

## THE STRUCTURE CONFORMAL VECTOR FIELDS ON A SASAKIAN MANIFOLD II

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ABSTRACT. The concept of the structure conformal vector field  $C$  on a Sasakian manifold  $M$  is defined. The existence of such a  $C$  on  $M$  is determined by an exterior differential system in involution. In this case  $M$  is a foliate manifold and the vector field  $C$  enjoys the property to be exterior concurrent. This allows to prove some interesting properties of the Ricci tensor and Obata's theorem concerning isometries to a sphere. Different properties of the conformal Lie algebra induced by  $C$  are also discussed.

### 0. Introduction

Let  $M(\Phi, \eta, \xi, g)$  be a  $(2m+1)$ -dimensional Sasakian manifold with soldering form  $dp \in \Gamma Hom(\Lambda^q TM, TM)$  ( $dp$ : canonical vector-valued 1-form) where  $\Phi, \eta, \xi$  and  $g$  are the  $(1,1)$ -tensor field, the structure 1-form, the structure vector field and the metric tensor of  $M$ , respectively. Since one may write  $\nabla\xi = \Phi dp$ , we give the following definition : Any vector field  $C$  such that

$$(0.1) \quad \nabla C = \rho dp + \lambda \nabla \xi ; \quad \rho, \lambda \in C^\infty M,$$

is defined as a *conformal vector field* ((0.1) implies  $\mathcal{L}_C g = 2\rho g$ )

In Section 2, it is proved that the existence of  $C$  on  $M(\Phi, \eta, \xi, g)$  is determined by an exterior differential system *in involution* (in the sence of É. Cartan [4]), and that any  $M$  which carries a vector field  $C$ , is foliated by *autoparallel* three-dimensional submanifold of scalar curvature  $+1$ , tangent to  $C, \Phi C$  and  $\xi$ . Besides such a Sasakian manifold

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possesses the remarkable property to be isometric to a unit sphere in a  $(2m+2)$ -dimensional Euclidean space [8].

Furthermore, any  $C$  is an *exterior concurrent* vector field (see [11], [9]) and of *conformal weight*  $\frac{2m+1}{m}$  [7].

Let  $\mathcal{D}_C$  be the autoparallel distribution tangent to  $C, \Phi C$  and  $\xi$ . In Section 3, we consider three conformal vector fields  $C_\nu \in \mathcal{D}_C, \nu = 1, 2, 3$  generated by  $C$  by means of the Lie bracket and agree to call them the *associated conformal vector fields of  $C$* . If  $\rho$  and  $\rho_\nu$  are the conformal scalars corresponding to  $C$  and  $C_\nu$  respectively, then it is proved that they define a 4-dimensional *eigenspace*  $\varepsilon^4(M)$  of the Laplacian  $\Delta$  of the eigenvalue  $2m + 1$ .

If  $r$  denotes the scalar curvature of  $M$ , then we may write

$$\mathcal{L}_C r = 2(2m(2m + 1) - r)\rho.$$

In this case when  $M$  is compact, some integral formulas of Watanabe type (see [13]) are obtained.

More generally, consider a  $K - contact$  manifold  $M(\Phi, \eta, \xi, g)$ , i.e. a contact metric manifold whose structure vector  $\xi$  is a Killing vector field [2].

We give the following definition : Any vector field  $X$  such that

$$(0.2) \quad \mathcal{L}_X \Omega = h\Omega + \gamma \wedge \eta$$

where  $\Omega = \frac{1}{2}d\eta, h \in C^\infty M, \gamma \in \Lambda^1 M$ , is called an *infinitesimal quasi-conformal contact transformation of  $\Omega$* . In particular, if  $X$  is a conformal vector field, then following [5], we will call it a *quasi-biconformal* vector field.

In Section 4, we discuss some infinitesimal transformations defined by  $C$  and its associated vector field  $C_\nu$ . We prove that if  $R$  denotes the curvature tensor, then the following formula holds:

$$(\mathcal{L}_C R)(Z, Z', Z'') = 2\rho(Z \wedge Z')Z'', \quad Z, Z', Z'' \in \mathcal{X}M$$

and the quality of  $C$  to be exterior concurrent (abbreviation : E.C.) implies

$$\mathcal{L}_C g(C, Z) = 2\rho g(C, Z).$$

This shows that  $C$  defines an infinitesimal conformal transformation of all the functions  $g(C, Z)$ .

We also prove that  $C$  and  $C_\nu$  are quasi-biconformal vector fields with respect to the pairing  $(\Omega, \eta)$ , and that  $\Omega$  is a *relatively integral invariant* of  $C_\nu$  (see [1]).

