

CRITICAL METRICS ON NEARLY KAEHLERIAN MANIFOLDS

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The study of the integral of the scalar curvature R , $J(g) := \int_M R dV_g$, as a function on the set of all Riemannian metrics with the same volume on a compact manifold, is now classical and it is well-known ([4]) that the critical points of $J(g)$ are the Einstein metrics; moreover, other functions of the curvature have been taken as integrands and studied in [1]-[6], etc. For example, M. Berger([1]) and Y. Mutō ([5]) investigated the critical point conditions of the functions

$$\int_M R^2 dV_g, \quad \int_M R_{ji} R^{ji} dV_g, \quad \int_M R_{kjih} R^{kjhi} dV_g,$$

where R_{ji} and R_{kjih} denote the Ricci and curvature tensors expressed in components, respectively. On the other hand D. E. Blair and S. Ianus ([3]) gave some results about the function $\int_M (R - R^*) dV_g$ on a symplectic manifold, where R^* is the $*$ -scalar curvature defined by $R^* := R_{kjih} J^{ji} J^{kh}$ for an almost complex structure J .

In this paper, we consider the function related with almost hermitian structure on a compact complex manifold. More precisely, on a $2n$ -dimensional complex manifold M admitting 2-form Ω of rank $2n$ everywhere, assume that M admits a metric g such that $g(JX, JY) = g(X, Y)$, that is, assume that g defines an hermitian structure on M admitting Ω as fundamental 2-form-the 'almost complex structure' J being determined by g and $\Omega : g(X, Y) = \Omega(X, JY)$. We consider the function $I(g) := \int_M \|N\|^2 dV_g$, where $\|N\|$ is the norm of Nijenhuis tensor N defined by (J, g) . Then, for this function restricted on the set of all nearly Kaehler metrics on M , we can show the following theorem;

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THEOREM. *Let M be a compact almost hermitian manifold and \mathcal{M} the set of nearly Kaehler metrics, Then a nearly Kaehler metric is a critical point of the function $I(g)$ on \mathcal{M} if and only if it is a Kaehler metric. Moreover, the critical metric is unique.*

1. Let M be a compact almost hermitian manifold with fundamental 2-form Ω and $\dim M = 2n$, i.e.,

$$g(X, JY) = \Omega(X, Y) \text{ and } g(X, Y) = g(JX, JY)$$

where J is an almost complex structure. We remark that g and J are constructed simultaneously by polarization of Ω evaluated on a local orthonormal basis of an arbitrary Riemannian metric. It is well-known that all hermitian metrics have the same volume element so that we may say $\int_M dV_g = 1$. The set \mathcal{A} of all hermitian metrics is infinite dimensional and is totally geodesic in the set of all Riemannian metrics on M . From now on we consider the function

$$I(g) := \int_M \|N\|^2 dV_g \text{ defined on } \mathcal{A}.$$

In the following we use local coordinates and a tensor is expressed in its components with respect to the natural frame. When we take a C^∞ curve $g(t)$ in \mathcal{A} , we get several tensor fields defined by

$$D_{ji} := \frac{\partial g_{ji}}{\partial t}, \quad D_i^h := D_{ik} g^{kh}, \quad D^{ih} := D_k^h g^{ki},$$

$$D_{ji}^h := \frac{1}{2}(\nabla_j D_i^h + \nabla_i D_j^h - \nabla^h D_{ji}), \quad D_{kji}^h := \nabla_k D_{ji}^h - \nabla_j D_{ki}^h,$$

where ∇ denotes the covariant differentiation with respect to the metric g and $\nabla^h = g^{hi} \nabla_i$. If $g(t)$ is a path in \mathcal{A} with $g(0) = g$, then we have

$$(1.1) \quad \frac{\partial J_{ji}}{\partial t} = 0, \quad \frac{\partial J_j^h}{\partial t} = -J_{ji} D^{ih}, \quad D_{ba} J_i^a D_j^b = -D_{ji}, \quad g^{ji} D_{ji} = 0.$$

In general, any symmetric tensor field D satisfying $DJ + JD = 0$ is tangent to some curve in \mathcal{A} (cf. [2]).

Using $\frac{\partial\{j^hi\}}{\partial t} = D_{ji}{}^h$ and (1.1), we have

$$(1.2) \quad \frac{\partial \nabla_j J_i^h}{\partial t} = -(\nabla_j J_{it})D^{th} - J_{it}(\nabla_j D^{th}) - D_{ji}{}^t J_t^h + D_{jt}{}^h J_i^t.$$

