

## A SUMMATION FORMULA ON KAMPÉ DE FÉRIET FUNCTION

By Manilal Shah

During the course of investigation, an attempt has been made by several workers to obtain the exact solution of many problems in quantum mechanics in terms of Appell's functions; in an extension of this work, the author establishes here a summation formula for Kampé de Fériet's function, which may prove to be useful. The formula is as follows:

$$\sum_{m=0}^n \frac{\Gamma(k+m+\frac{1}{2})(2yz)^m}{m!} F \left[ \begin{matrix} 1 & -k-m \\ 2 & \Delta(2, -m); \Delta(2, -m) \\ 2 & \Delta(2, -2k-2m) \\ 0 & -; - \end{matrix} \middle| \begin{matrix} y^{-2} \\ z^{-2} \end{matrix} \right]$$

$$= \frac{(2yz)^n \Gamma(k+n+\frac{1}{2})}{n!(y-z)} \left\{ y F \left[ \begin{matrix} 1 & -k-n \\ 2 & \Delta(2, -n-1); \Delta(2, -n) \\ 2 & \Delta(2, -2k-2n) \\ 0 & -; - \end{matrix} \middle| \begin{matrix} y^{-2} \\ z^{-2} \end{matrix} \right] \rightleftharpoons \left[ \begin{matrix} 1 & -k-n \\ 2 & \Delta(2, -n-1); \Delta(2, -n) \\ 2 & \Delta(2, -2k-2n) \\ 0 & -; - \end{matrix} \middle| \begin{matrix} z^{-2} \\ y^{-2} \end{matrix} \right] \right\}$$

where " $\rightleftharpoons$ " shows the presence of a similar term with  $y$  and  $z$  interchanged.

### 1. Introduction.

Kampé de Fériet J. [(1), p.150] has introduced a generalized hypergeometric function of two variables in the form

$$(1.1) \quad F \left[ \begin{matrix} m & a_m \\ l & b_l; b'_l \\ n & c_n \\ p & d_p; d'_p \end{matrix} \middle| \begin{matrix} x \\ y \end{matrix} \right] = \sum_{r=0}^{\infty} \sum_{s=0}^{\infty} \frac{\prod_{j=1}^m (a_j)_{r+s} \prod_{j=1}^l \{(b_j)_r (b'_j)_s\} x^r y^s}{r! s! \prod_{j=1}^n (c_j)_{r+s} \prod_{j=1}^p \{(d_j)_r (d'_j)_s\}}$$

where  $a_m$  stands for  $a_1, a_2, \dots, a_m$ ;  $\prod_{j=1}^m (a_j)_r$  represents the product  $(a_1)_r (a_2)_r \dots (a_m)_r$  and for absolute convergence of the series  $|x| < 1, |y| < 1, m+l \leq n+p+1$ .

The aim of this note is to establish a summation formula for Kampé de Fériet's function which may be utilized in the exact solution of a large number of

problems of Mathematics, both pure and applied, and in mathematical physics frequently expressed in terms of Kampé de Fériet's functions. In a recent paper [(3), p.9, (4.5)] the author has obtained the formal solution of a problem of heat conduction equation in terms of Kampé de Fériet's functions.

## 2. A Summation Formula.

We start from the known result [(3), p.5, (2.1) for  $n=0$ ]:

$$(2.1) \int_{-\infty}^{\infty} e^{-x^2} x^{2k} \left[ \left\{ x^l {}_2F_q \left( \Delta(2, -l), \frac{A_p}{B_q}; \mu x^{-2} \right) \right\} \right. \\ \left. \left\{ x^m {}_2F_q \left( \Delta(2, -m), \frac{a_p}{b_q}; \lambda x^{-2} \right) \right\} \right] dx \\ = \Gamma \left( k + \frac{1}{2}m + \frac{1}{2}l + \frac{1}{2} \right) F \left[ \begin{matrix} 1 \\ p+2 \\ 2 \\ q \end{matrix} \middle| \begin{matrix} -k - \frac{1}{2}m - \frac{1}{2}l \\ \Delta(2, -l), A_p; \Delta(2, -m), a_p \\ \Delta(2, -2k - m - l) \\ B_q; b_q \end{matrix} \right] \begin{matrix} -\mu \\ -\lambda \end{matrix} \right]$$

where  $l, m$  are positive integers and  $\Delta(m, n) = \frac{n}{m}, \frac{n+1}{m}, \dots, \frac{n+m-1}{m}$ .

