

異方性 Plasma 내에서 運動中인
Source에 의한 電磁界
(Electromagnetic Fields Due to Moving
Sources in Anisotropic Plasma)

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要 約

異方性 plasma 내를 等速運動中인 Source로 인한 電磁界理論을 一般的인 경우 즉 定磁界와 速度가 任意의 方向에 있을 경우에 대하여 論하였다. Minkowski의 關係式을 異方性分散媒体에 대해서 一般化했으며 媒体 parameter들의 相對論的變換公式를 誘導하여 速度比 β 의 多項式으로 展開함으로써 利用度를 높였다. 電磁界가 滿足하는 Helmholtz의 波動方程式을 tensor parameter로 特徵지워지는 媒体에 대해 一般化했으며 이를 行列形式으로 整理하였다. 波動方程式의 解는 波動演算子行列의 逆行列과 source函數 vector의 相乘積으로 表示되며 결국 電磁界는 波動演算子行列의 逆行列에 依해서 表示된다. 波動演算子行列의 逆行列을 求하고 그 結果를 具體的으로 表示하였다. 本 論文에서 誘導한 公式들은 一般的式들로서 特殊한 條件下에서 이들은 既知의 結果와 一致함을 例를 통해서 表示하였다.

Abstract

Fundamentals of electrodynamics of moving sources with constant velocity in an anisotropic plasma when the do magnetic field and the relative motion are oriented in arbitrary directions are presented. The well-known Minkowski's relations are generalized to accomodate anisotropic and dispersive media, and relativistic transformation formulae of constitutive parameters are derived and expanded into polynomials of the speed ratio β to increase the utility of the formulae. The helmholtz wave equation of electromagnetic fields is generalized to the media characterized by tensor parameters, and is solved in operator form. Also the solution of wave equation is expressed as a product of the inverse of the wave operator matrix and the source function vector, and the inverse of the wave operator matrix is presented in an explicit form. The equations and formulae derived in this paper are all general, and can be reduced to known and proven results upon imposing the restriction called for by specific situations.

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I. 序 論

相對論的電磁場的理論과 應用에 관한 關心은 主로 物理學의 本質的으로는 微視的分野에 限定되는 느낌이었었다. 近者에 이르러 人工衛星과 같은 더욱 빠르고 큰 人工飛行體의 出現과 相對運動中에 있는 無線局間의 信賴할 만한 通信路建設에 必須的인 高速 plasma wind下의 電磁波輻射特性에 대한 보다 많은 智識의 必要性때문에 本質的으로 巨視的인 이 問題가 새로운 關心을 惹起시키고 있다.¹¹⁻¹²

非等方性 plasma內에서의 相對運動中의 source에 의한 電磁場의 解析은 二重으로 複雜한바 그 理由는 場과 媒體가 tensor parameter를 거쳐서 相互作用하기 때문이며 또 한 理由는 電氣 vector와 磁氣 vector間에는 從來의 Faraday induction外에 새로운 型의 相互結合이 發生하기 때문이다. 이 複雜性때문에 大部分의 研究報告는 相互速度의 크기나 定磁界의 方向에 制約을 加한 比較的 簡單한 경우에 限定된 것들이다.

本 論文에서는 magneto-plasma內를 定速度로 運動中인 source로부터의 電磁界 問題를 可能한 限制約없이 가장 一般的인 형태로 解析하고자 努力하였다. 먼저 電磁場論과 Lorentz變換을 이용하여 媒體 parameter간의 Minkowsky의 關係式을 異方性媒體의 경우까지 包含되도록 一般化했으며 이를 土土로 媒體 parameter의 變換公式를 對稱型으로 誘導하였다. 다음에 電磁界가 滿足해야할 Helmholtz 波動方程式을 一般化하고 그 結果를 行列形式으로 整理하였다. 波動方程式의 解는 波動演算子行列과 source函數vector간의 相乘積으로 表示된다. Source 分布로부터 source函數vector는 곧 알 수 있으므로 電磁場을 求하는 問題는 결국 波動演算子行列의 逆行列을 求하는 問題로 歸着된다. V章에서는 이 逆行列을 구한 후 이를 利用하기 편리한 形으로 整理하였다.

II. 異方性媒體內的 電磁場論과 Lorentz 變換

電氣的 및 磁氣的 source $\underline{e}(\underline{r}, t)$ 및 $\underline{h}(\underline{r}, t)$ 가 異方性媒體內에서 等速度 \underline{v} 로 運動中일때의 電磁界를 생각해 보자. 媒體의 電磁氣的 特性은 tensor parameter $\underline{\epsilon}$ 및 $\underline{\mu}$ 로 표시된다. Source에 대한 靜止系를 S, 媒體에 대한 靜止系를 S'라 부르기로 하며 또 어떤 量을 표시하는데 있어서 S系에 대한 量은 그대로, S'系에 대한 것은 그 量에 記號를 添加 表示키로 한다.

S'系에 있어서 電磁界 $\underline{E}'(\underline{r}, t)$, $\underline{D}'(\underline{r}, t)$, $\underline{H}'(\underline{r}, t)$ 및 $\underline{B}'(\underline{r}, t)$ 가 만족하는 Maxwell方程式은

$$\nabla' \times \underline{E}' = -\frac{\partial \underline{B}'}{\partial t'} - \underline{h}'; \quad \nabla' \times \underline{H}' = \frac{\partial \underline{D}'}{\partial t'} + \underline{j}' \quad (1)$$

지금 Fourier 變換의 定義:

$$\underline{A}(\underline{k}, \omega) = \int_{-\infty}^{\infty} \underline{a}(\underline{r}, t) e^{-j\omega t - \underline{k} \cdot \underline{r}} d\underline{r} dt \quad (2)$$

$$\underline{a}(\underline{r}, t) = \frac{1}{(2\pi)^4} \int_{-\infty}^{\infty} \underline{A}(\underline{k}, \omega) e^{j\omega t - \underline{k} \cdot \underline{r}} d\underline{k} d\omega \quad (3)$$

에 따라 式(1)에 Fourier 變換을 가하면

$$-j\underline{k}' \times \underline{E}' = -j\omega' \underline{B}' - \underline{M}'; \quad -j\underline{k}' \times \underline{H}' = j\omega' \underline{D}' + \underline{J}' \quad (4)$$

여기서 $\underline{E}(\underline{k}, \omega)$, $\underline{D}(\underline{k}, \omega)$, $\underline{H}(\underline{k}, \omega)$ 및 $\underline{B}(\underline{k}, \omega)$ 는 각각 $\underline{E}(\underline{r}, t)$, $\underline{D}(\underline{r}, t)$, $\underline{H}(\underline{r}, t)$ 및 $\underline{B}(\underline{r}, t)$ 의 Fourier 變換을 표시한다. 이들 媒體 parameter간의 關係式은

$$\underline{D}' = \underline{\epsilon}_0 \underline{\epsilon}' \underline{E}' \quad (5)$$

$$\underline{B}' = \underline{\mu}_0 \underline{\mu}' \underline{H}' \quad (6)$$

여기서 ϵ_0 및 μ_0 는 自由空間 parameter들을 의미한다.

Maxwell方程式은 Lorentz 變換에 대해서 covariant property를 유지하므로 S系에 대한 Maxwell方程式은 S'系에 대한 것과 같은 形式 즉

$$\nabla \times \underline{E} = -\frac{\partial \underline{B}}{\partial t} - \underline{h}, \quad \nabla \times \underline{H} = \frac{\partial \underline{D}}{\partial t} + \underline{j} \quad (7)$$

이들의 Fourier 變換式은

$$-j\underline{k} \times \underline{E} = -j\omega \underline{B} - \underline{M}, \quad -j\underline{k} \times \underline{H} = j\omega \underline{D} + \underline{J} \quad (8)$$

그러나 S系로 變換했을 때의 式(5), (6)은 전혀 다른 形式을 취하게 된다.

S系와 S'系는 Lorentz變換에 의해서 密接히 結合되어 있으며 各量에 대한 兩系간의 變換公式만 確立된다면

어떤 現象을 한系에 대해서 풀었을때 다른 系에서 그 現象이 어떻게 나타나는지 豫告할 수 있다.

任意的 電磁界 vector 예컨대 $\underline{E}(\underline{r}, t)$ 와 그의 複素振幅vector 예컨대 $\underline{E}(\underline{k}, \omega)$ 에는 같은 形式의 Lorentz 變換이 적용된다.

電磁界 vector들의 Lorentz變換公式은 잘 알려진 바와 같이 ¹¹⁾

$$\underline{E}' = \gamma \underline{E} + (1 - \gamma) \frac{(\underline{E} \cdot \underline{V}) \underline{V}}{V^2} + \gamma \underline{V} \times \underline{B} \quad (9)$$

$$\underline{B}' = \gamma \underline{B} + (1 - \gamma) \frac{(\underline{B} \cdot \underline{V}) \underline{V}}{V^2} - \gamma \frac{\underline{V} \times \underline{E}}{C^2} \quad (10)$$

$$\underline{D}' = \gamma \underline{D} + (1 - \gamma) \frac{(\underline{D} \cdot \underline{V}) \underline{V}}{V^2} + \gamma \frac{\underline{V} \times \underline{H}}{C^2} \quad (11)$$

$$\underline{H}' = \gamma \underline{H} + (1 - \gamma) \frac{(\underline{H} \cdot \underline{V}) \underline{V}}{V^2} - \gamma \underline{V} \times \underline{D} \quad (12)$$

여기서

$$\gamma = (1 - \beta^2)^{-\frac{1}{2}} \quad (13)$$

$$\beta = \frac{V}{C} \quad (14)$$

$$C = 1/\sqrt{\epsilon_0 \mu_0} = \text{自由空間에서의 光速度} \quad (15)$$

한편 媒体 parameter에 대해서는 等方性媒体의 경우는 tensor parameter $\underline{\epsilon}$ 및 $\underline{\mu}$ 는 scalar量 ϵ' 및 μ' 로 縮退하며 S系에서의 該當量 ϵ, μ 는 Minkowski의 關係式^{15, 16)}

$$\underline{D} + \frac{1}{C^2} \underline{V} \times \underline{H} = \epsilon_0 \epsilon' (\underline{E} + \underline{V} \times \underline{B}) \quad (16)$$

$$\underline{B} - \frac{1}{C^2} \underline{V} \times \underline{E} = \mu_0 \mu' (\underline{H} - \underline{V} \times \underline{D}) \quad (17)$$

로부터 유도해낼 수 있다. 그러나 이 Minkowski의 關係式은 媒体parameter가 tensor가 되는 異方性媒体에 대해서는 그대로 적용할 수는 없다.

일반적으로 行列表記法에 의하면 2 vector간의 scalar 및 vector 積은 각각

$$\underline{A} \cdot \underline{B} = (A_1, A_2, A_3) \begin{vmatrix} B_1 \\ B_2 \\ B_3 \end{vmatrix} = \underline{A}^T \underline{B} \quad (18)$$

$$\underline{A} \times \underline{B} = \begin{vmatrix} 0 & -A_3 & A_2 \\ A_3 & 0 & -A_1 \\ -A_2 & A_1 & 0 \end{vmatrix} \begin{vmatrix} B_1 \\ B_2 \\ B_3 \end{vmatrix} = \underline{A} \underline{B} \quad (19)$$

여기서

$$\underline{A}^T = \text{行列 } \underline{A} \text{의 轉置行列} \quad (20)$$

$$\underline{A} \underline{B} = \begin{vmatrix} 0 & -A_3 & A_2 \\ A_3 & 0 & -A_1 \\ -A_2 & A_1 & 0 \end{vmatrix} = \text{Vector積 } \underline{A} \times \underline{B} \text{에서의 premultiplier } \underline{A} \text{의 等價行列} \quad (21)$$

이므로 式(9)~(12)를 行列表記할 수 있다. 예컨대 式(9)는

$$\begin{aligned} \underline{E}' &= \gamma \underline{E} + (1 - \gamma) \frac{\underline{V} (\underline{V}^T \underline{E})}{V^2} + \gamma \underline{V} \underline{B} \\ &= \gamma \underline{1} + \frac{1 - \gamma}{\gamma} \frac{\underline{\beta} \underline{\beta}^T}{\beta^2} \underline{1} \underline{E} + C \gamma \underline{V} \underline{B} \\ &= \gamma \underline{U} \underline{E} + C \gamma \underline{V} \underline{B} \end{aligned} \quad (22)$$

여기서

$$\underline{\beta} = \underline{V}/C, \quad \underline{V} = \underline{V}/C \quad (23)$$

$$\underline{U} = \underline{1} + \frac{1 - \gamma}{\gamma} \frac{\underline{\beta} \underline{\beta}^T}{\beta^2} \quad (24)$$

같은 방법으로 式(10), (11) 및 (12)는 각각

$$\underline{B}' = \gamma \underline{U} \underline{B} - \frac{1}{C} \gamma \underline{\beta} \underline{E} \tag{25}$$

$$\underline{D}' = \gamma \underline{U} \underline{D} + \frac{1}{C} \gamma \underline{\beta} \underline{H} \tag{26}$$

$$\underline{H}' = \gamma \underline{U} \underline{H} - C \gamma \underline{\beta} \underline{D} \tag{27}$$

이들을 式 (5), (6)에 代入하면

$$\underline{U} \underline{D} + \frac{1}{C} \underline{\beta} \underline{H} = \epsilon_0 \underline{\epsilon}' (\underline{U} \underline{E} + C \underline{\beta} \underline{B}) \tag{28}$$

$$\underline{U} \underline{B} - \frac{1}{C} \underline{\beta} \underline{E} = \mu_0 \underline{\mu}' (\underline{U} \underline{H} - C \underline{\beta} \underline{D}) \tag{29}$$

式 (28), (29)가 異方性媒体에 대해서도 成立되는 一般化된 Minkowski의 關係式이다.

III. 媒体 Parameter의 變換公式

S'系에서는 媒体 parameter $\underline{\epsilon}'$ 및 $\underline{\mu}'$ 는 각각 電氣 vector \underline{D}' 와 다른 電氣 vector \underline{E}' 와를, 또는 磁氣 vector \underline{B}' 와 다른 磁氣 vector \underline{H}' 와를 關係지을뿐 parameter를 통한 電氣 및 磁氣 vector간의 相互結合은 存在하지 않는다. 그러나 運動中에 있는 source에 대한 靜止系 S에서는 사정이 다르게 나타난다.