### 1. Preliminaries

Let  $(M, g)$  be an orientable  $C^\infty$ -Riemannian manifold and let  $\nabla$  be the covariant differential operator defined by the metric tensor  $g$ .

Let  $\Gamma(TM)$  be the set of sections of the tangent bundle  $TM$  and  $b : TM \rightarrow T^*M$  be the *musical isomorphism* [10] defined by  $g$ .

If, following [10], we denote by

$$A^q(M, TM) = \Gamma\text{Hom}(\Lambda^q TM, TM)$$

the set of vector-valued  $q$ -forms,  $q < \dim M$ , then

$$d^\nabla : A^q(M, TM) \rightarrow A^{q+1}(M, TM)$$

means the *exterior covariant derivative* operator with respect to  $\nabla$ . It should be noticed that generally  $d^{\nabla^2} = d^\nabla \circ d^\nabla \neq 0$ , unlike  $d^2$ .

If  $dp \in A^1(M, TM)$  denotes the soldering form of  $M$ , any vector field  $X$  such that

$$(1.1) \quad d^\nabla(\nabla X) = \nabla^2 X = \pi \wedge dp \in A^2(M, TM),$$

is defined as *exterior concurrent* (abbreviation :E.C.)(see [11],[9] ).

It has been proved [9] that  $\pi$  is necessarily given by

$$(1.2) \quad \pi = \nu b(X); \quad \nu \neq 0,$$

where  $\nu \in C^\infty M$  is the conformal scalar associated with  $X$ .

If  $\mathcal{R}$  denotes the *Ricci tensor* of  $\nabla$ , it follows from (1.1) and (1.2) that

$$(1.3) \quad \mathcal{R}(X, Z) = -(n - 1) \nu g(X, Z) \Rightarrow \nu = -\frac{1}{n - 1} \text{Ric} X,$$

where  $Z \in \mathcal{X}M$  and  $\dim M=n$ .

Let  $T \in \mathcal{X}M$  be any *conformal vector field* on  $M$  (or conformal Killing vector field), that is

$$(1.4) \quad \mathcal{L}_T g = 2\rho g \Leftrightarrow \langle \nabla_Z T, Z' \rangle + \langle \nabla_{Z'} T, Z \rangle = 2\rho \langle Z, Z' \rangle$$

where  $\rho \in C^\infty M; Z, Z' \in \mathcal{X}M$  and

$$(1.5) \quad \rho = \frac{\operatorname{div} T}{n}$$

If  $R$  and  $r$  denotes the curvature tensor and the scalar curvature of  $M$  respectively, we recall the following basic formulas (see [3], [13]):

$$(1.6) \quad \mathcal{L}_T \flat(Z) = 2\rho \flat(Z) + \flat[T, Z], \quad Z \in \mathcal{X}M,$$

$$(1.7) \quad \begin{aligned} (\mathcal{L}_T R)(Z, Z', Z'') &= -(\operatorname{Hess}_{\nabla \rho})(Z', Z'')Z + (\operatorname{Hess}_{\nabla \rho})(Z, Z'')Z' \\ &\quad - g(Z', Z'')(\nabla \operatorname{grad} \rho)Z + g(Z, Z'')(\nabla \operatorname{grad} \rho)Z', \end{aligned}$$

where  $(\operatorname{Hess}_{\nabla \rho})$  is a covariant and symmetric 2-tensor defined by  $(\operatorname{Hess}_{\nabla \rho})(Z, Z') = g(Z, H_\rho Z')$ ;  $H_\rho Z' = \nabla_{Z'} \operatorname{grad} \rho$ ; see also [3],

$$(1.8) \quad \mathcal{L}_T \mathcal{R}(Z, Z') = (\Delta \rho)g(Z, Z') - (n - 2)(\operatorname{Hess}_{\nabla \rho})(Z, Z')$$

and

$$(1.9) \quad \mathcal{L}_T r = (n - 1)\Delta \rho - r\rho,$$

where  $Z, Z', Z'' \in \mathcal{X}M$  and  $[, ]$  is the Lie bracket.

If  $T_1$  and  $T_2$  are conformal vector fields, then  $[T_1, T_2]$  is a conformal vector field with (see [3])

$$(1.10) \quad \rho_{[T_1, T_2]} = [\mathcal{L}_{T_1}, \rho_{T_2}] + [\rho_{T_1}, \mathcal{L}_{T_2}].$$

We also recall (see [13]) that if  $M$  is a compact manifold and its Ricci curvature is negative definite, then a conformal vector field other than zero vector field does not exist.

Any vector field  $X$  such that

$$(1.11) \quad \mathcal{L}_X \mathfrak{b}(X) = c(\operatorname{div} X)\mathfrak{b}(X); \quad c = \text{const}$$

is defined as a *self conformal* vector field.

