

# ON THE SEMI-SIMPLE GROUP SPACE WITH A KAEHLERIAN METRIC

By Sang-Seup Eum

## 1.

Let  $V$  be an orientable manifold of  $C^\omega$ -class, of  $2n$  dimensions. We consider a set of transformations of  $V$ , which are in one to one correspondence with the points of a space  $M$ , and we confine ourselves to the case in which  $M$  is a manifold of  $2r$  dimensions ( $r \geq n$ ). Furthermore we assume that the set of transformations form an  $2r$ -parameter compact semi-simple group, and that the following conditions are satisfied. Let  $(u_1, \dots, u_{2r})$  be a local coordinate system on  $N$  valid in some neighbourhood, and let  $A$  be any point of this neighbourhood. The transformations  $T_A$  transform the points of a neighbourhood  $N$  of  $V$  into points of a neighbourhood  $N'$  of  $V$ . If  $(x_1, \dots, x_{2n})$  is a coordinate system valid in  $N$ , and  $(x'_1, \dots, x'_{2n})$  is a coordinate system valid in  $N'$ , the transformation  $T_A$  where  $A$  has coordinates  $(u_1, \dots, u_{2r})$ , transforms the point  $P$ , whose coordinates are  $(x_1, \dots, x_{2n})$  into the point  $P'$  whose coordinates  $(x'_1, \dots, x'_{2n})$  be given by

$$(1.1) \quad x'_i = \mathcal{P}_i(x_1, \dots, x_{2n}; u_1, \dots, u_{2r})$$

the functions are real analytic functions of  $(x_1, \dots, x_{2n})$  and of  $(u_1, \dots, u_{2r})$ , and the determinant  $|\partial \mathcal{P}_i / \partial x_j|$  different from zero at any point of  $N$  for all positions of  $A$ . [1]

Now, if we put  $\bar{\alpha} = n + \alpha$ ,  $\bar{i} = r + i$  and

$$z_\alpha = x_\alpha + \sqrt{-1}x_{\bar{\alpha}}, \quad \bar{z}_\alpha = x_\alpha - \sqrt{-1}x_{\bar{\alpha}} \quad (\alpha = 1, \dots, n)$$

$$s_i = u_i + \sqrt{-1}u_{\bar{i}}, \quad \bar{s}_i = u_i - \sqrt{-1}u_{\bar{i}} \quad (i = 1, \dots, r)$$

then we get following relations instead of (1.1)

$$(1.2) \quad z'_X = \mathcal{P}_X(z_\alpha, \bar{z}_\alpha; s_i, \bar{s}_i) \quad (X = 1, \dots, n, \bar{1}, \dots, \bar{n})$$

If we eliminate  $(z_1, \dots, z_n, \bar{z}_1, \dots, \bar{z}_n)$  from equations (1.2) and the equations

$$\frac{\partial z'_X}{\partial s_I} = \frac{\partial}{\partial s_I} \mathcal{P}_X(z, \bar{z}; s, \bar{s}) \quad (I = 1, \dots, r, \bar{1}, \dots, \bar{r})$$

we obtain

$$(1.3) \quad \frac{\partial z'^x}{\partial s_I} = \xi^x_A(z, \bar{z}) A^A_I(s, \bar{s}) \quad (A, I=1, \dots, r, \bar{1}, \dots, \bar{r})$$

If we apply the conditions of integrability of (1.3), we obtain the equations

$$(1.4) \quad \xi^Y_A \frac{\partial \xi^X_B}{\partial z_Y} - \xi^Y_B \frac{\partial \xi^X_A}{\partial z_Y} = C^C_{AB} \xi^X_C$$

where

$$(1.5) \quad C^C_{AB} = B_A^I B_B^J \left( \frac{\partial A^C_I}{\partial s_J} - \frac{\partial A^C_J}{\partial s_I} \right), \quad \text{where } A^C_I B_B^J = \delta_B^C$$

and  $z$  was used instead of  $z'$  for convenience.