1. 變換公式

式 (28), (29)에서 \underline{B} 또는 \underline{D} 를 消去하면

$$\underline{D} = (\underline{U} + \underline{\epsilon}' \underline{\beta} \underline{U}^{-1} \underline{\mu}' \underline{\beta})^{-1} \epsilon_0 \underline{\epsilon}' (\underline{U} + \underline{\beta} \underline{U}^{-1} \underline{\beta}) \underline{E} + \frac{1}{C} (\underline{\epsilon}' \underline{\beta} \underline{U}^{-1} \underline{\mu}' \underline{U} - \underline{\beta}) \underline{H} \tag{30}$$

$$\underline{B} = (\underline{U} + \underline{\mu}' \underline{\beta} \underline{U}^{-1} \underline{\epsilon}' \underline{\beta})^{-1} \frac{1}{C} (\underline{\beta} - \underline{\mu}' \underline{\beta} \underline{U}^{-1} \underline{\epsilon}' \underline{U}) \underline{E} + \mu_0 \underline{\mu}' (\underline{U} + \underline{\beta} \underline{U}^{-1} \underline{\beta}) \underline{H} \tag{31}$$

지금

$$\underline{D} = \epsilon_0 \underline{\epsilon} \underline{E} + \underline{\kappa} \underline{H} \tag{32}$$

$$\underline{B} = \underline{y} \underline{E} + \mu_0 \underline{\mu} \underline{H} \tag{33}$$

라 놓으면 式 (32), (33)과 式 (30), (31)을 比較하므로써

$$\underline{\epsilon} = (\underline{U} + \underline{\epsilon}' \underline{\beta} \underline{U}^{-1} \underline{\mu}' \underline{\beta})^{-1} \underline{\epsilon}' (\underline{U} + \underline{\beta} \underline{U}^{-1} \underline{\beta}) \tag{34}$$

$$\underline{\kappa} = \frac{1}{C} (\underline{U} + \underline{\epsilon}' \underline{\beta} \underline{U}^{-1} \underline{\mu}' \underline{\beta})^{-1} (\underline{\epsilon}' \underline{\beta} \underline{U}^{-1} \underline{\mu}' \underline{U} - \underline{\beta}) \tag{35}$$

$$\underline{y} = \frac{1}{C} (\underline{U} + \underline{\mu}' \underline{\beta} \underline{U}^{-1} \underline{\epsilon}' \underline{\beta})^{-1} (\underline{\beta} - \underline{\mu}' \underline{\beta} \underline{U}^{-1} \underline{\epsilon}' \underline{U}) \tag{36}$$

$$\underline{\mu} = (\underline{U} + \underline{\mu}' \underline{\beta} \underline{U}^{-1} \underline{\epsilon}' \underline{\beta})^{-1} \underline{\mu}' (\underline{U} + \underline{\beta} \underline{U}^{-1} \underline{\beta}) \tag{37}$$

式 (34)~(37)은 媒体 parameter에 대한 靜止系와 運動系간의 一般變換公式이며 Lee와 Lo^{14, 15}의 非對稱公式과는 달리 對稱形인 점에 주의하라.

式 (24)의 行列 \underline{U} 를 다시 쓰면

$$\underline{U} = \underline{1} - (1 - \sqrt{1 - \beta^2}) \frac{\underline{\beta} \underline{\beta}^T}{\beta^2} = \underline{1} - B_2 \underline{m} \underline{m}^T \tag{38}$$

여기서

$$\underline{\beta} = \beta \underline{m} \tag{39}$$

$$\underline{m} = (m_1, m_2, m_3)^T = \underline{\beta} \text{의 方向餘弦} \tag{40}$$

$$B_2 = 1 - \sqrt{1 - \beta^2} = \frac{1}{2} \beta^2 (1 + \frac{1}{4} \beta^2 + \frac{1}{8} \beta^4 + \dots) \tag{41}$$

따라서

$$\underline{U}^{-1} = [\underline{1} - B_2 \underline{m} \underline{m}^T]^{-1} = \frac{1}{1 - B_2} \begin{bmatrix} 1 - B_2 (m_1^2 + m_2^2) & B_2 m_1 m_2 & B_2 m_1 m_3 \\ B_2 m_1 m_2 & 1 - B_2 (m_1^2 + m_2^2) & B_2 m_2 m_3 \\ B_2 m_1 m_3 & B_2 m_2 m_3 & 1 - B_2 (m_1^2 + m_2^2) \end{bmatrix}$$

$$\begin{aligned}
 &= \frac{1}{1-B_s} \{ (1-B_s) \underline{\underline{1}} + B_s \underline{\underline{m}} \underline{\underline{m}}^T \} \\
 &= \underline{\underline{1}} + \frac{B_s}{1-B_s} \underline{\underline{m}} \underline{\underline{m}}^T
 \end{aligned} \tag{42}$$

式(41)의 B_s 는 β^2 정도의 크기이므로 대부분의 경우 式(38) 및 (42)의 第2項은 第1項에 비해서 무시해도 무방하다. 따라서 式(34)~(37)은

$$\underline{\underline{\epsilon}} = (\underline{\underline{1}} + \underline{\underline{\epsilon}}' \underline{\underline{\beta}} \underline{\underline{\mu}}' \underline{\underline{\beta}})^{-1} \underline{\underline{\epsilon}}' (\underline{\underline{1}} + \underline{\underline{\beta}} \underline{\underline{\beta}}) \tag{43}$$

$$\underline{\underline{\kappa}} = \frac{1}{C} (\underline{\underline{1}} + \underline{\underline{\epsilon}}' \underline{\underline{\beta}} \underline{\underline{\mu}}' \underline{\underline{\beta}})^{-1} (\underline{\underline{\epsilon}}' \underline{\underline{\beta}} \underline{\underline{\mu}}' - \underline{\underline{\beta}}) \tag{44}$$

$$\eta = \frac{1}{C} (\underline{\underline{1}} + \underline{\underline{\mu}}' \underline{\underline{\beta}} \underline{\underline{\epsilon}}' \underline{\underline{\beta}})^{-1} (\underline{\underline{\beta}} - \underline{\underline{\mu}}' \underline{\underline{\beta}} \underline{\underline{\epsilon}}') \tag{45}$$

$$\underline{\underline{\mu}} = (\underline{\underline{1}} + \underline{\underline{\mu}}' \underline{\underline{\beta}} \underline{\underline{\epsilon}}' \underline{\underline{\beta}})^{-1} \underline{\underline{\mu}}' (\underline{\underline{1}} + \underline{\underline{\beta}} \underline{\underline{\beta}}) \tag{46}$$

Magneto-plasma 또는 gyromagnetic plasma에 있어서는 $\underline{\underline{\mu}}' = \underline{\underline{1}}$ 이므로 式(43)~(46)은 더욱 간단해져서

$$\underline{\underline{\epsilon}} = (\underline{\underline{1}} + \underline{\underline{\epsilon}}' \underline{\underline{\beta}} \underline{\underline{\beta}})^{-1} \underline{\underline{\epsilon}}' (\underline{\underline{1}} + \underline{\underline{\beta}} \underline{\underline{\beta}}) \tag{47}$$

$$\underline{\underline{\kappa}} = \frac{1}{C} (\underline{\underline{1}} + \underline{\underline{\epsilon}}' \underline{\underline{\beta}} \underline{\underline{\beta}})^{-1} (\underline{\underline{\epsilon}}' - \underline{\underline{1}}) \underline{\underline{\beta}} \tag{48}$$

$$\underline{\underline{\eta}} = \frac{1}{C} (\underline{\underline{1}} + \underline{\underline{\beta}} \underline{\underline{\epsilon}}' \underline{\underline{\beta}})^{-1} \underline{\underline{\beta}} (\underline{\underline{1}} - \underline{\underline{\epsilon}}') \tag{49}$$

$$\underline{\underline{\mu}} = (\underline{\underline{1}} + \underline{\underline{\beta}} \underline{\underline{\epsilon}}' \underline{\underline{\beta}})^{-1} (\underline{\underline{1}} + \underline{\underline{\beta}} \underline{\underline{\beta}}) \tag{50}$$

2. 變換公式의 多項式展開

β 의 값은 1보다 작은 것이 보통이므로 式(47)~(50)을 β 의 多項式으로 展開할 수 있다면 편리할 것이다. 式(47)~(50)을 展開하기 전에 式들에 나타난 逆行行列들의 要素를 계산하였으며 그 결과는 다음과 같다.

지금

$$({}_1\Delta)^{-1} \equiv (\underline{\underline{1}} + \underline{\underline{\epsilon}}' \underline{\underline{\beta}} \underline{\underline{\beta}})^{-1} = \frac{\text{adj}({}_1\Delta)}{\det({}_1\Delta)} = \frac{[\text{cof}({}_1\Delta_{ij})]}{\det({}_1\Delta)} \tag{51}$$

라 놓면 다음과 같이 된다.

$$\text{cof}({}_1\Delta_{ii}) = \beta^4 m_i m^T (E_{ii}, E_{ii}, E_{ii})^T + \beta^2 \{ m_j^2 \epsilon_{jj} + m_k^2 \epsilon_{kk} + m_i m_j \epsilon_{ji} + m_i m_k \epsilon_{ki} + m_j m_k (\epsilon_{jk} + \epsilon_{kj}) - (\epsilon_{jj} + \epsilon_{kk}) \} + 1, \tag{52}$$

$$\text{cof}({}_1\Delta_{ij}) \ (i \neq j) = \beta^4 m_i m^T (E_{ij}, E_{ij}, E_{ij})^T - \beta^2 \{ m_i^2 \epsilon_{ji}' + m_i m_j \epsilon_{jj}' + m_i m_k \epsilon_{jk}' \} - \epsilon_{ji}' \{, \tag{53}$$

$$\det({}_1\Delta) = m^T \left\{ \beta^4 \begin{bmatrix} E_{11} & E_{12} & E_{13} \\ E_{21} & E_{22} & E_{23} \\ E_{31} & E_{32} & E_{33} \end{bmatrix} + \beta^2 \begin{bmatrix} -(\epsilon_{22}' + \epsilon_{33}') & \epsilon_{12}' & \epsilon_{13}' \\ \epsilon_{21}' & -(\epsilon_{33}' + \epsilon_{11}') & \epsilon_{23}' \\ \epsilon_{31}' & \epsilon_{32}' & -(\epsilon_{11}' + \epsilon_{22}') \end{bmatrix} + \underline{\underline{1}} \right\} m \tag{54}$$

여기서 $E_{ij} = \det(\underline{\underline{\epsilon}})$ 의 要素 ϵ_{ij} 에 관한 餘因자

式(52), (53)으로부터

$$\begin{aligned}
 \text{adj}({}_1\Delta) &= \beta^4 m m^T \begin{bmatrix} E_{11} & E_{12} & E_{13} \\ E_{12} & E_{22} & E_{23} \\ E_{13} & E_{23} & E_{33} \end{bmatrix} + \beta^2 \left\{ \begin{bmatrix} -(\epsilon_{22}' + \epsilon_{33}') & \epsilon_{12}' & \epsilon_{13}' \\ \epsilon_{21}' & -(\epsilon_{33}' + \epsilon_{11}') & \epsilon_{23}' \\ \epsilon_{31}' & \epsilon_{32}' & -(\epsilon_{11}' + \epsilon_{22}') \end{bmatrix} \right\} \\
 &= \begin{bmatrix} -m^T \begin{bmatrix} 0 & 0 & 0 \\ \epsilon_{31}' & \epsilon_{32}' & \epsilon_{33}' \\ \epsilon_{21}' & \epsilon_{22}' & \epsilon_{23}' \end{bmatrix} m & m^T \begin{bmatrix} 0 & 0 & 0 \\ \epsilon_{11}' & \epsilon_{12}' & \epsilon_{13}' \\ 0 & 0 & 0 \end{bmatrix} m & m^T \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \epsilon_{11}' & \epsilon_{12}' & \epsilon_{13}' \end{bmatrix} m \\
 -m^T \begin{bmatrix} \epsilon_{31}' & \epsilon_{32}' & \epsilon_{33}' \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} m & -m^T \begin{bmatrix} \epsilon_{11}' & \epsilon_{12}' & \epsilon_{13}' \\ 0 & 0 & 0 \\ \epsilon_{21}' & \epsilon_{22}' & \epsilon_{23}' \end{bmatrix} m & m^T \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \epsilon_{21}' & \epsilon_{22}' & \epsilon_{23}' \end{bmatrix} m \\
 m^T \begin{bmatrix} \epsilon_{31}' & \epsilon_{32}' & \epsilon_{33}' \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} m & m^T \begin{bmatrix} 0 & 0 & 0 \\ \epsilon_{31}' & \epsilon_{32}' & \epsilon_{33}' \\ 0 & 0 & 0 \end{bmatrix} m & -m^T \begin{bmatrix} \epsilon_{11}' & \epsilon_{12}' & \epsilon_{13}' \\ \epsilon_{21}' & \epsilon_{22}' & \epsilon_{23}' \\ 0 & 0 & 0 \end{bmatrix} m
 \end{bmatrix} \tag{56}
 \end{aligned}$$

또

$$({}_2\Delta)^{-1} \equiv (\underline{\underline{1}} + \underline{\underline{\beta}} \underline{\underline{\epsilon}}' \underline{\underline{\beta}})^{-1} = \frac{\text{adj}({}_2\Delta)}{\det({}_2\Delta)} = \frac{[\text{cof}({}_2\Delta_{ij})]}{\det({}_2\Delta)} \tag{57}$$

라 놓면 여기서

$$\text{cof}({}_i\Delta_{ii}) = \beta^i m_i^2 \underline{m}^T \begin{bmatrix} E_{11} & E_{12} & E_{13} \\ E_{21} & E_{22} & E_{23} \\ E_{31} & E_{32} & E_{33} \end{bmatrix} \underline{m} - \beta^2 |m_i^2 (\epsilon'_{jj} + \epsilon'_{kk}) + (m_j^2 + m_k^2) \epsilon'_{ii} - m_i m_j (\epsilon'_{ij} + \epsilon'_{ji}) - m_i m_k (\epsilon'_{ik} + \epsilon'_{ki})| + 1 \tag{58}$$

$$\text{cof}({}_i\Delta_{ij}) (i \neq j) = \beta^i m_i m_j \underline{m}^T \begin{bmatrix} E_{11} & E_{12} & E_{13} \\ E_{21} & E_{22} & E_{23} \\ E_{31} & E_{32} & E_{33} \end{bmatrix} \underline{m} - \beta^2 |m_k^2 \epsilon'_{ij} + m_i m_j \epsilon'_{kk} - m_i m_k \epsilon'_{kj} - m_j m_k \epsilon'_{ik}| \tag{59}$$

$$\det({}_i\Delta) = \underline{m}^T \left\{ \beta^i \begin{bmatrix} E_{11} & E_{12} & E_{13} \\ E_{21} & E_{22} & E_{23} \\ E_{31} & E_{32} & E_{33} \end{bmatrix} + \beta^2 \begin{bmatrix} -(\epsilon'_{22} + \epsilon'_{33}) & \epsilon'_{12} & \epsilon'_{13} \\ \epsilon'_{21} & -(\epsilon'_{33} + \epsilon'_{11}) & \epsilon'_{23} \\ \epsilon'_{31} & \epsilon'_{32} & -(\epsilon'_{11} + \epsilon'_{22}) \end{bmatrix} + \underline{1} \right\} \underline{m} \tag{60}$$

式 (58), (59)로부터

$$\text{adj}({}_i\Delta) = \beta^i \underline{m} \underline{m}^T \left\{ \underline{m}^T \begin{bmatrix} E_{11} & E_{12} & E_{13} \\ E_{21} & E_{22} & E_{23} \\ E_{31} & E_{32} & E_{33} \end{bmatrix} \underline{m} \right. \\ \left. + \beta^2 \begin{bmatrix} \underline{m}^T \begin{bmatrix} -(\epsilon'_{22} + \epsilon'_{33}) & \epsilon'_{12} & \epsilon'_{13} \\ \epsilon'_{21} & -(\epsilon'_{33} + \epsilon'_{11}) & \epsilon'_{23} \\ \epsilon'_{31} & \epsilon'_{32} & -(\epsilon'_{11} + \epsilon'_{22}) \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} 0 & 0 & 0 \\ -\epsilon'_{33} & 0 & \epsilon'_{31} \\ \epsilon'_{23} & 0 & -\epsilon'_{21} \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} 0 & 0 & 0 \\ \epsilon'_{32} & -\epsilon'_{31} & 0 \\ -\epsilon'_{22} & \epsilon'_{21} & 0 \end{bmatrix} \underline{m} \\ \underline{m}^T \begin{bmatrix} 0 & -\epsilon'_{33} & \epsilon'_{32} \\ 0 & 0 & 0 \\ 0 & \epsilon'_{13} & -\epsilon'_{12} \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} -\epsilon'_{22} & \epsilon'_{12} & 0 \\ \epsilon'_{21} & -(\epsilon'_{33} + \epsilon'_{11}) & \epsilon'_{23} \\ 0 & \epsilon'_{32} & -\epsilon'_{22} \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} -\epsilon'_{32} & \epsilon'_{31} & 0 \\ 0 & \epsilon'_{31} & 0 \\ \epsilon'_{12} & -\epsilon'_{11} & 0 \end{bmatrix} \underline{m} \\ \underline{m}^T \begin{bmatrix} 0 & \epsilon'_{33} & -\epsilon'_{32} \\ 0 & -\epsilon'_{33} & \epsilon'_{13} \\ 0 & 0 & 0 \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} -\epsilon'_{23} & 0 & \epsilon'_{21} \\ \epsilon'_{13} & 0 & -\epsilon'_{11} \\ 0 & 0 & 0 \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} -\epsilon'_{33} & 0 & \epsilon'_{13} \\ 0 & -\epsilon'_{33} & \epsilon'_{23} \\ \epsilon'_{31} & \epsilon'_{32} & -(\epsilon'_{11} + \epsilon'_{22}) \end{bmatrix} \underline{m} \end{bmatrix} \right. \\ \left. + \underline{1} \right\} \tag{61}$$

