

$$G_l := \{x \in G \mid \operatorname{Ad}^*(x)l = l\} \quad (1.2)$$

the stabilizer of the coadjoint action Ad^* of G on \mathfrak{g}^* at l . Since G_l is a closed subgroup of G , G_l is a Lie subgroup of G . We denote by \mathfrak{g}_l the Lie subalgebra of \mathfrak{g}

corresponding to G_t . For all $X \in \mathfrak{g}_t$, $Y \in \mathfrak{g}$ and $t \in \mathbb{R}$, we have

$$\begin{aligned} \langle Y, l \rangle &= \langle Y, \text{Ad}^*(\exp tX)l \rangle \\ &= \langle \text{Ad}(\exp(-tX))Y, l \rangle \\ &= \langle e^{-t \text{ad} X}(Y), l \rangle \\ &= \langle X - t[X, Y] + \frac{t^2}{2}[X, [X, Y]] + \dots, l \rangle. \end{aligned}$$

Taking the derivative with respect to t at $t = 0$ yields

$$\langle [X, Y], l \rangle = 0 \quad \text{for all } X \in \mathfrak{g}_t \text{ and } Y \in \mathfrak{g}. \quad (1.3)$$

In order to interpret this identity geometrically, we first define the skew-symmetric \mathbb{R} -bilinear form $B_l \in \Lambda^2(\mathfrak{g}^*)$ associated with l on \mathfrak{g} by

$$B_l(X, Y) := \langle [X, Y], l \rangle, \quad X, Y \in \mathfrak{g}. \quad (1.4)$$

Lemma 1.1. *Let $\text{rad } B_l$ be the radical of B_l in \mathfrak{g} , i.e.,*

$$\text{rad } B_l := \{X \in \mathfrak{g} \mid B_l(X, Y) = 0 \text{ for all } Y \in \mathfrak{g}\}.$$

Then

$$\text{rad } B_l = \mathfrak{g}_l = \{X \in \mathfrak{g} \mid \text{ad}^*(X)l = 0\}.$$

Here ad^* denotes the differential of $\text{Ad}^* : G \rightarrow GL(\mathfrak{g}^*)$.

Proof. By (1.3), $\mathfrak{g}_l \subseteq \text{rad } B_l$. Conversely, suppose X is an element of \mathfrak{g} such that $B_l(X, Y) = 0$ for all $Y \in \mathfrak{g}$. Then

$$\begin{aligned} \langle Y, l \rangle &= \langle e^{-\text{ad}(tX)}(Y), l \rangle & (*) \\ &= \langle \text{Ad}(\exp(-tX))Y, l \rangle \\ &= \langle Y, \text{Ad}^*(\exp(tX))l \rangle \end{aligned}$$

holds for all $Y \in \mathfrak{g}$. Therefore $\text{Ad}^*(\exp(tX))l = l$. Thus $X \in \mathfrak{g}_l$. It remains to prove the identity (*). Differentiating with respect to t at $t = 0$,

$$\left. \frac{d}{dt} \right|_{t=0} \langle e^{-\text{ad}(tX)}Y, l \rangle = - \langle [X, Y], l \rangle = 0.$$

Thus $\langle e^{-\text{ad}(tX)}Y, l \rangle$ is constant for t and taking $t = 0$, we get the identity (*). For all $X, Y \in \mathfrak{g}$, we have the identities

$$\begin{aligned} \langle Y, \text{ad}^*(X)l \rangle &= \left. \frac{d}{dt} \right|_{t=0} \langle Y, \text{Ad}^*(\exp tX)l \rangle \\ &= \left. \frac{d}{dt} \right|_{t=0} \langle \text{Ad}(\exp(-tX))Y, l \rangle \\ &= - \langle \text{ad}(X)Y, l \rangle \\ &= - \langle [X, Y], l \rangle = -B_l(X, Y). \end{aligned}$$