If we assume that

$$g_{ab} = C_{ad}^i C_{eb}^j = 0 \quad (a, b, \dots = 1, \dots, r) \text{ (conj.)}$$

and put<sup>(\*)</sup>

$$g_{a\bar{b}} = C_{a\bar{c}}^i C_{b\bar{d}}^j$$

then, for a semi-simple group the rank of the matrix  $(g_{AB})$  is  $2n$  and, since the space is compact, the quadratic form  $g_{a\bar{b}} u^a \bar{u}^b$  is positive definite. Thus, denoting by  $(g^{a\bar{b}})$  the inverse of the matrix  $(g_{a\bar{b}})$ , we can use  $g^{a\bar{b}}$  and  $g_{a\bar{b}}$  for raising up and lowering down the indicies.

If we put

$$g^{a\bar{b}} = \xi^a_{\alpha} \xi^{\bar{b}}_{\bar{\beta}} g^{\alpha\bar{\beta}}$$

and

$$\xi^a_{\alpha} = g^{a\bar{b}} g_{\alpha\bar{b}} \xi^{\bar{b}}_{\bar{\beta}}$$

where  $(g_{\alpha\bar{\beta}})$  is the inverse matrix of  $(g^{\alpha\bar{\beta}})$ , then we have

$$(1.6) \quad g_{\alpha\bar{\beta}} = \xi^a_{\alpha} \xi^{\bar{b}}_{\bar{\beta}} g_{a\bar{b}}, \quad \xi^a_{\alpha} \xi^{\bar{b}}_{\bar{\beta}} = \delta^{\bar{b}}_{\bar{\beta}}$$

If  $r=n$ , we have  $\xi^a_{\alpha} \xi^{\bar{b}}_{\bar{\alpha}} = \delta^{\bar{b}}_{\bar{\alpha}}$ ,

but if  $r > n$ ,  $\xi^a_{\alpha} \xi^{\bar{b}}_{\bar{\alpha}} \neq \delta^{\bar{b}}_{\bar{\alpha}}$

We assume that the functions  $\xi^x_A$  are complex analytic in this section, i. e.  $\xi^a_{\alpha}$  and  $\xi^{\bar{a}}_{\bar{\alpha}}$  are functions of  $z^{\alpha}$  only,  $\xi^{\bar{a}}_{\bar{\alpha}}$  and  $\xi^a_{\alpha}$  are functions of  $\bar{z}^{\bar{\alpha}}$  only, then we have

$$g^{\gamma\bar{\delta}} \frac{\partial g^{\alpha\bar{\beta}}}{\partial z_{\gamma}} - g^{\gamma\bar{\beta}} \frac{\partial g^{\alpha\bar{\delta}}}{\partial z_{\gamma}} = \xi^{\bar{\delta}}_{\bar{\epsilon}} \xi^{\bar{\beta}}_{\bar{\epsilon}} g^{\epsilon\bar{\delta}} g^{\alpha\bar{\epsilon}} \left( \xi^{\gamma}_{\epsilon} \frac{\partial \xi^{\alpha}_{\epsilon}}{\partial z_{\gamma}} - \xi^{\gamma}_{\epsilon} \frac{\partial \xi^{\alpha}_{\epsilon}}{\partial z_{\gamma}} \right)$$

but on the other hand from (1.4) we get

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(\*) In this paper we assume the self-adjointness on the all indicies

$$\xi^\beta_a \frac{\partial \xi^\alpha_b}{\partial z_\beta} - \xi^\beta_b \frac{\partial \xi^\alpha_a}{\partial z_\beta} = C_{ab}^c \xi^\alpha_c + C_{ab}^{\bar{c}} \xi^{\alpha_{\bar{c}}}$$

$$\xi^\beta_a \frac{\partial \xi^{\alpha_{\bar{b}}}}{\partial z_\beta} - \xi^\beta_{\bar{b}} \frac{\partial \xi^\alpha_a}{\partial z_\beta} = C_{a\bar{b}}^c \xi^\alpha_c + C_{a\bar{b}}^{\bar{c}} \xi^{\alpha_{\bar{c}}}$$

