Laminar Lifted Methane Jet Flames in Co-flow Air

Narayan P. Sapkal*, Won June Lee, Jeong Park**, Byeong Jun Lee**, Oh Boong Kwon

ABSTRACT

The Laminar lifted methane jet flames diluted with helium and nitrogen in co-flow air have been investigated experimentally. The chemiluminescence intensities of OH^* and CH_2O^* radicals and the radius of curvature for tri-brachial flame were measured using an intensified charge coupled device (ICCD) camera, monochromator and digital video camera. The product of OH^* and CH_2O^* is used as a excellent proxy of heat release rate. These methane jet flames could be lifted in buoyancy and jet dominated regimes despite the Schmidt number less than unity. Lifted flames were stabilized due to buoyancy induced convection in buoyancy-dominated regime. It was confirmed that increased OH^* and CH_2O^* concentration caused an increase of edge flame speed via enhanced chemical reaction in buoyancy dominated regime. In jet momentum dominated regime lifted flames were observed even for nozzle exit velocities much higher than stoichiometric laminar flame speed. An increase in radius of curvature in addition to the increased OH^* and CH_2O^* concentration stabilizes such lifted flames.

Key Words: Schmidt number, Richardson number, Buoyancy effect, Chemical effect, Edge flame speed

1. Introduction

Laminar lifted flames in non-premixed jets have been studied to clarify the characteristics of flame stabilization[1-3]. Such laminar lifted flames have tri-brachial flame structure: a lea n premixed flame, a rich premixed flame and a trailing diffusion flame, all extending from a single location. Based on cold jet similarity so lution, experimentally it was shown that propa ne and n-butane fuels (Sc < 1) exhibited stab le lifted flame, while no stable lifted flame we re observed for methane and ethane fuels (Sc < 1) in free jets [1]. Concurrently, for the lift ed flame stabilization in hot co-flow environm ent, the important chemical role of intermediat e species [6] such as, OH*, CH2O* in laminar l ifted flame stabilization has been investigated. Also it was perceived that product of OH* an d CH₂O* can be used as a proxy for heat rele ase rate [11]. Furthermore, the triple flame sp eed could be dependent upon many factors su

* 연락저자, jeongpark@pknu.ac.kr TEL: (051)629-7911 ch as, mixture strength, heat loss, buoyancy, Lewis number[10]. In this regard, present stud y is to explore why laminar lifted methane jet flames diluted with helium and nitrogen (Sc < 1) can be stabilized. To confirm effect of buo yancy, Richardson number was evaluated and chemiluminescence intensities of OH^* and CH_2 O^* and radius of curvature were measured.

2. Experimental set-up

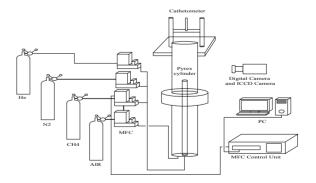


Fig. 1 Schematic diagram of experimental set-up

Experimental set-up consists of a co-flow b

^{*} Pukyong National University, Department of Mechanical Engineering, Combustion and Propulsion Lab

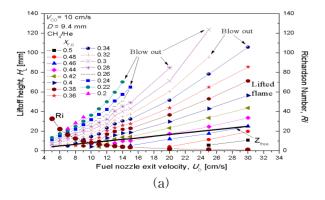
^{**} narayanpsapkal@gmail.com

urner, a flow control system, visualization syst em. Two co-flow burners used had a central f uel nozzles with 9.4 and 0.95 mm inner diamet ers and length is 100 times of the inner diame ter for the flow inside to be fully developed. T he co-flow air was supplied to the co-axial no zzle with 90.4 mm i.d. passing through glass b eads and aa honeycomb for the velocity to be uniform. The fuel was a pure grade methane diluted with helium and nitrogen and compres sed air was used for the co-flow. The flow ra tes were controlled by mass flow controller. Di gital video camera used to capture image of st ationary lifted flame. Liftoff height was measu red by the cathetometer. ICCD camera and mo nochromator were used to visualize the chemil uminescence and behaviors of lifted flame.

