

## Spin-glass behavior in (A,B)-site deficient manganese perovskites

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In the past years, a giant magnetoresistance (GMR) effect found in perovskite-like structured materials has attracted considerable attention among scientists and manufacturers, since, a practical point of view, the capacity of producing magnetic and sensing sensors [1]. In a stream of this interest, further efforts to understand the underlying mechanism that leads to the GMR effect relative to the correlation between transport and magnetic properties, have been extensively devoted [2-4]. In these cases, spin-glass-like behaviors are ascribed to the frustration of random competing exchange interactions, namely the ferromagnetic double-exchange interaction between  $\text{Co}^{3+}$  (or  $\text{Mn}^{3+}$ ) and  $\text{Co}^{4+}$  (or  $\text{Mn}^{4+}$ ) and the antiferromagnetic one like spins. Noticeably, the distinction of spin-glass region from cluster-glass one, involved in the remarkable changes in transport and magnetic properties at a critical value of doping concentration, was observed [2,3]. Magnetic anomalies in zero-field-cooled (ZFC) magnetization as well as ac magnetic susceptibility below Curie temperature  $T_C$  and the charge/orbital fluctuation were also realized [4]. In this work, we present a study of magnetic properties of a deficient manganese perovskites system of  $\text{La}_{0.6}\text{Sr}_x\text{MnTi}_y\text{O}_3$ , and particularly provide its new magnetic phase diagram.

$\text{La}_{0.6}\text{Sr}_x\text{MnTi}_y\text{O}_3$  ( $x=0.15, 0.2, 0.25, 0.3, 0.35, 0.4$ , and  $y=0.4-x$ , respectively) polycrystalline samples were prepared by a conventional solid-state reaction method. A detailed procedure of sample preparation will be described elsewhere. X-ray diffraction data and scanning electron microscope confirmed the quality of the samples. The magnetic measurements were performed using a Quantum Design MPMS-5 SQUID magnetometer or a PPMS-7 magnetometer.

In a primary material of  $\text{LaMnO}_3$  with a cubic structure antiferromagnetic and insulator characteristics are present. Substitution of divalent elements such as  $\text{Ca}^{2+}$ ,  $\text{Sr}^{2+}$ ,  $\text{Ba}^{2+}$ , etc. for  $\text{La}^{3+}$  in  $\text{LaMnO}_3$  introduces a paramagnetic to ferromagnetic transition accompanying an insulator to metal one [3,4]. The correlation between metallic conduction and ferromagnetism results from the Zener's double-exchange interaction between  $\text{Mn}^{3+}$  (electronic configuration  $t_{2g}^3 e_g^1$ ,  $S = 2$ ) and  $\text{Mn}^{4+}$  ( $t_{2g}^3 e_g^0$ ,  $S = 3/2$ ) via oxygen 2p orbital. The nature of magnetic ordering in the compositional range strongly depends on the relative concentrations of  $\text{Mn}^{3+}$  and  $\text{Mn}^{4+}$  ions as well as the  $\langle \text{Mn-O-Mn} \rangle$  angle and  $\langle \text{Mn-O} \rangle$  bond distance. A complex magnetic phase might arise due to the competing ferromagnetic  $\text{Mn}^{3+}(e_g)-\text{O}(2p_\sigma)-\text{Mn}^{4+}(e_g)$  double-exchange and antiferromagnetic  $\text{Mn}^{3+}(t_{2g})-\text{O}(2p_\pi)-\text{Mn}^{4+}(t_{2g})$  super-exchange. To the present work, the deficiency at both A-site (=La,Sr) and B-site (=Mn,Ti) led to a decrease in  $T_C$  as the titanium and strontium concentration increases and decreases, respectively (see Fig. 1). The drop in  $T_C$  is attributed to the disorder induced by the random substitution of  $\text{Ti}^{4+}$  in the manganese sublattice, expected to result in

regions of different  $T_c$ . Besides, the decrease of Sr-deficient contents here also causes an additional drop in  $T_c$  arising from the decrease of A-site average ion radius. In addition, the magnetization measured at  $T=5\text{K}$  and at a maximum field of 5T decreases as the Ti-deficient content,  $y$ , is increased between 0 and 0.25 (and, the Sr-deficient content,  $x$ , is decreased between 0.4 and 0.15, respectively). The magnetization is far from saturation at a magnetic field of 5T, but a considerable value of spontaneous magnetization is still observed. The decrease of magnetization here should be related to the dilution of the Mn sublattice, as the  $\text{Mn}^{4+}$  content does not change much in comparison with  $y=0.05$ . This dilution is assumed to result in weakening of the effective exchange interaction between the Mn magnetic moments with a reduced number of magnetic next neighbors. More interestingly, a prominent divergence between zero-field-cooled (ZFC) and field-cooled (FC) magnetization measured at  $H=50$  Oe below  $T_c$  for all compositions was observed. An appearance of a sharp peak of ZFC magnetization curves below  $T_c$  for  $y \geq 0.1$  prefer to a spin-glass-like behavior, rather than the case of the samples with  $y=0$  and 0.5. This cusp is the so-called spin-glass freezing temperature,  $T_g$  [2,3]. It is evidently confirmed by ac magnetic susceptibility measurements where a peak of this curve is coincident with that of the ZFC curve. In Fig. 1, a magnetic phase diagram of  $\text{La}_{0.6}\text{Sr}_x\text{MnTi}_y\text{O}_3$  where SG, FM, and PM denote “spin-glass”, “ferromagnetic”, and “paramagnetic”, respectively, is depicted.

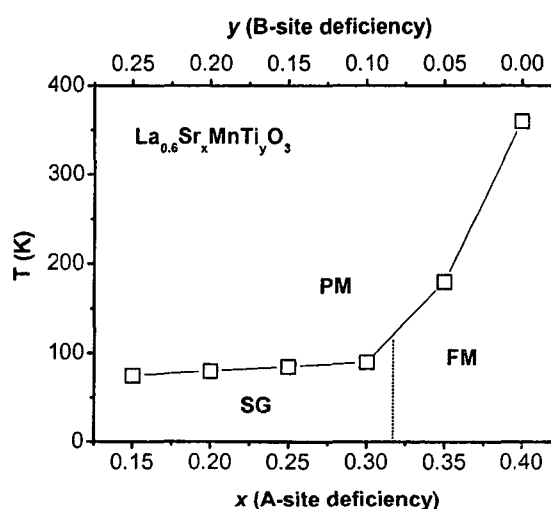


Fig. 1. The magnetic phase diagram of  $\text{La}_{0.6}\text{Sr}_x\text{MnTi}_y\text{O}_3$  as a function of Sr-deficient content (or Ti-deficient content).

**In summary**, the magnetic properties of deficient manganese perovskites of  $\text{La}_{0.6}\text{Sr}_x\text{MnTi}_y\text{O}_3$  have been studied by magnetization measurements in low magnetic fields. It is very instructive that there exists a spin-glass region for  $y \geq 0.1$  and  $x \leq 0.3$  but disappears at  $y=0$  and 0.05 (or  $x=0.4$  and 0.35, respectively). Nevertheless, the clarification in the correlation between transport and magnetic properties needs further study, and will be published elsewhere.

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