therefore we obtain

$$g^{\gamma\bar{\delta}} \frac{\partial g^{\alpha\bar{\beta}}}{\partial z_\gamma} - g^{\gamma\bar{\beta}} \frac{\partial g^{\alpha\bar{\delta}}}{\partial z_\gamma} = \xi^{\bar{\delta}}_{\bar{e}} \xi^{\bar{\beta}}_{\bar{b}} g^{c\bar{e}} g^{a\bar{b}} (C_{ab}^c \xi^\alpha_c + C_{ab}^{\bar{c}} \xi^{\alpha_{\bar{c}}})$$

and the following : [3]

**THEOREM 1. 1** *When  $r=n$ , a necessary and sufficient condition that (1. 6) is a Kaehlerian metric tensor is  $C_{ab}^c=0$  and  $C_{ab}^{\bar{c}}=0$ , and when  $r>n$  if  $C_{ab}^c=0$  and  $C_{ab}^{\bar{c}}=0$  then the metric tensor (1. 6) is a Kaehlerian.*

Under the Kaehlerian condition, the Christoffel symbols are given by

$$\Gamma_{\beta\gamma}^\alpha = \xi^\alpha_a \frac{\partial \xi^\beta_a}{\partial z_\gamma} = -\xi^\beta_b \frac{\partial \xi^\alpha_a}{\partial z_\gamma} \quad \text{when } r=n$$

$$\Gamma_{\beta\gamma}^\alpha = \xi^\alpha_a \xi^{\bar{e}}_{\bar{c}} \frac{\partial \xi^\beta_b}{\partial z_\gamma} \xi^{\bar{c}}_{\bar{e}} g^{a\bar{c}} g_{b\bar{e}} \quad \text{when } r>n$$

We consider only the case of  $r=n$  in this section, we can easily see that

$$\xi^\alpha_{a;\gamma} = \frac{\partial \xi^\alpha_a}{\partial z_\gamma} + \xi^\beta_a \Gamma_{\beta\gamma}^\alpha = 0$$

and

$$\xi^\alpha_{a;\gamma;\bar{\delta}} - \xi^\alpha_{a;\bar{\delta};\gamma} = \xi^\beta_a R^\alpha_{\beta\gamma\bar{\delta}} = 0$$

where ; indicates the covariant derivative w.r.t.  $\Gamma_{\beta\gamma}^\alpha$  and  $R^\alpha_{\beta\gamma\bar{\delta}}$  is the curvature tensor constructed by  $\Gamma_{\beta\gamma}^\alpha$

Let  $z^\alpha = z^\alpha(s)$  is a curve in  $V$ , and put

$$\frac{dz^\alpha}{ds} = e^a \xi^\alpha_a$$

where  $e^a$  are constants, then we can obtain [2]

$$(1. 7) \quad \frac{d^2 z^\alpha}{ds^2} + \Gamma_{\beta\gamma}^\alpha \frac{dz^\beta}{ds} \frac{dz^\gamma}{ds} = 0$$

and we shall call (1. 7) is the equation of geodesic. Hence we have the following :

**THEOREM 1. 2** *Under the assumption that  $\xi^x_A$  are complex analytic functions of  $z$ , and that  $r=n$ , if metric tensor (1. 6) satisfies the Kaehlerian condition then the following properties are satisfied.*

- (i)  $\xi^\alpha_a$  be parallel,  
(ii)  $V$  is a flat Kaehlerian manifold, ( $R^\alpha_{\beta\gamma\bar{\delta}}=0$ )  
(iii) Curve  $z^\alpha=z^\alpha(s)$  whose tangential vector is  $e^a\xi^\alpha_a$  is geodesic,  
where  $e^a$  are constants.

