

## Low roll-off of efficiency with increasing current density in phosphorescent OLEDs

**Jae-Wook Kang, Se-Hyung Lee, Hyung-Dol Park, Won-Ik Jeong, Kyung-Mo Yoo, Young-Seo Park, and Jang-Joo Kim\*\***

School of Materials Science and Engineering, and OLED Center, Seoul National University, Seoul 151-744, Korea

TEL:82-2-880-9216, e-mail: [jjkim@snu.ac.kr](mailto:jjkim@snu.ac.kr)

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### Abstract

*We demonstrate that the reduction of quantum efficiency with increasing current density in phosphorescent light emitting diodes (PhOLEDs) is related to the formation of excitons in hole transporting layer based on the analysis of emission spectra and exciton formation zone. By employing dual emitting layer we could achieve maintaining quantum efficiency at high current density up to 10 000 cd/m<sup>2</sup> as 13.1% compared to the devices with single emitting layer (S-EML) ( $\eta_{\text{ext}} = 6.9\%$  at 10 000 cd/m<sup>2</sup>).*

### 1. Introduction

Phosphorescent organic lighting emitting diodes (PhOLEDs) have received considerable attention due to their ability of highly efficient emission compared with fluorescent OLEDs.<sup>1-4</sup> Through harvesting both singlet and triplet excitons, the external quantum efficiency ( $\eta_{\text{ext}}$ ) of PhOLEDs has reached above 20 % by using the optimized material systems,<sup>5-7</sup> p-i-n structures<sup>8-9</sup> and microcavity structures.<sup>10-11</sup> However, the efficiency roll-off (the decrease of efficiency with increasing current density) occurs at much lower luminance than required in displays or solid-state lighting. The roll-off of the quantum efficiency is one of the most significant problems facing electrophosphorescent devices and its origin was attributed to the triplet-triplet annihilation coming from long lifetime of triplet excitons,<sup>12</sup> electric field induced dissociation of excitons,<sup>13</sup> and triplet-polaron annihilation.<sup>12</sup>

In this letter, we report that the roll-off of the quantum efficiency with increasing current density is related to the exciton formation in the hole transporting layer

(HTL). Analysis of the steady state emission spectra indicated that the significant portion of the efficiency reduction is originated from the more and more exciton formation in the HTL with increasing current density. Based on the results, we fabricated devices with double emitting layers (D-EMLs) in order to confine the exciton formation in the emitting layers. The external quantum efficiency ( $\eta_{\text{ext}}$ ) was maintained constant up to 10,000 cd/m<sup>2</sup> by adopting the D-EMLs in Ir(ppy)<sub>3</sub> based PhOLEDs, resulting in high external quantum efficiency at high luminescence compared to the devices with single emitting layer (S-EML). The OLEDs with D-EMLs show significantly lower roll-off of efficiency ( $\eta_{\text{ext}} = 13.1\%$  at 10 000 cd/m<sup>2</sup>) than conventional S-EML OLEDs ( $\eta_{\text{ext}} = 6.9\%$  at 10 000 cd/m<sup>2</sup>).

### 2. Experimental

The OLEDs were fabricated by thermal evaporation onto a cleaned glass substrate precoated with indium tin oxide (ITO) without breaking the vacuum. Prior to organic layer deposition, ITO substrates were exposed to UV-ozone flux for 10 min following degreasing in acetone and isopropylalcohol. All organic layers were grown by thermal evaporation at the base pressure of  $<5 \times 10^{-8}$  Torr in the following order: hole transporting layer (HTL)/emitting layer (EML)/hole blocking layer (HBL)/electron transporting layer (ETL)/cathode. 40-nm-thick 4'-bis[N-(1-naphthyl)-N-phenyl-amino]biphenyl (NPB) was used as the HTL, 30-nm-thick N,N'-dicarbazoly1-4-4'-biphenyl (CBP) doped with 6 wt.% Ir(ppy)<sub>3</sub> as the S-EML, 10-nm thick 2,9-dimethyl-4,7-diphenyl-1,10-phenanthroline (BCP) as

the HBL and 40-nm-thick tris-(8-hydroxyquinoline) aluminum ( $\text{Alq}_3$ ) as the ETL, respectively. For the device with D-EMLs, 20-nm-thick 4,4',4''-tris(N-carbazolyl)-triphenylamine (TCTA) and 10-nm-thick CBP doped with 6 wt.%  $\text{Ir}(\text{ppy})_3$ , respectively, were used as the emitting layer. Finally, the cathode consisting of a 1-nm-thick LiF and a 100-nm-thick layer of Al were deposited onto the sample surface. Figure 1 shows the structure of the devices and the materials used in this study. Current density-voltage-luminescence (J-V-L) characteristics of the OLEDs were measured simultaneously using a Keithley 2400 programmable source meter and SpectraScan PR650 (Photo Research).