式 (51) ~ (61)을 式 (47) ~ (50)에 代入하므로써 S系에서의 媒体 parameter들을 다음과 같이 β 의 多項式으로 展開할 수 있다.

$$= \frac{\underline{1}}{\det({}_i\Delta)} \left\{ \underline{\epsilon}' - \beta^2 \left\{ \underline{\epsilon}' (\underline{1} - \underline{m} \underline{m}^T) + \begin{bmatrix} E_{22} + E_{33} & -E_{21} & -E_{31} \\ -E_{12} & E_{22} + E_{11} & -E_{32} \\ -E_{13} & -E_{23} & E_{11} + E_{22} \end{bmatrix} \right. \right. \\ \left. + \begin{bmatrix} -\underline{m}^T \begin{bmatrix} 0 & 0 & 0 \\ 0 & E_{33} & E_{32} \\ 0 & -E_{23} & E_{22} \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} 0 & 0 & 0 \\ E_{33} & 0 & -E_{31} \\ -E_{23} & 0 & E_{21} \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} 0 & 0 & 0 \\ -E_{32} & E_{31} & 0 \\ E_{32} & -E_{21} & 0 \end{bmatrix} \underline{m} \\ \underline{m}^T \begin{bmatrix} 0 & E_{33} & -E_{32} \\ 0 & 0 & 0 \\ 0 & -E_{13} & E_{12} \end{bmatrix} \underline{m} & -\underline{m}^T \begin{bmatrix} E_{33} & 0 & -E_{31} \\ 0 & 0 & 0 \\ -E_{13} & 0 & E_{11} \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} E_{32} & -E_{31} & 0 \\ 0 & 0 & 0 \\ -E_{12} & E_{11} & 0 \end{bmatrix} \underline{m} \\ \underline{m}^T \begin{bmatrix} 0 & -E_{23} & E_{22} \\ 0 & E_{13} & -E_{12} \\ 0 & 0 & 0 \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} E_{23} & 0 & -E_{21} \\ -E_{13} & 0 & E_{11} \\ 0 & 0 & 0 \end{bmatrix} \underline{m} & -\underline{m}^T \begin{bmatrix} E_{32} & -E_{21} & 0 \\ -E_{12} & E_{11} & 0 \\ 0 & 0 & 0 \end{bmatrix} \underline{m} \end{bmatrix} \right. \\ \left. + \beta^4 \left\{ \det(\underline{\epsilon}') \underline{m} \underline{m}^T + \begin{bmatrix} E_{22} + E_{33} & -E_{21} & -E_{31} \\ -E_{12} & E_{22} + E_{11} & -E_{32} \\ -E_{13} & -E_{23} & E_{11} + E_{22} \end{bmatrix} (\underline{1} - \underline{m} \underline{m}^T) \right. \right. \\ \left. + \begin{bmatrix} -\underline{m}^T \begin{bmatrix} 0 & 0 & 0 \\ 0 & E_{33} & -E_{32} \\ 0 & -E_{23} & E_{22} \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} 0 & 0 & 0 \\ E_{33} & 0 & -E_{31} \\ -E_{23} & 0 & E_{21} \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} 0 & 0 & 0 \\ -E_{32} & E_{31} & 0 \\ E_{32} & -E_{21} & 0 \end{bmatrix} \underline{m} \\ \underline{m}^T \begin{bmatrix} 0 & E_{33} & -E_{32} \\ 0 & 0 & 0 \\ 0 & -E_{13} & E_{12} \end{bmatrix} \underline{m} & -\underline{m}^T \begin{bmatrix} E_{33} & 0 & -E_{31} \\ 0 & 0 & 0 \\ -E_{13} & 0 & E_{11} \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} E_{32} & -E_{31} & 0 \\ 0 & 0 & 0 \\ -E_{12} & E_{11} & 0 \end{bmatrix} \underline{m} \\ \underline{m}^T \begin{bmatrix} 0 & -E_{23} & E_{22} \\ 0 & E_{13} & -E_{12} \\ 0 & 0 & 0 \end{bmatrix} \underline{m} & \underline{m}^T \begin{bmatrix} E_{23} & 0 & -E_{21} \\ -E_{13} & 0 & E_{11} \\ 0 & 0 & 0 \end{bmatrix} \underline{m} & -\underline{m}^T \begin{bmatrix} E_{32} & -E_{21} & 0 \\ -E_{12} & E_{11} & 0 \\ 0 & 0 & 0 \end{bmatrix} \underline{m} \end{bmatrix} \right. \\ \left. - \beta^4 \det(\underline{\epsilon}') (\underline{1} - \underline{m} \underline{m}^T) \right\} \tag{62}$$

$$\underline{\underline{\mu}} = \frac{1}{\det({}_i\Delta)} \left[\underline{\underline{1}} - \beta' (\underline{\underline{1}} - \underline{\underline{m}} \underline{\underline{m}}^T) \right]$$

$$+ \left[\begin{array}{l} \underline{\underline{m}}^T \left[\begin{array}{ccc} \epsilon_{22} + \epsilon_{33} & -\epsilon_{12} & -\epsilon_{13} \\ -\epsilon_{21} & \epsilon_{11} & 0 \\ -\epsilon_{31} & 0 & \epsilon_{11} \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} 0 & 0 & 0 \\ \epsilon_{33} & 0 & -\epsilon_{31} \\ -\epsilon_{23} & 0 & \epsilon_{21} \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} 0 & 0 & 0 \\ -\epsilon_{32} & \epsilon_{31} & 0 \\ e_2 & -\epsilon_{21} & 0 \end{array} \right] \underline{\underline{m}} \\ \underline{\underline{m}}^T \left[\begin{array}{ccc} 0 & \epsilon_{33} & -\epsilon_{32} \\ -\epsilon_{13} & \epsilon_{12} & 0 \\ 0 & -\epsilon_{23} & \epsilon_{22} \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} \epsilon_{22} & -\epsilon_{12} & 0 \\ -\epsilon_{21} & \epsilon_{33} + \epsilon_{11} & -\epsilon_{23} \\ 0 & -\epsilon_{32} & \epsilon_{22} \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} \epsilon_{32} & -\epsilon_{31} & 0 \\ 0 & 0 & 0 \\ -\epsilon_{12} & \epsilon_{11} & 0 \end{array} \right] \underline{\underline{m}} \\ \underline{\underline{m}}^T \left[\begin{array}{ccc} 0 & -\epsilon_{23} & \epsilon_{22} \\ 0 & \epsilon_{13} & -\epsilon_{12} \\ 0 & 0 & 0 \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} \epsilon_{23} & 0 & -\epsilon_{21} \\ -\epsilon_{21} & 0 & \epsilon_{11} \\ 0 & 0 & 0 \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} 0 & 0 & -\epsilon_{13} \\ -0 & \epsilon_{33} & -\epsilon_{23} \\ -\epsilon_{31} & -\epsilon_{32} & \epsilon_{11} + \epsilon_{22} \end{array} \right] \underline{\underline{m}} \end{array} \right]$$

$$+ \beta' \left[\begin{array}{l} \underline{\underline{m}}^T \left[\begin{array}{ccc} \epsilon_{22} + \epsilon_{33} & -\epsilon_{12} & -\epsilon_{13} \\ -\epsilon_{21} & \epsilon_{11} & 0 \\ -\epsilon_{31} & 0 & \epsilon_{11} \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} 0 & 0 & 0 \\ \epsilon_{33} & 0 & -\epsilon_{31} \\ -\epsilon_{23} & 0 & \epsilon_{21} \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} -\epsilon_{32} & \epsilon_{33} & 0 \\ \epsilon_{22} & -\epsilon_{32} & 0 \\ 0 & 0 & 0 \end{array} \right] \underline{\underline{m}} \\ \underline{\underline{m}}^T \left[\begin{array}{ccc} 0 & \epsilon_{33} & -\epsilon_{32} \\ 0 & 0 & 0 \\ 0 & -\epsilon_{13} & \epsilon_{12} \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} \epsilon_{22} & -\epsilon_{12} & 0 \\ -\epsilon_{21} & \epsilon_{33} + \epsilon_{11} & -\epsilon_{23} \\ 0 & -\epsilon_{32} & \epsilon_{22} \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} \epsilon_{32} & -\epsilon_{31} & 0 \\ 0 & 0 & 0 \\ -\epsilon_{12} & \epsilon_{11} & 0 \end{array} \right] \underline{\underline{m}} \\ \underline{\underline{m}}^T \left[\begin{array}{ccc} 0 & -\epsilon_{23} & \epsilon_{22} \\ 0 & \epsilon_{13} & -\epsilon_{12} \\ 0 & 0 & 0 \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} \epsilon_{23} & 0 & -\epsilon_{21} \\ -\epsilon_{21} & 0 & \epsilon_{11} \\ 0 & 0 & 0 \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} \epsilon_{33} & 0 & -\epsilon_{13} \\ 0 & \epsilon_{33} & -\epsilon_{23} \\ -\epsilon_{31} & -\epsilon_{32} & \epsilon_{11} + \epsilon_{22} \end{array} \right] \underline{\underline{m}} \end{array} \right] (I - \underline{\underline{m}} \underline{\underline{m}}^T)$$

$$+ \left[\underline{\underline{m}}^T \left[\begin{array}{ccc} E_{11} & E_{12} & E_{13} \\ E_{21} & E_{22} & E_{23} \\ E_{31} & E_{32} & E_{33} \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \right]$$

(63)

$$\underline{\underline{\xi}} = \frac{\beta/C}{\det({}_i\Delta)} \left[(\underline{\underline{\xi}}' - \underline{\underline{1}}) \left[\begin{array}{ccc} 0 & -m_3 & m_3 \\ m_3 & 0 & -m_1 \\ -m_3 & m_1 & 0 \end{array} \right] \right.$$

$$+ \beta' \left[\begin{array}{l} \left[\begin{array}{ccc} m_2 D_{31} - m_3 D_{21} & m_3 D_{11} - m_1 D_{21} & m_1 D_{21} - m_2 D_{11} \\ m_2 D_{32} - m_3 D_{22} & m_3 D_{12} - m_1 D_{22} & m_1 D_{22} - m_2 D_{12} \\ m_2 D_{33} - m_3 D_{23} & m_3 D_{13} - m_1 D_{23} & m_1 D_{23} - m_2 D_{13} \end{array} \right] - \left[\begin{array}{ccc} m_2 A_{33} - m_3 A_{23} & m_3 C_{31} - m_1 B_{21} & m_1 B_{21} - m_2 C_{31} \\ m_2 B_{32} - m_3 C_{32} & m_3 A_{31} - m_1 A_{12} & m_1 C_{12} - m_2 B_{12} \\ m_2 C_{33} - m_3 B_{23} & m_3 B_{13} - m_1 C_{13} & m_1 A_{13} - m_2 A_{31} \end{array} \right] \\ \underline{\underline{m}}^T \left[\begin{array}{ccc} 0 & 0 & 0 \\ E_{21} & E_{22} & E_{23} \\ -E_{21} & -E_{22} & -E_{23} \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} -E_{31} & -E_{32} & -E_{33} \\ 0 & 0 & 0 \\ E_{11} & E_{12} & E_{13} \end{array} \right] \underline{\underline{m}} \underline{\underline{m}}^T \left[\begin{array}{ccc} E_{21} & E_{22} & E_{23} \\ -E_{21} & -E_{22} & -E_{23} \\ 0 & 0 & 0 \end{array} \right] \underline{\underline{m}} \end{array} \right]$$

(64)

여기서

$$A_{ij} = (1 - m_k^2) E_{jk} + m_i m_k E_{ji} + m_j m_k E_{ij}$$

$$B_{ij} = m_i^2 (\epsilon_{ij}' + \epsilon_{ji}') - m_k^2 E_{ij} + m_i m_k (E_{kj} + \epsilon_{jk}' + \epsilon_{kj}') + m_j m_k E_{ik}$$

$$C_{ij} = m_i^2 (E_{kk} - \epsilon_{ii}') + m_j^2 (E_{kk} - \epsilon_{jj}') + m_k^2 (E_{ii} - \epsilon_{kk}') - m_i m_k (E_{ik} + E_{ki} + \epsilon_{ik}' + \epsilon_{ki}') \\ D_{ii} = (\epsilon_{ii}' - 1) (\epsilon_{jj}' + \epsilon_{kk}') - \epsilon_{ji}' - \epsilon_{ki}' \\ D_{ij} (i \neq j) = \epsilon_{ii}' (\epsilon_{ji}' - \epsilon_{ij}') + \epsilon_{ji}' \epsilon_{kk}' - \epsilon_{ki}' \epsilon_{kj}' + \epsilon_{ij}'$$

(65)

$$\underline{\underline{\xi}} = \frac{\beta/C}{\det({}_i\Delta)} \left[\left[\begin{array}{ccc} 0 & -m_3 & m_3 \\ m_3 & 0 & -m_1 \\ -m_3 & m_1 & 0 \end{array} \right] (\underline{\underline{1}} - \underline{\underline{\xi}}') \right.$$

$$+ \beta' \left[\begin{array}{l} \left[\begin{array}{ccc} m_2 F_{12} - m_3 F_{13} & -(m_2 I_{12} + m_3 G_{12}) & m_3 G_{12} + m_2 I_{12} \\ m_1 G_{21} + m_3 I_{31} & m_1 F_{22} - m_3 F_{21} & -(m_1 I_{22} + m_3 G_{22}) \\ -(m_2 I_{31} + m_3 G_{31}) & m_2 G_{22} + m_1 I_{22} & m_2 F_{21} - m_1 F_{32} \end{array} \right] \end{array} \right]$$

(66)

여기서

$$F_{ij} = m_i^2 E_{ij} - (m_j^2 + m_k^2) \epsilon_{ji}' + m_i m_j (E_{jj} + \epsilon_{jj}') + m_i m_k (E_{ki} + \epsilon_{ki}') \\ G_{ij} = -m_i^2 \epsilon_{kj}' + (m_j^2 + m_k^2) E_{jk} + m_i m_j (E_{ik} + \epsilon_{ki}') \\ I_{ij} = m_i^2 (E_{ii} - \epsilon_{kk}') + (m_j^2 + m_k^2) (E_{kk} - \epsilon_{ii}') + m_i m_j (E_{ji} + \epsilon_{ij}') + m_i m_k (E_{ik} + E_{ki} + \epsilon_{ik}' + \epsilon_{ki}')$$

(67)

3. 例 題

以上에서 얻은 公式들을 特殊한 경우에 대해서 적용해 보자. 第1例로서 source와 媒体간에 相對運動이 存在하지 않은 경우는 $\beta = 0$ 이므로 式(54), (60)로부터 $\det({}_i\Delta) = \det({}_i\Delta) = 1$. 따라서 式(62)~(67)로부터

$$\underline{\epsilon} = \underline{\epsilon}', \quad \underline{\mu} = \underline{1}, \quad \underline{\xi} = 0, \quad \underline{\eta} = 0 \tag{68}$$

즉 우리가 豫測했던대로이다.