Further setting  $p=q=0, \mu=-y^{-2}, \lambda=-z^{-2}, l=m$  etc, in (2.1), we get

$$(2.2) \int_{-\infty}^{\infty} e^{-x^2} x^{2k} H_m(yx) H_m(zx) dx = \Gamma \left( k + m + \frac{1}{2} \right) (4yz)^m \\ F \left[ \begin{matrix} 1 \\ 2 \\ 2 \\ 0 \end{matrix} \middle| \begin{matrix} -k - m \\ \Delta(2, -m); \Delta(2, -m) \\ \Delta(2, -2k - 2m) \\ -; - \end{matrix} \right] \begin{matrix} y^{-2} \\ z^{-2} \end{matrix} \right]$$

where  $H_n(x) = (2x)^n {}_2F_0 \left[ \Delta(2, -n), -x^2 \right]$  is the Hermite polynomial.

Therefore

$$\sum_{m=0}^n \frac{\Gamma \left( k + m + \frac{1}{2} \right) (2yz)^m}{m!} F \left[ \begin{matrix} 1 \\ 2 \\ 2 \\ 0 \end{matrix} \middle| \begin{matrix} -k - m \\ \Delta(2, -m); \Delta(2, -m) \\ \Delta(2, -2k - 2m) \\ -; - \end{matrix} \right] \begin{matrix} y^{-2} \\ z^{-2} \end{matrix} \right] \\ = \sum_{m=0}^n \frac{1}{2^m m!} \int_{-\infty}^{\infty} e^{-x^2} x^{2k} H_m(yx) H_m(zx) dx.$$

The change in order of summation and integration is easily justified and we have the R.H.S. as

$$= \int_{-\infty}^{\infty} e^{-x^2} x^{2k} \left\{ \sum_{m=0}^n \frac{H_m(yx) H_m(zx)}{2^m m!} \right\} dx.$$

Using, Christoffel-Darboux formula [(2), p. 193, (11)] :

$$\sum_{m=0}^n \frac{H_m(x)H_m(y)}{2^m m!} = \frac{H_{n+1}(x)H_n(y) - H_n(x)H_{n+1}(y)}{2^{n+1} n! (x-y)},$$

the R.H.S. becomes

$$= \frac{(y-z)^{-1}}{2^{n+1} n!} \int_{-\infty}^{\infty} e^{-x^2} x^{2k-1} [H_{n+1}(yx)H_n(zx) - H_n(yx)H_{n+1}(zx)] dx.$$

Now separating the R.H.S. as the difference of two integrals then in view of (2.1), we obtain

$$(2.3) \quad \sum_{m=0}^n \frac{\Gamma(k+m+\frac{1}{2})(2yz)^m}{m!} F \left[ \begin{matrix} 1 & -k-m \\ 2 & \Delta(2, -m); \Delta(2, -m) \\ 2 & \Delta(2, -2k-2m) \\ 0 & -; - \end{matrix} \middle| \begin{matrix} y^{-2} \\ z^{-2} \end{matrix} \right]$$

$$= \frac{(2yz)^n \Gamma(k+n+\frac{1}{2})}{n!(y-z)} \left\{ y F \left[ \begin{matrix} 1 & -k-n \\ 2 & \Delta(2, -n-1); \Delta(2, -n) \\ 2 & \Delta(2, -2k-2n) \\ 0 & -; - \end{matrix} \middle| \begin{matrix} y^{-2} \\ z^{-2} \end{matrix} \right] - \right\}$$

where “ $\rightleftharpoons$ ” is used to indicate the presence of a similar term with  $y$  and  $z$  interchanged. This is the summation formula.

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REFERENCES

- [1] Appell, Paul, and Kampé de Fériet, J., *Fonctions Hypergéométriques et Hypersphériques; Polynomes d’Hermite*. Paris: Gauthier-Villars, 1926.
- [2] Erdélyi, A., *Higher Transcendental Functions*, Vol. II, McGraw-Hill, New York, 1953.
- [3] Shah Manilal, *On applications of Hermite polynomials*, Journal of Pure and Applied Sciences, Middle East Technical University, Ankara-Turkey, (In press).