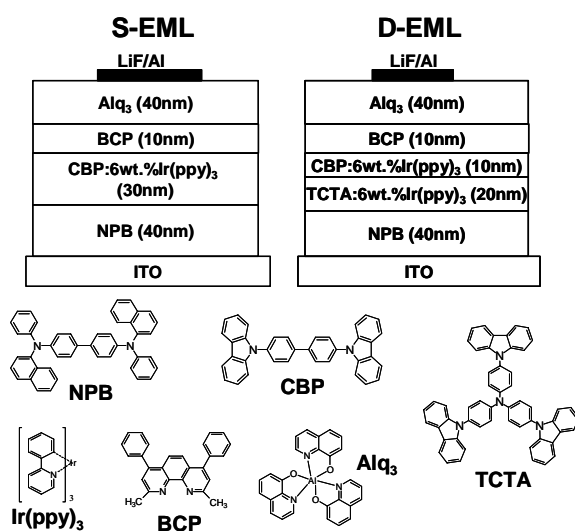


Figure 1 Structure of electronphosphorescent devices with single and double emitting layers and the chemical structure of the materials used in the devices.

### 3. Results and discussion

Current density-voltage-luminance (J-V-L) characteristics of the phosphorescent OLEDs are shown in Figure 2(a). The device with D-EMLs shows a lower driving voltage than S-EML. The driving voltage at  $1000 \text{ cd/m}^2$  was 5.4 and 7.3 V for D-EMLs and S-EML devices, respectively. Quantum efficiency and power efficiency of the devices are displayed in Figure 2(b). Initial quantum efficiency is almost the same between the S-EML and the D-EML devices. However, the roll-off of the quantum efficiency in the D-EML device is much smaller than the S-EML device. The quantum efficiency of the D-EML device was maintained almost the same with 13.5 % at  $10 \text{ cd/m}^2$ , 14.0 % at  $100 \text{ cd/m}^2$  and 13.1 % at  $10\,000 \text{ cd/m}^2$ . In contrast, the S-

EML device exhibited significant reduction of efficiency ( $\eta_{\text{ext}}=14.4 \%$  at  $10 \text{ cd/m}^2$ ,  $12.5 \%$  at  $100 \text{ cd/m}^2$  and  $6.9 \%$  at  $10\,000 \text{ cd/m}^2$ ).

Normalized electroluminescence (EL) spectra of the devices at different current densities are shown in Figure 3 in logarithmic scale in intensity. It is noteworthy that the emission around 440 nm grows with increasing current density in S-EML, and the emission begins to appear at the current density of  $0.1 \text{ mA/cm}^2$ , where the efficiency begins to be reduced.