We notice that in general (even if  $X$  is not a conformal vector field) if  $\dim M = n$ , then  $cn$  is called the *conformal weight of*  $\mathfrak{b}(X)$  (cf. [7]).

In addition, if  $M$  is a compact manifold, then the following integral formula holds (see [13]):

$$(1.12) \quad \int_M (\mathcal{R}(T, T - |\nabla T|^2 - \frac{n-2}{n}(\operatorname{div} T)^2) = 0.$$

We also recall the following theorem of M. Obata [8] (see also [3]): In order that a gradient vector field  $\operatorname{grad} f$  be an *infinitesimal concircular transformation* on an  $n$ -dimensional manifold  $M$ , it is necessary and sufficient that

$$(1.13) \quad \langle \nabla_Z \operatorname{grad} f, Z' \rangle = \nu \langle Z, Z' \rangle, \quad Z, Z' \in M,$$

where  $\nu$  is a non-vanishing scalar. If  $\nu = -c^2 f$ , then  $M$  is isometric to a sphere  $S^n$  of radius  $\frac{1}{c}$  in an  $(n + 1)$ -dimensional Euclidean space.

## 2. A structure conformal vector field on a Sasakian manifold

Let  $M(\Phi, \eta, \xi, g)$  be a  $(2m + 1)$ -dimensional contact metric manifold. In such a manifold the structure tensors  $\Phi, \eta$  and  $\xi$  satisfy (see [15]) the equations :

$$(2.1) \quad \begin{cases} \Phi\xi = 0, & \eta(\xi) = 1, \\ \Phi^2 = -I + \eta \otimes \xi, & \eta(Z) = g(\xi, Z), \\ g(\Phi Z, \Phi Z') = g(Z, Z') - \eta(Z)\eta(Z'), \\ g(\Phi Z, Z') = \frac{1}{2}d\eta(Z, Z'), & Z, Z' \in \mathcal{X}M. \end{cases}$$

The  $\Phi$ -Lie derivative is defined by

$$(2.2) \quad (\nabla\Phi)Z = \nabla\Phi Z - \Phi\nabla Z,$$

and it has been shown in [2] that  $\xi$  is a Killing vector fields if and only if  $\mathcal{L}_\xi \Phi$  vanishes. In this case M is called a *K-contact manifold*. A K-contact manifold for which one has

$$(2.3) \quad (\nabla_Z \Phi)Z' = -g(Z, Z')\xi + \eta(Z')Z,$$

is called a *Sasakian manifold*.

If M is a Sasakian manifold, then  $\xi$  is always E.C. and

$$(2.4) \quad \nabla^2 \xi = -\eta \wedge dp \Rightarrow \mathcal{R}(\xi, Z) = 2mg(\xi, Z)$$

(see [9]). Moreover, any E.C.vector field X satisfies

$$(2.5) \quad \nabla^2 X = -b(X) \wedge dp,$$

and the property of the exterior concurrency is invariant under the action of  $\Phi$  (i.e.  $\nabla^2 \Phi X = -b(\Phi X) \wedge dp$ ).

In the following we shall that

$$(2.6) \quad d\eta = 2\Omega$$

and agree to call  $\Omega$  the *fundamental presymplectic form* of  $M$  (abbreviation f.p.f.)

In the more general case when M is a K-contact manifold, we introduce the following two definitions

a) A vector field  $C$  on M such that

$$(2.7) \quad \nabla C = \rho dp + \lambda \nabla \xi; \quad \rho, \lambda \in C^\infty M,$$

is defined as a *structure conformal vector field*. Effectively, since  $\xi$  is a Killing vector field, it is easy to see that the equation (2.7) satisfies the conformal equation, that is (see (1.4)):

$$(2.8) \quad \mathcal{L}_C g = 2\rho g \Leftrightarrow \langle \nabla_Z C, Z' \rangle + \langle \nabla_{Z'} C, Z \rangle = 2\rho \langle Z, Z' \rangle,$$

where  $Z, Z' \in \mathcal{X}M$ , and this implies (see (1.5))

$$\operatorname{div} C = (2m + 1)\rho$$

b) Since we have set  $d\eta = 2\Omega$ , any vector field  $X$  such that

$$(2.10) \quad \mathcal{L}_X \Omega = \Phi \Omega + \gamma \wedge \eta; \quad \gamma \in \Lambda^1 M, \Phi \in C^\infty M,$$

is called an *infinitesimal quasi-conformal contact transformation of  $\Omega$*  (abbreviation : i.q.c.c.t)

Denote by  $\mu : TM \rightarrow T^*M, X \rightarrow i_X \Omega$  the bundle isomorphism defined by  $\Omega$ . If  $u$  is any 1-form on  $M(\Phi, \eta, \xi, g)$  such that  $du$  is equated by the second member of (2.10), then clearly  $\mu^{-1}(u)$  defines an i.q.c.c.t.