Therefore

$$\text{rad } B_l = \{X \in \mathfrak{g} \mid \text{ad}^*(X)l = 0\}.$$

□

Next we consider the smooth mapping $\Phi_l : G \rightarrow \mathfrak{g}^*$ defined by

$$\Phi_l(x) := \text{Ad}^*(x)l, \quad x \in G. \quad (1.5)$$

The mapping Φ_l is called the *coadjoint orbit mapping defined by* $l \in \mathfrak{g}^*$. Clearly the coadjoint orbit $\text{Ad}^*(G)l$ at l can be identified with the homogeneous manifold G/G_l . Let $\pi : G \rightarrow G/G_l$ be the canonical surjection. The bijective mapping

$$\dot{\Phi}_l : G/G_l \rightarrow \text{Ad}^*(G)l, \quad \dot{\Phi}_l(xG_l) := \text{Ad}^*(x)l, \quad x \in G \quad (1.6)$$

allows to identify G/G_l with $\text{Ad}^*(G)l$ and to equip $\text{Ad}^*(G)l$ with the structure of a C^∞ -submanifold of \mathfrak{g}^* in such a way that $\dot{\Phi}_l$ is a diffeomorphism. Since

$$\begin{aligned}\Phi_l(\exp tX) &= \text{Ad}^*(\exp tX)l \\ &= e^{\text{ad}^*(tX)}l \\ &= l + t \text{ad}^*(X)l + \frac{t^2}{2} \text{ad}^*(\text{ad}^*(X))l + \dots\end{aligned}$$

holds for all $X \in \mathfrak{g}$ and $t \in \mathbb{R}$,

$$\left. \frac{d}{dt} \right|_{t=0} \Phi_l(\exp tX) = \text{ad}^*(X)l.$$

Thus the differential $d\Phi_l(e) : \mathfrak{g} \rightarrow \mathfrak{g}^*$ of Φ_l at e is given by

$$d\Phi_l(e)(X) = \text{ad}^*(X)l, \quad X \in \mathfrak{g}. \quad (1.7)$$

Consequently, we have, for all $X, Y \in \mathfrak{g}$,

$$\begin{aligned}B_l(Y, X) &= \langle [Y, X], l \rangle \\ &= - \langle (\text{ad } X)Y, l \rangle \\ &= \langle Y, \text{ad}^*(X)l \rangle \\ &= \langle Y, d\Phi_l(e)(X) \rangle.\end{aligned}$$

We observe that if $X \in \mathfrak{g}_l$, $\langle Y, d\Phi_l(e)(X) \rangle = 0$ for all $Y \in \mathfrak{g}$ and hence $d\Phi_l(e)(X) = 0$. Therefore the image $d\Phi_l(e)(\mathfrak{g})$ is identified with the tangent space of the coadjoint orbit $\text{Ad}^*(G)l$ at the point $l \in \mathfrak{g}^*$. It is easy to show that the tangent space of $\text{Ad}^*(G)l$ at l is isomorphic to the quotient vector space $\mathfrak{g}/\text{rad } B_l$ (over \mathbb{R}).

Let \dot{B}_l denote the non-degenerate alternating \mathbb{R} -bilinear form on the quotient vector space $\mathfrak{g}/\text{rad } B_l$ induced by B_l .

Lemma 1.2. *The tangent space of the coadjoint orbit $\text{Ad}^*(G)l$ at l is a symplectic vector space with respect to \tilde{B}_l . In particular, it has an even dimension over \mathbb{R} .*

Proof. It follows immediately from the previous argument. \square

Now we are ready to prove that the coadjoint orbit $\Omega := \text{Ad}^*(G)l$ is a symplectic manifold. We denote by \tilde{X} the vector field on \mathfrak{g}^* corresponding to $X \in \mathfrak{g}$. That means that we have a canonical map $\mathfrak{g} \rightarrow T_l(\Omega)$ sending $X \rightarrow \tilde{X}_l := \text{ad}^*(X)l$. We observe that

$$\tilde{X}_l = \text{ad}^*(X)l = \left. \frac{d}{dt} \right|_{t=0} \text{Ad}^*(\exp tX)l.$$

According to Lemma 1.1, \mathfrak{g}_l is precisely the kernel of this map. We may define a 2-form B_Ω on Ω by

$$B_\Omega(\tilde{X}, \tilde{Y}) = B_\Omega(\text{ad}^*(X)l, \text{ad}^*(Y)l) := B_l(X, Y), \quad (1.8)$$

where $X, Y \in \mathfrak{g}$.