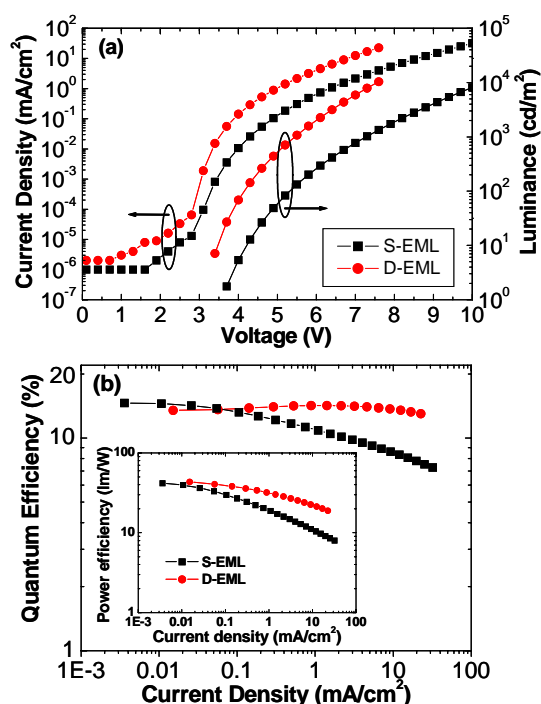
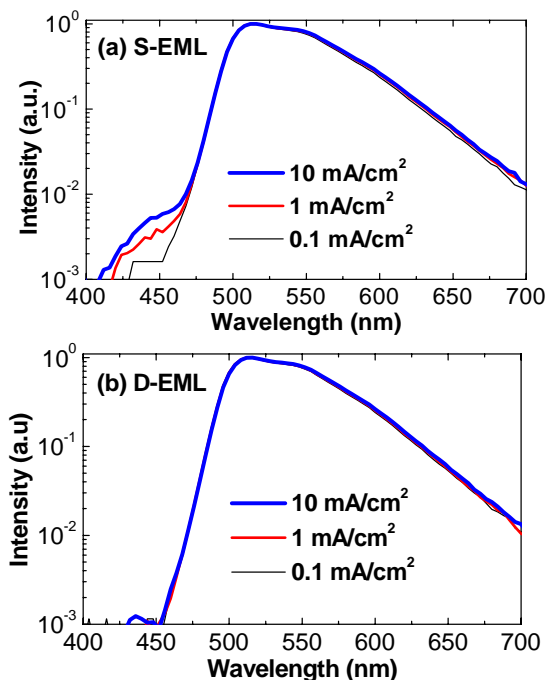


Figure 2 (a) Current density-voltage (J-V) and luminescence-voltage (L-V) characteristics, and (b) the external quantum efficiency and power efficiency of OLEDs for S-EML and D-EML.

Since the emission is consistent with the emission from NPB, the roll-off of the quantum efficiency with increasing current density in the S-EML device seems to be related to the exciton formation in the NPB, the HTL. On the contrary, however, the D-EMLs device shows no emission of NPB layer up to  $10 \text{ mA/cm}^2$ , resulting in the high quantum efficiency at high luminance, showing negligible efficiency reduction with increasing current density up to  $10\,000 \text{ cd/m}^2$ . Above the current density of  $10 \text{ mA/cm}^2$ , the emission from NPB layer is responsible for the lower quantum efficiency of  $\eta_{\text{ext}}=10.2 \%$  at  $50\,000 \text{ cd/m}^2$ . These results imply that recombination zone extends

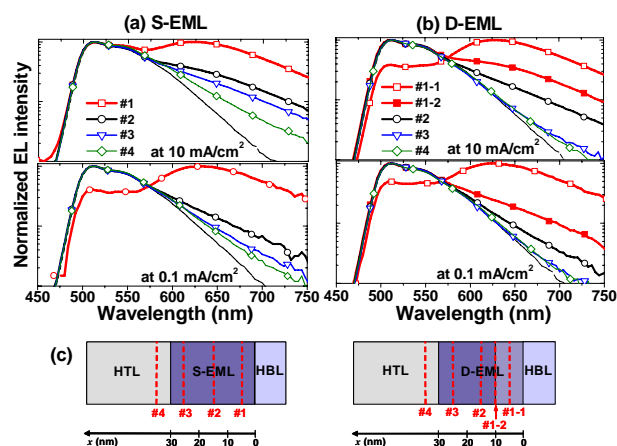
to HTL with increasing current and the degree of the extension is much larger in the S-EML device than D-EML. This can be understood based on the study of the location of the emission zones with increasing current density. To probe the exciton distribution in the devices, a 1-nm-thick 4-(dicyanomethylene)-2-(*t*-butyl)6-methyl-4H-pyran (DCJTB) was inserted as the sensing layer<sup>14</sup> at the various positions in the EML and HTL as shown in Figure 4(c). DCJTB was chosen as the sensing layer because it emits a red color around 630 nm which is easily distinguishable from the CBP:Ir(ppy)<sub>3</sub> emission. Even though the insertion of the layer perturbs the charge transport in the devices, still we expect that the resulting emission spectra provide the qualitative picture of the variation of the exciton distribution with increasing current density.