第 2 例로 相對運動은 存在하나 運動速度가 작아서 β' 以上の 高次項을 무시할 수 있을 경우에는 式 (54), (60) 로 부터 $\det(\underline{\Delta}) = \det(\underline{\Delta}') = 1$ 이므로 式 (62) ~ (67) 로부터

$$\left. \begin{aligned} \underline{\epsilon} &= \underline{\epsilon}' \\ \underline{\mu} &= \underline{1} \\ \underline{\xi} &= \frac{\beta}{C} (\underline{\epsilon}' - \underline{1}) \begin{bmatrix} 0 & -m_2 & m_1 \\ m_2 & 0 & -m_1 \\ -m_2 & m_1 & 0 \end{bmatrix} = \epsilon_0 \mu_0 \underline{\epsilon}' \underline{V} \times -\frac{1}{C^2} \underline{V} \times \\ \underline{\eta} &= \frac{\beta}{C} \begin{bmatrix} 0 & -m_2 & m_1 \\ m_2 & 0 & -m_1 \\ -m_2 & m_1 & 0 \end{bmatrix} (\underline{1} - \underline{\epsilon}') = -\epsilon_0 \mu_0 \underline{V} \times \underline{\epsilon}' + \frac{1}{C^2} \underline{V} \times \end{aligned} \right\} \tag{69}$$

(69) 式은 Tai²⁰⁾가 얻은 一次近似理論의 結果와 一致한다.

IV. 運動中の Source로 인한 電磁界

(\underline{k} , ω) 領域에서의 Maxwell 方程式은 式 (8) 에 표시한 바와 같이

$$-j\underline{k} \underline{E} = -j\omega \underline{B} - \underline{M} \tag{70}$$

$$-j\underline{k} \underline{H} = j\omega \underline{D} + \underline{J} \tag{71}$$

式 (32), (33) 을 式 (70), (71) 에 代入整理하면

$$\underline{H} = j \frac{1}{\omega \mu_0} \underline{\mu}^{-1} \underline{M} + \frac{1}{\omega \mu_0} \underline{\mu}^{-1} (\underline{k} - \omega \underline{\xi}) \underline{E} \tag{72}$$

$$\underline{E} = j \frac{1}{\omega \epsilon_0} \underline{\epsilon}^{-1} \underline{J} - \frac{1}{\omega \epsilon_0} \underline{\epsilon}^{-1} (\underline{k} + \omega \underline{\xi}) \underline{H} \tag{73}$$

式 (72), (73) 에서 $\underline{H}(\underline{k}, \omega)$ 또는 $\underline{E}(\underline{k}, \omega)$ 를 消去하면 (\underline{k}, ω) 領域에서 다음과 같은 Helmholtz 的 波動方程式을 얻는다.

$$\underline{W}_E \underline{E} = \underline{S}_E \tag{74}$$

$$\underline{W}_M \underline{H} = \underline{S}_M \tag{75}$$

여기서 波動演算子行列 (wave operator matrix) $\underline{W}_E(\underline{k}, \omega)$, $\underline{W}_M(\underline{k}, \omega)$ 는 각각

$$\underline{W}_E = \underline{k}_0^2 \underline{\epsilon} + (\underline{k} + \omega \underline{\xi}) \underline{\mu}^{-1} (\underline{k} - \omega \underline{\xi}) \tag{76}$$

$$\underline{W}_M = \underline{k}_0^2 \underline{\mu} + (\underline{k} - \omega \underline{\xi}) \underline{\epsilon}^{-1} (\underline{k} + \omega \underline{\xi}) \tag{77}$$

source 函數 vector $\underline{S}_E(\underline{k}, \omega)$, $\underline{S}_M(\underline{k}, \omega)$ 는 각각

$$\underline{S}_E = j\omega \mu_0 \underline{J} - j(\underline{k} + \omega \underline{\xi}) \underline{\mu}^{-1} \underline{M} \tag{78}$$

$$\underline{S}_M = j\omega \epsilon_0 \underline{M} + j(\underline{k} - \omega \underline{\xi}) \underline{\epsilon}^{-1} \underline{J} \tag{79}$$

式 (74), (75) 를 $\underline{E}(\underline{k}, \omega)$ 또는 $\underline{H}(\underline{k}, \omega)$ 에 대해서 풀면

$$\underline{E} = \underline{W}_E^{-1} \underline{S}_E \tag{80}$$

$$\underline{H} = \underline{W}_M^{-1} \underline{S}_M \tag{81}$$

만약 Source vector $\underline{S}_E(\underline{k}, \omega)$ 가 주어졌을 때는 우선 電界 $\underline{E}(\underline{k}, \omega)$ 를 式 (80) 로부터 求한다음 磁界 $\underline{H}(\underline{k}, \omega)$ 는 式 (72) 으로 求할 수 있다. 또 만약 source vector $\underline{S}_M(\underline{k}, \omega)$ 가 주어졌다면 磁界 $\underline{H}(\underline{k}, \omega)$ 를 式 (81) 에서 求한 다음 電界 $\underline{E}(\underline{k}, \omega)$ 는 式 (73) 으로 求할 수 있다.

V. 波動演算子行列과 逆行列

Source 函數가 주어졌을 때에는 式 (80), (81) 에 의하여 電磁界는 波動演算子行列의 逆行列과 source 函數 vector 的 相乘積으로 표시되므로 電磁界를 求하는 문제는 波動行列의 逆行列을 求하는 문제로 귀착된다.

1. 波動演算子行列

式 (76) 에 式 (62) ~ (67) 을 代入하여 整理하면

$$\underline{W}_E(\underline{k}) = (W_{Eij}) = \frac{1}{\det(\underline{\underline{M}})} \left[k^2 \underline{\underline{J}} + k\omega \left[\begin{array}{ccc} n_2 K_{21} - n_1 K_{21} & n_2 K_{22} - n_1 K_{22} & n_2 K_{23} - n_1 K_{23} \\ n_2 K_{11} - n_1 K_{11} & n_2 K_{12} - n_1 K_{12} & n_2 K_{13} - n_1 K_{13} \\ n_2 K_{31} - n_1 K_{31} & n_2 K_{32} - n_1 K_{32} & n_2 K_{33} - n_1 K_{33} \end{array} \right] \right. \\ \left. + \left[\begin{array}{ccc} n_2 L_{11} - n_1 L_{11} & n_2 L_{12} - n_1 L_{12} & n_2 L_{13} - n_1 L_{13} \\ n_2 L_{21} - n_1 L_{21} & n_2 L_{22} - n_1 L_{22} & n_2 L_{23} - n_1 L_{23} \\ n_2 L_{31} - n_1 L_{31} & n_2 L_{32} - n_1 L_{32} & n_2 L_{33} - n_1 L_{33} \end{array} \right] - \omega^2 \underline{\underline{N}} \right] + k_0^2 \underline{\underline{E}} \quad (82)$$

$$J_{ii} = -n_i^2 M_{kk} - n_k^2 M_{jj} + n_i n_k (M_{jk} + M_{kj}) \quad (83)$$

$$J_{ij} (i \neq j) = n_k^2 M_{ij} + n_i n_j M_{kk} - n_i n_k M_{kj} - n_j n_k M_{ik} \quad (84)$$

$$K_{ij} = (M_{1i}, M_{2i}, M_{3i}) (\eta_{1j}, \eta_{2j}, \eta_{3j})^T \quad (86)$$

$$L_{ij} = (\zeta_{1i}, \zeta_{2i}, \zeta_{3i}) (M_{j1}, M_{j2}, M_{j3})^T \quad (87)$$

$$N_{ij} = (\zeta_{1i}, \zeta_{2i}, \zeta_{3i}) (M)^T (\eta_{1j}, \eta_{2j}, \eta_{3j})^T \quad (88)$$

$$(M) = (M_{ij}) = \begin{bmatrix} M_{11} & M_{12} & M_{13} \\ M_{21} & M_{22} & M_{23} \\ M_{31} & M_{32} & M_{33} \end{bmatrix} \quad (89)$$

$$M_{ij} = \text{cofactor of the element } \mu_{ij} \text{ of } \det(\underline{\underline{M}}) \quad (90)$$

2. 電氣波動行列의 逆行列

波動行列의 逆行列를 求하는 과정은 대단히 複雜하고도 지루하므로 여기서는 그 結果만을 표시하겠다.

지금

$$\underline{W}_E^{-1} = \frac{\text{adj}(\underline{W}_E)}{\det(\underline{W}_E)} = \frac{[\text{cof}(W_{Eij})]}{\det(\underline{W}_E)} \quad (91)$$

라 농면 각 除因 子는 다음과 같은 k의 四次多項式이 된다.

$$\text{cof}(W_{Eij}) = k^4 \left(\frac{1}{\det(\underline{\underline{M}})} \right)^2 n_i n_j n^T (\mathcal{M}) \underline{n} + k^3 \omega \left(\frac{1}{\det(\underline{\underline{M}})} \right)^2 |P_{ij} + Q_{ij}| + k^2 \left(\frac{1}{\det(\underline{\underline{M}})} \right) \left\{ \frac{\omega^2}{\det(\underline{\underline{M}})} R_{ij} + k_0^2 S_{ij} \right\} \\ + k \omega \left(\frac{1}{\det(\underline{\underline{M}})} \right) \left\{ \frac{\omega^2}{\det(\underline{\underline{M}})} U_{ij} + k_0^2 V_{ij} \right\} + \omega^2 \left(\frac{1}{\det(\underline{\underline{M}})} \right) \left\{ \frac{\omega^2}{\det(\underline{\underline{M}})} (\mathcal{R}_{ij}, \mathcal{R}_{2j}, \mathcal{R}_{3j}) (\mathcal{M}) (Z_{1i}, Z_{2i}, Z_{3i})^T + k_0^2 T_{ij} \right\} \\ + k_0^2 E_{ij} \quad (92)$$

여기서

$$(E) = (E_{ij}) = (\text{cofactor of the element } \varepsilon_{ij} \text{ of } \det(\underline{\underline{E}}))$$

$$(M) = (M_{ij}) = (\text{cofactor of the element } \mu_{ij} \text{ of } \det(\underline{\underline{M}}))$$

$$(\varepsilon) = (\varepsilon_{ij}) = (\text{cofactor of the element } \varepsilon_{ij} \text{ of } \det(\underline{\underline{E}}))$$

$$(\mathcal{M}) = (\mathcal{M}_{ij}) = (\text{cofactor of the element } M_{ij} \text{ of } \det(\underline{\underline{M}}))$$

$$(\mathcal{R}) = (\mathcal{R}_{ij}) = (\text{cofactor of the element } \eta_{ij} \text{ of } \det(\underline{\underline{\eta}}))$$

$$(Z) = (Z_{ij}) = (\text{cofactor of the element } \zeta_{ij} \text{ of } \det(\underline{\underline{\zeta}})) \quad (93)$$

$$P_{ii} = n_i (n_i n_j n_k) \left[\begin{array}{ccc} 0 & 0 & 0 \\ \eta_{ij} & \eta_{jj} & \eta_{ki} \\ \eta_{ik} & \eta_{jk} & \eta_{kk} \end{array} \right] \left[\begin{array}{ccc} 0 & 0 & 0 \\ -\mu_{ki} & -\mu_{kj} & -\mu_{kk} \\ \mu_{ji} & \mu_{jj} & \mu_{jk} \end{array} \right] + \left[\begin{array}{ccc} 0 & 0 & 0 \\ \eta_{ik} & \eta_{jk} & \eta_{kk} \\ 0 & 0 & 0 \end{array} \right] \left[\begin{array}{ccc} -\mu_{ii} & -\mu_{ij} & \mu_{ik} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{array} \right] \\ + \left[\begin{array}{ccc} 0 & 0 & 0 \\ \eta_{ii} & \eta_{ji} & \eta_{ii} \\ 0 & 0 & 0 \end{array} \right] \left[\begin{array}{ccc} 0 & 0 & 0 \\ \mu_{ki} & \mu_{kj} & \mu_{kk} \\ -\mu_{ji} & -\mu_{jj} & -\mu_{jk} \end{array} \right] \left[\begin{array}{c} n_i \\ n_j \\ n_k \end{array} \right] \quad (94)$$

$$P_{ij} (i \neq j) = \varepsilon_{ijk} n_i (n_i n_j n_k) \left[\begin{array}{ccc} \eta_{ii} & \eta_{ji} & \eta_{ki} \\ 0 & 0 & 0 \\ \eta_{ik} & \eta_{jk} & \eta_{kk} \end{array} \right] \left[\begin{array}{ccc} \mu_{ki} & \mu_{kj} & \mu_{kk} \\ 0 & 0 & 0 \\ -\mu_{ii} & -\mu_{ij} & -\mu_{ik} \end{array} \right] + \left[\begin{array}{ccc} \eta_{ii} & \eta_{jj} & \eta_{kk} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{array} \right] \left[\begin{array}{ccc} -\mu_{ki} & -\mu_{kj} & -\mu_{kk} \\ 0 & 0 & 0 \\ \mu_{ii} & \mu_{ij} & \mu_{ik} \end{array} \right] \\ + \left[\begin{array}{ccc} \eta_{ik} & \eta_{jk} & \eta_{kk} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{array} \right] \left[\begin{array}{ccc} \mu_{ji} & \mu_{jj} & \mu_{jk} \\ -\mu_{ii} & -\mu_{ij} & -\mu_{ik} \\ 0 & 0 & 0 \end{array} \right] \left[\begin{array}{c} n_i \\ n_j \\ n_k \end{array} \right] \quad (95)$$

$$\begin{aligned}
 Q_{ii} = n_i(n_i n_j n_k) & \left[\begin{array}{c} \left[\begin{array}{ccc|ccc} 0 & -\mu_{ik} & \mu_{ij} & 0 & \zeta_{ji} & \zeta_{ki} \\ 0 & -\mu_{jk} & \mu_{jj} & 0 & \zeta_{jj} & \zeta_{kj} \\ 0 & -\mu_{rk} & \mu_{ki} & 0 & \zeta_{jk} & \zeta_{rk} \end{array} \right] + \left[\begin{array}{ccc|ccc} -\mu_{ik} & 0 & \mu_{ii} & \zeta_{ji} & 0 & 0 \\ -\mu_{jk} & 0 & \mu_{jj} & \zeta_{jj} & 0 & 0 \\ -\mu_{rk} & 0 & \mu_{ki} & \zeta_{jk} & 0 & 0 \end{array} \right] \\ + \left[\begin{array}{ccc|ccc} \mu_{ij} & -\mu_{ii} & 0 & \zeta_{ki} & 0 & 0 \\ \mu_{jj} & -\mu_{ji} & 0 & \zeta_{kj} & 0 & 0 \\ \mu_{kj} & -\mu_{ki} & 0 & \zeta_{rk} & 0 & 0 \end{array} \right] \left[\begin{array}{c} n_i \\ n_j \\ n_k \end{array} \right] \end{array} \right] \tag{96}
 \end{aligned}$$

$$\begin{aligned}
 Q_{ij} (i \neq j) = \epsilon_{ijk} n_j(n_i n_j n_k) & \left[\begin{array}{c} \left[\begin{array}{ccc|ccc} 0 & -\mu_{ik} & \mu_{ij} & 0 & \zeta_{ji} & \zeta_{ki} \\ 0 & -\mu_{jk} & \mu_{jj} & 0 & \zeta_{jj} & \zeta_{kj} \\ 0 & -\mu_{rk} & \mu_{ki} & 0 & \zeta_{jk} & \zeta_{rk} \end{array} \right] + \left[\begin{array}{ccc|ccc} -\mu_{ik} & 0 & \mu_{ii} & \zeta_{ji} & 0 & 0 \\ -\mu_{jk} & 0 & \mu_{jj} & \zeta_{jj} & 0 & 0 \\ -\mu_{rk} & 0 & \mu_{ki} & \zeta_{jk} & 0 & 0 \end{array} \right] \\ + \left[\begin{array}{ccc|ccc} \mu_{ij} & -\mu_{ii} & 0 & \zeta_{ki} & 0 & 0 \\ \mu_{ji} & -\mu_{ji} & 0 & \zeta_{kj} & 0 & 0 \\ \mu_{kj} & -\mu_{ki} & 0 & \zeta_{rk} & 0 & 0 \end{array} \right] \left[\begin{array}{c} n_i \\ n_j \\ n_k \end{array} \right] \end{array} \right] \tag{97}
 \end{aligned}$$