From now on we shall be concerned with Sasakian manifold carrying a structure conformal vector field  $C$ .

Next let

$$\mathcal{O} = \text{vect.}\{e_i, \Phi e_i = e_{i^*}, e_0 = \xi \mid i = 1, \dots, m; i^* = i + m\}$$

be an adapted local field of orthonormal frames on  $M$  and let

$$\mathcal{O}^* = \text{covect.}\{w^i, w^{i^*}, w^0 = \eta\}$$

be its associated coframe field.

Then the soldering form  $dp$  and É. Cartan's structure equations are :

$$(2.11) \quad dp = w^A \otimes e_A; \quad A \in \{i, i^*, 0\}$$

and

$$(2.12) \quad \nabla e = \theta \otimes e,$$

$$(2.13) \quad d\omega = -\Theta \otimes \omega,$$

$$(2.14) \quad d\theta = -\theta \wedge \theta + \Theta,$$

respectively.

In the above equation  $\theta$  (respectively  $\Theta$ ) is the local connection form in the bundle  $\mathcal{O}(M)$ . (respectively the curvature forms on  $M$ ), and in terms of  $\omega$ , the f.p.f.  $\Omega$  is expressed by

$$(2.15) \quad \Omega = \sum_i \omega^i \wedge \omega^{i^*},$$

Further since  $M$  is Sasakian, by (2.1),(2.3) and (2.12), one has

$$(2.16) \quad \theta_j^i = \theta_j^{i*}, \quad \theta_j^{i*} = \theta_i^{j*}$$

and by (2.16), we may check the formula  $\nabla^2 \Phi X = -b(\Phi X) \wedge dp(X$  is an E.C. vector field) and

$$R(Z, Z') = -\Phi R(Z, Z')\Phi + Z \wedge Z' - \Phi Z \wedge \Phi Z'$$

( $R$  is the curvature tensor).

Now in order to make simplifications, we set

$$(2.17) \quad \| C \|^2 = 2l, \quad b(C) = \alpha, \quad b(\Phi C) = \beta = i_C \Omega$$

and notice that one has

$$(2.18) \quad \beta = - \langle C, \nabla \xi \rangle$$

Next, with the help of (2.1) and (2.12), we obtain from (2.7) that

$$(2.19) \quad dl = \rho\alpha - \lambda\beta,$$

$$(2.20) \quad d\eta(C) = \rho\eta - \beta$$

and

$$(2.21) \quad d\alpha = 2\lambda\Omega \Rightarrow \lambda = \text{const..}$$

By (2.20), one gets at once

$$(2.22) \quad d\beta = d\rho \wedge \eta + 2\rho\Omega,$$

and by (2.17) the equation (2.22) implies

$$(2.23) \quad \mathcal{L}_C \Omega = 2\rho\Omega + d\rho \wedge \eta.$$



Since  $C$  is conformal vector field, then following the definition (0.2), we find the equation (2.23) that  $C$  is a quasi-biconformal vector field.

On the other hand, taking account of (2.4), one derives from (2.7) by covariant differentiation

$$(2.24) \quad \nabla^2 C = -(\lambda\eta - d\rho) \wedge p$$

The equation (2.24) proves that any structure conformal vector field on a Sasakian manifold is E.C.

Using (2.5), we find

$$(2.25) \quad \alpha = \flat(C) = \lambda\eta - d\rho,$$

and we notice that the equation (2.25) is consistent with (2.21).

Denote now by  $\sum$  the exterior differential system which defines the structure conformal vector field  $C$ . Then, by (2.19),(2.20),(2.21),(2.22) and (2.25), we see that the *characteristic number* of  $\sum$  (see [4]) are  $r = 5, s_0 = 3, s_1 = 2$ . Consequently, following É.Cartan's test [4], we conclude that  $\sum$  is in *involution* and depends on two arbitrary functions of one argument. Further, by (2.3) and (2.7), one derives

$$(2.26) \quad \nabla\Phi C = (\eta(C) - \lambda)d\rho + \rho\nabla\xi + d\rho \otimes \xi$$