If we define  $g_{i\bar{j}}$  on  $M$  by the relations

$$(1.8) \quad g_{i\bar{j}} = \sum_c A^c_i A^c_{\bar{j}}$$

and further assume that  $A^c_i$  also are complex analytic, i. e.  $A^c_i$  and  $A^c_{\bar{i}}$  are functions of  $s_i$  only,  $A^c_{\bar{i}}$  and  $A^c_i$  are functions of  $\bar{s}_i$  only, then we obtain

$$\begin{aligned} \frac{\partial g_{i\bar{j}}}{\partial s_k} - \frac{\partial g_{k\bar{j}}}{\partial s_i} &= A^c_j \left( \frac{\partial A^c_i}{\partial s_k} - \frac{\partial A^c_k}{\partial s_i} \right) \\ \frac{\partial g_{i\bar{j}}}{\partial \bar{s}_k} - \frac{\partial g_{i\bar{k}}}{\partial \bar{s}_j} &= A^c_i \left( \frac{\partial A^c_{\bar{j}}}{\partial \bar{s}_k} - \frac{\partial A^c_{\bar{k}}}{\partial \bar{s}_j} \right) \end{aligned}$$

and we have from (1.5)

$$\begin{aligned} C^c_{ab} &= B_a^i B_b^j \left( \frac{\partial A^c_i}{\partial s_j} - \frac{\partial A^c_j}{\partial s_i} \right) + B_a^{\bar{i}} B_b^{\bar{j}} \left( \frac{\partial A^c_{\bar{i}}}{\partial \bar{s}_j} - \frac{\partial A^c_{\bar{j}}}{\partial \bar{s}_i} \right) \\ C^c_{ab} &= B_a^i B_b^j \left( \frac{\partial A^c_i}{\partial s_j} - \frac{\partial A^c_j}{\partial s_i} \right) + B_a^{\bar{i}} B_b^{\bar{j}} \left( \frac{\partial A^c_{\bar{i}}}{\partial \bar{s}_j} - \frac{\partial A^c_{\bar{j}}}{\partial \bar{s}_i} \right) \end{aligned}$$

therefore if  $C^c_{ab}=C^c_{ba}=0$  then the Kaehlerian condition is satisfied, and in this case we can easily see that there exist the functions  $\varphi^c(s)$ ,  $\psi^c(\bar{s})$  satisfying

$$g_{i\bar{j}} = \sum_c \frac{\partial \varphi^c(s)}{\partial s_i} \frac{\partial \psi^c(\bar{s})}{\partial \bar{s}_j}$$

Further-more by putting

$$g^{i\bar{j}} A^c_i = A^c_{\bar{j}}, \quad g^{k\bar{j}} A^c_j = A^c_k$$

we have the following relations

$$g^{i\bar{j}} = \sum_c A^c_i A^c_{\bar{j}}$$

and

$$\Gamma^i_{jk} = g^{i\bar{l}} \frac{\partial g_{\bar{l}j}}{\partial s_k} = A^c_i \frac{\partial A^c_j}{\partial s_k}$$

and

$$\begin{aligned} A^c_{j:k} &= \frac{\partial A^c_j}{\partial s_k} - A^c_i \Gamma^i_{jk} = 0 \\ A^c_{j:k;\bar{l}} - A^c_{j;\bar{l}:k} &= -A^c_i R^i_{jk\bar{l}} = 0 \end{aligned}$$

where  $:$  indicates the covariant derivative w. r. t.  $\Gamma^i_{jk}$  and  $R^i_{jk\bar{l}}$  is the



curvature tensor constructed by  $\Gamma_{jk}^i$

Let  $s^i = s^i(t)$  is a curve in  $M$ , and put

$$\frac{ds^i}{dt} = e^a A_a^i$$

then we can obtain

$$(1.9) \quad \frac{d^2 s^i}{dt^2} + \Gamma_{jk}^i \frac{ds^j}{dt} \frac{ds^k}{dt} = 0$$

and we shall call it is geodesic.