Hence we have

$$(1.3) \quad \begin{aligned} \frac{dI(g)}{dt} &= \int_M \left[\frac{\partial(N_{jk}{}^i N_{bc}{}^a g^{bj} g^{ck} g_{ai})}{\partial t} + \frac{1}{2} \|N_{kj}{}^i\|^2 g^{ba} D_{ba} \right] dV_g \\ &= \int_M N_{bc}{}^a \left\{ 2 \frac{\partial N_{jk}{}^i}{\partial t} g_{ai} g^{bj} g^{ck} - 2 N_{jk}{}^i D^{ck} g^{bj} g_{ai} \right. \\ &\quad \left. + N_{jk}{}^i D_{ai} g^{bj} g^{ck} \right\} dV_g \\ &= \int_M N_{bc}{}^a \{ -2 D_{ai} J^{bt} \nabla_t J^{ci} + 2 J_{th} D_a{}^h \nabla^b J^{ct} - 4 J^{bt} J^{ci} \nabla_t D_{ai} \\ &\quad - 4 J_s{}^b D^{st} \nabla_t J_a{}^c + 4 J_{at} D^{st} \nabla^b J_s{}^c + 4 J_{at} J_s{}^c \nabla^b D^{st} \\ &\quad - 2 D^{ck} J^{bt} \nabla_t J_{ka} - 2 D^{ck} J_k{}^t \nabla_t J_a{}^b - 2 D^{ck} J_{at} \nabla^b J_k{}^t \\ &\quad + 2 D^{ck} J_{at} \nabla_k J^{bt} \} dV_g \\ &= \int_M [8 D_{ai} J_s{}^a J^{bt} (\nabla_b J_c{}^s) (\nabla_t J^{ci}) + 8 J_b{}^a (\nabla_t D_{ai}) (\nabla^i J^{tb}) \\ &\quad - 8 J^{ic} (\nabla_t D_{ai}) (\nabla_c J^{ta}) + 4 D^{ht} (\nabla_h J_c{}^a) (\nabla_t J_a{}^c) \\ &\quad - 8 D^{ht} (\nabla_t J_c{}^s) (\nabla_s J_h{}^c) + 4 D^{ht} J_s{}^a J_h{}^b (\nabla_t J_a{}^c) (\nabla_b J_c{}^s) \\ &\quad + 8 D^{ht} (\nabla^b J_h{}^c) (\nabla_b J_{ct})] dV_g. \end{aligned}$$

2. In this section we will consider only nearly Kaehler structure. By definition a nearly Kaehler structure is an almost hermitian structure (J, g) which satisfies

$$(\nabla_X J)Y + (\nabla_Y J)X = 0$$

for any vector fields X and Y on M .

If (J, g) is a nearly Kaehler structure, then we have

$$J_b^a(\nabla_t D_{ai})(\nabla^i J^{tb}) - J^{ic}(\nabla_t D_{ai})(\nabla_c J^{ta}) = 0,$$

and consequently from (1.3)

$$\begin{aligned} \frac{dI(g)}{dt} &= \int_M D_{ht} [8J_s^h J^{ba}(\nabla_b J_c^s)(\nabla_a J^{ct}) + 4(\nabla^h J_c^a)(\nabla^t J_a^c) \\ &\quad - 8(\nabla^t J^{sc})(\nabla_s J_c^h) + 4J_s^a J^{hb}(\nabla^t J_a^c)(\nabla_b J_c^s) \\ &\quad + 8(\nabla^b J^{hc})(\nabla_b J_c^t)] dV_g \\ (2.1) \quad &= -32 \int_M D_{ht}(\nabla^s J^{ct})(\nabla_s J_c^h) dV_g. \end{aligned}$$

On the other hand D_{ji} must satisfy the only condition

$$(2.2) \quad \int_M D_{ji} g^{ji} dV_g = 0.$$

Let \bar{g} be a nearly Kaehler metric such that, for every C^∞ curve $g(t)$ of M satisfying $g(0) = \bar{g}$, $I(g(t))$ has vanishing derivative at $t = 0$. Then from (2.1) and (2.2) it follows that

$$(2.3) \quad -32(\nabla^s J^{ct})(\nabla_s J_c^h) = cg^{ht}, \quad c = \text{const.}$$

for this metric \bar{g} and consequently

$$2nc = -32(\nabla^s J^{ct})(\nabla_s J_{ct}).$$

By the way, transvecting $J_h^e J_{et}$ to (2.3) and using the fact that (J, \bar{g}) is nearly Kaehlerian, we can obtain

$$c = -32(\nabla^s J^{ct})(\nabla_s J_{ct})$$

and so $c = 0$, that is, $\nabla_s J_{ct} = 0$ for this metric \bar{g} . Hence we can assert that the formal part of main theorem stated in introduction is true.

Next, we differentiate again (2.1) with respect to the parameter t . Taking account of (J, g) being nearly Kaehlerian, we can easily find

$$\begin{aligned} \frac{d^2 I(g)}{dt^2} &= -32 \int_M \left[\frac{\partial D_{hi}}{\partial t} (\nabla^s J^{ci}) (\nabla_s J_c^h) + D_{ht} \{ -D^{si} (\nabla_i J^{ct}) (\nabla_s J_c^h) \right. \\ &\quad + D_i^{sc} J^{it} (\nabla_s J_c^h) + D_i^{st} J^{ci} (\nabla_s J_c^h) \\ &\quad + (\nabla^s J^{ct}) (\nabla_s D_{cb}) J^{bh} + (\nabla^s J^{ct}) (\nabla_s J^{bh}) D_{cb} \\ &\quad \left. - D_{sc}{}^b J_b{}^h (\nabla^s J^{ct}) + D_{sb}{}^h J_c{}^b (\nabla^s J^{ct}) \right] dV_g \\ &= 32 \int_M D_{ht} D_{kc} (\nabla_s J^{hk}) (\nabla^s J^{tc}) dV_g. \end{aligned}$$

If we put $D_{ji} = fg_{ji}$, where f is a C^∞ function on M such that $\int_M f dV_g = 0$, then $\frac{d^2 I(g)}{dt^2} > 0$. This means that the last main theorem is also true.

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