## Perpendicular Magnetic Anisotropy of CoFeB/Pd Multilayers

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Magnetic tunnel junctions (MTJs) with perpendicular magnetic anisotropy (PMA) have the important advantage of a smaller critical current density for current induced magnetization switching (CIMS) over MTJs with an in-plane anisotropy (in-plane MTJs). This is because the current-induced spin transfer torque (STT) in the in-plane MTJs must overcome the large additional  $2\pi$ Ms factor that arises from the demagnetizing field of a magnetic thin film [1]. Since the CoFeB/MgO/CoFeB MTJs were shown to have large values of TMR, and most of the results were observed to be in the in-plane geometry, it would be highly desirable for the perpendicular MTJs structure to contain a CoFeB layer. In this study, an effort is made in this direction by investigating CoFeB/Pd multilayer system.

All the layers were deposited using a magnetron sputtering system that had two separate chambers with differentbase pressures of  $1 \times 10^{-8}$  Torr and  $2 \times 10^{-9}$  Torr. The M-H hysteresis loops, which are mainly used to characterize the PMA properties, were measured with a vibrating sample magnetometer at room temperature.

The effects of *N* were investigated and the results are shown in Figs. 1(a) and (b). The multilayer structure is shown in the inset of Fig. 1(a). The m-H hysteresis loops measured under the perpendicular applied field, shown in Fig. 1(a), indicate that all the samples have a good PMA. It can be seen in Fig. 1(b) that the coercivity ( $H_c$ ) remains nearly constant up to N=12 and is followed by a steep decrease at N=20. The results for the two critical fields, the saturation field ( $H_s$ ) and the nucleation field (or switching field) ( $H_n$ ), are also shown in the lower panel of Fig. 1(b), together with their definition in the inset. The value of  $H_s$  increases monotonically from 114 Oe to 1360 Oe, as *N* increases from 3 to 20. On the other hand, the value of  $H_n$  is nearly independent of *N* in the low *N* range ( $N=3\sim9$ ). In this range, the value of  $H_n$  is negative. However,  $H_n$  is nearly zero (at N=12) or even positive at the larger *N* values. At the largest *N* value (N=20), for example, the value of  $H_n$  is quite high (+390 Oe), indicating that the irreversible switching starts to occur even though a large value of  $H_{ext}$  is applied in the same (positive) direction. The value of *m* decreases quite steeply at  $H_n$  and it approaches zero at the remanent state. This behavior, which has been observed in the past, is explained by the formation of multi-domains with the PMA at large *N* values and the motion of the domain walls [2]. This behavior is also confirmed by an MFM image for the N=20 sample measured at the remanent state, as shown in the inset of Fig. 1(a).

## References

- [1] S. Mangin et al., Nat. Mater. 5, 210 (2006).
- [2] B. Kaplan et al., J. Magn. Magn. Mater. 128, 111 (1993).





Figure 1. (a) The *m*-*H* hysteresis loops of the CoFeB/Pd multilayers with different *N* values ( $N=3\sim20$ ). The insets show the investigated multilayer structure and an MFM image for the N=20 multilayer at the remanent state.

(b) The three characteristic fields (the coercivity  $(H_c)$ , the saturation field  $(H_s)$ , and the nucleation field  $(H_n)$ ) as a function of *N*. Shown in the inset is the definition of  $H_s$  and  $H_n$ .