Theorem 1.3. *B_Ω is non-degenerate and closed. So the coadjoint orbit $\Omega := \text{Ad}^*(G)l$ at l is a symplectic manifold.*

Proof. Let \tilde{X} be a vector field on \mathfrak{g}^* corresponding to $X \in \mathfrak{g}$ such that

$$B_\Omega(\tilde{X}, \tilde{Y}) = 0 \quad \text{for all } \tilde{Y} \text{ with } X \in \mathfrak{g}.$$

Since $B_\Omega(\tilde{X}, \tilde{Y}) = \langle [X, Y], l \rangle = 0$ for all $Y \in \mathfrak{g}$, according to Lemma 1.1, $X \in \mathfrak{g}_l$. Thus $\tilde{X} = 0$. Hence B_Ω is non-degenerate. If $\tilde{X}_1, \tilde{X}_2, \tilde{X}_3$ are three vector fields on Ω ($X_1, X_2, X_3 \in \mathfrak{g}$), then

$$\begin{aligned} & dB_\Omega(\tilde{X}_1, \tilde{X}_2, \tilde{X}_3) \\ &= \tilde{X}_1(B_\Omega(\tilde{X}_2, \tilde{X}_3)) - \tilde{X}_2(B_\Omega(\tilde{X}_1, \tilde{X}_3)) + \tilde{X}_3(B_\Omega(\tilde{X}_1, \tilde{X}_2)) \\ &\quad - B_\Omega([\tilde{X}_1, \tilde{X}_2], \tilde{X}_3) + B_\Omega([\tilde{X}_1, \tilde{X}_3], \tilde{X}_2) - B_\Omega([\tilde{X}_2, \tilde{X}_3], \tilde{X}_1) \\ &= - \langle [[X_1, X_2], X_3] - [[X_1, X_3], X_2] + [[X_2, X_3], X_1], l \rangle = 0 \end{aligned}$$

Therefore B_Ω is closed. Then (Ω, B_Ω) is a symplectic manifold. \square

2. General Facts of Symplectic Geometry

On a symplectic manifold (M, ω) there is an isomorphism between vector and covector fields. Denoting the former space by $\text{Vect}(M)$ and the latter by $A^1(M)$, we have

$$\text{Vect}(M) \ni \xi \leftrightarrow \omega(\xi, \cdot) = \iota_\xi \omega \in A^1(M). \quad (2.1)$$

Let L denote the Lie derivative and let $\text{Vect}(M, \omega)$ denote the set of Hamiltonian vector fields, that is,

$$\text{Vect}(M, \omega) := \{\xi \in \text{Vect}(M) \mid L_\xi \omega = 0\}. \quad (2.2)$$

Theorem 2.1. *A vector field ξ is Hamiltonian if and only if the corresponding covector field $\iota_\xi \omega$ is closed.*

Proof.

$$L_\xi \omega = \iota_\xi \circ d\omega + d\iota_\xi \circ \omega = 0 \Leftrightarrow d\iota_\xi \omega = 0. \quad \square$$

If $\iota_\xi \omega$ is not only closed but exact, then we say that ξ is *strictly Hamiltonian*, and define its *Hamiltonian* f_ξ by :

$$\iota_\xi \omega = -df_\xi. \quad (2.3)$$

Conversely, if we start with a function f we obtain a Hamiltonian vector field ξ_f defined by : $\omega(\xi_f, \cdot) = -df$.