Next, taking account of (2.25) and  $\text{div } Z = \text{tr } \nabla Z$ , one finds

$$(2.27) \quad \text{div}\Phi C = 2m(\eta(C) - \lambda)$$

We will outline the following property connected with this subject. First, by (2.25), the equation (2.22) becomes

$$(2.28) \quad d\beta = \eta \wedge \alpha + 2\rho\Omega,$$

and by (2.1) , one has

$$(2.29) \quad i_{\Phi C}\Omega = \flat(\Phi^2 C) = \left(\frac{2m+1}{m}\right)\eta(C)\eta - \alpha$$

Then, taking account of (2.27), one may write

$$(2.30) \quad \mathcal{L}_{\Phi C}\beta = \frac{1}{m}\beta(\text{div}\Phi C) = \frac{2m+1}{m} \frac{(\text{div}\Phi C)\beta}{\text{dim}M}$$

Hence, by definition (1.11), the equation (2.30) proves the following salient property: The structure conformal vector field  $C$  on a  $(2m + 1)$ -dimensional Sasakian manifold  $M$ , turns out, under the action of  $\Phi$ , to a self-conformal vector field of conformal weight  $\frac{2m+1}{m}$

Next by (2.5),  $\Phi C$  is also an E.C. vector field, i.e.

$$(2.31) \quad \nabla^2 \Phi C = -b(\Phi C) \wedge d\rho = -\beta \wedge d\rho.$$

Therefore, according to the general formula

$$S_{Z' \wedge Z} = \frac{\langle R(Z', Z)Z, Z' \rangle}{\|Z'\|^2 \cdot \|Z\|^2 - \langle Z', Z \rangle^2}; \quad Z, Z' \in \mathcal{X}M,$$

and by (2.24) and (2.31), one finds that the sectional curvature  $S_{\Phi C \wedge C}$  of  $\Phi C$  and  $C$  is given by

$$S_{\Phi C \wedge C} = \frac{\langle R(\Phi C, C), \Phi C \rangle}{\|C\|^2 \cdot \|\Phi C\|^2} = +1$$

Denote now by  $\mathcal{D}_C = \{C, \Phi C, \xi\}$  the  $\mathcal{D}$ -distribution defined by  $C, \Phi C$  and  $\xi$ . Then, if  $X_C, X'_C \in \mathcal{D}_C$  are any vector fields of  $\mathcal{D}_C$ , it is easy to see by (2.1), (2.7) and (2.26), that one has  $\nabla_{X'_C} X_C \in \mathcal{D}_C$  which expresses the fact that  $\mathcal{D}_C$  is an *autoparallel foliation* (cf. [6]). On the other hand, since  $\xi, C$  and  $\Phi C, \xi$  and E.C. vector fields, it follows, by linearity that any vector field  $X_C$  of  $\mathcal{D}_C$  is E.C.

As a consequence of this fact and the results of [9], we conclude that the leaf  $M_C$  of  $\mathcal{D}_C$  is an autoparallel submanifold of scalar curvature  $+1$  of the Sasakian manifold  $M(\Phi, \eta, \xi, g)$  under consideration.

Next, from (2.25) it follows

$$(2.32) \quad \text{grad } \rho = \lambda \xi - C$$

and taking account of (2.7) one gets at once

$$(2.33) \quad \nabla \text{grad } \rho = -\rho dp$$

which shows that  $\text{grad } \rho$  is a *concurrent vector field* [14]. From (2.33) one gets instantly

$$(2.34) \quad \langle \nabla_Z \text{grad } \rho, Z' \rangle = -\rho \langle Z, Z' \rangle.$$

Applying Obata's theorem (see (1.13)), we obtain that the Sasakian manifold under consideration enjoys the remarkable property to be isometric to a unit sphere in a  $(2m+2)$ -dimensional Euclidean space.

Thus, we proved the following theorem :

**Theorem 2.1** Any Sasakian manifold  $M(\Phi, \eta, \xi, g)$  which carries a structure conformal vector field  $C$  is foliated by autoparallel 3-dimensional submanifolds of scalar curvature  $+1$  tangent to  $C, \Phi C$  and  $\xi$  and is isometric to a unit sphere in a  $(2m + 2)$ -dimensional Euclidean space. Further, one has the following properties :

- (i) The existence of  $C$  is determined by an exterior differential system in involution.
- (ii) Any  $C$  is an E.C. vector field and defines an infinitesimal quasi-conformal contact transformation of  $\Omega$
- (iii) The vector field  $\Phi C$  is self-conformal of conformal weight  $\frac{2m+1}{m}$

### 3. Associated conformal vector fields of $C$

From (2.33) we get

$$(3.1) \quad \operatorname{div} \operatorname{grad} \rho = -(2m + 1)\rho$$

Applying the general formula

$$\Delta \mu = -\operatorname{div} \operatorname{grad} \mu; \quad \mu \in C^\infty M,$$

one finds

$$(3.2) \quad \Delta \rho = (2m + 1)\rho$$

This proves that the conformal scalar  $\rho$  is an *eigenfunction* of  $\Delta$  and has  $2m + 1$  as its associated *eigenvalue*. As a consequence of this fact, by (1.9), we may write

$$(3.3) \quad \mathcal{L}_C r = 2(2m(2m + 1) - r)\rho,$$

where  $r$  denotes the scalar curvature of  $M(\Phi, \eta, \xi, g)$ .