Therefore we may now conclude as follows:

**THEOREM 1.3** *Under the assumption that  $A^c_i$  are complex analytic functions of  $s$ , if metric tensor (1.8) satisfies the Kaehlerian condition then the following properties are satisfied.*

- (i)  $A^c_i$  is a parallel gradient vector
- (ii)  $M$  is a flat Kaehlerian manifold. ( $R^i_{jki} = 0$ )
- (iii) Curve  $s^i = s^i(t)$  whose tangential vector is  $e^a A_a^i$  is geodesic.
- (iv) Metric tensor of  $V$  which is defined by (1.6) also satisfies the Kaehlerian condition for all  $r$  such that  $r \geq n$ , therefore  $V$  holds all properties of Theorem 1.2.

## 2.

Take a compact semi-simple group space with Maurer-Cartan equations \*\*

$$(2.1) \quad h^b_c \frac{\partial h^a_c}{\partial z_b} - h^b_c \frac{\partial h^a_b}{\partial z_c} = C^a_{bc} h^a_c$$

where

$$(2.2) \quad C^a_{bc} = A_b^i A_c^j \left( \frac{\partial A^a_i}{\partial s_j} - \frac{\partial A^a_j}{\partial s_i} \right), \quad \text{where} \quad A^a_i A^j_a = \delta^j_i$$

and

$$\begin{aligned} C^a_{ic} &= -C^a_{ci} \\ C^a_{ab} C^f_{ce} + C^e_{bc} C^f_{ae} + C^e_{ca} C^f_{be} &= 0 \end{aligned}$$

Now, if we put

$$\frac{1}{2}(x_\alpha + \sqrt{-1}s_\alpha) = z_\alpha, \quad \frac{1}{2}(x_\alpha - \sqrt{-1}s_\alpha) = \bar{z}_\alpha$$

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(\*\*) In this section we assume that the all indices take the values 1, 2, ..., n unless otherwise stated.

then

$$(2.3) \quad x_\alpha = z_\alpha + \bar{z}_\alpha, \quad s_\alpha = \frac{1}{\sqrt{-1}}(z_\alpha - \bar{z}_\alpha)$$

and  $h^\alpha_a$  and  $A^a_i$  are functions of  $z^\alpha$  and  $\bar{z}^\alpha$ , and

$$\frac{\partial h^\alpha_a}{\partial x_\gamma} = \frac{\partial h^\alpha_a}{\partial z_\gamma}, \quad \frac{\partial A^a_i}{\partial s_\gamma} = \sqrt{-1} \frac{\partial A^a_i}{\partial z_\gamma}$$

Here we shall write  $\alpha, \beta, \gamma, \dots$  instead of  $i, j, k, \dots$  and put

$$\overline{h^\alpha_a(z, \bar{z})} = h^{\bar{\alpha}}_{\bar{a}}(z, \bar{z}), \quad \overline{A^a_\alpha(z, \bar{z})} = A^{\bar{a}}_{\bar{\alpha}}(z, \bar{z}), \quad (\bar{\alpha}, \bar{a} = \bar{1}, \dots, \bar{n})$$

then

$$h^x_A = (h^\alpha_a, 0, 0, h^{\bar{\alpha}}_{\bar{a}}), \quad A^A_x = (A^a_\alpha, 0, 0, A^{\bar{a}}_{\bar{\alpha}})$$

thus, we may obtain pure contravariant vector  $h^x_A (A=1, \dots, n, \bar{1}, \dots, \bar{n})$  and pure covariant vector  $A^A_x (A=1, \dots, n, \bar{1}, \dots, \bar{n})$ , and we also get the following relations

$$(2.4) \quad \begin{aligned} \frac{\partial h^\alpha_a}{\partial \bar{z}_\gamma} &= \frac{\partial h^\alpha_a}{\partial z_\gamma}, & \frac{\partial h^{\bar{\alpha}}_{\bar{a}}}{\partial z_\gamma} &= \frac{\partial h^{\bar{\alpha}}_{\bar{a}}}{\partial \bar{z}_\gamma} \\ \frac{\partial A^a_\alpha}{\partial \bar{z}_\gamma} &= -\frac{\partial A^a_\alpha}{\partial z_\gamma}, & \frac{\partial A^{\bar{a}}_{\bar{\alpha}}}{\partial z_\gamma} &= -\frac{\partial A^{\bar{a}}_{\bar{\alpha}}}{\partial \bar{z}_\gamma} \end{aligned}$$