The space $C^\infty(M)$ is a Lie algebra with the *Poisson* bracket :

$$\{f, g\} = \omega(\xi_f, \xi_g) = -df(\xi_g) = -\xi_g f = \xi_f g.$$

Since

$$\xi_{\{f,g\}} = [\xi_f, \xi_g]$$

$f \rightarrow \xi_f : C^\infty(M) \rightarrow \text{Vect}(M)$ is a Lie algebra homomorphism. Suppose a Lie group G acts on M and preserves ω . There is a homomorphism $\mathfrak{g} \rightarrow \text{Vect}(M, \omega) : X \rightarrow \xi_X$. Suppose ξ_X is strictly Hamiltonian. If there is a Lie algebra homomorphism $\mathfrak{g} \rightarrow C^\infty(M) : X \rightarrow h_X$ so that the following diagram is commutative, the action is a *Poisson* action :

$$\begin{array}{ccc} & \text{Vect}(M, \omega) & \\ & \uparrow & \\ \mathfrak{g} & & C^\infty(M). \end{array}$$

Theorem 2.2. *Suppose G and M are connected and G acts transitively on M and preserves ω . There exists :*

- (1) *A covering $\tilde{M} \rightarrow M$*
- (2) *A central extension $\tilde{G} \rightarrow G$*
- (3) *A Poisson action of \tilde{G} on \tilde{M}*

so that the diagram below is commutative.

$$\begin{array}{ccc} \tilde{G} \times \tilde{M} & \longrightarrow & \tilde{M} \\ \downarrow & & \downarrow \\ G \times M & \longrightarrow & M. \end{array}$$

Here the horizontal arrows denote actions and vertical arrows the natural projections. \square

This leads to a characterization of homogeneous symplectic manifolds.

Theorem 2.3. *Each G -homogeneous manifold M is locally isomorphic to a coadjoint orbit of G , or of a central extension \tilde{G} of G .*

In view of the preceding Lemma one need only to consider Poisson actions. For these we introduce the *momentum* map $\mu : M \rightarrow \mathfrak{g}^*$ which is defined by :

$$\langle \mu(m), X \rangle = h_X(m). \quad (2.4)$$

Theorem 2.3 follows from

Theorem 2.4. *If G is connected,*

$$\mu : M \rightarrow \mathfrak{g}^*$$

is covariant, that is, $\mu(gm) = \text{Ad}^*(g)\mu(m)$.

Proof. Since G is connected it is enough to show infinitesimal covariance

$$-\xi_Y \mu(m) = \text{Ad}^*(Y)\mu(M)$$

all $Y \in \mathfrak{g}$. Applying both sides to $X \in \mathfrak{g}$ we see that the desired equality becomes

$$-\xi_Y h_X = -h_{[Y, X]}$$

which is a consequence of the definition of the Poisson bracket. \square

Momentum maps are also interesting in the infinite-dimensional situation. Here are some examples.

- (1) $M = G = \mathbb{R}^{2n}$, $\tilde{M} = M$, $\tilde{G} = H_n$.
- (2) $M = \mathbb{R}^{2n} \setminus \{0\}$, $G = Sp(2n, \mathbb{R})$. \mathfrak{g} may be identified with the space of symmetric matrices $Sym(2n, \mathbb{R})$ which is isomorphic to \mathfrak{g}^* in such a way that $\langle F, X \rangle = \text{tr}(FX)$. In this case $\tilde{M} = M$, $\tilde{G} = G$, and $\mu(v) = vv^t$, v a column vector.

- (3) $M = \mathbb{C}\mathbb{P}^n$, $G = U(n+1)$. If $(z_0 : z_1 : \cdots : z_n)$ are homogeneous coordinates of z , then

$$\mu(z) = (h_{ij}), \quad h_{ij} = \frac{z_i \bar{z}_j}{\|z\|^2}.$$

- (4) $M = T^*(N)$, $G = \text{Diff}(N)$, $\mathfrak{g} = \text{Vect}(N)$. Here μ is defined by :

$$\langle \mu(A), \xi \rangle = \langle A, \xi(p(A)) \rangle,$$

where $p : M \rightarrow N$ is the natural projection.