In the view of the further discussions, we agree to set

$$(3.4) \quad \text{grad}\rho = C_1,$$

and since  $C_1$  is concurrent, it also enjoys the properties to be conformal and E.C.

Therefore, by (1.10), the bracket  $[C, C_1]$  is also a conformal vector field: By (2.7) and (2.33), one has

$$(3.5) \quad [C, C_1] = \lambda(\Phi C - \rho\xi).$$

Noticing that

$$(3.6) \quad \Phi C - \rho\xi = [C, \xi]$$

and  $\lambda = \text{const.}$ , we shall denote by  $C_2$  the conformal vector field  $[C, \xi]$ .

Then with the help of (2.26) one derives

$$(3.7) \quad \nabla C_2 = \rho_2 dp,$$

where we have set (see (2.27))

$$(3.8) \quad \rho_2 = \eta(C) - \lambda = \frac{\text{div}\Phi C}{2m}$$

for the conformal scalar associated with  $C_2$ . By (2.20) and similar devices, one gets

$$(3.9) \quad \text{grad}\rho_2 = [\xi, C] = \rho\xi - \Phi C,$$

$$(3.10) \quad d\rho_2 = \rho\eta - \beta$$

and

$$(3.11) \quad \nabla\rho_2 = (2m + 1)\rho_2,$$

which shows that  $\rho_2$  enjoys the same properties as  $\rho$  and  $\rho_1 = -\rho$ , i.e. it is an eigenfunction of  $\nabla$  with  $2m + 1$  as the associated eigenvalue. One

also easily finds  $\nabla \text{grad} \rho_2 = -\rho_2 dp$  which matches Obata's theorem (see (2.34)).

Next consider the conformal vector field

$$(3.12) \quad [C, C_2] = C_3 = -\rho C_2 - \eta(C)C_1$$

By (2.20), (2.25), (2.33) and (3.7), the covariant differentiation of  $C_3$  is expressed by

$$(3.13) \quad \nabla C_3 = \lambda \rho dp + \Phi C \lambda + \rho(C \wedge \xi) + \lambda(\xi \wedge \Phi C).$$

Since  $\lambda = \text{const.}$ , the equation (3.13) shows that the conformal scalar  $\rho_3 = \lambda \rho$  associated with  $C_3$  satisfies also equation (3.2).

Finally consider the vector field  $[C_1, C_2]$  and denote by  $\rho_{[C_1, C_2]}$  its associated conformal scalar. Using computations, similar to those performed above, one finds  $\rho_{[C_1, C_2]} = 0$ , and this proves that the conformal vector field  $[C_1, C_2]$  is a *Killing vector field*, i.e.  $\mathcal{L}_{[C_1, C_2]} g = 0$ . Effectively, setting  $K = [C_1, C_2]$ , using (3.7) and applying a straightforward calculation, one finds

$$(3.14) \quad \nabla K = -\Phi C \wedge C' - \rho(C \wedge \xi).$$

It is easy to obtain from the equation (3.14) that the Killing equation is satisfied:  $\mathcal{L}_K g = 0$ . Besides, comparison of (3.13) and (3.14) gives

$$\nabla C_3 + \nabla K = \lambda \rho dp.$$

Therefore one may say that the Killing vector field  $K$  is homologous exterior concurrent to the conformal vector field  $C_3$ .

In the following we agree to call  $C_\nu, \nu = 1, 2, 3$  (respectively  $K$ ), the *associated conformal vector fields* (respectively the *associated Killing vector field*) of the structure vector field  $C$ .

As can be seen from (2.33), (3.2), (3.11) and (3.13), the conformal scalars  $\rho, \rho_1, \rho_2$  and  $\rho_3$  associated with these vector fields are eigenfunctions of  $\Delta$  with the same eigenvalue equal to  $\dim M$ . Therefore one may say that these eigenfunctions define a 4-dimensional *eigenspace*  $\varepsilon^4(M)$  of the eigenvalue  $\dim M$  [12].

Further we will need some properties of conformal vector fields  $C, C_1$  and  $C_3$ . We will discuss them now.