$$\begin{aligned}
 R_{ii} = (n_i n_j n_k) & \left[\begin{array}{c} \left[\begin{array}{ccc|ccc} \mathcal{R}_{ii} & \mathcal{R}_{ji} & \mathcal{R}_{ki} & \mu_{ii} & \mu_{ij} & \mu_{ik} \\ 0 & 0 & 0 & \mu_{ji} & \mu_{jj} & \mu_{jk} \\ 0 & 0 & 0 & \mu_{ki} & \mu_{kj} & \mu_{rk} \end{array} \right] + \left[\begin{array}{ccc|ccc} Z_{ii} & 0 & 0 \\ Z_{ij} & 0 & 0 \\ Z_{ik} & 0 & 0 \end{array} \right] \\ + \left[\begin{array}{c} \left(\eta_{ij} \eta_{jj} \eta_{kj} \right) \left[\begin{array}{ccc|ccc} \mu_{rk} & 0 & -\mu_{ki} & \zeta_{ji} \\ 0 & 0 & 0 & \zeta_{jj} \\ -\mu_{ik} & 0 & \mu_{ii} & \zeta_{jk} \end{array} \right] \left(\eta_{ij} \eta_{jj} \eta_{ki} \right) \left[\begin{array}{ccc|ccc} 0 & -\mu_{rk} & -\mu_{kj} & \zeta_{ji} \\ 0 & 0 & 0 & \zeta_{jj} \\ 0 & -\mu_{ik} & \mu_{ij} & \zeta_{jk} \end{array} \right] \\ \left(\eta_{ij} \eta_{jj} \eta_{kj} \right) \left[\begin{array}{ccc|ccc} 0 & 0 & 0 & \zeta_{ji} \\ \mu_{rk} & 0 & -\mu_{ki} & \zeta_{jj} \\ -\mu_{jk} & 0 & \mu_{ji} & \zeta_{jk} \end{array} \right] \left(\eta_{ij} \eta_{jj} \eta_{ki} \right) \left[\begin{array}{ccc|ccc} 0 & 0 & 0 & \zeta_{ji} \\ 0 & \mu_{rk} & -\mu_{kj} & \zeta_{jj} \\ 0 & -\mu_{jk} & \mu_{jj} & \zeta_{jk} \end{array} \right] \\ 0 & & & 0 & & 0 \end{array} \right] \\ + \left[\begin{array}{c} \left(\eta_{ij} \eta_{jj} \eta_{kj} \right) \left[\begin{array}{ccc|ccc} -\mu_{kj} & \mu_{ki} & 0 & \zeta_{ki} \\ 0 & 0 & 0 & \zeta_{kj} \\ \mu_{ij} & -\mu_{ii} & 0 & \zeta_{rk} \end{array} \right] \left(\eta_{ij} \eta_{jj} \eta_{ki} \right) \left[\begin{array}{ccc|ccc} 0 & \mu_{rk} & -\mu_{kj} & \zeta_{ki} \\ 0 & - & 0 & \zeta_{kj} \\ 0 & -\mu_{ik} & \mu_{ij} & \zeta_{rk} \end{array} \right] \\ \left(\eta_{ij} \eta_{jj} \eta_{kj} \right) \left[\begin{array}{ccc|ccc} 0 & 0 & 0 & \zeta_{ki} \\ -\mu_{ki} & \mu_{ki} & 0 & \zeta_{kj} \\ \mu_{jj} & -\mu_{ji} & 0 & \zeta_{rk} \end{array} \right] \left(\eta_{ij} \eta_{jj} \eta_{ki} \right) \left[\begin{array}{ccc|ccc} 0 & 0 & 0 & \zeta_{ki} \\ 0 & \mu_{rk} & -\mu_{kj} & \zeta_{kj} \\ 0 & -\mu_{jk} & \mu_{jj} & \zeta_{rk} \end{array} \right] \\ 0 & & & 0 & & 0 \end{array} \right] \\ + \left[\begin{array}{c} \left(\eta_{ik} \eta_{jk} \eta_{rk} \right) \left[\begin{array}{ccc|ccc} -\mu_{jk} & 0 & \mu_{jj} & \zeta_{ji} \\ \mu_{ik} & 0 & -\mu_{ii} & \zeta_{jj} \\ 0 & 0 & 0 & \zeta_{jk} \end{array} \right] \left(\eta_{ik} \eta_{jk} \eta_{rk} \right) \left[\begin{array}{ccc|ccc} 0 & -\mu_{jk} & \mu_{jj} & \zeta_{ji} \\ 0 & \mu_{ik} & -\mu_{ii} & \zeta_{jj} \\ 0 & 0 & 0 & \zeta_{jk} \end{array} \right] \\ 0 & & & 0 & & 0 \\ \left(\eta_{ik} \eta_{jk} \eta_{rk} \right) \left[\begin{array}{ccc|ccc} 0 & 0 & 0 & \zeta_{ji} \\ \mu_{rk} & 0 & -\mu_{ki} & \zeta_{jj} \\ -\mu_{jk} & 0 & \mu_{ji} & \zeta_{jk} \end{array} \right] \left(\eta_{ik} \eta_{jk} \eta_{rk} \right) \left[\begin{array}{ccc|ccc} 0 & 0 & 0 & \zeta_{ji} \\ 0 & \mu_{rk} & -\mu_{ki} & \zeta_{jj} \\ 0 & -\mu_{jk} & \mu_{jj} & \zeta_{jk} \end{array} \right] \\ 0 & & & 0 & & 0 \end{array} \right] \end{array} \right]
 \end{aligned}$$

$$\begin{aligned}
S_{ii} = & (n_i n_j n_k) \begin{bmatrix} 0 & -M_{jj} & -M_{jk} \\ 0 & M_{ij} & M_{ik} \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \varepsilon_{ij} & 0 & 0 \\ \varepsilon_{jj} & 0 & 0 \\ \varepsilon_{kj} & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & -M_{kj} & M_{kk} \\ 0 & 0 & 0 \\ 0 & M_{ij} & M_{ik} \end{bmatrix} \begin{bmatrix} \varepsilon_{ik} & 0 & 0 \\ \varepsilon_{jk} & 0 & 0 \\ \varepsilon_{kk} & 0 & 0 \end{bmatrix} \\
& + \begin{bmatrix} 0 & 0 & 0 \\ \varepsilon_{ji} & \varepsilon_{jj} & \varepsilon_{jk} \\ \varepsilon_{kj} & \varepsilon_{kj} & \varepsilon_{kk} \end{bmatrix} \begin{bmatrix} 0 & 0 & 0 \\ M_{ji} & 0 & 0 \\ M_{ki} & 0 & 0 \end{bmatrix} - M_{ii} \begin{bmatrix} 0 & 0 & 0 \\ 0 & \varepsilon_{jj} & \varepsilon_{jk} \\ 0 & \varepsilon_{kj} & \varepsilon_{kk} \end{bmatrix} \begin{bmatrix} n_i \\ n_j \\ n_k \end{bmatrix} \\
S_{ij} (i \neq j) = & (n_i n_j n_{12}) \begin{bmatrix} 0 & M_{jj} & M_{jk} \\ 0 & -M_{ij} & -M_{ik} \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \varepsilon_{ii} & 0 & 0 \\ \varepsilon_{ji} & 0 & 0 \\ \varepsilon_{kj} & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 \\ 0 & -M_{kj} & -M_{1212} \\ 0 & M_{jj} & M_{j12} \end{bmatrix} \begin{bmatrix} \varepsilon_{ik} & 0 & 0 \\ \varepsilon_{jk} & 0 & 0 \\ \varepsilon_{kk} & 0 & 0 \end{bmatrix} \\
& + \begin{bmatrix} 0 & 0 & 0 \\ \varepsilon_{ji} & \varepsilon_{jj} & \varepsilon_{jk} \\ \varepsilon_{kj} & \varepsilon_{kj} & \varepsilon_{kk} \end{bmatrix} \begin{bmatrix} 0 & M_{ii} & 0 \\ 0 & 0 & 0 \\ 0 & M_{ki} & 0 \end{bmatrix} - M_{ji} \begin{bmatrix} 0 & 0 & 0 \\ \varepsilon_{ji} & 0 & \varepsilon_{jk} \\ \varepsilon_{kj} & 0 & \varepsilon_{kk} \end{bmatrix} \begin{bmatrix} n_i \\ n_j \\ n_k \end{bmatrix}
\end{aligned} \tag{100}$$

$$\begin{aligned}
U_{ij} = & n_i \left\{ \begin{bmatrix} \mathcal{K}_{ii} & \mathcal{K}_{ji} & \mathcal{K}_{ki} \end{bmatrix} \begin{bmatrix} -\mu_{ik} & 0 & M_{ii} \\ -\mu_{jk} & 0 & M_{ji} \\ -\mu_{kk} & 0 & M_{kk} \end{bmatrix} \begin{bmatrix} \xi_{ji} \\ \xi_{jj} \\ \xi_{jk} \end{bmatrix} + (H_{ii} H_{ji} H_{ki}) \begin{bmatrix} \mu_{ij} & -\mu_{ii} & 0 \\ \mu_{jj} & -\mu_{ji} & 0 \\ \mu_{kj} & -\mu_{ki} & 0 \end{bmatrix} \begin{bmatrix} \zeta_{ki} \\ \zeta_{kj} \\ \zeta_{kk} \end{bmatrix} \right. \\
& + (\eta_{ij} \eta_{jj} \eta_{ki}) \begin{bmatrix} -\mu_{ki} & -\mu_{kjc} & -\mu_{kk} \\ 0 & 0 & 0 \\ \mu_{ii} & \mu_{ij} & \mu_{ik} \end{bmatrix} \begin{bmatrix} Z_{ii} \\ Z_{ij} \\ Z_{ik} \end{bmatrix} + (\eta_{ik} \eta_{jk} \eta_{kk}) \begin{bmatrix} \mu_{ji} & \mu_{jj} & \mu_{jk} \\ -\mu_{ii} & -\mu_{ij} & -\mu_{ik} \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} Z_{ii} \\ Z_{ij} \\ Z_{ik} \end{bmatrix} \\
& + n_j \left\{ \begin{bmatrix} \mathcal{K}_{ii} & \mathcal{K}_{ji} & \mathcal{K}_{ki} \end{bmatrix} \begin{bmatrix} 0 & -\mu_{ik} & \mu_{ij} \\ 0 & -\mu_{jk} & \mu_{jj} \\ 0 & -\mu_{kk} & \mu_{kj} \end{bmatrix} \begin{bmatrix} \zeta_{ji} \\ \zeta_{jj} \\ \zeta_{jk} \end{bmatrix} + (\eta_{ij} \eta_{jj} \eta_{ki}) \begin{bmatrix} 0 & 0 & 0 \\ -\mu_{ki} & -\mu_{kj} & -\mu_{kk} \\ \mu_{ii} & \mu_{ij} & \mu_{jk} \end{bmatrix} \begin{bmatrix} Z_{ii} \\ Z_{ij} \\ Z_{ik} \end{bmatrix} \right. \\
& + n_k \left\{ \begin{bmatrix} \mathcal{K}_{ii} & \mathcal{K}_{ji} & \mathcal{K}_{ki} \end{bmatrix} \begin{bmatrix} 0 & -\mu_{ik} & \mu_{ij} \\ 0 & -\mu_{jk} & \mu_{jj} \\ 0 & -\mu_{kk} & \mu_{ij} \end{bmatrix} \begin{bmatrix} \zeta_{ki} \\ \zeta_{kj} \\ \zeta_{kk} \end{bmatrix} + (\eta_{ik} \eta_{jk} \eta_{kk}) \begin{bmatrix} 0 & 0 & 0 \\ -\mu_{ki} & -\mu_{kj} & -\mu_{kk} \\ \mu_{ji} & \mu_{jj} & \mu_{jk} \end{bmatrix} \begin{bmatrix} Z_{ii} \\ Z_{ij} \\ Z_{ik} \end{bmatrix} \right\}
\end{aligned} \tag{101}$$

$$\begin{aligned}
U_{ij} (i \neq j) = & \varepsilon_{ijk} \left[n_i \left\{ \begin{bmatrix} \mathcal{K}_{ii} & \mathcal{K}_{ji} & \mathcal{K}_{ki} \end{bmatrix} \begin{bmatrix} -\mu_{ik} & 0 & \mu_{ii} \\ -\mu_{jk} & 0 & \mu_{ji} \\ -\mu_{kk} & 0 & \mu_{ii} \end{bmatrix} \begin{bmatrix} \zeta_{ji} \\ \zeta_{jj} \\ \zeta_{jk} \end{bmatrix} + \begin{bmatrix} \mathcal{K}_{ij} & \mathcal{K}_{jj} & \mathcal{K}_{kj} \end{bmatrix} \begin{bmatrix} \mu_{ij} & 0 \\ \mu_{jj} & 0 \\ \mu_{kj} & 0 \end{bmatrix} \begin{bmatrix} \zeta_{ki} \\ \zeta_{kj} \\ \zeta_{kk} \end{bmatrix} \right. \right. \\
& + (\eta_{ii} \eta_{ji} \eta_{ki}) \begin{bmatrix} \mu_{ki} & \mu_{kj} & \mu_{kk} \\ 0 & 0 & 0 \\ -\mu_{ii} & -\mu_{ij} & -\mu_{ik} \end{bmatrix} \begin{bmatrix} Z_{ii} \\ Z_{ij} \\ Z_{ik} \end{bmatrix} + n_j \left\{ \begin{bmatrix} \mathcal{K}_{ii} & \mathcal{K}_{ji} & \mathcal{K}_{ki} \end{bmatrix} \begin{bmatrix} 0 & -\mu_{ij} \\ 0 & -\mu_{jj} \\ 0 & -\mu_{kj} \end{bmatrix} \begin{bmatrix} \zeta_{ji} \\ \zeta_{jj} \\ \zeta_{jk} \end{bmatrix} \right. \\
& + (\eta_{ii} \eta_{jj} \eta_{ki}) \begin{bmatrix} 0 & 0 & 0 \\ \mu_{ki} & \mu_{kj} & \mu_{kk} \\ -\mu_{ji} & -\mu_{jj} & -\mu_{ik} \end{bmatrix} \begin{bmatrix} Z_{ii} \\ Z_{ij} \\ Z_{ik} \end{bmatrix} + (\eta_{ik} \eta_{jk} \eta_{kk}) \begin{bmatrix} \mu_{ji} & \mu_{jj} & \mu_{j12} \\ -\mu_{ii} & -\mu_{ij} & -\mu_{i12} \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} Z_{ii} \\ Z_{ij} \\ Z_{ik} \end{bmatrix} \\
& \left. \left. + n_{12} \left\{ \begin{bmatrix} \mathcal{K}_{ij} & \mathcal{K}_{jj} & \mathcal{K}_{kj} \end{bmatrix} \begin{bmatrix} 0 & -\mu_{ik} & \mu_{ij} \\ 0 & -\mu_{jk} & \mu_{jj} \\ 0 & -\mu_{kk} & \mu_{kj} \end{bmatrix} \begin{bmatrix} \zeta_{ki} \\ \zeta_{kj} \\ \zeta_{kk} \end{bmatrix} + (\eta_{ik} \eta_{jk} \eta_{kk}) \begin{bmatrix} \mu_{ki} & \mu_{kj} & \mu_{kk} \\ 0 & 0 & 0 \\ -\mu_{ii} & -\mu_{ij} & -\mu_{ik} \end{bmatrix} \begin{bmatrix} Z_{ii} \\ Z_{ij} \\ Z_{ik} \end{bmatrix} \right\} \right]
\end{aligned} \tag{102}$$

$$V_{ii} = n_i \left\{ -(\eta_{ij} \eta_{jj} \eta_{ki}) \begin{bmatrix} 0 & M_{ij} & M_{ik} \\ 0 & M_{jj} & M_{jk} \\ 0 & M_{kj} & M_{kk} \end{bmatrix} \begin{bmatrix} \varepsilon_{ik} \\ \varepsilon_{jk} \\ \varepsilon_{kk} \end{bmatrix} + (\eta_{ik} \eta_{jk} \eta_{kk}) \begin{bmatrix} 0 & M_{ij} & M_{ik} \\ 0 & M_{jj} & M_{jk} \\ 0 & M_{kj} & M_{kk} \end{bmatrix} \begin{bmatrix} \varepsilon_{ii} \\ \varepsilon_{jj} \\ \varepsilon_{kj} \end{bmatrix} \right\}$$

$$+ (\eta_{ik}\eta_{jk}\eta_{kk}) \begin{bmatrix} M_{ii} & M_{ij} & M_{ik} \\ M_{ji} & M_{jj} & M_{jk} \\ M_{ki} & M_{kj} & M_{kk} \end{bmatrix} \begin{bmatrix} 0 & \zeta_{ki} & -\zeta_{ji} \\ 0 & \zeta_{kj} & -\zeta_{jj} \\ 0 & \zeta_{kk} & -\zeta_{jk} \end{bmatrix} \begin{bmatrix} \varepsilon_{ii} \\ \varepsilon_{jj} \\ \varepsilon_{ki} \end{bmatrix}$$