3. Polarization

Let \mathfrak{g} be a Lie algebra over a field K . Let \mathfrak{g}^* be the dual space of \mathfrak{g} . We recall that a Lie subalgebra \mathfrak{h} of \mathfrak{g} is said to be *subordinate* to $l \in \mathfrak{g}^*$ if \mathfrak{h} forms a totally isotropic vector space of \mathfrak{g} relative to the alternating K -bilinear form $B_l : \mathfrak{g} \times \mathfrak{g} \rightarrow K$ given by $B_l(X, Y) := \langle [X, Y], l \rangle$ with $X, Y \in \mathfrak{g}$ associated with l on \mathfrak{g} , i.e.,

$$B_l|_{\mathfrak{h} \times \mathfrak{h}} = \langle [\mathfrak{h}, \mathfrak{h}], l \rangle = 0.$$

Definition 3.1. A Lie subalgebra \mathfrak{h} of \mathfrak{g} subordinate to $l \in \mathfrak{g}^*$ is said to be a *K-polarization* of \mathfrak{g} for l if \mathfrak{h} is maximal among the totally isotropic vector subspaces of \mathfrak{g} relative to B_l . In other words, if \mathfrak{h} is a vector subspace of \mathfrak{g} such that $\mathfrak{h} \subseteq P$ and $B_l|_{P \times P} = 0$, then we have $\mathfrak{h} = P$. In particular, each K -polarization \mathfrak{h} of \mathfrak{g} for $l \in \mathfrak{g}^*$ is subordinate to l and contains $\text{rad } B_l = \mathfrak{g}_l$. Thus we have the inclusions,

$$\mathcal{Z} \hookrightarrow \text{rad } B_l \hookrightarrow \mathfrak{h}, \tag{3.1}$$

where \mathcal{Z} denotes the center of \mathfrak{g} .

Remark 3.2. A maximal totally isotropic vector subspace of \mathfrak{g} relative to B_l need not to be a K -polarization of \mathfrak{g} for $l \in \mathfrak{g}^*$. If \mathfrak{g} is finite dimensional over K , a Lie subalgebra \mathfrak{h} of \mathfrak{g} subordinate to $l \in \mathfrak{g}^*$ of maximal dimension over K is not necessarily a K -polarization of \mathfrak{g} for l . Moreover, it is not true that there exist K -polarizations of general Lie algebras \mathfrak{g} over K for all K -linear forms $l \in \mathfrak{g}^*$. However, if \mathfrak{g} is a nilpotent real Lie algebra, then there exist real polarizations of \mathfrak{g} for arbitrary \mathbb{R} -linear forms l on \mathfrak{g} and the subalgebras \mathfrak{h} of \mathfrak{g} subordinate to $l \in \mathfrak{g}^*$ of maximal dimension over \mathbb{R} are exactly the real polarizations of \mathfrak{g} for l .

Let E be a finite dimensional vector space over a field K and $B : E \times E \rightarrow K$ an alternating K -bilinear form on E . For any subset F of E , we let the vector subspace

$$F^\perp := \{x \in E \mid B(x, y) = 0 \text{ for all } y \in F\}$$

of E be the *orthogonal* subspace of E for F relative to B . In particular, $E^\perp = \text{rad } B$.

Suppose that F is a vector subspace of E and define the K -linear mapping $f_F : E \rightarrow \left(F / F \cap \text{rad } B\right)^*$ by

$$f_F(x) := B(x, \cdot)|_F, \quad x \in E. \quad (3.2)$$

Then $\ker f_F = F^\perp$. Consequently we have

$$\dim_K F - \dim_K(F \cap \text{rad } B) + \dim_K F^\perp = \dim_K E. \quad (3.3)$$

The vector subspace F of E is said to be *isotropic* or *coisotropic* relative to B if $F \subseteq F^\perp$ or $F^\perp \subseteq F$ respectively. In the *isotropic case* we have

$$2 \dim_K F \leq \dim_K E + \dim_K(F \cap \text{rad } B) \quad (3.4)$$

and in the *coisotropic case*

$$2 \dim_K F \geq \dim_K E + \dim_K(F \cap \text{rad } B). \quad (3.5)$$

Since $E^\perp = \text{rad } B$ is an isotropic vector subspace of E , a *maximal* vector subspace F among the isotropic vector subspaces of E satisfies $\text{rad } B \subseteq F$ and hence

$$2 \dim_K F \leq \dim_K E + \dim_K \text{rad } B. \quad (3.6)$$

Consequently the maximality property of F suggests that $\dim_K F$ actually attains the upper bound.