**Figure 3** Normalized electroluminescent spectra of (a) S-EML and (b) D-EML as a function of current density.

Figure 4 shows a series of EL emission spectra at two different current densities at each position of DCJTB layer for the S-EML and D-EMLs based devices. At low current density ( $0.1\text{mA}/\text{cm}^2$ ), the red emission close to HBL (position of #1) was dominant and there was a little and negligible emission at the position #3 and #4 for the S-EML and D-EML devices, respectively, indicating that the recombination zone is

close to HBL for both devices. At high current density ( $10\text{ mA}/\text{cm}^2$ ), however, the S-EML device showed significant red emission close to and inside HTL (position of #3 and #4), demonstrating that the recombination zone penetrates the HTL. The portion of exciton formation near the hole blocking BCP layer is reduced at the high current density as indicated by the relative emission intensity ratio of Ir(ppy)<sub>3</sub> and DCJTB at the position 1. The D-EMLs based device, in contrast, shows most of the emission in the CBP:Ir(ppy)<sub>3</sub> layer and increasing emission in the TCTA:Ir(ppy)<sub>3</sub> near the interface between the TCTA and CBP layers at the current density of  $10\text{ mA}/\text{cm}^2$ . Negligible emission at the position #3 and #4 was observed even at the high current density. These results indicate the confinement of excitons inside of CBP:Ir(ppy)<sub>3</sub> layer without the exciton formation close to HTL. More excitons are formed in TCTA:Ir(ppy)<sub>3</sub> layer with increasing current, but little excitons are formed near or inside of the HTL. Analysis of the emission spectra confirms that the significant portion of the efficiency reduction of S-EML based device is originated from the more and more exciton formation in the HTL with increasing current density.



**Figure 4** Normalized electroluminescent spectra of (a) S-EML and (b) D-EML having 1-nm-thick DCJTB exciton sensing layer at various positions (c) in the devices. The spectra were taken at different current densities. The solid line is the emission spectrum without DCJTB layer.

The exciton confinement inside the CBP:Ir(ppy)<sub>3</sub> layer in the D-EMLs devices<sup>8,11,15</sup> can be understood based on the charge mobility in the consisting layers. From the time-of-flight (TOF) measurement, CBP has

high electron and hole mobility of  $3 \times 10^{-4}$  and  $2 \times 10^{-3}$   $\text{cm}^2\text{V}^{-1}\text{s}^{-1}$  at the applied field of 0.5 MV/cm, respectively. However, TCTA has the hole mobility of  $3 \times 10^{-4}$   $\text{cm}^2\text{V}^{-1}\text{s}^{-1}$  at the applied field of 0.5 MV/cm and the electron mobility not measurable by the TOF method ( $<10^{-8}$   $\text{cm}^2\text{V}^{-1}\text{s}^{-1}$ .) By doping Ir(ppy)<sub>3</sub> in TCTA, transporting of electrons injected from the CBP:Ir(ppy)<sub>3</sub> layer through TCTA is effectively blocked in the layer. Electrons can only be transported by hopping between Ir(ppy)<sub>3</sub> dopant molecules. Therefore, the excitons can be mostly confined inside the CBP:Ir(ppy)<sub>3</sub> layer to prevent them from formation in HTL.

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#### 4. Summary

In summary, we have demonstrated that the reduction of EL quantum efficiency with increasing current density in single EML CBP:Ir(ppy)<sub>3</sub> devices is related to the formation of excitons in the NPB HTL based on the analysis of EL spectra and exciton formation zone. By adopting D-EML structure, we achieved very low roll-off of the quantum efficiency with increasing current density in PhOLEDs, in which hole transporting TCTA is used as one of the hosts. The external quantum efficiency was maintained constant up to 10 000  $\text{cd}/\text{m}^2$  by adopting the D-EMLs in Ir(ppy)<sub>3</sub> based PhOLEDs, resulting in high external quantum efficiency at high luminescence ( $\eta_{\text{ext}}=13.1$  % at 10 000  $\text{cd}/\text{m}^2$ ) compared to the devices with S-EML ( $\eta_{\text{ext}}=6.9$  % at 10 000  $\text{cd}/\text{m}^2$ ). These results clearly indicate that phosphorescent OLEDs with less roll off in efficiency with increasing current can be realized by confining recombination zone inside the emitting layer and minimizing the exciton formation in hole transporting or hole blocking layers. This is a promising approach for display and solid-state lighting application.

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