한편 行列式을 예컨대 第 1 行에 따라 展開 하면

$$\det(\underline{W}_E) = W_{E11} \text{ cof}(W_{E11}) + W_{E12} \text{ cof}(W_{E12}) + W_{E13} \text{ cof}(W_{E13}) \tag{106}$$

波動行列 W_{E11} , W_{E12} , W_{E13} 는 모두 k 의 二次式이며 그들의 餘因子는 모두 k 의 四次式이므로 波動演算子行列의 行列式은 k 의 六次式이다. 그러나 긴 計算 끝에 最高次의 2 項이 消滅되어 式(106)의 行列式은 결국 다음과 같은 k 의 四次式이 됨을 알 수 있다.

$$\begin{aligned} \det(\underline{W}_E) &= \frac{k^4}{(\det(\underline{\mu}))^2} \left\{ \frac{\omega^2}{\det(\underline{\mu})} \Gamma + k_0^2 \Delta \right\} + \frac{k^3 \omega}{(\det(\underline{\mu}))^2} \left\{ \frac{\omega^2}{\det(\underline{\mu})} \Theta + k_0^2 \Lambda \right\} \\ &+ \frac{k^2}{\det(\underline{\mu})} \left\{ \frac{\omega^4}{(\det(\underline{\mu}))^2} \Xi + \frac{\omega^2 k_0^2}{\det(\underline{\mu})} O + k_0^4 \Pi \right\} \\ &+ \frac{k \omega}{\det(\underline{\mu})} \left\{ \frac{\omega^4}{(\det(\underline{\mu}))^2} \Sigma + \frac{\omega^2 k_0^2}{\det(\underline{\mu})} \Upsilon + k_0^4 \Phi \right\} \\ &+ \frac{\omega^2}{\det(\underline{\mu})} \left\{ \frac{\omega^4}{(\det(\underline{\mu}))^2} X + \frac{\omega^2 k_0^2}{\det(\underline{\mu})} \Psi + k_0^4 \Omega \right\} + k_0^6 \det(\underline{\varepsilon}) \end{aligned} \tag{107}$$

여기서 $\Gamma, \Delta, \Theta, \Lambda, \Xi, O, \Pi, \Sigma, \Upsilon, \Phi, X, \Psi$, 및 Ω 는 다음과 같다.

$$\begin{aligned} \Gamma &= (n_1 N_{11} + n_2 N_{12} + n_3 N_{13}) n_1 \underline{n}^T (\underline{M}) \underline{n} + (P_{11} + Q_{11}) \{ (n_2 K_{31} - n_3 K_{21}) + (n_2 L_{13} - n_3 L_{12}) \} \\ &+ (P_{12} + Q_{11}) \{ (n_2 K_{32} - n_3 K_{22}) + (n_3 L_{11} - n_1 L_{13}) \} + (P_{13} + Q_{13}) \{ (n_2 K_{33} - n_3 K_{23}) + (n_1 L_{12} - n_2 L_{11}) \} \\ &+ R_{11} J_{11} + R_{12} J_{12} + R_{13} J_{13} \end{aligned} \tag{108}$$

$$\begin{aligned} \Delta &= \underline{n}^T \left\{ \underline{\varepsilon} + (n_1^2 \varepsilon_{11} + n_2^2 \varepsilon_{22} + n_3^2 \varepsilon_{33}) \begin{bmatrix} 0 & \mathcal{M}_{12} & \mathcal{M}_{13} \\ \mathcal{M}_{21} & 0 & \mathcal{M}_{23} \\ \mathcal{M}_{31} & \mathcal{M}_{32} & 0 \end{bmatrix} \right. \\ &+ \left. \begin{bmatrix} n_3^2 (\varepsilon_{13} + \varepsilon_{31}) (M_{13} + M_{31}) & n_3^2 \left\{ \begin{array}{l} (\varepsilon_{13} + \varepsilon_{31}) (\mathcal{M}_{23} + \mathcal{M}_{23}) \\ + (\varepsilon_{23} + \varepsilon_{32}) (\mathcal{M}_{13} + \mathcal{M}_{31}) \end{array} \right\} & n_2^2 \left\{ \begin{array}{l} (\varepsilon_{12} + \varepsilon_{21}) (\mathcal{M}_{23} + \mathcal{M}_{32}) \\ + (\varepsilon_{23} + \varepsilon_{32}) (\mathcal{M}_{12} + \mathcal{M}_{21}) \end{array} \right\} \\ n_1^2 (\varepsilon_{12} + \varepsilon_{21}) (\mathcal{M}_{12} + \mathcal{M}_{21}) & & n_1^2 \left\{ \begin{array}{l} (\varepsilon_{12} + \varepsilon_{21}) (\mathcal{M}_{13} + \mathcal{M}_{31}) \\ + (\varepsilon_{13} + \varepsilon_{31}) (\mathcal{M}_{12} + \mathcal{M}_{21}) \end{array} \right\} \\ & & n_2^2 (\varepsilon_{23} + \varepsilon_{32}) (\mathcal{M}_{23} + \mathcal{M}_{32}) \end{bmatrix} \right\} \underline{n} \end{aligned} \tag{109}$$

$$\begin{aligned} \Theta &= - (P_{11} + Q_{11}) N_{11} - (P_{12} + Q_{12}) N_{12} - (P_{13} + Q_{13}) N_{13} + R_{11} \{ (n_2 K_{31} - n_3 K_{21}) + (n_2 L_{13} - n_3 L_{12}) \} \\ &+ R_{12} \{ (n_2 K_{32} - n_3 K_{22}) + (n_3 L_{11} - n_1 L_{13}) \} + R_{13} \{ (n_2 K_{33} - n_3 K_{23}) + (n_1 L_{12} - n_2 L_{11}) \} \\ &+ U_{11} J_{11} + U_{12} J_{12} + U_{13} J_{13} \end{aligned} \tag{110}$$

$$A = \underline{n}^T \left[n_1 \left[\begin{array}{l} (\mathcal{M}_{21} - \mathcal{M}_{11} O) \begin{bmatrix} \eta_{13} & 0 & -\eta_{11} \\ \eta_{23} & 0 & -\eta_{21} \\ \eta_{33} & 0 & -\eta_{31} \end{bmatrix} + (\mathcal{M}_{31} O - \mathcal{M}_{11}) \begin{bmatrix} -\eta_{12} & \eta_{11} & 0 \\ -\eta_{22} & \eta_{21} & 0 \\ -\eta_{32} & \eta_{31} & 0 \end{bmatrix} \\ (\mathcal{M}_{21} - \mathcal{M}_{11} O) \begin{bmatrix} 0 & \eta_{13} & -\eta_{12} \\ 0 & \eta_{23} & -\eta_{22} \\ 0 & \eta_{33} & -\eta_{32} \end{bmatrix} + (O \mathcal{M}_{31} - \mathcal{M}_{21}) \begin{bmatrix} -\eta_{12} & \eta_{11} & 0 \\ -\eta_{22} & \eta_{21} & 0 \\ -\eta_{32} & \eta_{31} & 0 \end{bmatrix} \\ (\mathcal{M}_{31} O - \mathcal{M}_{11}) \begin{bmatrix} 0 & \eta_{13} & -\eta_{12} \\ 0 & \eta_{23} & -\eta_{22} \\ 0 & \eta_{33} & -\eta_{32} \end{bmatrix} + (O \mathcal{M}_{31} - \mathcal{M}_{21}) \begin{bmatrix} -\eta_{13} & 0 & \eta_{11} \\ -\eta_{23} & 0 & \eta_{21} \\ -\eta_{33} & 0 & \eta_{31} \end{bmatrix} \right] \right]$$

$$\begin{aligned}
 & \left[\begin{array}{c} (\mathcal{M}_{22} - \mathcal{M}_{12} \ 0) \\ (\mathcal{M}_{22} - \mathcal{M}_{12} \ 0) \\ (\mathcal{M}_{32} \ 0 - \mathcal{M}_{12}) \\ (\mathcal{M}_{23} - \mathcal{M}_{13} \ 0) \\ (\mathcal{M}_{23} - \mathcal{M}_{13} \ 0) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \end{array} \right] \left[\begin{array}{c} \eta_{13} \ 0 \ -\eta_{11} \\ \eta_{23} \ 0 \ -\eta_{21} \\ \eta_{33} \ 0 \ -\eta_{31} \\ 0 \ \eta_{13} \ -\eta_{12} \\ 0 \ \eta_{23} \ -\eta_{22} \\ 0 \ \eta_{33} \ -\eta_{32} \end{array} \right] + \left[\begin{array}{c} (\mathcal{M}_{32} \ 0 - \mathcal{M}_{12}) \\ (0 \ \mathcal{M}_{32} - \mathcal{M}_{22}) \\ (0 \ \mathcal{M}_{32} - \mathcal{M}_{22}) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \\ (0 \ \mathcal{M}_{33} - \mathcal{M}_{23}) \end{array} \right] \left[\begin{array}{c} -\eta_{12} \ \eta_{11} \ 0 \\ -\eta_{22} \ \eta_{21} \ 0 \\ -\eta_{32} \ \eta_{31} \ 0 \\ -\eta_{12} \ \eta_{11} \ 0 \\ -\eta_{22} \ \eta_{21} \ 0 \\ -\eta_{32} \ \eta_{31} \ 0 \end{array} \right] \\
 & + n^2 \left[\begin{array}{c} (\mathcal{M}_{22} - \mathcal{M}_{12} \ 0) \\ (\mathcal{M}_{22} - \mathcal{M}_{12} \ 0) \\ (\mathcal{M}_{32} \ 0 - \mathcal{M}_{12}) \\ (\mathcal{M}_{23} - \mathcal{M}_{13} \ 0) \\ (\mathcal{M}_{23} - \mathcal{M}_{13} \ 0) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \end{array} \right] \left[\begin{array}{c} 0 \ \eta_{13} \ -\eta_{12} \\ 0 \ \eta_{23} \ -\eta_{22} \\ 0 \ \eta_{33} \ -\eta_{32} \\ 0 \ \eta_{13} \ -\eta_{12} \\ 0 \ \eta_{23} \ -\eta_{22} \\ 0 \ \eta_{33} \ -\eta_{32} \end{array} \right] + \left[\begin{array}{c} (0 \ \mathcal{M}_{32} - \mathcal{M}_{22}) \\ (0 \ \mathcal{M}_{32} - \mathcal{M}_{22}) \\ (0 \ \mathcal{M}_{32} - \mathcal{M}_{22}) \\ (0 \ \mathcal{M}_{33} - \mathcal{M}_{23}) \\ (0 \ \mathcal{M}_{33} - \mathcal{M}_{23}) \\ (0 \ \mathcal{M}_{33} - \mathcal{M}_{23}) \end{array} \right] \left[\begin{array}{c} -\eta_{12} \ \eta_{11} \ 0 \\ -\eta_{22} \ \eta_{21} \ 0 \\ -\eta_{32} \ \eta_{31} \ 0 \\ -\eta_{12} \ \eta_{11} \ 0 \\ -\eta_{22} \ \eta_{21} \ 0 \\ -\eta_{32} \ \eta_{31} \ 0 \end{array} \right] \\
 & + n^3 \left[\begin{array}{c} (\mathcal{M}_{23} - \mathcal{M}_{13} \ 0) \\ (\mathcal{M}_{23} - \mathcal{M}_{13} \ 0) \\ (\mathcal{M}_{23} - \mathcal{M}_{13} \ 0) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \end{array} \right] \left[\begin{array}{c} \eta_{13} \ 0 \ -\eta_{11} \\ \eta_{23} \ 0 \ -\eta_{21} \\ \eta_{33} \ 0 \ -\eta_{31} \\ 0 \ \eta_{13} \ -\eta_{12} \\ 0 \ \eta_{23} \ -\eta_{22} \\ 0 \ \eta_{33} \ -\eta_{32} \end{array} \right] + \left[\begin{array}{c} (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \\ (\mathcal{M}_{33} \ 0 - \mathcal{M}_{13}) \end{array} \right] \left[\begin{array}{c} -\eta_{12} \ 0 \ \eta_{11} \\ -\eta_{22} \ 0 \ \eta_{21} \\ -\eta_{32} \ 0 \ \eta_{31} \\ -\eta_{12} \ 0 \ \eta_{11} \\ -\eta_{22} \ 0 \ \eta_{21} \\ -\eta_{32} \ 0 \ \eta_{31} \end{array} \right] \left[\begin{array}{c} \epsilon_{11} \ \epsilon_{21} \ \epsilon_{31} \\ \epsilon_{12} \ \epsilon_{22} \ \epsilon_{32} \\ \epsilon_{13} \ \epsilon_{23} \ \epsilon_{33} \end{array} \right] n \\
 & + n_1^3 (\epsilon_{11} \epsilon_{21} \epsilon_{31}) \left[\begin{array}{c} \dots \zeta_{21} \ \dots \zeta_{22} \ \dots \zeta_{23} \\ \zeta_{11} \ \zeta_{12} \ \zeta_{13} \\ 0 \ 0 \ 0 \end{array} \right] \left[\begin{array}{c} \mathcal{M}_{13} \\ 0 \\ \mathcal{M}_{11} \end{array} \right] + \left[\begin{array}{c} \zeta_{31} \ \zeta_{32} \ \zeta_{33} \\ 0 \ 0 \ 0 \\ -\zeta_{11} \ -\zeta_{12} \ -\zeta_{13} \end{array} \right] \left[\begin{array}{c} \mathcal{M}_{12} \\ -\mathcal{M}_{11} \\ 0 \end{array} \right] \\
 & + n_2^3 (\epsilon_{12} \epsilon_{22} \epsilon_{32}) \left[\begin{array}{c} + \zeta_{11} \ -\zeta_{12} \ \dots \zeta_{13} \\ \zeta_{11} \ \zeta_{12} \ \zeta_{13} \\ 0 \ 0 \ 0 \end{array} \right] \left[\begin{array}{c} 0 \\ \mathcal{M}_{23} \\ -\mathcal{M}_{22} \end{array} \right] + \left[\begin{array}{c} \zeta_{31} \ \zeta_{32} \ \zeta_{33} \\ -\zeta_{21} \ -\zeta_{22} \ -\zeta_{23} \\ -\zeta_{21} \ -\zeta_{22} \ -\zeta_{23} \end{array} \right] \left[\begin{array}{c} \mathcal{M}_{22} \\ -\mathcal{M}_{21} \\ 0 \end{array} \right] \\
 & + n_3^3 (\epsilon_{13} \epsilon_{23} \epsilon_{33}) \left[\begin{array}{c} -\zeta_{31} \ -\zeta_{32} \ \dots \zeta_{33} \\ 0 \ 0 \ \zeta_{13} \\ \zeta_{11} \ \zeta_{12} \ 0 \end{array} \right] \left[\begin{array}{c} 0 \\ \mathcal{M}_{33} \\ -\mathcal{M}_{32} \end{array} \right] + \left[\begin{array}{c} \zeta_{31} \ \zeta_{32} \ \zeta_{33} \\ -\zeta_{21} \ -\zeta_{22} \ -\zeta_{23} \\ -\zeta_{21} \ -\zeta_{22} \ -\zeta_{23} \end{array} \right] \left[\begin{array}{c} \mathcal{M}_{33} \\ 0 \\ -\mathcal{M}_{31} \end{array} \right] \\
 & + n_1^2 n_2 (\epsilon_{11} \epsilon_{21} \epsilon_{31}) \left[\begin{array}{c} -\zeta_{21} \ -\zeta_{22} \ \dots \zeta_{23} \\ \zeta_{11} \ \zeta_{12} \ \zeta_{13} \\ 0 \ 0 \ 0 \end{array} \right] \left(\left[\begin{array}{c} \mathcal{M}_{23} \\ 0 \\ -\mathcal{M}_{21} \end{array} \right] + \left[\begin{array}{c} 0 \\ \mathcal{M}_{13} \\ -\mathcal{M}_{12} \end{array} \right] \right) + \left[\begin{array}{c} \zeta_{31} \ \zeta_{32} \ \zeta_{33} \\ 0 \ 0 \ 0 \\ -\zeta_{11} \ -\zeta_{12} \ -\zeta_{13} \end{array} \right] \left[\begin{array}{c} \mathcal{M}_{22} \\ -\mathcal{M}_{12} \\ 0 \end{array} \right] \\
 & + n_1 n_2^2 (\epsilon_{12} \epsilon_{22} \epsilon_{32}) \left[\begin{array}{c} \dots \zeta_{21} \ \dots \zeta_{22} \ \dots \zeta_{23} \\ \zeta_{11} \ \zeta_{12} \ 0 \\ 0 \ 0 \ 0 \end{array} \right] \left(\left[\begin{array}{c} \mathcal{M}_{23} \\ 0 \\ -\mathcal{M}_{21} \end{array} \right] + \left[\begin{array}{c} 0 \\ \mathcal{M}_{13} \\ -\mathcal{M}_{12} \end{array} \right] \right) + \left[\begin{array}{c} \zeta_{31} \ \zeta_{32} \ \zeta_{33} \\ -\zeta_{21} \ -\zeta_{22} \ -\zeta_{23} \\ -\zeta_{21} \ -\zeta_{22} \ -\zeta_{23} \end{array} \right] \left[\begin{array}{c} \mathcal{M}_{12} \\ -\mathcal{M}_{11} \\ 0 \end{array} \right] \quad (111)
 \end{aligned}$$