First of all, since we have seen that any vector field  $X_C$  of the distribution  $D_C$  is E.C., it follows, by (1.3), (2.17), (3.4) and (3.6) that

$$(3.15) \quad \begin{cases} \mathcal{R}(C, C) = 2m\|C\|^2 = 4ml, \\ \mathcal{R}(C_1, C_1) = 2m\|C_1\|^2 = 2m((\rho_2)^2 + 2l - \eta(C))^2, \\ \mathcal{R}(C_2, C_2) = 2m\|C_2\|^2 = 2m((\rho_1)^2 + 2l - \eta(C))^2, \end{cases}$$

where  $\rho_2$  is expressed by (3.8) and  $\rho_1 = -\rho$ .

Next, by (2.7), (2.33) and (3.7), a short calculation gives

$$(3.16) \quad \begin{cases} \|C\|^2 = (m + 1)\rho^2 + 2m\lambda^2, \\ \|C_1\|^2 = (m + 1)\rho_1^2, \\ \|C_2\|^2 = (m + 1)\rho_2^2. \end{cases}$$

Assume now that  $M(\Phi, \eta, \xi, g)$  is compact. Then, by reference to (1.12) and identifying  $\mathcal{T}$  with  $C, C_1$  and  $C_2$  respectively, from (3.15) and (3.16) one finds the following integral formulas:

$$(3.17) \quad \begin{cases} \int (\|C\|^2 + \lambda^2 - (1 + 2m)\rho^2) = 0, \\ \int (\|C\|^2 + \rho_2^2 - ((\eta(C))^2 - (1 + 2m)\rho^2)) = 0, \\ \int (\|C\|^2 + \rho^2 - ((\eta(C))^2 - (1 + 2m)\rho^2)) = 0. \end{cases}$$

The following theorem combines all results obtained in this section:

**THEOREM 3.1.** *Let  $C$  be a structure conformal vector field on a  $(2m + 1)$ -dimensional Sasakian manifold  $M(\Phi, \eta, \xi, g)$ . By means of the Lie bracket, the vector field  $C$  generates three other conformal vector fields  $C_\nu, \nu = 1, 2, 3$  called the associated conformal vector fields of  $C$ . The conformal scalars  $\rho, \rho_\nu$  corresponding to these vector fields, define a 4-dimensional eigenspace  $\varepsilon^4(M)$  of  $\Delta$  of the eigenvalue  $\dim M = 2m + 1$ . If  $\tau$  denotes the scalar curvature of  $M$ , then*

$$\mathcal{L}_C r = 2(2m(2m + 1) - r)\rho,$$

and if  $M$  is compact, the integral formulas (3.17) hold.

#### 4. Infinitesimal transformations defined by $C$ and its associated vector fields

In this section we shall consider some infinitesimal transformations defined by the structure conformal vector field  $C$ . If  $\rho$  is the conformal scalar associated with  $C$ , we recall that one has

$$(4.1) \quad (\text{Hess}_{\nabla} \rho)(Z, Z') = g(Z, H_{\rho} Z'); H_{\rho} Z' = \nabla_{Z'} \text{grad} \rho$$

and

$$(4.2) \quad \nabla \text{grad} \rho = -\rho dp.$$

First of all, by (1.7), (4.1) and (4.2), one finds after a short calculation that

$$(4.3) \quad (\mathcal{L}_C R)(Z, Z', Z'') = 2\rho(Z \wedge Z')Z''; Z, Z', Z'' \in \mathcal{X}M.$$

Next, since  $C$  is E.C. which implies (cf.(1.3))

$$(4.4) \quad \mathcal{R}(C, Z) = 2mg(C, Z),$$

by (1.8) and (3.2), one derives

$$(4.5) \quad \mathcal{L}_C g(C, Z) = 2\rho g(C, Z).$$

Therefore we may say that  $C$  defines as infinitesimal conformal transformation (i.c.t.) of all of the tensors  $g(C, Z)$ . By a routine matter we may prove that similar results hold for the associated conformal vector fields  $C_{\nu}$  of  $C$ .

Further, by the general formula (1.6) and (3.6), one readily finds

$$(4.6) \quad \mathcal{L}_C \eta = \rho \eta + \beta.$$

Now with the help of (4.6),(4.20) and the Lie derivative properties one writes

$$(4.7) \quad \mathcal{L}_{\eta(C)C} \eta = 2\eta(C)\rho\eta.$$

Hence the vector field  $Y'_C = \eta(C)C$  defines an (i.c.t) of the structure 1-form  $\eta$ . Further, since  $\text{div}C = (2m + 1)\rho$ , one finds by (2.20)

$$(4.8) \quad \text{div}Y'_C = 2(m + 1)\eta(C)\rho,$$

and so (4.7) moves to

$$(4.9) \quad \mathcal{L}_{Y'_C} \eta = \frac{2m + 1}{m + 1} \left( \frac{\text{div}Y'_C}{2m + 1} \right) \eta,$$

which shows that  $Y'_C$  is of conformal weight  $\frac{2m+1}{m+1}$  with respect to  $\eta$ [7].