From (2.1) and (2.2), we get

$$h^\beta_b(z, \bar{z}) \frac{\partial h^\alpha_c(z, \bar{z})}{\partial z_\beta} - h^\beta_c(z, \bar{z}) \frac{\partial h^\alpha_b(z, \bar{z})}{\partial z_\beta} = C^a_{bc} h^\alpha_a(z, \bar{z}) \text{ (conj.)}$$

$$C^a_{bc} = \sqrt{-1} A^b_\beta(z, \bar{z}) A^c_\gamma(z, \bar{z}) \left( \frac{\partial A^a_\beta(z, \bar{z})}{\partial z_\gamma} - \frac{\partial A^a_\gamma(z, \bar{z})}{\partial z_\beta} \right) \text{ (conj.)}$$

(i) By putting

$$(2.10) \quad \begin{aligned} g_{bc} &= -C^f_{bc} C^c_{cf} \\ b_{\beta\gamma} &= h_\beta^b h_\gamma^c g_{bc} \end{aligned}$$

where

$$h_\beta^b = g^{bc} g_{\beta\gamma} h_\gamma^c$$

we obtain

$$h^\alpha_a h_\beta^a = \delta^\alpha_\beta \quad \text{and} \quad h^\alpha_a h_\alpha^b = \delta^b_a$$

Further-more by putting

$$(2.11) \quad \Omega_{\beta\gamma}^\alpha = \frac{1}{2} C^a_{bc} h_\beta^b h_\gamma^c h^\alpha_a$$

$$(2.12) \quad E^\alpha_{\beta\gamma} = h^\alpha_a \frac{\partial h_\beta^a}{\partial z_\gamma}$$

we may obtain the following relations by the same way in [3] (pp. 90-92)

$$(b)\{\frac{\alpha}{\beta\gamma}\} = -\frac{1}{2}(E^{\alpha}_{\beta\gamma} + E^{\alpha}_{\gamma\beta})$$

$$\Omega_{\beta\gamma}^{\alpha} = \frac{1}{2}(E^{\alpha}_{\beta\gamma} - E^{\alpha}_{\gamma\beta})$$

$$(b) R^{\alpha}_{\beta\gamma\delta} = \Omega_{\gamma\delta}^{\rho} \Omega_{\rho\beta}^{\alpha}$$

$$(b) R_{\beta\gamma} = -\frac{1}{4}b_{\beta\gamma}$$

where  $(b)\{\frac{\alpha}{\beta\gamma}\}$  are the Christoffel symbols which are calculate usually from  $b_{\beta\gamma}$  and  $(b)R^{\alpha}_{\beta\gamma\delta}$  is the curvature tensor calculated from  $(b)\{\frac{\alpha}{\beta\gamma}\}$ .

(ii) If we put

$$(2. 20) \quad a_{\beta\gamma} = A^b_{\beta} A^c_{\gamma} g_{bc}$$

instead of (2. 10) then by putting

$$(2. 21) \quad \begin{aligned} L^{\alpha}_{\beta\gamma} &= A^{\alpha}_c \frac{\partial A^c_{\beta}}{\partial z_{\gamma}} \\ \phi_{\beta\gamma}^{\alpha} &= -\frac{1}{2}(L^{\alpha}_{\beta\gamma} - L^{\alpha}_{\gamma\beta}) \end{aligned}$$

we may obtain the following relations by the same way in [3] (pp. 90-92)

$$(2. 22) \quad \begin{aligned} (a) \{\frac{\alpha}{\beta\gamma}\} &= \frac{1}{2}(L^{\alpha}_{\beta\gamma} + L^{\alpha}_{\gamma\beta}) \\ (a) R^{\alpha}_{\beta\gamma\delta} &= \phi_{\gamma\delta}^{\sigma} \phi_{\sigma\beta}^{\alpha} \\ (a) R_{\alpha\beta\gamma\delta} &= -\phi_{\alpha\beta\sigma} \phi_{\gamma\delta}^{\sigma} \end{aligned}$$