3. Results and Discussion

3.1 Stationary lifted flame

Experiments were conducted using 9.4 mm i.d. nozzle co-flow burner with co-flow jet vel ocity $V_{\rm CO}$ = 10 cm/s, by varying the fuel nozzl e exit velocity $U_{\rm O}$, and fuel mole fraction $X_{\rm F,O}$ for methane jet flames diluted with helium is as shown in Fig. 2(a). Liftoff heights increase d non-linearly with U_0 and were in the range of several millimeters to 123.7 mm. For a refer ence, the developing region of a free jet, Z_{free} was marked by dotted line. This was obtained to be [8], Z_{free}/d = 0.0165 X Re_{d} . where Re_{d} is the Reynolds number defined as $U_0 d_0 / v$ and vis the kinematic viscosity. The variation in Z_{fre} e with Uo is nearly linear which substantiates two different stabilization modes in the develo ping and developed regions of jet flame. Stabil izatio n mechanism has to be the balance



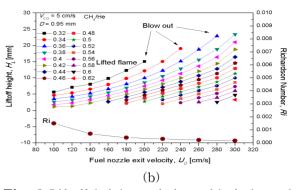


Fig. 2 Liftoff height variation with fuel nozzle exit velocity for methane diluted with helium (Sc < 1) at various fuel mole fraction for (a) 9.4 mm (b) 0.95 mm.

of edge flame speed to the local flow speed ir respective of the stabilized location in the dev eloped and developing regions. It was noted t hat, lifted flames are observed at smaller nozz le exit velocities less than stoichiometric unstr etched laminar flame speed and this was depi cted well by buoyancy effect. The effect of b uovancy was evaluated by Richardson numbe r, $Ri = \Delta \rho g d/\rho U o^{2}$, which was the ratio of buo yancy induced momentum to the jet momentu m, where g was the gravitational acceleration, ρ was the unburned density and Δρ was the d ensity difference between unburned and burne d gases. Ri number for 5-30 cm/s is in the r ange of 0.8848 < Ri < 32.86. The results sho ws that influence of buoyancy is more effective e at low Uo with high Ri number and hence the stationary lifted flames were observed. Bu t with an increase in the Uo the value of Ri number is decreasing precipitously and buoyan cy effect is suppressed significantly. Also to o btain the lifted flames at much higher Uo tha n stoichiometric un-stretched laminar flame sp eed experiments were performed using 0.95 m m i.d. nozzle with $V_{co} = 5$ cm/s. variation of $H_{\rm L}$ with $U_{\rm O}$ in laminar lifted methane jet flam es is shown in Fig. 2(b). Addition of helium d iluent and increase in the Uo causes an increa se in the $H_{\rm L}$ non-linearly until the flame is bl own out. The triple flame structure was attain ed at high H_L . Ri number also evaluated to o bserve buoyancy effect and it was in the rang e of 0.0015 < Ri < 0.0017. It is noted that, bu ovancy effect can be suppressed and lifted fla mes were obtained at larger Uo than stoichio metric unstretched non-adiabatic flam

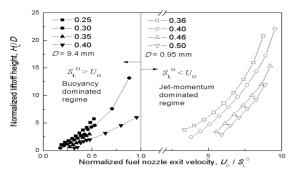


Fig. 3 Normalized liftoff height with fuel nozz le exit velocity considering stoichiometric un-s tretched non-adiabatic laminar burning velocit y for two different fuel nozzles.

e speed.

The lifted flames obtained from two different nozzle diameters co-flow burners are shown in Fig. 3. Using 9.4 mm diameter nozzle lifted flames are obtained at lower U0 than S_L° and with 0.95 mm i.d nozzle lifted flames are obtained at higher U0 than S_L° . This stoichiometric un-stretched non-adiabatic laminar burning velocity was achieved in a counter flow configuration through extrapolation of the linear relation of flame speed versus global strain rate.

Also, variation of $H_{\rm L}$ with $U_{\rm O}$ for methane diluted with nitrogen at various $X_{\rm f,o}$ is shown in Fig. 4. The flame edge has a triple flame s tructure even if it was not shown clearly. While nozzle attached flame edge shows the lean premixed flame wing is merged to the diffusion flame, such that the bi-brachial structure is exhibited. In this case, lifted flames were observed only in the developing region due to buo yancy effect at low $U_{\rm O}$ and it was confirmed by evaluating Ri number. It was in the range of 0.15 < Ri < 32.11. But at high $U_{\rm O}$, Ri number was decreased still lifted flames were observed.

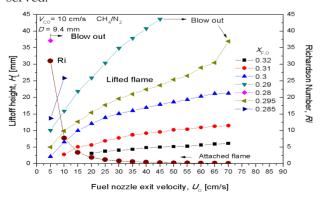
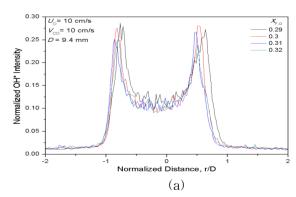


Fig. 4. Variation of H_L with U_O for methane

diluted with Nitrogen (Sc < 1) at various $X_{F,O}$ Further investigations may require to clarify the reason behind these lifted flames.