We shall now investigate whether there exists another vector field  $Y''_C \neq Y'_C$  of  $\mathcal{D}_C$ , which defines an (i.c.t) of  $\eta$ . With the help of (2.19) and (2.25) one finds by a straightforward calculation that

$$(4.10) \quad Y''_C = 2\lambda C + \rho\Phi C + 2l\xi,$$

and operating by  $\nabla$ , one gets

$$(4.11) \quad \nabla Y''_C = \rho(\lambda + \eta(C))dp + \lambda\Phi C \wedge \xi + (\lambda\rho\eta + \rho\alpha - \lambda\beta) \otimes \xi + (2\lambda^2 + \rho^2 + 2l)\nabla\xi.$$

From (4.11) one derives by taking  $\text{tr} \nabla Y''_C$

$$(4.12) \quad \text{div}Y''_C = 2(m + 1)(\lambda + \eta(C))\rho,$$

and a standard calculation gives

$$\mathcal{L}_{Y''_C} \eta = \frac{2m + 1}{m + 1} \left( \frac{\text{div}Y''_C}{2m + 1} \right) \eta.$$



Hence as  $Y'_C$ , the vector field  $Y''_C$  is of conformal weight  $\frac{2m+1}{m+1}$  with respect to  $\eta$ .

We shall now focus our attention on the distinguished conformal vector fields of  $\mathcal{D}_C$ . We recall that we have shown (see (2.33)) that  $C$  defines an infinitesimal quasi-conformal contact transformation of  $\Omega$ . From (2.32), (3.6) and (3.12) one derives

$$(4.13) \quad \begin{cases} \mathcal{L}_{C_1}\Omega = -\eta \wedge \alpha - 2\rho\Omega, \\ \mathcal{L}_{C_2}\Omega = -\eta \wedge \beta + 2(\eta(C) - \lambda)\Omega, \\ \mathcal{L}_{C_3}\Omega = -\lambda\eta \wedge \alpha + 2\lambda\rho\Omega. \end{cases}$$

which proves that all these vector fields are quasi-biconformal vector fields with respect to the pairing  $(\eta, \Omega)$ . Moreover, by a short calculation one derives from (2.23) and (4.13) that

$$d(\mathcal{L}_C\Omega) = 0, \quad d(\mathcal{L}_{C_\nu}\Omega) = 0, \quad \nu = 1, 2, 3.$$

Hence, following a well-known definition (see [1]), we may say that the structure 2-form  $\Omega$  is a *relatively integral invariant* of the distinguished conformal vector fields of  $\mathcal{D}_C$ .

Finally since the Killing vector field  $K = [C_1, C_2]$  can be expressed as

$$(4.14) \quad K = \rho_2 C_1 + \rho C_2; \quad \rho_2 = \eta(C) - \lambda,$$

one quickly finds

$$i_K\Omega = \eta(C)d\eta(C) - \lambda dl \Rightarrow \mathcal{L}_K\Omega = 0.$$

According one may say that  $K$  defines an infinitesimal automorphism of the 2-form  $\Omega$ . We arrived at the final theorem of our considerations:

**THEOREM 4.1.** *Let  $M(\Phi, \eta, \xi, g)$  be a  $(2m + 1)$ -dimensional Sasakian manifold carrying a structure conformal vector field  $C$  and let  $\mathcal{D}_C$  be the autoparallel distribution defined by  $C$ . Let  $C_\nu \in \mathcal{D}_C$ ,  $\nu = 1, 2, 3$  be the associated conformal vector fields of  $C$  and let  $K \in \mathcal{D}_C$  be the associated Killing vector field of  $C$  by the Lie bracket. Then*

(i) *There are two vector fields  $Y'_C, Y''_C \in \mathcal{D}_C$  which define infinitesimal conformal transformation of  $\eta$ , and both are of conformal weight  $\frac{2m+1}{m+1}$  with respect to  $\eta$ .*

(ii)  $C$  and  $C_\nu$  define infinitesimal conformal contact transformations of the fundamental 2-form  $\frac{1}{2}d\eta$ , and  $\frac{1}{2}d\eta$  is a relatively integral invariant of  $C$  and  $C_\nu$ .

(iii) The Killing vector field  $K$  defines an infinitesimal automorphism of  $\frac{1}{2}d\eta$ .

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