$$\begin{aligned}
 \Xi = & -R_{11}U_{11} - R_{12}U_{12} - R_{13}U_{13} + U_{11} (n_2 K_{31} - n_3 K_{21}) + (n_2 L_{13} - n_3 L_{12}) \{ + U_{12} (n_2 K_{32} - n_3 K_{22}) + (n_3 L_{11} - n_1 L_{13}) \} \\
 & + U_{13} (n_2 K_{33} - n_3 K_{23}) + (n_1 L_{12} - n_2 L_{11}) \{ + (\mathcal{R}_{11} \ \mathcal{R}_{21} \ \mathcal{R}_{31}) (\mathcal{M}) (Z_{11} \ Z_{12} \ Z_{13})^T J_{11} + (\mathcal{R}_{12} \ \mathcal{R}_{22} \ \mathcal{R}_{32}) (\mathcal{M}) (Z_{11} \ Z_{12} \ Z_{13})^T J_{12} \\
 & + (\mathcal{R}_{13} \ \mathcal{R}_{23} \ \mathcal{R}_{33}) (\mathcal{M}) (Z_{11} \ Z_{12} \ Z_{13})^T J_{13} \quad (112)
 \end{aligned}$$

$$O = n^T \underline{\mathbf{g}} \quad n \quad (113)$$

$$\begin{aligned}
 \mathbf{g}_{11} = & (\mathcal{M}_{11} \ \mathcal{M}_{21} \ \mathcal{M}_{31}) (\mathcal{R}) (\epsilon_{i1} \ \epsilon_{i2} \ \epsilon_{i3})^T + (\epsilon_{i1} \ \epsilon_{21} \ \epsilon_{31}) (Z) (\mathcal{M}_{i1} \ \mathcal{M}_{i2} \ \mathcal{M}_{i3})^T \\
 & + (\eta_{i1} \ \eta_{j1} \ \eta_{k1}) \left[\begin{array}{c} \mathcal{M}_{kk} \ 0 \ -\mathcal{M}_{ki} \\ 0 \ 0 \ 0 \\ -\mathcal{M}_{ik} \ 0 \ \mathcal{M}_{ii} \end{array} \right] \left[\begin{array}{c} -\zeta_{i1} \ \zeta_{i1} \ 0 \\ -\zeta_{j1} \ \zeta_{j1} \ 0 \\ -\zeta_{k1} \ \zeta_{k1} \ 0 \end{array} \right] \left[\begin{array}{c} \epsilon_{ij} \\ \epsilon_{jj} \\ \epsilon_{kj} \end{array} \right] \\
 & + (\eta_{i1} \ \eta_{j1} \ \eta_{k1}) \left[\begin{array}{c} \mathcal{M}_{kk} \ 0 \ -\mathcal{M}_{ki} \\ 0 \ 0 \ 0 \\ -\mathcal{M}_{ik} \ 0 \ \mathcal{M}_{ii} \end{array} \right] \left[\begin{array}{c} \zeta_{j1} \ -\zeta_{i1} \ 0 \\ \zeta_{j1} \ -\zeta_{i1} \ 0 \\ \zeta_{jk} \ -\zeta_{ik} \ 0 \end{array} \right] \left[\begin{array}{c} \epsilon_{i1} \\ \epsilon_{j1} \\ \epsilon_{k1} \end{array} \right]
 \end{aligned}$$

$$\begin{aligned}
 & + (\eta_{ii}\eta_{jj}\eta_{kk}) \begin{bmatrix} 0 & \mathcal{M}_{jk} & -\mathcal{M}_{jj} \\ 0 & -\mathcal{M}_{ik} & \mathcal{M}_{ij} \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \zeta_{ji} & -\zeta_{ii} & 0 \\ \zeta_{jj} & -\zeta_{ij} & 0 \\ \zeta_{jk} & -\zeta_{ik} & 0 \end{bmatrix} \begin{bmatrix} \varepsilon_{ik} \\ \varepsilon_{jk} \\ \varepsilon_{kk} \end{bmatrix} \\
 & + (\eta_{ii}\eta_{jk}\eta_{kk}) \begin{bmatrix} 0 & \mathcal{M}_{jk} & -\mathcal{M}_{jj} \\ 0 & -\mathcal{M}_{ik} & \mathcal{M}_{ij} \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} -\zeta_{ji} & \zeta_{ii} & 0 \\ -\zeta_{jj} & \zeta_{ij} & 0 \\ -\zeta_{jk} & \zeta_{ik} & 0 \end{bmatrix} \begin{bmatrix} \varepsilon_{ii} \\ \varepsilon_{ji} \\ \varepsilon_{ki} \end{bmatrix} \\
 & + (\eta_{ij}\eta_{jj}\eta_{kk}) \begin{bmatrix} \mathcal{M}_{ij} & -\mathcal{M}_{ji} & 0 \\ -\mathcal{M}_{ij} & \mathcal{M}_{ii} & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} -\zeta_{ki} & 0 & \zeta_{ii} \\ -\zeta_{kj} & 0 & \zeta_{ij} \\ -\zeta_{kk} & 0 & \zeta_{ik} \end{bmatrix} \begin{bmatrix} \varepsilon_{ik} \\ \varepsilon_{jk} \\ \varepsilon_{kk} \end{bmatrix} \\
 & + (\eta_{ik}\eta_{jk}\eta_{kk}) \begin{bmatrix} \mathcal{M}_{jj} & -\mathcal{M}_{ji} & 0 \\ -\mathcal{M}_{ij} & \mathcal{M}_{ii} & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \zeta_{ki} & 0 & -\zeta_{ii} \\ \zeta_{ij} & 0 & -\zeta_{ij} \\ \zeta_{ik} & 0 & -\zeta_{ik} \end{bmatrix} \begin{bmatrix} \varepsilon_{ii} \\ \varepsilon_{jj} \\ \varepsilon_{kj} \end{bmatrix} \\
 & + (\eta_{ij}\eta_{jj}\eta_{kj}) \begin{bmatrix} \mathcal{M}_{jk} & 0 & -\mathcal{M}_{ji} \\ -\mathcal{M}_{ik} & 0 & \mathcal{M}_{ii} \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \zeta_{ji} & -\zeta_{ii} & 0 \\ \zeta_{jj} & -\zeta_{ij} & 0 \\ \zeta_{jk} & -\zeta_{ik} & 0 \end{bmatrix} \begin{bmatrix} \varepsilon_{ik} \\ \varepsilon_{jk} \\ \varepsilon_{kk} \end{bmatrix} \\
 & + (\eta_{ik}\eta_{jk}\eta_{kk}) \begin{bmatrix} \mathcal{M}_{jk} & 0 & -\mathcal{M}_{ji} \\ -\mathcal{M}_{ik} & 0 & \mathcal{M}_{ii} \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} -\zeta_{ji} & \zeta_{ii} & 0 \\ -\zeta_{jj} & \zeta_{ij} & 0 \\ -\zeta_{jk} & \zeta_{ik} & 0 \end{bmatrix} \begin{bmatrix} \varepsilon_{ij} \\ \varepsilon_{jj} \\ \varepsilon_{kj} \end{bmatrix}
 \end{aligned} \tag{115}$$

$$\begin{aligned}
 \mathbb{H} = \underline{n}^T & \left[\begin{array}{ccc|ccc} \hline E_{22}+E_{33} & -E_{12} & -E_{13} & M_{22}E_{23}+M_{33}E_{22} & M_{12}E_{33} & M_{13}E_{22} \\ -E_{21} & E_{33}+E_{11} & -E_{32} & M_{21}E_{33} & M_{21}E_{31}+M_{13}E_{13} & M_{23}E_{11} \\ -E_{31} & -E_{32} & E_{11}+E_{22} & M_{31}E_{23} & M_{32}E_{11} & M_{13}E_{12}+M_{21}E_{21} \\ \hline 0 & M_{13}E_{23}+M_{33}E_{13} & M_{12}E_{32}+M_{33}E_{12} & & & \\ & +M_{31}E_{33}+M_{33}E_{31} & +M_{21}E_{23}+M_{33}E_{21} & & & \\ & 0 & M_{21}E_{31}+M_{31}E_{21} & & & \\ & & +M_{12}E_{13}+M_{13}E_{12} & & & \\ & & 0 & & & \\ \hline \end{array} \right] \underline{n}
 \end{aligned} \tag{116}$$

$$\begin{aligned}
 \Sigma = n_1 [& (\mathcal{M}_{22}-\mathcal{M}_{33}) \det(\underline{M}) \det(\underline{\zeta}) + (\zeta_{33}, -\zeta_{31}, 0) + (-\zeta_{33}, 0, \zeta_{21}) \} (\mathcal{M})^T (\mathcal{M})^T (\eta)^T (M) (\zeta_{11}, \zeta_{13}, \zeta_{12})^T \\ & + n_2 [(\mathcal{M}_{31}-\mathcal{M}_{13}) \det(\underline{M}) \det(\underline{\zeta}) + (0, -\zeta_{33}, \zeta_{22}) (\mathcal{M})^T (\mathcal{M})^T (\eta)^T (M) (\zeta_{11}, \zeta_{12}, \zeta_{12})^T \\ & + (Z_{11}, Z_{12}, Z_{13}) (\mathcal{M})^T (\mathcal{M})^T (\eta)^T (M_{11}, M_{22}, M_{33})^T] + n_3 [(\mathcal{M}_{12}-\mathcal{M}_{21}) \det(\underline{M}) \det(\underline{\zeta}) \\ & + (0, -\zeta_{33}, \zeta_{22}) (\mathcal{M})^T (\mathcal{M})^T (\eta)^T (M) (\zeta_{11}, \zeta_{12}, \zeta_{12})^T - (Z_{11}, Z_{12}, Z_{13}) (\mathcal{M})^T (\mathcal{M})^T (\eta)^T (M_{11}, M_{22}, M_{33})^T]
 \end{aligned} \tag{117}$$

$$\begin{aligned}
 \mathbb{T} = n_1 [& (\mathcal{M}_{11}\mathcal{M}_{12}\mathcal{M}_{13}) (Z)^T (\varepsilon) \} (\eta_{22}-\eta_{21}, 0)^T + (-\eta_{22}, 0, \eta_{21})^T \} + (\zeta_{33}-\zeta_{12}, 0) + (-\zeta_{33}, 0, \zeta_{12}) \} (\varepsilon) (\mathcal{M})^T (\mathcal{M}_{11}\mathcal{M}_{21}\mathcal{M}_{31})^T \\ & + (\mathcal{M}_{21}\mathcal{M}_{22}\mathcal{M}_{23}) (Z)^T (\varepsilon) (\eta_{12}, 0, -\eta_{11})^T + (\zeta_{31}, 0, -\zeta_{11}) (\varepsilon) (\mathcal{M})^T (\mathcal{M}_{12}\mathcal{M}_{22}\mathcal{M}_{32})^T \\ & + (\mathcal{M}_{31}\mathcal{M}_{32}\mathcal{M}_{33}) (Z)^T (\varepsilon) (-\eta_{12}, \eta_{11}, 0)^T + (-\zeta_{31}, \zeta_{11}, 0) (\varepsilon) (\mathcal{M})^T (\mathcal{M}_{13}\mathcal{M}_{23}\mathcal{M}_{33})^T \\ & + n_2 [(\mathcal{M}_{21}\mathcal{M}_{22}\mathcal{M}_{23}) (Z)^T (\varepsilon) \} (\eta_{22}-\eta_{21}, 0)^T + (0, \eta_{13}, -\eta_{12})^T \} + (\zeta_{22}-\zeta_{13}, 0) + (0, \zeta_{31}-\zeta_{21}) \} (\varepsilon) (\mathcal{M})^T (\mathcal{M}_{12}\mathcal{M}_{22}\mathcal{M}_{32})^T \\ & + (\mathcal{M}_{31}\mathcal{M}_{32}\mathcal{M}_{33}) (Z)^T (\varepsilon) (-\eta_{22}, \eta_{21}, 0)^T + (-\zeta_{22}, \zeta_{12}, 0) (\varepsilon) (\mathcal{M})^T (\mathcal{M}_{13}\mathcal{M}_{23}\mathcal{M}_{33})^T \\ & + (\mathcal{M}_{11}\mathcal{M}_{12}\mathcal{M}_{13}) (Z)^T (\varepsilon) (0, -\eta_{22}, \eta_{22})^T + (0, -\zeta_{22}, \zeta_{22}) (\varepsilon) (\mathcal{M})^T (\mathcal{M}_{11}\mathcal{M}_{21}\mathcal{M}_{31})^T \\ & + n_3 [(\mathcal{M}_{21}\mathcal{M}_{22}\mathcal{M}_{23}) (Z)^T (\varepsilon) \} (-\eta_{22}, 0, \eta_{21})^T + (0, \eta_{13}, -\eta_{12})^T \} + (-\zeta_{33}, 0, \zeta_{22}) + (0, \zeta_{31}-\zeta_{21}) \} (\varepsilon) (\mathcal{M})^T (\mathcal{M}_{13}\mathcal{M}_{23}\mathcal{M}_{33})^T \\ & + (\mathcal{M}_{11}\mathcal{M}_{12}\mathcal{M}_{13}) (Z)^T (\varepsilon) (0, -\eta_{22}, \eta_{22})^T + (0, -\zeta_{22}, \zeta_{22}) (\varepsilon) (\mathcal{M})^T (\mathcal{M}_{11}\mathcal{M}_{21}\mathcal{M}_{31})^T \\ & + (\mathcal{M}_{31}\mathcal{M}_{32}\mathcal{M}_{33}) (Z)^T (\varepsilon) (\eta_{22}, 0, -\eta_{21})^T + (\zeta_{33}, 0, -\zeta_{13}) (\varepsilon) (\mathcal{M})^T (\mathcal{M}_{13}\mathcal{M}_{23}\mathcal{M}_{33})^T
 \end{aligned} \tag{118}$$