where  $(a)\{\frac{\alpha}{\beta\gamma}\}$  are the Christoffel symbols which are calculated usully from  $a_{\beta\gamma}$ , and  $(a)R^{\alpha}_{\beta\gamma\delta}$  is the curvature tensor calculated from  $(a)\{\frac{\alpha}{\beta\gamma}\}$ , and

$$\phi_{\alpha\beta\sigma} = a_{\alpha\rho} \phi_{\beta\sigma}^{\rho}$$

Further-more from (2. 2), (2. 21) and the last of (2. 22) we may also obtain

$$(a) R_{\beta\gamma} = -\frac{1}{4}a_{\beta\gamma}$$

(iii) If we put

$$(2. 30) \quad g_{\alpha\bar{\beta}} = h^a_{\alpha} h^{\bar{b}}_{\bar{\beta}} g_{ab}$$

then we may obtain the follows by a straightforward calculation

$$\begin{aligned} g^{\alpha\bar{\beta}} &= h^{\alpha}_a h^{\bar{\beta}}_{\bar{b}} g^{ab} \\ \frac{\partial g_{\alpha\bar{\beta}}}{\partial z_{\gamma}} &= g_{\rho\bar{\beta}} E^{\rho}_{\alpha\gamma} + g_{\alpha\bar{\rho}} E^{\bar{\rho}}_{\bar{\beta}\gamma} \quad (E^{\bar{\rho}}_{\bar{\beta}\gamma} = \overline{E^{\rho}_{\beta\gamma}}) \end{aligned}$$

and the Kaehlerian condition

$$\frac{\partial g_{\alpha\bar{\beta}}}{\partial z_\gamma} = \frac{\partial g_{\gamma\bar{\beta}}}{\partial z_\alpha}$$

is equivalent to

$$g_{\rho\bar{\beta}} \Omega_{\alpha\gamma}{}^\rho = g_{\rho\bar{\gamma}} E^{\bar{\rho}}{}_{[\bar{\beta}|\alpha]}$$

where the right hand members indicate

$$g_{\rho\bar{\gamma}} E^{\bar{\rho}}{}_{\bar{\beta}\alpha} - g_{\rho\bar{\alpha}} E^{\bar{\rho}}{}_{\bar{\beta}\gamma}$$

If the above condition is satisfied then

$$\Gamma_{\bar{\beta}\gamma}^\alpha = g^{\alpha\bar{\epsilon}} \frac{\partial g_{\beta\bar{\epsilon}}}{\partial z_\gamma} = E^\alpha{}_{\gamma\beta} + g^{\alpha\bar{\epsilon}} g_{\gamma\bar{\rho}} E^{\bar{\rho}}{}_{\bar{\epsilon}\beta}$$

and contracting by  $\alpha=\gamma$  we get

$$\Gamma_{\bar{\beta}\alpha}^\alpha = E^\alpha{}_{\alpha\beta} + \overline{E^\alpha{}_{\alpha\beta}}$$

then from (2. 4) and (2. 12)

$$R_{\beta\bar{\gamma}} = -\frac{\partial \Gamma_{\bar{\beta}\alpha}^\alpha}{\partial \bar{z}_\gamma} = -\left( \frac{\partial}{\partial z_\gamma} E^\alpha{}_{\alpha\beta} + \overline{\frac{\partial}{\partial z_\gamma} E^\alpha{}_{\alpha\beta}} \right)$$

and we may now conclude as follows:

**THEOREM 2** *When we introduce the metric tensor (2. 30) in our semi-simple group space endowed with complex coordinates  $(z_\alpha, \bar{z}_\alpha)$  by (2. 3), if the Kaehlerian condition is satisfied then the Ricci tensor is real.*

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Mathematical Department,  
Liberal Arts and Science College  
Kyungpook University

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