Lemma 3.2. *A vector subspace F of the finite dimensional vector space E over a field K is maximal among the isotropic vector subspaces of E relative to B if and only if*

$$2 \dim_K F = \dim_K E + \dim_K \text{rad } B.$$

Proof. It will suffice to show that the dimension of each vector subspace F of E which is maximal among the isotropic subspaces of E relative to B attains the *upper bound*

$$\frac{1}{2}(\dim_K E + \dim_K \text{rad } B).$$

Let \dot{B} denote the non-degenerate alternating K -bilinear form induced by B on the quotient vector space $E/\text{rad } B$ over K . Then $(E/\text{rad } B, \dot{B})$ is a symplectic vector space over K . In particular, $\dim_K(E/\text{rad } B)$ is an even positive integer. Choose a *Lagrangian* vector subspace L of $E/\text{rad } B$. Since L is isotropic and coisotropic, it coincides with its orthogonal vector subspace of $E/\text{rad } B$ for L relative to the symplectic form \dot{B} , i.e., $L = L^\perp$. It follows that

$$\dim_K L = \frac{1}{2} \dim_K(E/\text{rad } B).$$

Denote by $\pi : E \rightarrow E/\text{rad } B$ be the canonical surjection and let $F = \pi^{-1}(L)$ denote the preimage of L relative to π . Then F is an isotropic subspace of E relative to B

such that $\text{rad } B \subseteq F$ satisfying the following

$$\begin{aligned} \dim_K F &= \frac{1}{2} \dim_K(E/\text{rad } B) + \dim_K(\text{rad } B) \\ &= \frac{1}{2} \dim_K E - \frac{1}{2} \dim_K(\text{rad } B) + \dim_K(\text{rad } B) \\ &= \frac{1}{2} \dim_K E + \frac{1}{2} \dim_K(\text{rad } B). \end{aligned}$$

Since for each isotropic vector subspace P of E relative to B containing $\text{rad } B$ the image $\pi(P)$ is isotropic in $E/\text{rad } B$ relative to \dot{B} , the proof is complete. \square

If we apply the preceding lemma to the case when E is a finite dimensional Lie algebra \mathfrak{g} over a field K and $B \in \Lambda^2(E^*)$ is the alternating K -bilinear form $B_l : \mathfrak{g} \times \mathfrak{g} \rightarrow K$ defined by $B(X, Y) := \langle [X, Y], l \rangle$, then we get the following characterization of the K -polarization of \mathfrak{g} for l .

Proposition 3.3. *Let $l \in \mathfrak{g}^*$ denote a K -linear form on the finite dimensional Lie algebra over a field K . For a Lie subalgebra \mathfrak{h} of \mathfrak{g} subordinate to l the following conditions mutually equivalent :*

- (1) \mathfrak{h} forms a K -polarization of \mathfrak{g} for l .
- (2) For any element $X \in \mathfrak{g}$ such that $B_l(X, Y) = 0$ holds for all $Y \in \mathfrak{h}$, we have $X \in \mathfrak{h}$.
- (3) The orthogonal vector subspace \mathfrak{h}^\perp of \mathfrak{h} relative to B_l is contained in \mathfrak{h} .
- (4) $\dim_K \mathfrak{h} = \frac{1}{2}(\dim_K \mathfrak{g} + \dim_K(\text{rad } B))$.

Remark 3.4. *For more details on symplectic manifolds, we refer to [A-B], [A-N] and [Au]. For instance, [A-B] discusses Floer homology, the moduli of pseudoholomorphic curves and the compactness of the moduli space, [A-N] deals with Legendre singularities and cobordisms, and [Au] deals with equivariant cohomology and toric manifolds.*

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