$$\begin{aligned}
 \Phi = \underline{n}^T & \left[\begin{array}{ccc} 0 & 0 & 0 \\ -E_{21} & -E_{22} & -E_{23} \\ E_{21} & E_{22} & E_{23} \end{array} \right] (\eta)^T \begin{bmatrix} M_{11} \\ M_{21} \\ M_{31} \end{bmatrix} + \left[\begin{array}{ccc} 0 & 0 & 0 \\ M_{21} & M_{22} & M_{23} \\ -M_{21} & -M_{22} & -M_{23} \end{array} \right] (\zeta)^T \begin{bmatrix} E_{11} \\ E_{21} \\ E_{31} \end{bmatrix} \\
 & + \left[\begin{array}{ccc} E_{31} & E_{32} & E_{33} \\ 0 & 0 & 0 \\ -E_{11} & -E_{12} & -E_{13} \end{array} \right] (\eta)^T \begin{bmatrix} M_{12} \\ M_{22} \\ M_{32} \end{bmatrix} + \left[\begin{array}{ccc} -M_{31} & -M_{32} & -M_{33} \\ 0 & 0 & 0 \\ M_{11} & M_{12} & M_{13} \end{array} \right] (\zeta)^T \begin{bmatrix} E_{12} \\ E_{22} \\ E_{32} \end{bmatrix} \\
 & + \left[\begin{array}{ccc} -E_{21} & -E_{22} & -E_{23} \\ E_{11} & E_{12} & E_{13} \\ 0 & 0 & 0 \end{array} \right] (\eta)^T \begin{bmatrix} M_{13} \\ M_{23} \\ M_{33} \end{bmatrix} + \left[\begin{array}{ccc} M_{21} & M_{22} & M_{23} \\ -M_{11} & -M_{12} & -M_{13} \\ 0 & 0 & 0 \end{array} \right] (\zeta)^T \begin{bmatrix} E_{13} \\ E_{23} \\ E_{33} \end{bmatrix} \tag{119}
 \end{aligned}$$

$$X = -(\zeta_{11}, \zeta_{12}, \zeta_{13}) (M)^T (\eta) (\mathcal{A})^T (\mathcal{A}) (Z_{11}, Z_{12}, Z_{13})^T \tag{120}$$

$$\Psi = (\epsilon_{11}, \epsilon_{12}, \epsilon_{13}) (\mathcal{A})^T (\mathcal{A}) (Z_{11}, Z_{12}, Z_{13})^T + (\epsilon_{21}, \epsilon_{22}, \epsilon_{23}) (\mathcal{A})^T (\mathcal{A}) (Z_{21}, Z_{22}, Z_{23})^T + (\epsilon_{31}, \epsilon_{32}, \epsilon_{33}) (\mathcal{A})^T (\mathcal{A}) (Z_{31}, Z_{32}, Z_{33})^T \tag{121}$$

$$\Omega = -(\eta_{11}, \eta_{12}, \eta_{13}) (M) (\zeta)^T (E_{11}, E_{21}, E_{31})^T - (\eta_{21}, \eta_{22}, \eta_{23}) (M) (\zeta)^T (E_{12}, E_{22}, E_{32})^T - (\eta_{31}, \eta_{32}, \eta_{33}) (M) (\zeta)^T (E_{13}, E_{23}, E_{33})^T \tag{122}$$

3. 電氣型 및 磁氣型問題간의 analogy

Maxwell方程式 (70)과 (71), Helmholtz의 波動方程式 (74)와 (75), 波動演算子行列 (76)과 (77), 및 source 函數 vector (78)과 (97) 등을 比較하면 電氣型 및 磁氣型電磁場問題를 公式化하는데 있어서 서로간에 다음과 같은 密接한 類似性 즉 analogy가 存在함을 알 수 있다.

$$\begin{aligned}
 \underline{W}_E(\underline{k}, \omega) & \leftrightarrow -\underline{W}_M(\underline{k}, \omega), & \underline{\epsilon} & \leftrightarrow -\underline{\mu} \\
 \underline{S}_E(\underline{k}, \omega) & \leftrightarrow -\underline{S}_M(\underline{k}, \omega), & \epsilon_0 & \leftrightarrow \mu_0, & \underline{n} & \leftrightarrow -\underline{\zeta} \\
 \underline{E}(\underline{k}, \omega) & \leftrightarrow \underline{H}(\underline{k}, \omega), & \omega & \leftrightarrow \omega, & \underline{E} & \leftrightarrow -\underline{M} \\
 \underline{D}(\underline{k}, \omega) & \leftrightarrow -\underline{E}(\underline{k}, \omega), & k_0 & \leftrightarrow k_0, & \underline{\mathcal{A}} & \leftrightarrow -\underline{Z} \\
 \underline{J}(\underline{k}, \omega) & \leftrightarrow -\underline{M}(\underline{k}, \omega), & \underline{k} & \leftrightarrow \underline{k}, & \underline{\mathcal{E}} & \leftrightarrow -\underline{\mathcal{E}}
 \end{aligned} \tag{123}$$

여기서 \underline{E} 와 \underline{M} , $\underline{\mathcal{A}}$ 와 \underline{Z} 및 $\underline{\mathcal{E}}$ 와 $-\underline{\mathcal{E}}$ 등은 式 (93)에서 定義된 行列이며 媒体 parameter의 逐次的除 因子 次數가 漸次로 높아지는 順序로 配列되어 있다.

本 論文에서 誘導한 어떤 한型的 表現式중의 모든 量은 위 analogy對照表上的 相應한 量으로 置換하면 다른 型的 相應하는 表現式을 얻을 수 있다. 이 原理은 波動演算子行列과 그 逆行列에도 적용되므로 이제 電氣型波動行列과 그 逆行列이 각각 式 (82)~(90) 및 式 (91)~(122)와 같이 알려졌으므로 磁氣型波動行列과 그 逆行列의 表現式을 求하는 일은 간단한 analogy法의 例에 불과하다.

VI. 相對論的公式의 應用例

II 3節에서 相對運動이 없을 때는 S系에 대한 媒体 parameter의 相對論的公式이 S'系에 대한 公式과 一致하게 된다는 것을 證明하였다. 波動演算子行列과 그 逆行列에 대해서도 같은 事實을 立證할 수 있다.

즉 相對運動이 없을 때는 $\beta=0$ 이므로 式 (68)로부터

$$\underline{\epsilon} = \underline{\epsilon}, \underline{\mu} = \underline{\mu}, \underline{\zeta} = \underline{\eta} = 0 \tag{124}$$

따라서 式 (76), (77)로부터 波動演算子行列은

$$\underline{W}_E = k_0^2 \underline{\epsilon} + \underline{k} \underline{k} \tag{125}$$

$$\underline{W}_M = k_0^2 \underline{\mu} + \underline{k} \underline{\epsilon}^{-1} \underline{k} \tag{126}$$

한편 式 (78), (79)로부터 source 函數 vector는

$$\underline{S}_E = j\omega\mu_0 \underline{J} - j\underline{k} \underline{M} \tag{127}$$

$$\underline{S}_M = j\omega\epsilon_0 \underline{M} + j\underline{k} \underline{J} \tag{128}$$

이 式들은 靜止系에서의 잘 알려진 式들과 一致한다.¹²¹⁾

式 (124)로부터

$$\det(\underline{\underline{L}}) = 1 \tag{129}$$

따라서 式 (93)으로부터

$$(M) = (M_{ij}) = \underline{\underline{1}} \tag{130}$$

$$(\mathcal{M}) = (\mathcal{M}_{ij}) = \underline{\underline{1}} \tag{131}$$

$$(\mathcal{X}) = (\mathcal{X}_{ij}) = 0 \tag{132}$$

$$(Z) = (Z_{ij}) = 0 \tag{133}$$

그러므로 式 (94) ~ (106)로부터

$$P_{ij} = Q_{ij} = R_{ij} = U_{ij} = V_{ij} = T_{ij} = 0 \tag{134}$$

$$S_{ii} = (n_i n_j n_k) \begin{vmatrix} -\epsilon_{jj} - \epsilon_{kk} & 0 & 0 \\ 0 & -\epsilon_{jj} & -\epsilon_{jk} \\ 0 & -\epsilon_{ki} & -\epsilon_{kk} \end{vmatrix} \begin{vmatrix} n_i \\ n_j \\ n_k \end{vmatrix} \\ = -n_i^2 (\epsilon_{jj} + \epsilon_{kk}) - n_j^2 \epsilon_{jj} - n_k^2 \epsilon_{kk} - n_j n_k (\epsilon_{jk} + \epsilon_{ki}) \tag{135}$$

$$S_{ij} (i \neq j) = (n_i n_j n_k) \begin{vmatrix} \epsilon_{ji} & 0 & 0 \\ -\epsilon_{kk} & \epsilon_{ji} & 0 \\ \epsilon_{jk} & \epsilon_{ki} & 0 \end{vmatrix} \begin{vmatrix} n_i \\ n_j \\ n_k \end{vmatrix} \\ = n_i^2 \epsilon_{ji} + n_j^2 \epsilon_{ji} - n_j n_i \epsilon_{kk} + n_k n_j \epsilon_{ki} + n_k n_i \epsilon_{jk} \tag{136}$$

式 (129) ~ (136)을 式 (90)에 代入하면

$$\text{cof}(W_{Eij}) = k^4 n_i n_j - k^2 k_o^2 S_{ij} + k_o^4 E_{ij} \tag{137}$$

한편 波動行列의 行列式에 대해서는 式 (108) ~ (122)로부터

$$\Gamma = \Theta = \Lambda = \Xi = 0 = \Sigma = \Upsilon = \Phi = X = \Psi = \Omega = 0 \tag{138}$$

$$\Delta = \underline{\underline{n}}^T \underline{\underline{\epsilon}} \underline{\underline{n}} \tag{139}$$

$$\Pi = -\underline{\underline{n}}^T \begin{vmatrix} E_{33} + E_{33} & -E_{13} & -E_{13} \\ -E_{31} & E_{33} + E_{11} & -E_{23} \\ -E_{31} & -E_{33} & E_{11} + E_{22} \end{vmatrix} \underline{\underline{n}} \tag{140}$$

式 (129)과 式 (139) ~ (140)을 式 (128)에 代入하면

$$\det(\underline{\underline{W}}_E) = k^4 k_o^2 \underline{\underline{n}}^T \underline{\underline{\epsilon}} \underline{\underline{n}} - k^2 k_o^4 \underline{\underline{n}}^T \begin{vmatrix} E_{33} + E_{33} & -E_{13} & -E_{13} \\ -E_{31} & E_{33} + E_{11} & -E_{23} \\ -E_{31} & -E_{33} & E_{11} + E_{22} \end{vmatrix} \underline{\underline{n}} + k_o^4 \det(\underline{\underline{\epsilon}}) \tag{139}$$

따라서 $\det(\underline{\underline{W}}_E) = 0$ 라 놓면

$$k^4 \underline{\underline{n}}^T \underline{\underline{\epsilon}} \underline{\underline{n}} - k^2 k_o^2 \underline{\underline{n}}^T \begin{vmatrix} E_{33} + E_{33} & -E_{13} & -E_{13} \\ -E_{31} & E_{33} + E_{11} & -E_{23} \\ -E_{31} & -E_{33} & E_{11} + E_{22} \end{vmatrix} \underline{\underline{n}} + k_o^4 \det(\underline{\underline{\epsilon}}) = 0 \tag{140}$$

式 (140)은 k^2 에 관한 二次式으로서 잘 알려진 Appleton-Hertree式^[22]의 가장 一般化된 表現式이다.

以上の 例로 表示된 바와 같이 本論文에서 誘導한 相對論的 一般公式들은 適當한 特殊條件下에서 既知의 結果와 一致한다.

相對論的公式系와 非相對論的公式系間의 差異點은 到處에서 볼 수 있으나 顯著한 한例로서 Appleton-Hartree 方程式 $\det(\underline{\underline{W}}_E) = 0$ 를 들 수 있다. 非相對論的 경우에는 前述한 바와 같이 式 (140)은 k 의 四次式이나 k^2 에 대해서는 二次式이므로 k^2 에 대해서 2個의 根이 存在하며 2根은 각각 ordinary wave와 extraordinary wave를 代表한다.

한편 相對論的 경우에는 式 (128)로부터 $\det(\underline{\underline{W}}_E) = 0$ 란 方程式은 k 의 四次式이기는 하나 k^2 의 二次式은 아니다. 따라서 一般의 4個의 相異한 根이 存在하므로써 非相對論的 경우의 波面에서 歪曲된 波面을 나타낸다. 그러

나 速度가 減少하여 零이되면 k 의 四次式은 式(140)에 表示한 바와 같이 k' 의 二次式으로 縮退하고 波面의 歪曲은 消滅된다.

VII. 結 論

Magneto-plasma內를 定速度로 運動中인 source에 依한 電磁界問題를 定磁界와 相對運動의 方向이 任意方向에 놓여있는 一般의인 경우에 대해서 論하였다.

Minkowski의 關係式을 異方性分散媒体에서도 適用할 수 있게끔 一般化했으며 誘導率等 媒体parameter의 相對論의 變換公式를 誘導하였다. 그 結果를 速度比 β 의 多項式으로 整理하므로써 相對速度의 影響을 評價하는데 便利하도록 하였다.

電磁界가 만족하는 Helmholtz의 波動方程式도 異方性媒体에 대해서 一般化했으며 이를 行列形式으로 整理하였다. 電磁界는 波動演算子行列의 逆行列과 source函數 vector의 相乘數으로 表示되므로 波動行列의 逆行列을 求하여 그 結果를 具體的으로 表示하였다.

本 論文에서 誘導한 公式들은 一般의 式들로서 既知의 特殊한 條件下에서는 잘 알려진 結果과 一致함을 例示하였다.

波動演算子行列의 決定方程式은 電磁波의 屈折率과 傳播特性을 左右하는 바 相對運動時는 非相對運動때와는 달리 決定方程式은 k 의 四次式이기는하나 k' 의 二次式은 아니기 때문에 波面에 歪曲을 招來한다.

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