3.2 Stabilization of lifted flames

In the previous discussions, lifted flames we re observed for methane jet diluted with heliu m and nitrogen in buoyancy and jet momentu m dominated regime. So how these lifted flam es were stabilized is questionable. By stabilizat ion mechanism, edge flame speed has to increa se even with the mole fractions of helium and nitrogen despite the reduction in mixture stren gth. To clarify it, flame images were captured by the ICCD camera and chemiluminescence in tensities of OH* radicals are measured at vario us conditions. Typical radial distribution of che miluminescence intensity is as shown in the Fi g.5(a). The OH* chemiluminescence intensities have maxima at the triple point, indicating a d ouble peak. We investigated the normalized ma ximum OH* intensity which was defined by th e ratio of the OH* intensity at triple point to t he saturated intensity. Fig. 5 (b) and (c) show s that normalized maximum OH* intensities de creased with increasing fuel mole fraction and increased with helium and nitrogen mole fracti on for D = 9.4 mm nozzle diameter, which cor responds to buoyancy dominated regime. This implies that, buoyancy induced convection incr eases the reactant fluxes to the edge flame, in creasing the reaction rate of edge flame and h ence edge flame speed. Similar investigations were observed D = 0.95 mm nozzle diameter a s shown in Fig. 5(d). Normalized maximum O H* intensity increases with helium mole fractio n as well. For such a high Uo, Richardson nu mber and hence buoyancy effect is negligible, meaning that there are some other reasons.



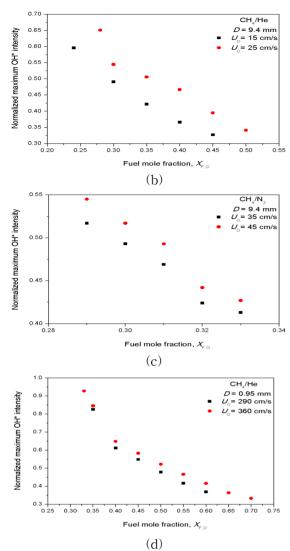


Fig. 5 (a) Typical Radial distribution of chemi luminescence intensity passing through the triple point at various $X_{\rm F,O}$. Measured OH* radic al intensities for (b) methane diluted with heli um at $U_0 = 15$ and 25 cm/s (c) methane diluted with nitrogen at $U_0 = 35$ and 45 cm/s, using 9.4 mm nozzle diameter (d) methane diluted with helium at $U_0 = 290$ and 360 cm/s using 0.95 mm nozzle inner diameter.

Also, the product of OH* and CH₂O* is obta ined using monochromatic experiment as show n in Fig. 6, which is the marker for heat rele ase rate as per previous studies [11]. It implies that, increase in the reaction rate and hence edge flame speed. It is to be noted that, edge flame speed is dependent upon mixture strength, buoyancy, fuel concentration gradient and thereby radius of curvature. Hence important role on edge flame speed enhancement may addressed to the radius of curvature. Fig. 7, sho

ws the appreciable increase in the radius of c urvature and thereby increasing edge flame sp eed.

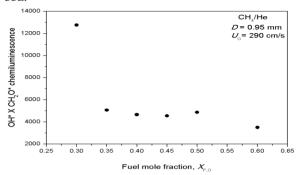


Fig. 6. Product of OH* and CH₂O* as marker for heat release rate.

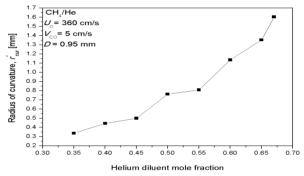


Fig. 7 Radius of curvature $r*_{cur}$, of lifted meth ane jet flame with helium dilution.

References

- [1] S.H. Chung, B.J. Lee, Combust. Flame 86 (1991) 62–72.
- [2] B.J. Lee, S.H. Chung, Combust. Flame 109 (1997) 163-172.
- [3] Y.S. Ko, S.H. Chung, Combust. Flame 118 (1999) 151-163.
- [4] S.H. Won, J.Kim, K.J. Hong, M.S.Cha, S.H. Chung, Proc. Combust. Inst. 30(2005) 339–347.
- [5] S.H. Won, S.H. Chung, M.S.Cha, B.J.Lee, Proc. Combust. Inst. 28 (2000) 2093–2099.
- [6] P.R.Medwell, D.L.Blunck, B.B. Dally, Combust. Flame 161 (2014) 465–474.
- [7] J.Bukmaster, N.Peters, Proc. Combust. Inst. 21(1986) 1829–1836.
- [8] D.S. Lee, K.D.Kihm,S.H. Chung, J.Fluids E ng. 119(3) 716–718 (1997).
- [9] G.R. Ruetsch, L. Vervisch, A.Linan, Phys. Fluids 7(6) 1447–1454 (1995).
- [10] S.H. Chung Proc. Combust. Inst.31 (2007) 877–892.
- [11] R.L. Gordon, A.R. Masri, E. Mastorakos, Combust. Theory Mod. 13(2009) 645–670.